Fermi Liquid Theory

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1 Introduction

Fermi Gas is an easy and interesting problem in statistical mechanics. In metals we can treat the electrons as Fermi Gas at $T = 0$K because fermi energy of common metal is much higher than room temperature. But this model neglected the Coulomb interaction between the electrons. Now we have to establish a theory to describe the interacting electron system by path integral.

2 Path Integral of Fermi Gas and Fermi Liquid

2.1 Fermi Gas

Fermion system is described by a Grassmannian Action. For free Fermions with dispersion $\epsilon = k^2/2m$, the imaginary-time Action is:

$$S[\bar{\psi}, \psi] = \int_0^\beta d\tau \int d^3x \, \bar{\psi}_\sigma(\mathbf{x}, \tau) \left(\partial_\tau - \frac{\nabla^2}{2m} - \mu \right) \psi_\sigma(\mathbf{x}, \tau) \tag{1}$$

Change the Action into the momentum-Mastubara (imaginary-time frequencies) representation, the result is:

$$S[\bar{\psi}, \psi] = \sum_p \bar{\psi}_{p\sigma} \left(-i\omega_n + \frac{\mathbf{p}^2}{2m} - \mu \right) \psi_{p\sigma} \tag{2}$$

By this action, we can get the Grand Canonical partition function is:

$$Z = \int D\bar{\psi} D\psi \, e^{-S[\bar{\psi}, \psi]} \tag{3}$$

Clearly this is a Gaussian Integral, we can obtain the result easily, and it is:

$$\ln Z = \ln \prod_{\mathbf{q}\omega_n} \beta^{-1} \left(-i\omega_n + \frac{\mathbf{p}^2}{2m} - \mu \right) = \sum_{\mathbf{p}, \omega_n} \ln \left[\beta \left(-i\omega_n + \frac{\mathbf{p}^2}{2m} - \mu \right) \right]$$

From the theory of thermodynamics, we can conclude that the particle number is the derivative of the free energy, so by this path integral method we can get the occupying number function for Fermions:

$$N = -\frac{\partial F}{\partial \mu} = -\sum_p \frac{\beta^{-1}}{-i\omega_n + \frac{\mathbf{p}^2}{2m} - \mu} = \sum_{\mathbf{p}} n_{\mathrm{F}}(\epsilon_{\mathbf{p}}) \tag{4}$$

The second equality can be obtained by a contour integral and it is easy to do, but difficult to think. For details, Altland and Simons' Book 4.2 will be a good reference. I will use this result directly in this note.

2.2 Fermi Liquid

Now we need to add interactions into the Action. Assumes that there is Coulomb interaction between electrons, we can represent the Action in 4-momentum space as:

$$S[\bar{\psi}, \psi] = \sum_p \bar{\psi}_{p\sigma} \left(-i\omega_n + \frac{\mathbf{p}^2}{2m} - \mu \right) \psi_{p\sigma} + \frac{T}{2L^3} \sum_{pp'q} V(\mathbf{q}) \bar{\psi}_{k+q\sigma} \bar{\psi}_{k'-q\sigma'} \psi_{k'\sigma'} \psi_{k\sigma} \tag{5}$$

in which $V(\mathbf{q}) = 4\pi e^2 / \mathbf{q}^2$ ($\mathbf{q} \neq 0$) is the Coulomb interaction. Obviously the partition function is:

$$Z = \int D\bar{\psi} D\psi \exp \left\{ - \left[\sum_p \bar{\psi}_{p\sigma} \left(-i\omega_n + \frac{\mathbf{p}^2}{2m} - \mu \right) \psi_{p\sigma} + \frac{T}{2L^3} \sum_{pp'q} V(\mathbf{q}) \bar{\psi}_{k+q\sigma} \bar{\psi}_{k'-q\sigma'} \psi_{k'\sigma'} \psi_{k\sigma} \right] \right\} \tag{6}$$

This is what we We can use this path integral to calculate the ensemble average of lots of quantities. For example, the electron propagator is the ensemble average of the product $\langle \psi_{p\sigma} \bar{\psi}_{p\sigma} \rangle$.

3 The Diagram Approach of the Jellium Model

3.1 Hartree-Fock Approximation

Now if we define Z_0 as the partition function of the free electron system, the partition function of the Jellium model will be:

$$Z = Z_0 \langle e^{-S_{\text{int}}[\bar{\psi}, \psi]} \rangle_0 \tag{7}$$

Use this formula we can expand the free energy to the first order of $V(\mathbf{q})$, say:

$$F = F_0 + \frac{T^2}{2L^3} \sum_{pp'q} V(\mathbf{q}) \langle \bar{\psi}_{k+q\sigma} \bar{\psi}_{k'-q\sigma'} \psi_{k'\sigma'} \psi_{k\sigma} \rangle_0 \tag{8}$$

The second term can be represented by bubble diagrams with one interaction line: Now let's find out how to calculate the

Figure 1: Feynman Diagrams of Hartree Fock Approximation

two diagrams. It is obvious that the photon line in the first diagram has a momentum $V(\mathbf{q} = 0)$. But our electrons are moving in a Jellium with positive charges, so it has no contribution to the free energy. The second diagram is called "Fock" term. It has a Fermion loop so there will be a minus sign in the contribution. It is easy to write down the contribution by free electrons Green function

$$G_p = \frac{1}{-i\omega_n + \frac{\mathbf{p}^2}{2m} - \mu}$$

So the contribution of the second diagram is:

$$F^{(1)} = -\frac{T^2}{L^3} \sum_{pp'} G_p G_{p'} V(\mathbf{p} - \mathbf{p}') = -\frac{1}{L^3} \sum_{\mathbf{pp'}} n_F(\epsilon_{\mathbf{p}}) n_F(\epsilon_{\mathbf{p}'}) V(\mathbf{p}' - \mathbf{p}) \tag{9}$$

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This summation can be done exactly[1] at $T = 0$:

$$F^{(1)} = -\frac{1}{L^3} \sum_{\epsilon_{\mathbf{p}}, \epsilon_{\mathbf{p}'} < \mu} \frac{4\pi e^2}{|\mathbf{p} - \mathbf{p}'|^2} \simeq -L^3 \int_{\epsilon_{\mathbf{p}} < \mu} \frac{d^3 p}{(2\pi)^3} \int_{\epsilon_{\mathbf{p}'} < \mu} \frac{d^3 p'}{(2\pi)^3} \frac{4\pi e^2}{|\mathbf{p} - \mathbf{p}'|^2} = -\frac{e^2 L^3 p_F^4}{4\pi^3} \tag{10}$$

But this result is not physical: we can prove that the excitation mode of electron has a zero density of states, say $g(\mu) = 0$. If this is true, the transport properties will be different from what we see in experiments. So if the interaction is screened, this problem will be solved. We will see how to solve this problem by an infinite expansion.

3.2 Random Phase Approximation by Feynman Diagrams

Before we talk about the RPA approximation, I would like to introduce the **linked-cluster theorem**: partition function is the summation of all the **Feynman diagrams** with no out legs; the logarithm is the summation of all the **connected Feynman diagrams** with no out legs[2]. So the free energy is the summation of all the connected diagrams. RPA liquid, is the summation of all the diagrams as shown:

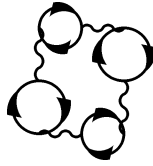
Figure 2: RPA diagram of order 4.

Now let's calculate the nth order contribution of RPA diagram to the free energy:

$$F^{(n)} = -T \frac{2^{n-1}(n-1)!}{n!} (-)^n (-)^n \left(\frac{T}{2L^3} \right)^n \sum_q \left(2V(\mathbf{q}) \sum_p G_p G_{p+q} \right)^n = -\frac{T}{2n} \sum_q \left(V(\mathbf{q}) \frac{2T}{L^3} \sum_p G_p G_{p+q} \right)^n \tag{11}$$

I will try to explain all of the coefficients in the first equality. $(-)^n n!^{-1}$ comes from the expansion of the $e^{-S_{\text{int}}}$; $2^{n-1}(n-1)!$ is the numbers of the ways you can get this diagram in the correlation function by Wick's theorem (each interacting line has two way to be put in the ring, and there are $(n-1)!$ ways to rearrange all the n interacting lines); another $(-)^n$ comes from the fact that there are n fermion loops in the diagram; the factor 2 in the summation comes from the spin of the electrons.

Now we can get the free energy of the RPA theroy:

$$F^{RPA} = \sum_n F^{(n)} = -\frac{T}{2} \sum_q \ln(1 - V(\mathbf{q})\Pi_q) \tag{12}$$

in which Π_q is defined as

$$\Pi_q = \frac{2T}{L^3} \sum_p G_p G_{p+q}$$

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[1]Hint: Use the spherical coordinate in momentum space to do the integral. It is tedious but not difficult.
[2]This theorem is proofed in Srednicki's book, p.74-76.

We will see that is the polarization operator which describes the vacuum polarization of photons and gives us a screened interaction. It's Feynman Diagram is:

$$\sim\!\!\sim\!\!\sim\!\!\sim\!\!\sim \quad = \quad \sim\!\!\sim\!\!\sim\!\!\sim \quad + \quad \sim\!\!\sim\!\!\bigcirc\!\!\sim\!\!\sim \tag{13}$$

That means:

$$V_{\text{eff}}(\mathbf{q}) = V(\mathbf{q}) + V(\mathbf{q})\Pi_q V_{\text{eff}}(\mathbf{q}) \quad \Rightarrow \quad V_{\text{eff}}(\mathbf{q}) = \frac{V_{\mathbf{q}}}{1 - V(\mathbf{q})\Pi_q}$$

At the low frequency limit, this will make the potential has the form of Yukawa interaction, say:

$$V(\mathbf{q}) = \frac{4\pi e^2}{\mathbf{q}^2 + \nu}$$

For more detailed discussion, Altland's book 5.2 is a good reference.