

**Stable cubic metal-semiconductor alloy clusters:  $X_4Y_4$  ( $X=\text{Cu,Ag,Au,Ti}$ ;  $Y=\text{C,Si}$ )**S. F. Li,<sup>1</sup> Xinlian Xue,<sup>1</sup> Yu Jia,<sup>1</sup> Gaofeng Zhao,<sup>2</sup> Mingfeng Zhang,<sup>1</sup> and X. G. Gong<sup>3</sup><sup>1</sup>*School of Physics and Engineering, Zhengzhou University, Zhengzhou-450052, People's Republic of China*<sup>2</sup>*Key Laboratory of Material Physics, Institute of Solid State Physics, Chinese Academy of Science, 230031-Hefei, People's Republic of China*<sup>3</sup>*Surface Physics Laboratory and Department of Physics, Fudan University, Shanghai-200433, People's Republic of China*

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*Ab initio* calculations have been performed on the electronic structures for some semiconductor-metal alloy clusters  $\text{Au}_n\text{Si}_m$  ( $n, m=1-4$ ), upon which very stable cubiclike clusters  $X_4Y_4$  ( $X=\text{Cu,Ag,Au,Ti}$ ;  $Y=\text{C,Si}$ ) have been found. The calculations show that relativistic effect as well as the ionic bonds may together play the key role in the cubic structures. The large highest occupied lowest unoccupied molecular orbital gaps (0.8–2.24 eV) make them attractive for optoelectronic applications and the perfect cubic structures make some of them desirable for cluster-assembled materials, such as one-dimensional nanotubes.

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**I. INTRODUCTION**

As silicon is the most important semiconducting element in microelectronics industry, Si clusters have been extensively investigated both experimentally<sup>1</sup> and theoretically<sup>2,3</sup> for their potential applications as building blocks to build up well-controlled nanostructures. Since the discovery of the very stable cage-like  $\text{C}_{60}$  fullerenes,<sup>4</sup> searching for free standing carbon and carbon-free clusters such as  $\text{C}_n$  and  $\text{Si}_n$  with cage-like structures has attracted much attention during the last decade. Unfortunately, the cage-like structures for pure Si clusters have never been found due to their unfavorable  $sp^2$  hybridization. However, by doping transition metal atoms, hexagonal prism  $W@Si_{12}$ ,<sup>5</sup> fullerene-like  $M@Si_{16}$ , cubiclike  $M@Si_{14}$  and Frank-Kasper tetrahedral  $Ti@Si_{16}$  and adodecahedral  $Zr@Si_{20}$  clusters<sup>6</sup> have been found. All these metal-encapsulated silicon clusters are stabilized by the interaction of the center doped metal atoms with the silicon cages. However, at present, the  $M@Si_n$  assembled nanomaterials are rarely constructed.<sup>7</sup>

Gold is a well known typical metal and shows novel properties in atomic scale. Though the  $d$  electrons may take part in the bonding and dominate the physical and chemical properties for  $\text{Au}_n$  clusters, the jellium model still successfully predicated that  $\text{Au}_n$  ( $n=2, 8, 18, 20, 34$ ) may be magic clusters, which has been confirmed experimentally by mass spectrometry analysis.<sup>8,9</sup> Studies show that due to relativistic effect,<sup>10</sup> which enhanced stronger  $sd$  hybridization and  $dd$  interaction in gold, up to 13 atoms, small anionic  $\text{Au}_n$  clusters favor planar structures.<sup>11</sup> And recent research predicates that the relativistic effect continues to play a role in the structures and bonding of large Au clusters. The configuration of  $\text{Au}_{20}$  (Ref. 12) was found to be a tetrahedral structure with all of the atoms located on the surfaces, which can be considered four triangle  $\text{Au}_{10}$  planes merged together making a quasi-two-dimensional structure. The latest studies show that  $\text{Au}_n$  ( $n=32, 33, 34, 35$ ) clusters can unexpectedly have cage-like structures.<sup>13</sup> The above results proposed a great challenge for the studies of silicon structures. Can we stabilize one class of silicon cage, such as cube, by appealing to gold?

As cubic structures are easily found in human life we expect that cubic structures may exist and even be universal in the microscope world. As the smallest special cage, cubic units may be even more suitable to build up cluster-assembled materials. In this paper, we report that  $\text{Si}_4$  and  $\text{C}_4$  clusters indeed can be stabilized to almost perfect cubic structure by alloying with four Au atoms. Other stable cubic metal-semiconductor alloy clusters  $X_4Y_4$  ( $X=\text{Cu,Ag,Au,Ti}$ ;  $Y=\text{C,Si}$ ) with large highest occupied molecular orbital-lowest unoccupied molecular orbital (HOMO-LUMO) gaps are also studied.

**II. CALCULATION DETAILS**

The calculations are based on the density functional theory (DFT) with spin-polarized generalized gradient approximation (GGA)<sup>14</sup> implemented in the VASP code.<sup>15</sup> The interactions of the valence electrons with the ionic cores are described by projected augmented wave (PAW)<sup>16</sup> method, with a scalar relativistic effect included. The wave functions are expanded in plane waves with an energy cutoff of  $\sim 400$  eV. We use a simple cubic cell of 15 Å edge length with a periodic boundary condition, and the  $\Gamma$  point approximation for Brillouin zone sampling. The ground states of clusters are obtained by comparing the energies of many random configurations followed with optimizations by conjugated gradient (CG)<sup>17</sup> and simulated annealing methods. In order to explicitly consider the relativistic effect for Au atoms, we have also used the ADF code<sup>18</sup> to do ZORA (zero order regular approximation) scalar relativistic calculations for solid support.

**III. RESULTS AND DISCUSSION**

First of all, we have studied the geometrical and electronic structures for  $\text{Au}_n\text{Si}_m$  ( $n, m=1-4$ ) clusters. The ground state structures, the binding energies defined with the form of  $E_b = -[E(X_nY_m) - nE(X) - mE(Y)]$ , the HOMO-LUMO gaps ( $E_{\text{gap}}$ ) and the bond lengths have been presented in Fig. 1. The calculated average Au-Si bond lengths in-

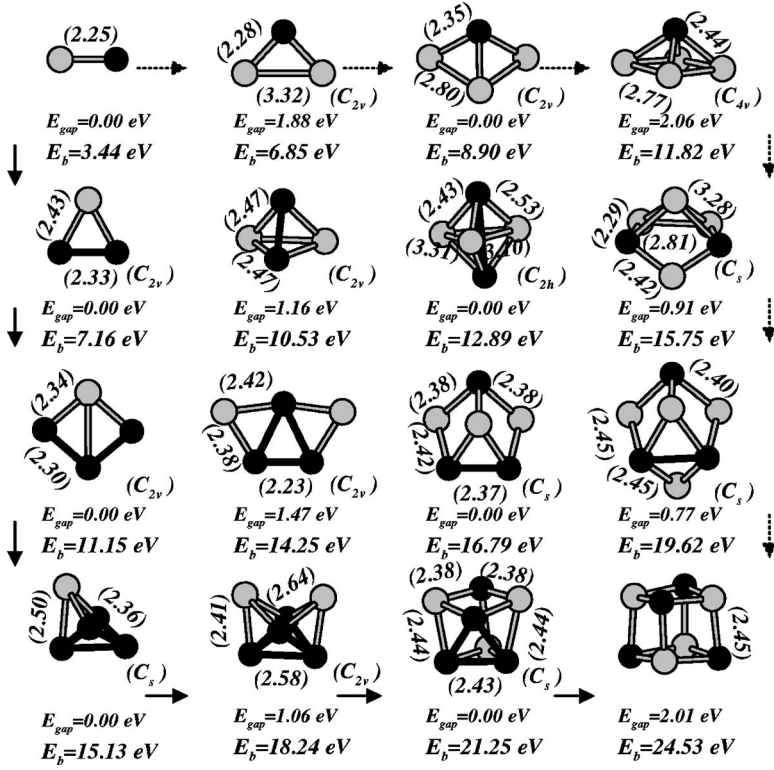


FIG. 1. The geometrical structures for  $\text{Au}_n\text{Si}_m$  ( $n, m=1-4$ ) clusters. The HOMO-LUMO gaps ( $E_{\text{gap}}$ ) and the main bond lengths are marked for all clusters. Two possible growth modes along the solid line and the dotted line are also shown, respectively. Black spheres: Si atoms, gray spheres: Au atoms.

crease from 2.25 Å for  $\text{AuSi}$  dimer to 2.45 Å for  $\text{Au}_4\text{Si}_4$  cluster. For every cluster comprising an odd number of gold atoms, we find  $1 \mu_B$  magnetic moment. For those clusters with an even number of gold atoms, large HOMO-LUMO gaps ranging from 0.77 to 2.06 eV have been obtained. These results indicate that the electronic structures of silicon clusters may be tunable by doping gold atoms with different stoichiometries. As shown in Fig. 1,  $\text{Au}_4\text{Si}$ ,  $\text{Au}_2\text{Si}_4$ , and  $\text{Au}_4\text{Si}_4$  clusters have 2.06, 1.06 and 2.01 eV HOMO-LUMO gaps, respectively, indicating their high stabilities, which has been confirmed, in the case of  $\text{Au}_4\text{Si}$ , by experimental work.<sup>19</sup> For the noble metal  $\text{Au}(\text{Ag}, \text{Cu})$ , electronic configuration  $s^1d^{10}$  for the PAW potential has been used in VASP calculation. If we regard that per gold (silicon) contributes one (four) valent electron(s),  $\text{Au}_4\text{Si}$ ,  $\text{Au}_2\text{Si}_4$ , and  $\text{Au}_4\text{Si}_4$  clusters have 8, 18, and 20 electrons, respectively, which happens to coincide with the magic numbers predicated by the jellium model for gold clusters. However, it seems that  $\text{Au}_2\text{Si}$ ,  $\text{Au}_2\text{Si}_2$ , and  $\text{Au}_2\text{Si}_3$  clusters with large HOMO-LUMO gaps carrying 6, 10, and 14 valence electrons, respectively, failed to be predicated as stable clusters by the jellium model.

It is interesting to find that the  $\text{Au}_4\text{Si}_4$  cluster prefers cubic structure with about 24.53 eV binding energy and a 2.01 eV HOMO-LUMO gap, indicating its high inertness. As presented in Fig. 1 and Table I, the bond length of Si-Au is 2.45 Å. The four Au atoms and the four Si atoms are just alternatively located on the eight apexes of one cube with little distortion. From Fig. 1, one can notice that there may be many growth paths from the reactive linear  $\text{AuSi}$  dimer to the high stable  $\text{Au}_4\text{Si}_4$  cluster, however, we assume that the two most possible growth paths may be along the dotted-line

path and the solid-line path, respectively. In fact, along the dotted-line path, the four Au atoms will first be grown up one by one and then the four Si atoms have been grown step by step on the  $\text{Au}_4$  core. On the contrary, for the solid-line path the four Si atoms will be first grown up. It is also worth mentioning that the geometrical structures for both  $\text{Si}_3\text{Au}_4$  and  $\text{Au}_3\text{Si}_4$  clusters are just like one cubic removed one apex. If one  $\text{Si}(\text{Au})$  atom is added to the removed apex of  $\text{Si}_3\text{Au}_4(\text{Au}_3\text{Si}_4)$  followed with one simple optimization, then

TABLE I. Calculated properties of  $X_4Y_4$  ( $X=\text{Cu}, \text{Ag}, \text{Au}, \text{Ti}; Y=\text{C}, \text{Si}$ ) clusters. In the brackets, C and  $T_d$  represent cubic and tetrahedral configurations, respectively.  $\bar{R}$ (Å) includes average bond lengths of  $X-X$ ,  $X-Y$  and  $Y-Y$ , respectively;  $E_{\text{gap}}$ (eV) is a HOMO-LUMO gap;  $E_b$ (eV) is the binding energy defined with  $E_b = -[E(X_4Y_4) - 4E(X) - 4E(Y)]$ ;  $X-C$  indicates the average charge flow from per  $X$  atom to per  $Y$  atom in  $X_4Y_4$  clusters, obtained from Voronoi deformation density (VDD)<sup>20</sup> analysis.

Clusters	$\bar{R}$ (Å)			$E_b$ (eV)	$E_{\text{gap}}$ (eV)	$X-C(e)$
	$X-X$	$X-Y$	$Y-Y$			
$\text{Cu}_4\text{C}_4$ (C)	2.69	1.91	2.71	26.74	2.10	0.35
$\text{Ag}_4\text{C}_4$ (C)	3.03	2.14	3.03	19.07	1.67	0.34
$\text{Au}_4\text{C}_4$ (C)	2.96	2.15	3.10	22.81	1.83	0.30
$\text{Cu}_4\text{Si}_4$ ( $T_d$ )	3.62	2.30	2.77	24.31	1.93	0.12
$\text{Ag}_4\text{Si}_4$ ( $T_d$ )	4.18	2.55	2.70	20.27	1.42	0.17
$\text{Au}_4\text{Si}_4$ (C)	3.30	2.45	3.61	24.53	2.01	0.07
$\text{Ti}_4\text{C}_4$ (C)	2.77	1.97	2.78	45.73	2.24	0.55
$\text{Ti}_4\text{Si}_4$ (C)	3.02	2.43	3.77	30.54	0.84	0.33

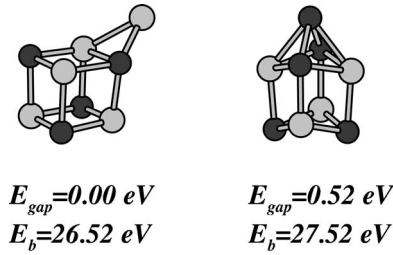


FIG. 2. The geometrical structures for  $\text{Au}_5\text{Si}_4$  and  $\text{Au}_4\text{Si}_5$  clusters. Black spheres: Si atoms, gray spheres: Au atoms.

the cubic  $\text{Si}_4\text{Au}_4$  cluster has been obtained. Recent study results declare that both  $\text{Si}_4$  and  $\text{Au}_4$  show much higher stability relative to their neighbors,<sup>21</sup> which supports our idea that  $\text{Si}_4$  and  $\text{Au}_4$  cores could be first grown and there may be competition between these two different growth paths. As presented in Fig. 2, the geometrical structure of  $\text{Au}_5\text{Si}_4$  ( $\text{Au}_4\text{Si}_5$ ) can be obtained by growing one Au(Si) atom on the edge (face) of the cubic  $\text{Au}_4\text{Si}_4$ , with a small distortion of the cube, which again predicates the stability of  $\text{Au}_4\text{Si}_4$  cluster.

In Fig. 3, we present the dissociation energy for one Au(Si) atom from  $\text{Au}_n\text{Si}_4$  ( $\text{Au}_4\text{Si}_m$ ) clusters ( $n, m=1-5$ ). Much higher stability for  $\text{Au}_4\text{Si}_4$  relative to its neighbors can be seen clearly. We have also further studied the interaction between two cubic  $\text{Au}_4\text{Si}_4$  clusters. The reacted product keeps these two cubic units almost intact with about 2.56 Å distance and about 1.66 eV reaction energy  $E_R$  [for two  $\text{Au}_4\text{Si}_4$  clusters,  $E_R = -[E((\text{Au}_4\text{Si}_4)_2) - 2E(\text{Au}_4\text{Si}_4)]$ ], as shown in the inserted panel in Fig. 3. The HOMO-LUMO gap still reaches about 1.0 eV, which suggests that stable self-organized assemblies of the cubic  $\text{Au}_4\text{Si}_4$  unit may be very possible.

The relativistic effect may indeed play an important role in the structure for the  $X_4\text{Si}_4$  cluster. As has been mentioned, Au clusters favor planar or cage-like structures due to the relativistic effect, then one may ask a question—does the relativistic effect still work in the binary alloy cubic  $\text{Au}_4\text{Si}_4$  cluster? In order to make this question clear, we have also studied the electronic structures for  $\text{Cu}_4\text{Si}_4$  and  $\text{Ag}_4\text{Si}_4$  clusters. The calculated results show that both  $\text{Cu}_4\text{Si}_4$  and  $\text{Ag}_4\text{Si}_4$  clusters prefer cubicle-like structures with relatively large distortion. In fact, the most stable structures for both  $\text{Cu}_4\text{Si}_4$  and  $\text{Ag}_4\text{Si}_4$  are inlaid tetrahedron-like, which can also be regarded as four metal atoms which adsorb on the four surfaces of the inner  $\text{Si}_4$  tetrahedron, constituting the outer tetrahedron. And the ratio of Cu-Cu (Ag-Ag) distance to that of the Si-Si is

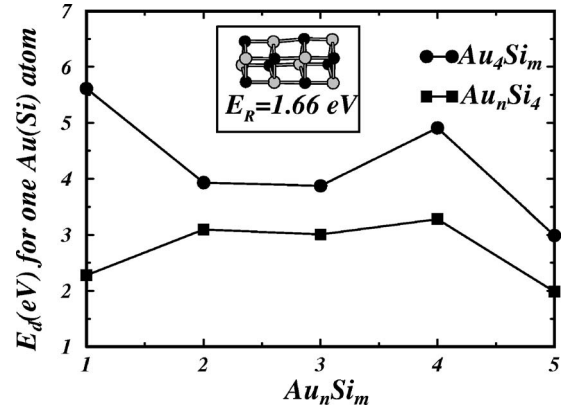


FIG. 3. The dissociation energy ( $E_d$ ),  $E_d(n) = \{E[\text{Au}(\text{Si})_n\text{Si}(\text{Au})_4] - E[\text{Au}(\text{Si})] - E[\text{Au}(\text{Si})_{(n-1)}\text{Si}(\text{Au})_4]\}$  for one Au(Si) atom from  $\text{Au}_n\text{Si}_4$  ( $\text{Au}_4\text{Si}_m$ ) cluster ( $n, m=1-5$ ). Much higher stability for  $\text{Au}_4\text{Si}_4$  relative to its neighbors can be seen clearly. In the inserted figure, the configuration for two reacted  $\text{Au}_4\text{Si}_4$  clusters is presented, which also indicates the extraordinary stability of the  $\text{Au}_4\text{Si}_4$  cluster.

about 1.31 Å (1.55 Å), which can be seen from Table I and Fig. 4. These results indicate that the relativistic effect does work in the cubic  $\text{Au}_4\text{Si}_4$  cluster. We have also made the comparison between relativistic and nonrelativistic calculations performed by both VASP and ADF codes for  $\text{Au}_4\text{Si}_4$ . The calculations performed by both these two codes show that, in a nonrelativistic level, the cubic structure is not stable for  $\text{Au}_4\text{Si}_4$ , which will distort to the inlaid-tetrahedron-like structure, however, in a relativistic calculation level, the inlaid-tetrahedron-like configuration will relax again to cubic structure. Our calculations also show that in the cubic structure (owing to the stabilization of  $s$  electrons of Au), the charge flow from per Au atom to per Si atom is only about 0.07  $e$  (see Table I). While in the inlaid-tetrahedron-like structure obtained with nonrelativistic calculations, the charge flow reaches to about 0.2  $e$ , however, the Au-Si bonding lengths have been lengthened to about 2.68 Å, making the total binding energy decrease. The relativistic effect may be very weak in  $\text{Cu}_4\text{Si}_4$  and  $\text{Ag}_4\text{Si}_4$  clusters, and these two clusters prefer inlaid-tetrahedron-like structures. These results indicate that the relativistic effect may really play a key role in the cubic structure, though we also guess that the size effect for Cu, Ag, and Au atoms may be important in the structural differences.

As the same important semiconducting element silicon, carbon has also attracted much attention for applications in

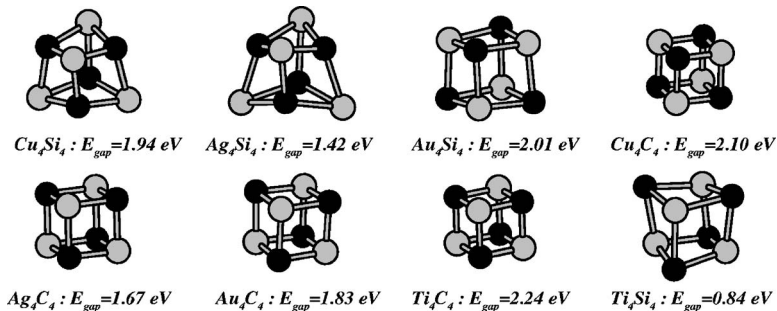


FIG. 4. The geometrical configurations for  $X_4Y_4$  ( $X=\text{Cu}, \text{Ag}, \text{Au}, \text{Ti}$ ;  $Y=\text{C}, \text{Si}$ ) clusters. The HOMO-LUMO gaps are also marked for each cluster. The black spheres and the gray spheres correspond to the semiconductor atoms and the metal atoms, respectively.



nanoscience. Can a  $C_4$  cluster also be stabilized to cubic structures by alloying with noble metal atoms? We are surprised to find that the most stable structures for all the  $X_4C_4$  ( $X=Cu, Ag, Au$ ) clusters are almost perfect cubic configurations (see Table I and Fig. 4). The bond length of  $Cu-C$  in  $Cu_4C_4$  cluster is 1.91 Å, which is about 10% smaller than the  $Ag-C$  and  $Au-C$  bond lengths (2.14 Å, 2.15 Å) in  $Ag_4C_4$  and  $Au_4C_4$  clusters, respectively. There are 2.10, 1.67, and 1.83 eV HOMO-LUMO gaps for  $X_4C_4$  ( $X=Cu, Ag, Au$ ) clusters respectively, indicating their high inertness and stability. As discussed above, due to the relativistic effect, the  $Au_4Si_4$  cluster can be stabilized to cubic structure. So it is reasonable for  $Au_4C_4$  to favor cubic configuration. However, it seems difficult to understand the almost perfect cubic structures for  $X_4C_4$  ( $X=Cu, Ag$ ) clusters. In fact, the cubic structures for  $X_4C_4$  ( $X=Cu, Ag$ ) clusters may be understood from the point of view of  $X-C$  ionic bonds. As shown in Table I, the charge flow from  $X$  ( $X=Cu, Ag, Au$ ) to  $C$  atom is about 0.30–0.35  $e$ , which is obviously larger than that in  $X_4Si_4$  ( $X=Cu, Ag$ ) clusters with inlaid-tetrahedronlike structures. In the mean time, as all the  $X_4C_4$  ( $X=Cu, Ag, Au$ ) clusters are almost perfect cubic configurations, so we regard that the size effect does not work in a cubic-noncubic alternative in  $X_4Si_4$  ( $X=Cu, Ag, Au$ ) clusters. The results above encourage us to make one conclusion that both the relativistic effect and the ionic bonding may dominate the cubic structures.

In order to know whether the semiconductor  $Y_4$  ( $Y=C, Si$ ) clusters can be stabilized to cubic structures by alloying with other transitional metal elements, we have also studied the electronic structures of  $Ti_4Y_4$  ( $Y=C, Si$ ) clusters. (It is well known that  $Ti$  can easily alloy with carbon and silicon.) As presented in Fig. 4, the  $Ti_4C_4$  cluster also prefers perfect cubic structure with about 1.97 Å  $Ti-C$  bond length. The 2.24 eV (see Table I) large HOMO-LUMO gap and the 45.73 eV large binding energy indicate its extraordinary stability. In fact, intensive work has been done for  $Ti_xC_y$  clusters in both theoretical and experimental fields<sup>22</sup> following the original discovery of the metallocarbohedrenes (Met-Cars)  $Ti_8C_{12}$  by Castleman and co-workers.<sup>23</sup> Wang *et al.*<sup>22</sup> found that  $Ti_3C_8^-$ ,  $Ti_4C_8^-$ ,  $Ti_6C_{13}^-$ ,  $Ti_7C_{13}^-$ ,  $Ti_9C_{15}^-$ , and  $Ti_{13}C_{22}^-$  clusters show abundant (magic) negative ion peaks in mass spectra, and from  $Ti_4C_8$  these clusters prefer cubic-like growth pathways. It is interesting that the structure for the  $Ti_4C_8$  cluster can just be regarded as putting an additional four  $C$  atoms on the apexes of the cubic  $Ti_4C_4$  cluster forming four  $C-C$  dimers. The results above strongly imply that the  $Ti_4C_4$  cluster may be the smallest cubic unit in  $Ti_xC_y$  growths.  $Ti_4Si_4$  cluster is found to favor cubiclike structure with about 0.8 ratio of  $Ti-Ti$  distance to  $Si-Si$  distance. Despite the large binding energies,  $Ti_4Si_4$  has a relatively smaller HOMO-LUMO gap of 0.84 eV, which could be partially attributed to the distortions from perfect cubic structures. The electronic analysis shows that the average charge flow from per  $Ti$  atom to per  $Y$  ( $Y=C, Si$ ) atom is 0.55  $e$  and 0.33  $e$ , respectively, indicating again that the strength of the ionic bonds of  $Ti-Y$  ( $Y=C, Si$ ) may indeed play an important role in a cubic-noncubic alternative. Here, we simply state that in our calculations Mulliken analysis predicates obvi-

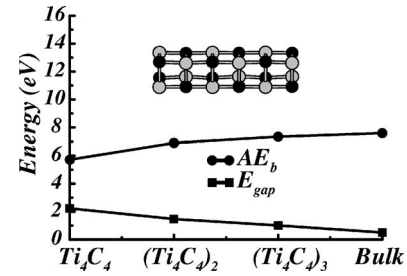


FIG. 5. The average binding energies per atom ( $AE_b$ ) and the HOMO-LUMO gaps ( $E_{gap}$ ) in eV for finite nanotubes of  $(Ti_4C_4)_n$  ( $n=1, 2, 3$ ) and for the infinite  $Ti_4C_4$  nanotube. The structure for the infinite nanotube has also been shown.

ously larger charge flow for  $Au_4Si_4$ ,  $Ti_4C_4$ , and  $Ti_4Si_4$  clusters than VDD does, for example, about one electron charge flow from per  $Ti$  atom to per  $C$  atom has been obtained with Mulliken analysis for  $Ti_4C_4$ , whereas for other clusters with the  $Au$  and  $Si$  components, both methods lead to similar results.

Recent studies show that both finite and infinite hexagonal silicon nanotubes can be stabilized by the interactions of transition metal atom with hexagonal  $M@Si_{12}$  cluster units,<sup>7</sup> and by using high-resolution transmission electron microscopy, a stable 3 Å in diameter armchair carbon nanotube can be grown inside a multiwalled carbon-carbon nanotube.<sup>24</sup> However, here we simply report that  $Ti_4C_4$  units can be self-organized to stable finite and infinite quadrate nanotubes with only about 1.99 Å diameter by the strong  $Ti-C$  ionic bonding. (For the infinite nanotube,  $7 \times 1 \times 1$   $k$  points have been used.) As presented in Fig. 5, in the nanotubes, the  $Ti_4C_4$  units still keep almost perfect cubic symmetry with about 5% bonding lengths increasing along the tube axis. The average binding energy ( $AE_b$ ) increases from about 5.7 eV for the  $Ti_4C_4$  cluster to about 7.3 eV for the infinite  $Ti_4C_4$  nanotube. For the infinite nanotube, the HOMO-LUMO gap has been decreased to about 0.5 eV, indicating its semiconducting character. Constant-temperature (600 K) molecular dynamic simulations performed by VASP show that the infinite (finite) nanotubes still keep almost perfect quadrate symmetry in a time duration of 1 ps with the time step of 1 fs, showing that the  $Ti_4C_4$ -organized nanotubes are indeed very stable. (The details for the cubic nanotubes will be published separately.)

#### IV. SUMMARY

In summary, we have shown that, due to the relativistic effect and the ionic bonding, semiconducting  $Y_4$  ( $Y=C, Si$ ) clusters can be stabilized to cubic structures by alloying with transition metal elements  $X$  ( $X=Cu, Ag, Au, Ti$ ). The metal-semiconductor alloy cubic clusters  $Au_4Si_4$ ,  $X_4C_4$  ( $X=Cu, Ag, Au$ ) and  $Ti_4C_4$  clusters have large binding energies and large HOMO-LUMO gaps, making them desirable for self-organized microelectronic devices and other nanoscale device applications. And we also hope such cubic clusters deserve further experimental investigations.

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