Cosmology Calculator

In[34]:= Clear["Global`*"];
pdConv[f_]:=TraditionalForm[f/.Derivative[inds__][g_][vars__]:→Apply[Defer[D[g[vars],##]]&

$$z \equiv 1/a - 1$$

In[36]:= (*scale factor in terms of the Redshift *)

scalfactor[z]:=
$$\frac{1}{(1+z)}$$
;
redshift[a]:= $\frac{1}{a}$ -1;

In[38]:=

 $(*Function to compute conformal time for a given redshift*)\\ tau[zs_,H0_,\Omega m_,\Omega d_,w_,\Omega r_]:=NIntegrate\Big[\frac{H0^{-1}}{\sqrt{\Omega d\left(1+z\right)^{3}(1+w)}+\Omega m_-(1+z)^{3}+\Omega r_-(1+z)^{4}}},\{z,zs,\infty\}\Big]$

In[39]:= (*Function to compute Physical time for a given redshift = $\frac{a'(tau)}{a}$ *)

Friedmann equation: (FROM https://en.wikipedia.org/wiki/Hubble%27s_law)

If the universe is both matter-dominated and dark energy-dominated, then the above equation for the Hubble parameter will also be a function of the equation of state of dark energy. So now:

$$ho =
ho_m(a) +
ho_{de}(a),$$

where ho_{de} is the mass density of the dark energy. By definition, an equation of state in cosmology is $P=w
ho c^2$, and if this is substituted into the fluid equation, which describes how the mass density of the universe evolves with time, then

$$\dot{\rho} + 3\frac{\dot{a}}{a}\left(\rho + \frac{P}{c^2}\right) = 0;$$

$$rac{d
ho}{
ho} = -3rac{da}{a}(1+w).$$

If w is constant, then

$$\ln \rho = -3(1+w) \ln a;$$

implying:

In[41]:=

$$a = a^{-3(1+w)}$$

Therefore, for dark energy with a constant equation of state w, $\rho_{dc}(a) = \rho_{de0}a^{-3(1+w)}$. If this is substituted into the Friedman equation in a similar way as before, but this time set k=0, which assumes a snatially flat universe, then (see shape of the universe)

ln[40]:= (*Physical Hubble Function at a given redshift = $\frac{\dot{a}}{a}$ *)

 $Physical Hubble [z_{-}, H0_{-}, \Omega m_{-}, \Omega d_{-}, w_{-}, \Omega r_{-}] := H0 \ \sqrt{\Omega d \left(1+z\right)^{3 \, (1+w)} \ + \Omega m \ \left(1+z\right)^{3} + \Omega r \left(1+z\right)^{4}}$

(*Conformal Hubble Function at a given redshift *)

 ${\sf Conformal Hubble} \big[{\sf z_, H0_, \Omega m_, \Omega d_, w_, \Omega r_} \big] := {\sf H0} \ \sqrt{\Omega d \, (1+z)^{\, (1+3 \ w)}} \ + \Omega m \ (1+z) + \Omega r \, (1+z)^{\, 2}$

(*Physical Hubble Function at a given redshift = $\frac{\dot{a}}{a}$ *)

 $Conformal Hubble Prime \\ \left[z_, H0_, \Omega m_, \Omega d_, w_, \Omega r_\right] := D \\ \left[Conformal Hubble \\ \left[z_, H0_, \Omega m_, \Omega d_, w_, \Omega r\right]_\right] \\ * \left(-(1+z)^{2} + (1+z)^{2} + (1+z)^$

ln[50]:= (*(*Physical Hubble Function at a given redshift = $\frac{\dot{a}}{a}$ *)

 $Conformal Hubble Prime [z_, H0_, \Omega m_, \Omega d_, w_, \Omega r_] := D [Conformal Hubble Prime [z_, H0_, \Omega m_, \Omega d_, w_, \Omega r]] := D [Conformal Hubble Prime [z_, H0_, \Omega m_, \Omega d_, w_, \Omega r]] := D [Conformal Hubble Prime [z_, H0_, \Omega m_, \Omega d_, w_, \Omega r]] := D [Conformal Hubble Prime [z_, H0_, \Omega m_, \Omega d_, w_, \Omega r]] := D [Conformal Hubble Prime [z_, H0_, \Omega m_, \Omega d_, w_, \Omega r]] := D [Conformal Hubble Prime [z_, H0_, \Omega m_, \Omega d_, w_, \Omega r]] := D [Conformal Hubble Prime [z_, H0_, \Omega m_, \Omega d_, w_, \Omega r]] := D [Conformal Hubble Prime [z_, H0_, \Omega m_, \Omega d_, w_, \Omega r]] := D [Conformal Hubble Prime [z_, H0_, \Omega m_, \Omega d_, w_, \Omega r]] := D [Conformal Hubble Prime [z_, H0_, \Omega m_, \Omega d_, w_, \Omega r]] := D [Conformal Hubble Prime [z_, H0_, \Omega m_, \Omega d_, w_, \Omega r]] := D [Conformal Hubble Prime [z_, H0_, \Omega m_, \Omega d_, w_, \Omega r]] := D [Conformal Hubble Prime [z_, H0_, \Omega m_, \Omega d_, w_, \Omega r]] := D [Conformal Hubble Prime [z_, H0_, \Omega m_, \Omega d_, w_, \Omega r]] := D [Conformal Hubble Prime [z_, H0_, \Omega m_, \Omega d_, w_, \Omega r]] := D [Conformal Hubble Prime [z_, H0_, \Omega m_, \Omega d_, w_, \Omega r]] := D [Conformal Hubble Prime [z_, H0_, \Omega m_, \Omega d_, W]] := D [Conformal Hubble Prime [z_, H0_, \Omega m_, \Omega$

The equation to find a"(\tau) which must be consistent with what we find previously. We don't need to do it, instead we can just

To find Ψ and Ψ' we define $N = \ln a$ as a new time coordinate, where we have,

$$\mathcal{H}^2 = H_0^2 \left[\Omega_m e^{-N} + \Omega_r e^{-2N} + \Omega_{kess} e^{-(1+3w)N} \right]$$
 (1.6)

$$\partial_{\tau}^{2}\phi + 3\frac{\dot{a}}{a}\partial_{\tau}\phi - \left(8\pi Ga^{2}\bar{p} + 2K\right)\phi = 4\pi Ga^{2}\delta p.$$

Here we neglect pressure from radiation, meaning that we don't go deep to the radiation domination and we start our equations from z = 1000 for example

In[51]:=

(*Here we neglect pressure from radiation, meaning that we don't go deep to the radiation domin

$$\begin{split} & \text{dNlnH}\big[N_-, \text{HO}_-, \Omega\text{m}_-, \Omega\text{d}_-, \text{w}_-, \Omega\text{r}_-\big] := \ D\Big[\ \text{Log}\Big[\text{ConformalHubble}\Big[\frac{1}{\text{E}^N} - 1, \text{HO}_+, \Omega\text{m}_+, \Omega\text{d}_+, \text{w}_+, \Omega\text{r}_-\big] := \ \psi''[N] \ + \left(3 + \ \text{dNlnH}\big[N_+, \text{HO}_+, \Omega\text{m}_+, \Omega\text{d}_+, \text{w}_+, \Omega\text{r}_-\big] \right) \psi'[N] + \left(2 - \frac{2}{2} \ \Omega\text{m} + \frac{2}{2} \ \Omega\text$$

$$(*s = NDSolve \Big[\Big\{ \Big(\Big(a'[tau] - a[tau]H0\sqrt{\Omega m} \ a[tau]^{-1} + (1-\Omega m)a[tau]^{-(1+3 \ w)} = 0 \Big) / \Omega m \rightarrow 3/10 / N \rightarrow -0.9 / (1+3 \ w) \Big] = 0 \Big]$$

RESULTS

The formula for \HH(z)

In[53]:=

ConformalHubble $[z, H0, \Omega m, \Omega d, w, \Omega r] //pdConv$

Out[53]//TraditionalForm=

H0
$$\sqrt{\Omega d(z+1)^{3w+1} + \Omega m(z+1) + \Omega r(z+1)^2}$$

The formula for d \HH/dtau (z)

In[55]:=

ConformalHubblePrime $[z,H0,\Omega m,\Omega d,w,\Omega r]//pdConv$

Out[55]//TraditionalForm=

$$-\frac{1}{2} H0^{2} (z+1) \left((3 w+1) \Omega d (z+1)^{3 w} + \Omega m + 2 \Omega r (z+1) \right)$$

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(*(* Cosmology setting *)

\Omega\Lambda=0.687;

H0=0.0691023;

\Omegam=0.31;

\Omegar=9 10<sup>-5</sup>;

w=-1;*)
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