

Our equation is

$$\begin{aligned}
\left(K_X + \frac{\pi'^2}{a^2} K_{XX}\right) \pi'' = & \\
& a^2 \left(K_\pi - \frac{\pi'^2}{a^2} K_{\pi X}\right) - \mathcal{H} \left(2K_X - \frac{\pi'^2}{a^2} K_{XX}\right) \pi' \\
& + K_X \nabla^2 \pi + K_{\pi X} \partial_i \pi \partial^i \pi + 2 \frac{K_{XX}}{a^2} \pi' \partial_i \pi \partial^i \pi' - \mathcal{H} \frac{K_{XX}}{a^2} \pi' \partial_i \pi \partial^i \pi \\
& - \frac{K_{XX}}{a^2} \partial_i \pi \partial_j \pi \partial^j \partial^i \pi
\end{aligned} \tag{0.1}$$

where

$$\begin{aligned}
K &= -g_0 + g_2(X - \hat{X})^2 + g_4(X - \hat{X})^4 \\
K_X &= 2g_2(X - \hat{X}) + 4g_4(X - \hat{X})^3 \\
K_{XX} &= 2g_2 + 12g_4(X - \hat{X})^2 \\
K_\pi &= 0 \\
K_{\pi X} &= 0
\end{aligned}$$

where

$$X = \frac{1}{2a^2} (\pi'^2 - \partial_i \pi \partial^i \pi)$$

## 1 Rewriting the PDE

We write down the PDE in the following form to simplify our numerics implementation,

$$\begin{aligned}
\pi'' = \frac{1}{\left(K_X + \frac{\pi'^2}{a^2} K_{XX}\right)} & \left[ -\mathcal{H} \left(2K_X - \frac{\pi'^2}{a^2} K_{XX}\right) \pi' + K_X \nabla^2 \pi + K_{\pi X} \partial_i \pi \partial^i \pi \right. \\
& \left. + 2 \frac{K_{XX}}{a^2} \pi' \partial_i \pi \partial^i \pi' - \mathcal{H} \frac{K_{XX}}{a^2} \pi' \partial_i \pi \partial^i \pi - \frac{K_{XX}}{a^2} \partial_i \pi \partial_j \pi \partial^j \partial^i \pi \right]
\end{aligned} \tag{1.1}$$

## 2 Numerical implementation

### 2.1 Euler method

$$\pi_{(i,j,k)}^{n+1} = \pi_{(i,j,k)}^n + \pi'_{(i,j,k)}{}^n \Delta\tau \tag{2.1}$$

$$\pi'_{(i,j,k)}{}^{n+1} = \pi'_{(i,j,k)}{}^n + \pi''_{(i,j,k)}{}^n \Delta\tau \tag{2.2}$$

$$\pi_{(i,j,k)}^{n+2} = \pi_{(i,j,k)}^{n+1} + \pi'_{(i,j,k)}{}^{n+1} \Delta\tau \tag{2.3}$$

where the superscript  $n$  and subscript  $(i, j, k)$  shows respectively the time step and the position on the lattice, i.e  $\pi'_{(i,j,k)}{}^n$  is the field  $\pi'$  at discrete time step  $(n)$  and point  $((i, j, k))$  on the lattice. To find  $\pi''_{(i,j,k)}{}^n$  we discretize our equation as

$$\begin{aligned}
\pi''_{(i,j,k)} = & \frac{1}{\left(K_X + \frac{(\pi'^n)^2}{a^2} K_{XX}\right)} \left[ a^2 \left( K_\pi - \frac{(\pi'^n)^2}{a^2} K_{\pi X} \right) - \mathcal{H} \left( 2K_X - \frac{(\pi'^n)^2}{a^2} K_{XX} \right) \pi'^n \right. \\
& + K_X \nabla^2 \pi^n + K_{\pi X} \partial_i \pi^n \partial^i \pi^n + 2 \frac{K_{XX}}{a^2} \pi'^n \partial_i \pi^n \partial^i \pi'^n - \mathcal{H} \frac{K_{XX}}{a^2} \pi'^n \partial_i \pi^n \partial^i \pi^n \\
& \left. - \frac{K_{XX}}{a^2} \partial_i \pi^n \partial_j \pi^n \partial^j \partial^i \pi^n \right]
\end{aligned} \tag{2.4}$$

## 2.2 The leap-frog method

We use the Newton-Stormer-Verlet-leapfrog method [?] to solve the two first order partial differential equations for the linear  $k$ -essence scalar field on the lattice,

$$\pi'^{n+\frac{1}{2}}_{(i,j,k)} = \pi'^{n-\frac{1}{2}}_{(i,j,k)} + \pi''^n_{(i,j,k)} \Delta\tau \tag{2.5}$$

$$\pi^{n+1}_{(i,j,k)} = \pi^n_{(i,j,k)} + \pi'^{n+\frac{1}{2}}_{(i,j,k)} \Delta\tau \tag{2.6}$$

It is important to note that in our scheme we have split the background from perturbations, as a result we have access to the  $\mathcal{H}^{n+\frac{1}{2}}$  independently of the value of the fields. Moreover we compute  $\pi^{n+\frac{1}{2}}_{(i,j,k)}$  as,

$$\pi^n_{(i,j,k)} = \pi'^{n-\frac{1}{2}}_{(i,j,k)} + \pi''^n_{(i,j,k)} \frac{\Delta\tau}{2} \tag{2.7}$$

To do this in the code we need to also define the  $\pi''^n_{(i,j,k)}$  as a new field to make sure that we update the  $\pi'$  field correctly.