

1 Non-linear contrubution

Martin: We can look at $k = 10$

Just note that, if we dont run the GEvolution from enough large scale, we lose the the long mode modulation and the pertubration on quadratic variables would change like $\chi \dots$, and on the other hand also we have the large power from non linearities of scalar field Here we write the full equations and try to solve them numerically in Gevolution, and if we got something intreesting or strange we need to solve them in mathematica to see if we did not have make a mistake !

To solve in mathematica we must solve with Non -linear solve command and to exactly compare with Gevolution we can get some symmetric situations. The discussion with Julian is as following,

I just have a problem with the last part of the note where the time solution of discrete Poisson equation is written
It looked strange to me that for any modes the time dependence of the growing mode decreases!

No, epsilon can have either sign. It vanishes for $K=\text{Laplace}$ which is the continuum limit. For $K>\text{Laplace}$ epsilon is negative, while it is positive for $K<\text{Laplace}$

So I solved the Poisson equation myself and took the real part, the solution in mathematica suggest different relation than is written in the note. At the end depending on the K and Δ we get different behaviour for growing mode. Did I make a mistake somewhere? Please simplify your expressions and you'll see that you got exactly what I wrote.

$P_{\phi}(\phi)/\chi H(k) \sim P_{\phi}(\phi)(k) * 6.675 * (k/k_{Ny})^4$
where k_{Ny} is the Nyqvist wavenumber in your simulation
Can you please give me more information about where this formula come from?

The first (and least trivial) step is to compute the explicit expressions for K and Laplace in discrete Fourier space, given the CIC and NGP weight functions and discrete derivative operators used in the code. This allows you to compute the growth rate of each mode individually.

The next step is to do a Taylor series expansion for small k/k_{Ny} , for which the first term gives the $(k/k_{Ny})^4$ factor - this factor is easy to understand because in the continuum limit $K=\text{Laplace}$, and therefore

$$K/\text{Laplace} = 1 + c(n) (k/k_{Ny})^2 + \dots$$

is inevitable. $c(n)$ here is a function of the direction of the k -vector, and due to the symmetries of the grid it can only be a monopole plus an $l=4$ spherical harmonic. Explicit computation gives

$$c(n) = -(4/5) \pi^{5/2} [Y_{00}(n) - (1/3) Y_{40}(n) - \sqrt{5/126} Y_{44}(n) - \sqrt{5/126} Y_{4-4}(n)]$$

So this already gives you the modified growth as a function of n and k/k_{Ny} for $k/k_{Ny} \ll 1$. The power spectrum estimator is constructed by integrating over the directions. This gives the numerical constant in my final formula, which I quoted as 6.675 but which more precisely is $12 \pi^4 / 175$.

Could you maybe quickly plot this curve to see where this effect becomes relevant?
The result for different redshifts is attached, I assumed the $K_{Ny} = \pi * N_{grid}/\text{Boxsize}$.

This is excellent. It shows that your resolution was good enough that this effect only becomes important right at the Nyqvist wavenumber, and the ϕ you get from gevolution should be quite safe from discretization errors (at first order). Had you chosen $N_{grid}=1024$ or even 512, the smallest scales would have been quite contaminated, as the curves would be shifted upwards by a factor of 16 or 256, respectively.

Let me know if you have further questions.

Maybe we can discuss this in the skype meeting. I think one testbed could be to consider plane symmetric configurations first. For these it turns out that $K=\text{Laplacian}$ (as an identity) for any k/k_{Ny} - in other words, my function $c(n)$ vanishes along the principal axes. These modes will therefore have the correct linear growth despite the discretization. Its a very good idea. Also it looks straightforward to me. I'll try it.

Best,
Farbod

On 26 Apr 2018, at 22:28, Julian Adamek <julian.adamek@qmul.ac.uk> wrote:

Dear Farbod,

Exactly. But unfortunately for my local tests at most I can run with $N_{\text{grid}}=256$, which the interesting scales are contaminated by the error!

Maybe we can discuss this in the skype meeting. I think one testbed could be to consider plane symmetric configurations first. For these it turns out that $K=\text{Laplacian}$ (as an identity) for any k/k_{Ny} - in other words, my function $c(n)$ vanishes along the principal axes. These modes will therefore have the correct linear growth despite the discretization.

Another advantage would be that you can solve the evolution in Mathematica and compare to your implementation in gevolution.

Best wishes,
Julian

$$\begin{aligned} \pi'' + \mathcal{H}(1-3w)\pi' + 3\mathcal{H}\left(-c_s^2 + w\right)\Psi - \Psi' - 3c_s^2\Phi' + \left(3\mathcal{H}^2(c_s^2 - w) + \mathcal{H}'(1-3c_s^2)\right)\pi \\ - c_s^2\nabla^2\pi - 2c_s^2\Phi\nabla^2\pi + (1-c_s^2)\Psi\nabla^2\pi + 3c_s^2\mathcal{H}(1+w)\pi\nabla^2\pi - (1-c_s^2)(\mathcal{H}\pi + \pi')\nabla^2\pi \\ + c_s^2\nabla\Phi.\nabla\pi - (2c_s^2 - 1)\nabla\Psi.\nabla\pi + \frac{\mathcal{H}}{2}\left(2 + 3w + c_s^2\right)\nabla\pi.\nabla\pi - 2(1-c_s^2)\nabla\pi.\nabla(\mathcal{H}\pi + \pi') = 0 \end{aligned} \quad (1)$$

$$\begin{aligned} T_0^0(Gev) &= \Omega_{kess}^0 a^{-3w} \left[1 + \frac{1+w}{c_s^2} \left(-3\mathcal{H}c_s^2\pi - \Psi + (\pi' + \mathcal{H}\pi) - (1-2c_s^2)\frac{(\vec{\nabla}\pi)^2}{2} \right) \right] \\ T_0^i(Gev) &= -\Omega_{kess}^0 a^{-3w}(1+w) \left[1 - \left(\frac{1}{c_s^2} - 1\right)\frac{(\vec{\nabla}\pi)^2}{2} \right] \partial_i\pi \\ T_j^i(Gev) &= w\Omega_{kess}^0 a^{-3w} \left(1 + \frac{1+w}{w} \left[-3\mathcal{H}w\pi - \Psi + (\pi' + \mathcal{H}\pi) - \frac{(\vec{\nabla}\pi)^2}{2} \right] \delta_j^i + \frac{1+w}{w} \delta^{ik} \partial_k\pi \partial_j\pi \right) \end{aligned} \quad (2)$$

we take π and ζ defined as following as new set of variables,

$$\zeta \doteq -\Psi + \pi' + \mathcal{H}\pi, \quad (3)$$

After substitution $\pi' \rightarrow \zeta + \Psi - \mathcal{H}\pi$ we get the following expression for stress tensor according to mathematica, which is very clear,

$$\begin{aligned} &= T_{00} = \frac{1+w}{cs^2} \left(-3\mathcal{H}cs^2\pi\pi[x, t] - \Psi + \pi\text{prime} + \mathcal{H}\pi\pi[x, t] - (1-2cs^2)(\partial_x\pi\pi[x, t]\partial_x\pi\pi[x, t]) \right); \\ &= \pi\text{prime} = \zeta + \Psi - \mathcal{H}\pi\pi[x, t] \\ &= \zeta + \Psi - \mathcal{H}\pi\pi[x, t] \\ &= T_{00} \\ &= \frac{(1+w) \left(\zeta - 3cs^2\mathcal{H}\pi\pi[x, t] - (1-2cs^2)\pi\pi^{(1,0)}[x, t]^2 \right)}{cs^2} \end{aligned}$$

So the stress tensor reads,

$$\begin{aligned}
T_0^0(Gev) &= \Omega_{kess}^0 a^{-3w} \left[1 + \frac{1+w}{c_s^2} \left(\zeta - 3\mathcal{H}c_s^2\pi - (1-2c_s^2)\frac{(\vec{\nabla}\pi)^2}{2} \right) \right] \\
T_0^i(Gev) &= -\Omega_{kess}^0 a^{-3w}(1+w) \left[1 - \left(\frac{1}{c_s^2} - 1 \right) \frac{(\vec{\nabla}\pi)^2}{2} \right] \partial_i \pi \\
T_j^i(Gev) &= w \Omega_{kess}^0 a^{-3w} \left(1 + \frac{1+w}{w} \left[-3\mathcal{H}w\pi + \zeta - \frac{(\vec{\nabla}\pi)^2}{2} \right] \delta_j^i + \frac{1+w}{w} \delta^{ik} \partial_k \pi \partial_j \pi \right) \quad (4)
\end{aligned}$$

After substitution in mathematica we get,

$$\begin{aligned}
&\zeta' - (\mathcal{H}\pi)' + \Psi' + \mathcal{H}(1-3w)(\zeta - \mathcal{H}\pi + \Psi) + 3\mathcal{H}(-c_s^2 + w)\Psi - \Psi' - 3c_s^2\Phi' + (3\mathcal{H}^2(c_s^2 - w) + \mathcal{H}'(1-3c_s^2))\pi \\
&- c_s^2\nabla^2\pi - 2c_s^2\Phi\nabla^2\pi + (1-c_s^2)\Psi\nabla^2\pi + 3c_s^2\mathcal{H}(1+w)\pi\nabla^2\pi - (1-c_s^2)(\zeta + \Psi)\nabla^2\pi \\
&+ c_s^2\nabla\Phi.\nabla\pi - (2c_s^2 - 1)\nabla\Psi.\nabla\pi + \frac{\mathcal{H}}{2}(2+3w+c_s^2)\nabla\pi.\nabla\pi - 2(1-c_s^2)\nabla\pi.\nabla(\zeta + \Psi) = 0 \quad (5)
\end{aligned}$$

The linear part simplification gives the same equation we had before and we can rewrite the equation as following,

$$\begin{aligned}
&\zeta' - 3w\mathcal{H}\zeta + 3c_s^2(\mathcal{H}^2 - \mathcal{H}')\pi - 3c_s^2(\Phi' + \mathcal{H}\Psi) - c_s^2\nabla^2\pi \\
&- 2c_s^2\Phi\nabla^2\pi + (1-c_s^2)\Psi\nabla^2\pi + 3c_s^2\mathcal{H}(1+w)\pi\nabla^2\pi - (1-c_s^2)(\zeta + \Psi)\nabla^2\pi \\
&+ c_s^2\nabla\Phi.\nabla\pi - (2c_s^2 - 1)\nabla\Psi.\nabla\pi + \frac{\mathcal{H}}{2}(2+3w+c_s^2)\nabla\pi.\nabla\pi - 2(1-c_s^2)\nabla\pi.\nabla(\zeta + \Psi) = 0 \quad (6)
\end{aligned}$$

The other equation is like before,

$$\pi' = \zeta + \Psi - \mathcal{H}\pi \quad (7)$$

We actually did nothing except substituting the π' and in non-linear part the substitution is very straightforward. Just note that in Gevolution to compute $\nabla\pi.\nabla(\zeta + \Psi)$ we use the symmetric derivative as following,

$$\nabla\pi.\nabla(\zeta + \Psi) = \frac{1}{4dx^2} \sum_i \left[\pi(x_i + 1) - \pi(x_i - 1) \right] \left[(\zeta(x_i+1) + \Psi(x_i+1)) - (\zeta(x_i-1) + \Psi(x_i-1)) \right] \quad (8)$$

where "1/4" coefficient appear since we are using symmetric derivative and using points with distance two. " x_i " is the lattice coordinate.

1.1 Solving the equation by leap-frog method

Here we explain how we solve the two equations using leap-frog method,

1.1.1 π equation

for the π equation we have,

$$\pi_{n+1} = \pi_n + \pi'_{n+\frac{1}{2}} \Delta\tau \quad (9)$$

$$\pi'_{n+\frac{1}{2}} = \zeta_{n+\frac{1}{2}} - \mathcal{H}_{n+\frac{1}{2}}\pi_{n+\frac{1}{2}} + \Psi_{n+\frac{1}{2}} \quad (10)$$

According to the definition, since the equation for updating π is linear we do not have any trouble with non-linear terms and this part is done like linear equations.

The background part is better to be updated before this step to have a_{kess} at $(n+1/2)$ and then having $\mathcal{H}(n+1/2)$. Just note that making ζ updating at half steps help us because we need all the terms in $\zeta'(n)$ at integer steps.

Since the scalar field Stress energy tensor must be synchronized with particles stress tensor, we need to have all the variables at the same step which is n , so we need to write all the terms at step $n + \frac{1}{2}$ in terms of the values at step n and $n + 1$ as following, of course except ζ which we have it at $n + 1/2$. The easiest model to calculate $F_{n+\frac{1}{2}}$ is by taking average of the next and last step, so

$$\pi_{n+\frac{1}{2}} = \frac{\pi_{n+1} + \pi_n}{2} \quad (11)$$

and the same for all the other variables at step $n + \frac{1}{2}$.

For π we have,

$$\pi_{n+1} = \pi_n + \Delta\tau \left[\zeta_{n+\frac{1}{2}} - \mathcal{H}_{n+\frac{1}{2}} \left(\frac{\pi_{n+1} + \pi_n}{2} \right) + \Psi_{n+\frac{1}{2}} \right] \quad (12)$$

$$\pi_{n+1} = \frac{1}{1 + \mathcal{H}_{n+\frac{1}{2}}\Delta\tau/2} \left[\pi_n + \Delta\tau \left[\zeta_{n+\frac{1}{2}} - \mathcal{H}_{n+\frac{1}{2}} \frac{\pi_n}{2} + \Psi_{n+\frac{1}{2}} \right] \right] \quad (13)$$

As it is clear from the formula we don't have access to the $\Psi_{n+\frac{1}{2}}$ and Ψ_{n+1} , so we use the extrapolation to have them in next half steps,

$$\Psi_{n+\frac{1}{2}} = \Psi_n + \Psi'_n \frac{d\tau}{2} \quad (14)$$

Moreover to have Ψ'_n we use the following formula by saving Ψ at two different steps!

$$\Psi'_n = \frac{\Psi_n - \Psi_{n-1}}{d\tau} \quad (15)$$

1.1.2 ζ equation

We chose to update ζ at half steps since then we need its derivative at integer steps and make the computation simplified. Moreover first π is updated to have it at step $n+1$, then ζ is updated to get the values at $n+3/2$ and from taking the average of the two steps $(n+1/2)$ and $(n+3/2)$ we get ζ at step n which is synchronized with particles stress tensor at step n .

$$\zeta_{n+\frac{3}{2}} = \zeta_{n+\frac{1}{2}} + \zeta'_{n+1}\Delta\tau \quad (16)$$

Just note that in the last step π is updated, so at the current state we have $\zeta_{n+1/2}$, π_{n+1} and ζ_n . $\zeta_{n+1/2}$ which is obtained in the previous loop, ζ_n which in the previous loop was the average of $\zeta_{n-1/2}$ and $\zeta_{n+1/2}$, and π_n was recently updated.

ζ'_n in the equation 16 reads from the differential equation as following,

$$\zeta'_{n+1} = 3\mathcal{H}_{n+1}(w\zeta_{n+1} + c_s^2\Psi_{n+1}) - c_s^2(3\mathcal{H}_{n+1}^2 - 3\mathcal{H}'_{n+1})\pi_{n+1} + 3c_s^2\Psi'_{n+1} + c_s^2\nabla^2\pi_{n+1} \quad (17)$$

$$\begin{aligned} & + 2c_s^2\Phi_{n+1}\nabla^2\pi_{n+1} - (1 - c_s^2)\Psi_{n+1}\nabla^2\pi_{n+1} - 3c_s^2\mathcal{H}_{n+1}(1 + w)\pi_{n+1}\nabla^2\pi_{n+1} + (1 - c_s^2)(\zeta_{n+1} + \Psi_{n+1})\nabla^2\pi_{n+1} \\ & - c_s^2\nabla\Phi_{n+1}\cdot\nabla\pi_{n+1} + (2c_s^2 - 1)\nabla\Psi_{n+1}\cdot\nabla\pi_{n+1} - \frac{\mathcal{H}_{n+1}}{2}(2 + 3w + c_s^2)\nabla\pi_{n+1}\cdot\nabla\pi_{n+1} + 2(1 - c_s^2)\nabla\pi_{n+1}\cdot\nabla(\zeta_{n+1} + \Psi_{n+1}) \end{aligned} \quad (18)$$

Since we need to have ζ_n at ζ'_{n+1} so we write $\zeta_{n+1} = \frac{\zeta_{n+3/2} + \zeta_{n+1/2}}{2}$
We get,

$$\begin{aligned} \zeta_{n+\frac{3}{2}} = & \zeta_{n+\frac{1}{2}} + \Delta\tau \left[3\mathcal{H}_{n+1} \left(w \frac{\zeta_{n+\frac{3}{2}} + \zeta_{n+\frac{1}{2}}}{2} + c_s^2 \Psi_{n+1} \right) - c_s^2 (3\mathcal{H}_{n+1}^2 - 3\mathcal{H}'_{n+1}) \pi_{n+1} + 3c_s^2 \Phi'_{n+1} + c_s^2 \nabla^2 \pi_{n+1} \right. \\ & + 2c_s^2 \Phi_{n+1} \nabla^2 \pi_{n+1} - (1 - c_s^2) \Psi_{n+1} \nabla^2 \pi_{n+1} - 3c_s^2 \mathcal{H}_{n+1} (1 + w) \pi_{n+1} \nabla^2 \pi_{n+1} + (1 - c_s^2) \left(\frac{\zeta_{n+3/2} + \zeta_{n+1/2}}{2} + \Psi_{n+1} \right) \nabla^2 \pi_{n+1} \\ & - c_s^2 \nabla \Phi_{n+1} \cdot \nabla \pi_{n+1} + (2c_s^2 - 1) \nabla \Psi_{n+1} \cdot \nabla \pi_{n+1} - \frac{\mathcal{H}_{n+1}}{2} (2 + 3w + c_s^2) \nabla \pi_{n+1} \cdot \nabla \pi_{n+1} \\ & \left. + 2(1 - c_s^2) \nabla \pi_{n+1} \cdot \nabla \left(\frac{\zeta_{n+3/2} + \zeta_{n+1/2}}{2} + \Psi_{n+1} \right) \right] \end{aligned} \quad (19)$$

As one can see the term $\nabla \left(\frac{\zeta_{n+3/2} + \zeta_{n+1/2}}{2} + \Psi_{n+1} \right)$ is somehow problematic, since we cannot factorize the $\zeta_{n+3/2}$ and $\zeta_{n+1/2}$ simply. Here we use the predictor-corrector method to solve this equation.

As the first guess we take $\zeta_{n+1} = \zeta_{n+1/2}$, which means we have neglected the term $\zeta'_{n+1/2} d\tau/2$, then we get the following equation,

$$\begin{aligned} \zeta_{n+\frac{3}{2}} = & \zeta_{n+\frac{1}{2}} + \Delta\tau \left[3\mathcal{H}_{n+1} \left(w \frac{\zeta_{n+\frac{3}{2}} + \zeta_{n+\frac{1}{2}}}{2} + c_s^2 \Psi_{n+1} \right) - c_s^2 (3\mathcal{H}_{n+1}^2 - 3\mathcal{H}'_{n+1}) \pi_{n+1} + 3c_s^2 \Phi'_{n+1} + c_s^2 \nabla^2 \pi_{n+1} \right. \\ & + 2c_s^2 \Phi_{n+1} \nabla^2 \pi_{n+1} - (1 - c_s^2) \Psi_{n+1} \nabla^2 \pi_{n+1} - 3c_s^2 \mathcal{H}_{n+1} (1 + w) \pi_{n+1} \nabla^2 \pi_{n+1} + (1 - c_s^2) \left(\frac{\zeta_{n+3/2} + \zeta_{n+1/2}}{2} + \Psi_{n+1} \right) \nabla^2 \pi_{n+1} \\ & - c_s^2 \nabla \Phi_{n+1} \cdot \nabla \pi_{n+1} + (2c_s^2 - 1) \nabla \Psi_{n+1} \cdot \nabla \pi_{n+1} - \frac{\mathcal{H}_{n+1}}{2} (2 + 3w + c_s^2) \nabla \pi_{n+1} \cdot \nabla \pi_{n+1} \\ & \left. + 2(1 - c_s^2) \nabla \pi_{n+1} \cdot \nabla \left(\zeta_{n+1/2} + \Psi_{n+1} \right) \right] \end{aligned} \quad (20)$$

Factorizing $\zeta_{n+3/2}$ gives,

$$\begin{aligned} \zeta_{n+\frac{3}{2}} \left[1 - 3\mathcal{H}_{n+1} w \Delta\tau/2 - (1 - c_s^2) \nabla^2 \pi_{n+1} \Delta\tau/2 \right] = & \zeta_{n+\frac{1}{2}} + \Delta\tau \left[3\mathcal{H}_{n+1} \left(w \frac{\zeta_{n+\frac{1}{2}}}{2} + c_s^2 \Psi_{n+1} \right) - c_s^2 (3\mathcal{H}_{n+1}^2 - 3\mathcal{H}'_{n+1}) \pi_{n+1} + 3c_s^2 \Phi'_{n+1} \right. \\ & + c_s^2 \nabla^2 \pi_{n+1} + 2c_s^2 \Phi_{n+1} \nabla^2 \pi_{n+1} - (1 - c_s^2) \Psi_{n+1} \nabla^2 \pi_{n+1} - 3c_s^2 \mathcal{H}_{n+1} (1 + w) \pi_{n+1} \nabla^2 \pi_{n+1} + (1 - c_s^2) \left(\frac{\zeta_{n+1/2}}{2} + \Psi_{n+1} \right) \nabla^2 \pi_{n+1} \\ & - c_s^2 \nabla \Phi_{n+1} \cdot \nabla \pi_{n+1} + (2c_s^2 - 1) \nabla \Psi_{n+1} \cdot \nabla \pi_{n+1} - \frac{\mathcal{H}_{n+1}}{2} (2 + 3w + c_s^2) \nabla \pi_{n+1} \cdot \nabla \pi_{n+1} \\ & \left. + 2(1 - c_s^2) \nabla \pi_{n+1} \cdot \nabla \left(\zeta_{n+1/2} + \Psi_{n+1} \right) \right] \end{aligned} \quad (21)$$

Finally we can write the below equation,

$$\zeta_{n+\frac{3}{2}} = \frac{1}{1 - 3\mathcal{H}_{n+1}w\Delta\tau/2 - (1 - c_s^2)\nabla^2\pi_{n+1}\Delta\tau/2} \left[\zeta_{n+\frac{1}{2}} + \Delta\tau \left(3\mathcal{H}_{n+1} \left(w\frac{\zeta_{n+\frac{1}{2}}}{2} + c_s^2\Psi_{n+1} \right) - c_s^2(3\mathcal{H}_{n+1}^2 - 3\mathcal{H}'_{n+1})\pi_{n+1} + 3c_s^2\Phi'_{n+1} \right. \right. \\ + c_s^2\nabla^2\pi_{n+1} + 2c_s^2\Phi_{n+1}\nabla^2\pi_{n+1} - (1 - c_s^2)\Psi_{n+1}\nabla^2\pi_{n+1} - 3c_s^2\mathcal{H}_{n+1}(1 + w)\pi_{n+1}\nabla^2\pi_{n+1} + (1 - c_s^2)\left(\frac{\zeta_{n+1/2}}{2} + \Psi_{n+1}\right)\nabla^2\pi_{n+1} \\ - c_s^2\nabla\Phi_{n+1}\cdot\nabla\pi_{n+1} + (2c_s^2 - 1)\nabla\Psi_{n+1}\cdot\nabla\pi_{n+1} - \frac{\mathcal{H}_{n+1}}{2}(2 + 3w + c_s^2)\nabla\pi_{n+1}\cdot\nabla\pi_{n+1} \\ \left. \left. + 2(1 - c_s^2)\nabla\pi_{n+1}\cdot\nabla(\zeta_{n+1/2} + \Psi_{n+1}) \right) \right] \quad (22)$$

which at linear order we get the same equation as we had before,

$$\zeta_{n+\frac{3}{2}} = \frac{1}{1 - 3\mathcal{H}_{n+1}w\Delta\tau/2} \left[\zeta_{n+\frac{1}{2}} + \Delta\tau \left(3\mathcal{H}_{n+1} \left(\frac{w\zeta_{n+\frac{1}{2}}}{2} + c_s^2\Psi_{n+1} \right) - c_s^2(3\mathcal{H}_{n+1}^2 - 3\mathcal{H}'_{n+1})\pi_{n+1} + 3c_s^2\Phi'_{n+1} + c_s^2\nabla^2\pi_{n+1} \right) \right] \quad (23)$$

It is very important to check everything for the first loop specifically! For the first loop we first compute $\zeta^{(-1/2)}$ by $\zeta_{-1/2} = \zeta_0 - \zeta'_0 d\tau/2$ and then ζ is updated to have $\zeta^{(1/2)}$ from the values of $\Psi^{(0)}$ but we take $\Phi'^{(0)} = 0$ which is an approximation in our scheme.

Then we go through the procedure which is the same for all loops (first loop included), which we update π to have it at step 1 and then we update ζ to have it at step 3/2 during all of these procedures, Φ is assumed to be constant and we use $\Phi^{(1/2)} = \Phi^{(0)}$ since in the first loop Φ' is zero! So this is an approximation that in the first loop we take $\Psi' = 0$ and use the same Φ at 1/2 and 0 step the same. But in the other loops everything seems correct as following,

To compute $\zeta_{-1/2}$ we just use the linear equation for ζ'_0 and also neglect Φ'_0 .

The whole updating loop is as following,

$$\zeta^{n+1/2}, \zeta^n, \pi^n, \Psi^n, \Psi'^n \longrightarrow a_{kess}^{n+1/2}, \Psi^{n+1/2}, \zeta^{n+1/2}, \pi^{n+1} \longrightarrow a_{kess}^{n+1}, \pi^{n+1}, \Psi^{n+1}, \zeta^{n+3/2}. \quad (24)$$

Approximations:

The important approximations here are:

- 1- At the first loop we take $\Phi' = 0$
- 2- To compute $\zeta_{-1/2}$ we neglect non-linear terms, but it should not be important, since we use predictor corrector method in other loops, which correct the values.
- 3- The value of $\Phi^{(n+1)} = \Phi^{(n)} + \Phi'^{(n)} \times d\tau$ to compute $\zeta^{n+3/2}$ which again at the first loop we assume $\Phi^{(1)} = \Phi^{(0)}$
- 4- Since we do not have $\Phi'^{(n+1)}$ we assume that it does not change fast, so we take $\Phi'^{(n+1)} = \Phi'^{(n)}$ or equivalently $\Phi''^{(n)} \approx 0$
- All the top approximations can be suppressed by increasing the number of kessence update or decreasing the time stepping of Gevolution.
- 5- We use predictor-corrector method for couple of times in each loop to correct our bad guess of non-linear term $\nabla\zeta_{n+1}$

Before talking about predictor-corrector method, we mention the important parts in the Gevolution which is changed for non-linear implementation,

The parameters are the same as linear part, so we just need to check the updatings in the main loop which is as following and the same as linear updating!

```

//*****
//Kessence - LeapFrog:START
//*****
double a_kess=a;
for (i=0;i<sim.nKe_numsteps;i++)
{
//First we update zeta to have it at 0-1/2 just in the first loop
if(cycle==0)
{
for (x.first(); x.test(); x.next())
{
//computing zeta(-1/2)
zeta_half(x)=zeta_integer(x) - 0.5 * dtau * ( 3. * Hconf(a_kess, fourpiG, cosmo) * ( cosmo.
w_kessence * zeta_integer(x) + cosmo.cs2_kessence * phi(x) - chi(x) ) - cosmo.cs2_kessence *
(3. * Hconf(a_kess, fourpiG, cosmo) * Hconf(a_kess, fourpiG, cosmo) - 3. * Hconf_prime(a_kess,
fourpiG, cosmo) ) * pi_k(x));
//Approximations: 1-The linear definition of derivative
// 2-phi_prime = 0
// 3- cs^2 Laplace pi =0
// 4- neglecting the non-linear terms for computing zeta(-1/2)
//Phi_prime is omitted since in the first loop is zero
// We also have neglected Laplace term since Laplace itself is small and is multiplied to cs^2
// which is very suppressed!
// TO MAKE SURE THE APPROXIMATIONS WORK WELL, WE NEED TO INCREASE THE PRECISION AND SEE THE
IMPROVEMENTS!
}
//Updating zeta to get zeta(1/2) and zeta(0) just in the first loop
// In sum: zeta(1/2) = zeta(-1/2)=zeta(0) + zeta'(0) dtau for the first loop
update_zeta(dtau/ sim.nKe_numsteps, dx, a_kess, phi, phi_old, chi, chi_old, pi_k, zeta_half,
zeta_integer, cosmo.Omega_kessence, cosmo.w_kessence, cosmo.cs2_kessence, Hconf(a_kess, fourpiG,
cosmo), Hconf_prime(a_kess, fourpiG, cosmo));
zeta_half.updateHalo();
zeta_integer.updateHalo();
}

//Since we have pi(n+1)=pi(n) + pi'(n+1/2), and in pi'(n+1/2) we have H(n+1/2) we update the
background before updating the pi to have H(n+1/2)
rungekutta4bg(a_kess, fourpiG, cosmo, dtau / sim.nKe_numsteps / 2.0);
//First we update pi to have it at n+1 (at first loop from the value at (0) and the value of zeta at
1/2 and H(n+1/2) we update pi at (1))
update_pi_k(dtau/ sim.nKe_numsteps, dx, a_kess, phi, phi_old, chi, chi_old, pi_k, zeta_half,
zeta_integer, cosmo.Omega_kessence, cosmo.w_kessence, cosmo.cs2_kessence, Hconf(a_kess, fourpiG,
cosmo),Hconf_prime(a_kess, fourpiG, cosmo)); // H_old is updated here in the function
pi_k.updateHalo();
// Now we have pi(n+1) and a_kess(n+1/2) so we update background by halfstep to have a_kess(n+1) to
calculate zeta'(n+1) to have zeta(n+1/2)=zeta(n-1/2) + zeta'(n+1) dtau
rungekutta4bg(a_kess, fourpiG, cosmo, dtau / sim.nKe_numsteps / 2.0 );
//Now from the values of zeta at step (1/2) we calculate zeta(3/2) and then we calculate zeta(1) which
is synched with pi(1)
update_zeta(dtau/ sim.nKe_numsteps, dx, a_kess, phi, phi_old, chi, chi_old, pi_k, zeta_half,
zeta_integer, cosmo.Omega_kessence, cosmo.w_kessence, cosmo.cs2_kessence, Hconf(a_kess, fourpiG,
cosmo), Hconf_prime(a_kess, fourpiG, cosmo));
zeta_half.updateHalo();
zeta_integer.updateHalo();
}
}
#endif BENCHMARK
kessence_update_time += MPI_Wtime() - ref_time;
ref_time = MPI_Wtime();
#endif
//*****
//Kessence - LeapFrog: End
//*****

```

The most important changes compared to linear equations occur in the Gevolution.hpp, since we need to add, non-linear terms in the stress tensor correctly, adding non-linear terms in the updating equations and also predictor-corrector method should be added.

The stress tensor of kessence is modified as following,

```

//////////
template <class FieldType>
void projection_Tmunu_kessence( Field<FieldType> & T00, Field<FieldType> & T0i, Field<FieldType> & Tij,
double dx, double a, Field<FieldType> & phi,Field<FieldType> & phi_old, Field<FieldType> & chi, Field<
FieldType> & pi_k, Field<FieldType> & zeta_integer, double Omega_fld, double w, double cs2, double
Hcon, double fourpiG, int method )
{
Site xField(phi.lattice());
double coeff1, coeff2, coeff3, Hdot, psi;
double gradientpi_squared, Dx_pi_Dx_pi, Dx_pi_Dy_pi, Dx_pi_Dz_pi, Dy_pi_Dy_pi, Dy_pi_Dz_pi,
Dz_pi_Dz_pi;
Site x(phi.lattice());
double gradient_pi2;
coeff1=Omega_fld*pow(a,-3.*w)*(1.+w)/(cs2);
coeff2=Omega_fld*pow(a,-3.*w)*(1.+w);

for (xField.first(); xField.test(); xField.next())
{
//*****
// (D_i pi)^2
//*****
gradientpi_squared =0.25*(pi_k(xField+0) - pi_k(xField-0))* (pi_k(xField+0) - pi_k(xField-0))/(dx*
dx);
gradientpi_squared+=0.25*(pi_k(xField+1) - pi_k(xField-1))* (pi_k(xField+1) - pi_k(xField-1))/(dx*
dx);
gradientpi_squared+=0.25*(pi_k(xField+1) - pi_k(xField-2))* (pi_k(xField+1) - pi_k(xField-2))/(dx*

```

```

    dx);
// *****
// (X,X) ::::::: > Dx_pi_Dx_pi = GradX(pi).GradX(pi)
// *****
Dx_pi_Dx_pi = 0.25*(pi_k(xField+0) - pi_k(xField-0))* (pi_k(xField+0) - pi_k(xField-0))/(dx*dx);
// *****
// (X,Y) ::::::: > Dx_pi_Dy_pi = GradX(pi).GradY(pi) = GradY(pi).GradX(pi)
// *****
Dx_pi_Dy_pi = 0.25*(pi_k(xField+0) - pi_k(xField-0))* (pi_k(xField+1) - pi_k(xField-1))/(dx*dx);
// *****
// (X,Z) ::::::: > Dx_pi_Dz_pi = GradX(pi).GradZ(pi)
// *****
Dx_pi_Dz_pi = 0.25*(pi_k(xField+0) - pi_k(xField-0))* (pi_k(xField+2) - pi_k(xField-2))/(dx*dx);
// *****
// (Y,Y) ::::::: > Dy_pi_Dy_pi = GradY(pi).GradY(pi)
// *****
Dy_pi_Dy_pi = 0.25*(pi_k(xField+1) - pi_k(xField-1))* (pi_k(xField+1) - pi_k(xField-1))/(dx*dx);
// *****
// (Y,Z) ::::::: > Dy_pi_Dz_pi = Grady(pi).Gradz(pi)
// *****
Dy_pi_Dz_pi = 0.25*(pi_k(xField+1) - pi_k(xField-1))* (pi_k(xField+2) - pi_k(xField-2))/(dx*dx);
// *****
// (Z,Z) ::::::: > Dz_pi_Dz_pi = Gradz(pi).Gradz(pi)
// *****
Dz_pi_Dz_pi = 0.25*(pi_k(xField+2) - pi_k(xField-2))* (pi_k(xField+2) - pi_k(xField-2))/(dx*dx);
// *****
psi= phi(xField) - chi(xField);

// *****
//STRESS TENSOR COMPONENTS
// *****
// 0-0-component: (Time,Time)
T00(xField) = - coeff1 * ( -3. * cs2 * Hcon * pi_k(xField) + zeta_integer(xField)
/*Non-linear*/ - (1. - 2. * cs2) * gradientpi_squared / 2. );
// *****
// 1-1-component: (X,X)
Tij(xField, 0, 0) = + coeff2 * (-3.* w * Hcon* pi_k(xField) + zeta_integer(xField)
/*Non-linear*/ - gradientpi_squared / 2. + Dx_pi_Dx_pi );
// *****
// 2-2-component: (Y,Y)
Tij(xField, 1, 1) = + coeff2 * (-3.* w * Hcon* pi_k(xField) + zeta_integer(xField)
/*Non-linear*/ - gradientpi_squared / 2. + Dy_pi_Dy_pi );
// *****
// 3-3-component: (Z,Z)
Tij(xField, 2, 2) = + coeff2 * (-3.* w * Hcon* pi_k(xField) + zeta_integer(xField)
/*Non-linear*/ - gradientpi_squared / 2. + Dz_pi_Dz_pi );
// *****
// 1-2-component: (X,Y)
Tij(xField, 0, 1) = + coeff2 * ( /*Non-linear*/ Dx_pi_Dy_pi );
// *****
// 1-3-component: (X,Z)
Tij(xField, 0, 2) = + coeff2 * ( /*Non-linear*/ Dx_pi_Dz_pi );
// *****
// 2-3-component: (Y,Z)
Tij(xField, 1, 2) = + coeff2 * ( /*Non-linear*/ Dy_pi_Dz_pi );
// *****

// *****
//In the case of Vector parabolic
// *****
if(method==1) // method=1 Turn on vector elliptic
{
    // T01:(Time,X)
    T0i(xField, 0) = -coeff2 * (1. - /*Non-linear*/ (1./cs2 -1.) * gradientpi_squared / 2.) *
(pi_k(xField + 0) - pi_k(xField - 0)) / (2. * dx);
// *****
    // T02:(Time,Y)
    T0i(xField, 1) = -coeff2 * (1. - /*Non-linear*/ (1./cs2 -1.) * gradientpi_squared / 2.) *
(pi_k(xField + 1) - pi_k(xField - 1)) / (2. * dx);
// *****
    // T03:(Time,Z)
    T0i(xField, 2) = -coeff2 * (1. - /*Non-linear*/ (1./cs2 -1.) * gradientpi_squared / 2.) *
(pi_k(xField + 2) - pi_k(xField - 2)) / (2. * dx);
// *****
}
}
}

```

The updating π equation does not change, since by definition is linear,

```

void update_pi_k( double dtau, double dx, double a, Field<FieldType> & phi, Field<FieldType> &
phi_old, Field<FieldType> & chi, Field<FieldType> & chi_old, Field<FieldType> & pi_k, Field<
FieldType> &, Field<FieldType> & zeta_half, double Omega fld, double w, double cs2, double
Hcon, double H_prime)
{
    double psi, psi_prime, psi_half;
    double Coeff1 = 1./(1. + Hcon * dtau/2.);
    Site x(phi.lattice());
    for (x.first(); x.test(); x.next())
    {
        psi=phi(x) - chi(x);
        psi_prime= ((phi(x) - chi(x))-(phi_old(x) - chi_old(x)))/dtau;
        psi_half= psi + psi_prime * dtau/2.; //psi_half (n+1/2) = psi(n) + psi_prime'(n) dtau/2
// *****
//pi Updating which is linear by definition
// *****
        pi_k(x)=Coeff1 * (pi_k(x) + dtau * ( zeta_half(x) - Hcon * pi_k(x)/2. + psi_half )); //
        pi_k(n+1)
    }
}

```



```

// *****
}
}

```

The most challenging part is implementing non-linear terms in the updating ζ equation and since we use some approximations we first introduce the predictor-corrector idea and then put the C++ code.

1.1.2.1 Predictor-corrector method

This is an important part which should be done to have the errors propagation under the control! As we have explained already in the first loop of updating we take $\nabla\zeta_{n+1} = \nabla\zeta_{n+1/2}$ which is just prediction and then we calculate $\zeta_{n+3/2}$ from the predictor and then we compute ζ_{n+1} by averaging the two ζ at half steps.

So at the end we have ζ_{n+1} which we did not have in the start, so we are going to use ζ_{n+1} as new value to compute $\zeta_{n+3/2}$ and then again we calculate new corrected ζ_{n+1} and we take it as initial value to compute the corrected $\zeta_{n+3/2}$ and ζ_{n+1} .

The best way to do it is first to check that ζ_n from the next order corrector is near to original one (like the relative error $2\frac{\zeta_1 - \zeta_2}{\zeta_1 + \zeta_2} < 0.01$, which gives less than 1% error).

The code for predictor corrector method in C++ is written in below, note that all the lines have comments make the code easier to follow,

```

// We use predictor corrector method to calculate \zeta precisely for non-linear case. (for linear
// equation is does not make better)
template <class FieldType>
void update_zeta(double dtau, double dx, double a, Field<FieldType> & phi, Field<
FieldType> & phi_old, Field<FieldType> & chi, Field<FieldType> & chi_old, Field<
FieldType> & pi_k, Field<FieldType> & zeta_half, Field<FieldType> &
zeta_integer, double Omega_fld, double w, double cs2, double Hcon, double
H_prime )
{
double Gradphi_Gradpi, Gradpsi_Gradpi, Gradpi_Gradpi, GradPsiZeta_Gradpi, Dx_psi, Dy_psi, Dz_psi;
double C1, C2, C3, psi, psi_old, psi_prime, phi_prime, Laplacian_pi,
zeta_old_half;
// Since a_kess is at (n+1) so H_prime is at (n+1) which is needed to calculate zeta(n+1/2)
// *****
// Coefficient two, H(n+1), H_prime(n+1) since BG already updated
// *****
C2 = cs2 * (3. * Hcon * Hcon - 3. * H_prime);
// *****
// Coefficient three, H(n+1), H_prime(n+1) since BG already updated
// *****
C3 = (2. + 3. * w + cs2) * Hcon/2.;

Site x(phi.lattice());
for (x.first(); x.test(); x.next())
{
// *****
// Laplace pi, pi(n+1) since pi updated before zeta and a_kess(n+1)
// *****
Laplacian_pi = pi_k(x-0) + pi_k(x+0) - 2. * pi_k(x);
Laplacian_pi += pi_k(x+1) + pi_k(x-1) - 2. * pi_k(x);
Laplacian_pi += pi_k(x+2) + pi_k(x-2) - 2. * pi_k(x);

Laplacian_pi = Laplacian_pi / (dx * dx);
// *****
// psi(n), APPROXIMATION: we take psi(n+1) = psi(n)
// *****
psi = phi(x) - chi(x);

// *****
// Coefficient one, H( at n+1)
// *****
C1 = 1. / (1. - 3. * Hcon * w * dtau/2. - (1. - cs2) * Laplacian_pi * dtau/2.);
// *****
// phi'(n), APPROXIMATION: we take phi'(n+1) = phi'(n)
// *****
phi_prime = (phi(x) - phi_old(x)) / dtau; // phi_prime(n+1) since we
// want to use it to compute zeta (n+3/2)
// NOTE: We dont have phi'(n+1) since it will be updated by particles later, but we remains
// constant and phi_prime(n) = phi_prime(n+1) or take the second derivative for this period
// approximately to zero!
// *****
// psi'(n)
// psi(n+1) = psi(n) + psi'(n) * dtau
// *****
psi_prime = ((phi(x) - chi(x)) - (phi_old(x) - chi_old(x))) / dtau;
psi = psi + psi_prime * dtau;
// *****
// Grad_i Psi
// *****
Dx_psi = ((phi(x + 0) - chi(x + 0)) - (phi(x - 0) - chi(x - 0)));

```



```

/*Non-linear(5,6,7)*/ - cs2 * Gradphi_Gradpi + (2. * cs2 -1.) * Gradpsi_Gradpi - C3 *
Gradpi_Gradpi
/*Non-linear(8) */ + 2.* (1. - cs2) * GradPsiZeta_Gradpi;
//*****
//Computing the zeta(n+3/2) from the new value of zeta'(n+1) and zeta(n+1/2)
//*****
zeta_predictor_half_n1 = zeta_old_half + zeta_prime_int_n0 * dtau;
//*****
//Computing the corrected zeta(n+1) from the new value of zeta(n+3/2) and zeta(n+1/2)
//*****
zeta_predictor_int_n0 = (zeta_predictor_half_n1 + zeta_old_half)/2.;
//*****
// Now we must check the relative error between the two corrected and predicted ones
zeta_integer(x)=zeta_predictor_int_n0 and the previous step: zeta_half(x), zeta_integer(x)
//If the error has reached less than 10^-6 it breaks the loop
//*****
if (2.* abs(zeta_integer(x)-zeta_predictor_int_n0)/(zeta_integer(x)+zeta_predictor_int_n0)<1.e
-8)break;
zeta_half(x)=zeta_predictor_half_n1; // zeta(n+3/2) after correction
zeta_integer(x)= (zeta_half(x) + zeta_old_half)/2.; //zeta(n+1) after correction
numerator++;
}
if (numerator==n_correcor_steps) cout << "\033[1;31mbold WARNING: PRECISION ERROR ON KESSENCE
FIELD ZETA, More than 1% Error\033[0m\n" << '\n';
}
}

```

As we know although the implementation is checked well but it is very important to design some tests to check the terms as precise as possible.

2 Non-linear terms tests: **How?**

Here we aim to check the non-linear terms alone, just to make sure everything is correctly computed or we did not forgot any trivial point. What we are going to do is solving the below two equations in mathematica and Gevolution to see if we get reasonable results! We believe that these two equations are enough to be solved, since all other terms look like one of these two terms and also we have already checked that ζ and π up to a great order agree in mathematica and Gevolution, but maybe it would be good check to put all the possible terms together to see if we can solve in mathematica?

$$\pi' = \zeta + \Psi - \mathcal{H}\pi \quad (25)$$

$$\zeta' = \mathcal{H}\nabla\pi.\nabla\pi \quad (26)$$

$$\zeta' = c_s^2\Phi(\nabla\pi)^2 \quad (27)$$

Unfortunately, I do not know how to solve in mathematica to compare with the result of Gevolution!

2.1 Predictor-corrector method

Here we try to measure the effect of non-linear terms. Before doing it we first check if predictor-corrector method works well?

To do, we measure how many times it goes through predictor-corrector method to converge? After doing some experiments, it seems that it works well! So for demanding precision like 10^{-10} which is actually the relative improvement, it just go through the loop once mostly or twice and some times specially low redshifts it must go 10 times in the loop to reach the precision, so it hits precision error in low redshifts.

As an example when we ask for the the relative precision of 10^{-8} and ask if the numerator was more than 6 print the precision we get the following result in the terminal,

```

numerator: 7 The rel diff: 7.00482e-10
numerator: 7 The rel diff: 8.28997e-10
numerator: 8 The rel diff: 6.95646e-10
numerator: 7 The rel diff: 6.65728e-10
numerator: 7 The rel diff: 3.09033e-09
numerator: 7 The rel diff: 1.9231e-09
numerator: 7 The rel diff: 3.23853e-09
numerator: 8 The rel diff: 6.69198e-10
numerator: 7 The rel diff: 1.74146e-09
numerator: 7 The rel diff: 6.87903e-10
numerator: 7 The rel diff: 7.36749e-10
numerator: 7 The rel diff: 2.34831e-09
numerator: 8 The rel diff: 6.397e-10
numerator: 7 The rel diff: 3.06782e-09
numerator: 7 The rel diff: 1.82822e-09
numerator: 7 The rel diff: 7.01182e-10

```

The numerator is the number of times which go through corrector loop to improve the variable! For the more realistic case we consider getting near to the previous one to 10^{-4} relative difference which is less than 1%, we see that it goes at most 4 times over the loop and if we ask for printing the values of numerator more than 4 we get,

```

cycle 90, background information: z = 1.85694, average T00 = 0.312045, background model = 0.312046
cycle 90, time integration information: max |v| = 0.00339536 (cdm Courant factor = 0.0648959), time step
/ Hubble time = 0.04
numerator: 5 The rel diff: 1.0313e-05
numerator: 5 The rel diff: 2.32914e-05
writing power spectra at z = 0.996995 (cycle 99), tau/boxsize = 18.0884
cycle 100, background information: z = 0.918918, average T00 = 0.312044, background model = 0.312046
cycle 100, time integration information: max |v| = 0.00415543 (cdm Courant factor = 0.087711), time step
/ Hubble time = 0.04
numerator: 5 The rel diff: 1.71626e-05
cycle 110, background information: z = 0.285983, average T00 = 0.312043, background model = 0.312046
cycle 110, time integration information: max |v| = 0.00598645 (cdm Courant factor = 0.124585), time step
/ Hubble time = 0.04
numerator: 5 The rel diff: 6.9317e-06
numerator: 5 The rel diff: 1.67579e-05
numerator: 5 The rel diff: 1.45638e-05

```

Which is just four times and at low redshifts! We also do not see numerator more than 5!

3 Non-linear contribution in the code

To measure the non-linear terms contribution, we have couple of ways. One is comparing the the linear versus non-linear contribution plot on different variables...

The other one is just internally look at the relative importance of non-linear terms in comparison with linear terms, so in ζ' we just measure the relative between non-linear terms and linear ones $\frac{\text{all non-linear contributions} - \text{all linear contributions}}{\text{all linear contributions}}$, which is as following,

$$\frac{\left(+ 2c_s^2 \Phi_{n+1} \nabla^2 \pi_{n+1} - (1 - c_s^2) \Psi_{n+1} \nabla^2 \pi_{n+1} \right) \dots \left(- 3\mathcal{H}_{n+1} \left(w \frac{\zeta_{n+1}}{2} + c_s^2 \Psi_{n+1} \right) - c_s^2 (3\mathcal{H}_{n+1}^2 - 3\mathcal{H}'_{n+1}) \pi_{n+1} + \dots \right)}{3\mathcal{H}_{n+1} \left(w \frac{\zeta_{n+1}}{2} + c_s^2 \Psi_{n+1} \right) - c_s^2 (3\mathcal{H}_{n+1}^2 - 3\mathcal{H}'_{n+1}) \pi_{n+1} + 3c_s^2 \Phi'_{n+1} + c_s^2 \nabla^2 \pi_{n+1}} \quad (28)$$

what we get is somehow strange! So in different points of lattice we get the following terminal output for the quantity which shows the importance of non-linear terms!

```

NL_REL: -2.36611
NL_REL: -2.38774
NL_REL: -2.29923
NL_REL: -2.22916
NL_REL: -2.18889
NL_REL: -2.15739
NL_REL: -2.31122
NL_REL: -2.48264
NL_REL: -2.54007
NL_REL: -2.4611
NL_REL: -2.47478
NL_REL: -2.38359
NL_REL: -2.35794
NL_REL: -2.38377
NL_REL: -2.4141
NL_REL: -2.33127
NL_REL: -2.29721
NL_REL: -2.33786
NL_REL: -2.33795
NL_REL: -2.35654
NL_REL: -2.30935

```

```

NL_REL: -2.24938
NL_REL: -2.30449
NL_REL: -2.4566
NL_REL: -2.37959
NL_REL: -2.64718
NL_REL: -2.40469
NL_REL: -2.30812
NL_REL: -2.44066
NL_REL: -2.42957
NL_REL: -2.24973
NL_REL: -2.48907
NL_REL: -2.41429
NL_REL: -2.40282

```

4 Linear and non-linear comparisons with class results

Here we want to compare linear and non-linear results versus class for different variables, ζ , π and δ_{kess} to see what is the effect of non-linearities and does it make sense at all?

5 Sensitivity to initial conditions

What would be error if we set the initial condition at $z=100$ to zero?

6 The effect of non linearities on gravitational potential at some different redshifts Ψ and matter power spectrum

7 Solve the solution to non-linear terms in mathematica with Gevolution

8 Measuring the average of field in Gevolution to see the backreaction