

The Laniakea supercluster of galaxies

R. Brent Tully¹, Hélène Courtois², Yehuda Hoffman³ & Daniel Pomarède⁴

Galaxies congregate in clusters and along filaments, and are missing from large regions referred to as voids. These structures are seen in maps derived from spectroscopic surveys^{1,2} that reveal networks of structure that are interconnected with no clear boundaries. Extended regions with a high concentration of galaxies are called ‘super-clusters’, although this term is not precise. There is, however, another way to analyse the structure. If the distance to each galaxy from Earth is directly measured, then the peculiar velocity can be derived from the subtraction of the mean cosmic expansion, the product of distance times the Hubble constant, from observed velocity. **The peculiar velocity is the line-of-sight departure from the cosmic expansion and arises from gravitational perturbations; a map of peculiar velocities can be translated into a map of the distribution of matter³.** Here we report a map of structure made using a catalogue of peculiar velocities. We find locations where peculiar velocity flows diverge, as water does at watershed divides, and we trace the surface of divergent points that surrounds us. Within the volume enclosed by this surface, the motions of galaxies are inward after removal of the mean cosmic expansion and long range flows. We define a supercluster to be the volume within such a surface, and so we are defining the extent of our home supercluster, which we call Laniakea.

The distribution of matter can be determined by two independent methods: either based on surveys of the distribution of galaxies in projection and redshift, or from the motions of galaxies. With the former, using galaxy redshift surveys, the assumption is required that the galaxy ‘lighthouses’ and mass distribution are strongly correlated, a condition that requires confirmation if, as is suspected, only a minor fraction of matter is baryonic. Moreover, with the former there is a stringent demand that the survey be complete, or at least that its incompleteness be well understood. With the latter, studies of galaxy motions, sparse sampling is acceptable (indeed inevitable) but dealing with errors is a challenge. Except for the very closest galaxies, uncertainties in distance measurements translate into uncertainties in the peculiar velocities of galaxies that are larger in amplitude than the actual peculiar velocities. Many measurements are required for suitable averaging and care must be taken to avoid systematic errors. Overall, the two paths to determining the distribution of matter are in good agreement, a consequence that represents a considerable success for the standard model of structure formation via gravitational instability^{4–7}.

The path from velocities to mass distributions benefits from the coherence in velocities on large scales. Multipole components in the velocity field can point to tidal influences beyond the survey region. The current all-sky redshift surveys and distance measurement surveys reach similar depths, but the latter probe structure to greater distances because of sensitivity to uncharted attractors and repellers. Coherence in motions on large scales means that signals can be measured by averaging over data in circumstances where individual contributions are very noisy.

Details about the actual measurement of galaxy distances and the derivation of peculiar velocities are given in Methods. These two parameters are available for more than 8,000 galaxies, affording extremely detailed information locally, degrading outward to increasingly coarse coverage. We use this material to reconstruct the large-scale structure of the near-by Universe⁷.

We obtain the underlying three-dimensional velocity and density fields by the Wiener filter algorithm^{8,9}, assuming the standard model of cosmology as a Bayesian prior. Large-scale structure is assumed to develop from gravitational instabilities out of primordial random Gaussian fluctuations. The developing density and velocity fields retain their Gaussian properties as long as the growth is in the linear regime. It has been shown⁸ that with a random Gaussian field, the optimal Bayesian estimator of the field given the data is the Wiener filter minimal variance estimator. At the present epoch, large-scale structure has become nonlinear on small scales. However, it is an attractive feature of the velocity field that the break from linearity is only on scales of a few megaparsecs, an order of magnitude smaller in scale than the deviations from linearity for the density field. In any event, the present discussion concerns structure on scales of tens to hundreds of megaparsecs, comfortably in the linear regime.

The Wiener filter result is determined by the ratio of power to power + noise. Hence, the large-scale structure is strongly constrained nearby, where uncertainties are small and the coverage is extensive. At large distances, where the data become more sparse and noisy, the Wiener filter attenuates the recovered density and velocity fields to the null field that is expected in the absence of data. However in the linear regime there is coherence in galaxy flows on much larger scales than seen in density fluctuations. Tidal influences from beyond the surveyed regions can be manifested in cosmic flows on scales that exceed the coverage in measured distances by a factor of two (ref. 10).

The ultimate goal is to map the velocity field to a radius that completely encompasses the sources of the motion of the Local Group (of 631 km s^{-1}) reflected in the cosmic microwave background dipole¹¹. However, our knowledge of flows on large scales remains inevitably modulated by the extent of the data. Our analysis of the data in the Cosmicflows-2 galaxy catalogue (see Methods for details) tells us that a coherent flow extends across the full extent of the region that we can map, reaching the Shapley concentration¹². It is clear that we do not yet have a sufficiently extensive compendium of distances to bound the full source of our deviant motion from the cosmic expansion.

For the present discussion, we focus on intermediate-scale flow patterns. **The standard model of cosmology predicts that on the scale we are considering the flow is irrotational: namely, the velocity field \mathbf{v} is the gradient of a potential ϕ , $\mathbf{v} = -\nabla\phi$.** The local minima and maxima of the potential (attractors and repellers respectively) are the drivers of the large-scale flow. We can define a ‘basin of attraction’ as the volume containing all points whose flow lines converge at a given attractor. The large-scale structure can be characterized on scales of a few megaparsecs and above by attractors and their basins of attraction.

The Wiener filter provides a straightforward way of decomposing the velocity field into a local component that is induced by the distribution within a zone and a tidal residual¹³. In the linear regime, the velocity and density fields, \mathbf{v} and δ , are directly related: $\nabla \cdot \mathbf{v} = -H_0 f(\Omega_m, \Omega_\Lambda) \delta$ where f depend on the cosmological matter and vacuum energy densities characterized by Ω_m and Ω_Λ . We specify a centre and radius, and the density field within the defined volume is assumed to be a cut-out from the full Wiener filter density map. Then the Poisson-like equation between velocity and density can be solved to derive the velocity

¹Institute for Astronomy, University of Hawaii, Honolulu, Hawaii 96822, USA. ²Université Claude Bernard Lyon I, Institut de Physique Nucléaire, Université Lyon I, CNRS/IN2P3, Lyon 69622, France. ³Racah Institute of Physics, Hebrew University, Jerusalem 91904, Israel. ⁴Institut de Recherche sur les Lois Fondamentales de l'Univers, CEA/Saclay, 91191 Gif-sur-Yvette, France.

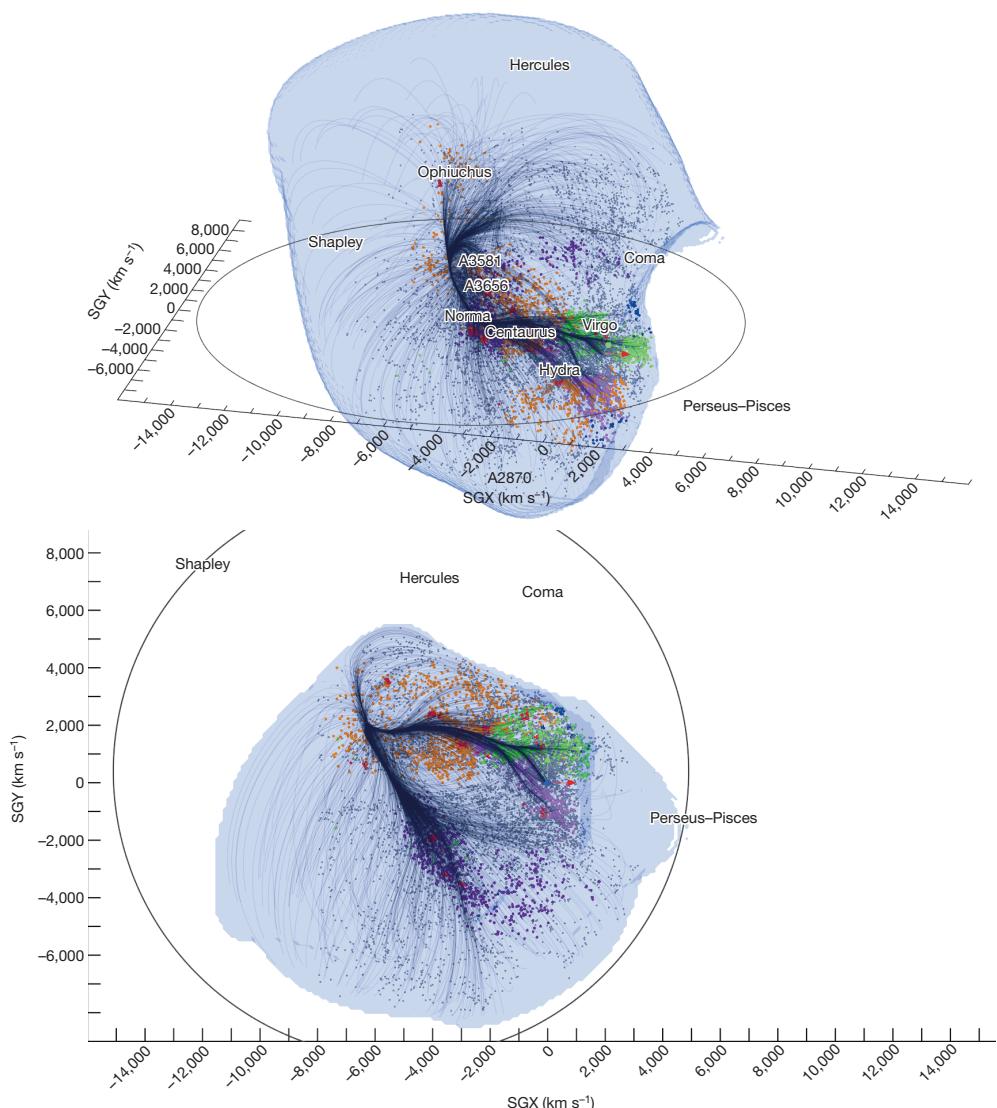


Figure 1 | Two views of the Laniakea supercluster. The outer surface (blue) demarcates the limits of local velocity flows. The plot is in supergalactic coordinates with axes SGX and SGY shown in the plane at SGZ = 0. Our Milky Way galaxy is at the origin. Units of $1,000 \text{ km s}^{-1}$ in velocity correspond to roughly 13.3 Mpc. Velocity streamlines are shown in black and terminate in the vicinity of the Norma cluster. Individual galaxies from a redshift catalogue are given colours to distinguish major components within the Laniakea supercluster: the historical Local supercluster in green, the Great Attractor region in orange, the Pavo-Indus filament in purple, and structures including the Antlia wall and Fornax-Eridanus cloud in magenta. Several major entities are named. Norma, Hydra, Centaurus, Virgo, Ophiuchus, A2870, A3581 and A3656 are individual clusters of galaxies embedded within the Laniakea supercluster. Shapley, Hercules, Coma and Perseus–Pisces are complexes of galaxies outside Laniakea. The outer black circle defines the domain used to separate between local and tidal flows. The panels provide two perspectives of the same scene.

field responding to just the matter within the prescribed volume. The vector subtraction of the local velocity component from the full flow gives the external component of the velocity field. The residual component is responsible for the bulk motion of the zone under consideration and for a quadrupole component within the zone. The decomposition allows us to probe the local velocity field, with the tidal field induced by distant structures filtered out. Relative attractors and their basins of attraction are defined with respect to that local field.

One more useful tool is now mentioned before we turn to results. At each position in space, the three eigenvalues of the velocity shear tensor can be calculated. If these eigenvalues are ordered from most positive to most negative, then a threshold can be set that captures four possibilities. Flows can be inward on all three axes, the condition of a cluster, inward on two axes and outward on the third, the condition of a filament, inward on one axis and outward on two, hence a sheet, or outward on all three axes, hence a void. Boundaries can be created around contiguous regions with the same shear properties, and the contours outline the cosmic web as reconstructed by the V-web algorithm¹⁴.

We now consider the full presentation of our results given in the Supplementary Video and in Extended Data Figs 1–5. We pay attention in turn to smaller (but still in the linear regime) scales to examine the separation of local and tidal flows and to isolate local basins of attraction. Particular attention has been given to locations where there are local divergences. In co-moving coordinates and with the removal of long-range

flows, there are places where relatively neighbouring galaxies can be found to be moving in opposing directions towards separate local basins of attraction. Voids are usually the demarcations between attraction basins, but divergences can occur along filaments and sheets. We note the very nearby case between our home basin of attraction and the Perseus–Pisces complex^{15,16}, where we find a particularly impressive example of velocity flows in an apparent bridge between attractors that are diverging. This divergence occurs in the feature called the Arch in Extended Data Fig 2. Similar structures abound on close inspection. Velocity information reveals the locations of divergence along filaments between high density regions. This dissipation of the cosmic web is expected with the accelerated expansion of the Universe. We emphasize that peculiar velocity information can reveal details that are otherwise hard to discern.

The particular interest in the present discussion is with the largest structure that can be circumscribed within the currently available distance and peculiar-velocity data: the structure schematically illustrated in Figs 1 and 2. The region includes 13 Abell clusters (with the Virgo cluster). Local flows within the region converge towards the Norma and Centaurus clusters, in good approximation to the location of what has been called the ‘Great Attractor’¹⁷. This volume includes the historical Local and Southern superclusters¹⁸, the important Pavo-Indus filament, an extension to the Ophiuchus cluster, the Local Void, and the Sculptor and other bounding voids. This region of inflow towards a local basin

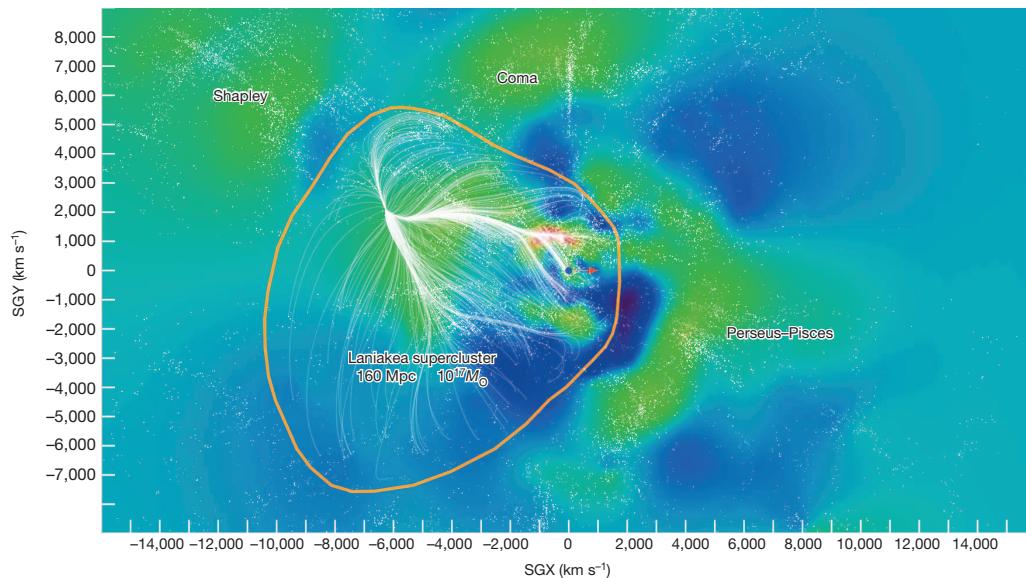


Figure 2 | A slice of the Laniakea supercluster in the supergalactic equatorial plane. Shaded contours represent density values within the equatorial slice, with red at high densities, green at intermediate densities and blue in voids. Our Milky Way galaxy is located at the black dot at the origin of the supergalactic coordinates system: a red arrow points right from the black dot toward increasing SGX and a green arrow points up toward increasing SGY. Individual galaxies from a redshift catalogue are given as white dots. Velocity flow streams within the Laniakea basin of attraction are shown in white. The orange contour encloses the outer limits of these streams. This domain has an extent of $\sim 12,000 \text{ km s}^{-1}$ ($\sim 160 \text{ Mpc}$ diameter) and encloses $\sim 10^{17}$ solar masses, M_{\odot} .

of attraction can be reasonably called a supercluster. The region, if approximated as round, has a diameter of $12,000 \text{ km s}^{-1}$ in units of the cosmic expansion or 160 megaparsecs, and encompasses $\sim 10^{17}$ solar masses. We propose that this region be named the Laniakea supercluster of galaxies (from the Hawaiian; *lani*, heaven, and *akea*, spacious, immeasurable).

Online Content Methods, along with any additional Extended Data display items and Source Data, are available in the online version of the paper; references unique to these sections appear only in the online paper.

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Supplementary Information is available in the online version of the paper.

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Author Contributions R.B.T. guided the project, was involved in the data acquisition, interacted closely in the development of the ideas that are presented here, and wrote most of the Letter. H.C. was involved in the observing programme, was instrumental in coordinating activities, and was involved in all facets. Y.H. took responsibility for the theoretical analysis, including the Wiener filter, the cosmic web and the Malmquist bias correction. D.P. developed and applied visualization tools useful to this research.

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METHODS

The present discussion draws on a new catalogue of galaxy distances and peculiar velocities¹⁹, one that extends to recession velocities of $30,000 \text{ km s}^{-1}$ (redshift $z = 0.1$) and with 8,161 entries provides a density of coverage previously unknown. The new catalogue is called Cosmicflows-2 and the six main methodologies for the distance estimates rely on the characteristics of Cepheid star pulsations, the luminosity terminus of stars at the tip of the red giant branch, surface brightness fluctuations of the ensemble of stars in elliptical galaxies, the standard candle nature of supernovae of type Ia, the adherence by elliptical galaxies to a fundamental plane in luminosity, radius, and velocity dispersion, and the correlation between the luminosities of spirals and their rates of rotation. Each of the methodologies has strengths and weaknesses. The Cepheid and tip of the red giant branch techniques provide high precision distances but only to very nearby galaxies. The elliptical fundamental plane and spiral luminosity–rotation methods provide individually less accurate distances but can be used to acquire samples of thousands of galaxies in a large volume. Type Ia supernovae are excellent distance indicators and can be seen far away but they arise serendipitously and current samples are small. Jointly the sky coverage is now substantial within $\sim 100 \text{ Mpc}$, with only spotty coverage out to $\sim 400 \text{ Mpc}$.

When the Cosmicflows-2 compendium was being compiled, considerable effort went into ensuring that the six independent methodologies were on a common scale. This scale determines the value of the Hubble constant. In Cosmicflows-2, this parameter is determined from the velocities and distances of type Ia supernovae at redshifts $0.03 < z < 0.5$, in a domain where peculiar velocities should be a negligible fraction of observed velocities. We found²⁰ $H_0 = 75.2 \pm 3.0 \text{ km s}^{-1} \text{ Mpc}^{-1}$.

The value of H_0 remains somewhat contentious, especially with the value of $67.3 \pm 1.2 \text{ km s}^{-1} \text{ Mpc}^{-1}$ claimed by the Planck collaboration²¹. Other determinations that actually measure distances and velocities are more compatible with our value^{22,23}. Variations of the assumed H_0 within the Cosmicflows-2 analysis introduce a monopole term: infall if H_0 is increased and outflow if H_0 is decreased. Our value minimizes the monopole term. Tests with the current distances demonstrate unphysical monopole terms if H_0 is varied by more than $\pm 1.5 \text{ km s}^{-1} \text{ Mpc}^{-1}$ from the fiducial value. The Planck determination at face value introduces a massive outflow. It is to be emphasized that, for a study of galaxy flows such as ours, the value assumed for H_0 needs to be internally consistent with the scaling of distance measures. A rescaling, say by a re-evaluation of the distances to the nearest galaxies, leaves peculiar velocities unchanged. The extremely low H_0 value found by the Planck collaboration is only plausible in the face of the monopole implication if distances to nearby calibrator galaxies are significantly greater than currently measured.

The Wiener filter has been used before to reconstruct the underlying flow field from sparse and noisy velocity surveys^{24,25}. Here it is used to overcome the Malmquist bias introduced by errors in distance that scatter galaxies out of regions containing more galaxies towards emptier regions²⁶. This correction is done in two steps. In the first, the Wiener filter is applied to the raw data. New estimated distances are defined by $d_{\text{WF}} = (V_{\text{obs}} - V_{\text{WF}})/H_0$ where V_{WF} is the radial component of the Wiener filter velocity field at the measured position. This adjustment brings the estimated distances to be much closer to their position in redshift space. The Wiener filter is then applied to the revised catalogue to yield the reconstructed flow that is presented here. The Wiener filter correction of the Malmquist bias has been tested against mock catalogues drawn from N -body simulations constrained by the Cosmicflows-2 data to reproduce the actual Universe. The Wiener filter strongly suppresses the spurious (negative) monopole term introduced by the Malmquist bias.

The Wiener filter provides the Bayesian mean field, given the data, its uncertainties and an assumed prior model. The robustness of the Wiener filter reconstruction is gauged by sampling the distribution of the residual from the mean field, assuming that the underlying field is Gaussian²⁷. To meet this requirement, an ensemble of 20 constrained realizations of the underlying flow has been constructed²⁵. A local/tidal decomposition has been applied to a sphere of radius $6,000 \text{ km s}^{-1} = 80 \text{ Mpc}$ centred on the putative centroid of the attractor region ($[-4,700, 1,300, -500] \text{ km s}^{-1}$). Each one of these realizations has been found to harbour a distinct monolithic over-dense structure that dominates that sphere. Extending the radius of the sphere to a radius of $7,000 \text{ km s}^{-1}$ and above we find the breakdown of the monolithic structure and the emergence of more than one attractor within the sphere. The ensemble of constrained realizations can be used to assess the statistical significance of our claim for the existence of a basin of attraction on a scale of $6,000 \text{ km s}^{-1} = 80 \text{ Mpc}$. All of the 20 realizations, each consistent with the data and the assumed prior Λ CDM model, reproduce such an attractor.

A superiority of the recovery of structure from peculiar velocities with the Wiener filter over reconstructions from redshift surveys is the sensitivity to structures in regions that cannot be observed, in particular in the heavily obscured equatorial band of the plane of our Galaxy. Indeed a large part of our watershed is hidden by obscuration. The connections between the important Pavo-Indus and Centaurus

components of Laniakea, the extension up to Ophiuchus cluster, the filament from the Antlia Wall that ultimately connects to the Perseus–Pisces complex, and the continuation of the Perseus–Pisces filament beyond Perseus are all examples of hidden structures manifested in flow patterns. The two dominant attractors in Shapley and Lepus are both at low galactic latitudes. A further nice example is the identification of the previously unknown Arch.

The Supplementary Video and the figures presented in this paper were created using SDvision²⁸, an interactive visualization software developed within the framework of IDL Object Graphics. In Supplementary Information we give the transcript of the dialogue in versions that include a sound track. The transcript is also provided in the closed caption version (transcript of the dialogue embedded in the video).

The Supplementary Video begins with a half-rotation display of galaxies from the 2MASS Extended Source Catalog²⁹ with distances given by measured velocities assuming the Hubble law. This display gives a general idea of the structure in the nearby Universe although there is increasing incompleteness with distance from our position at the centre and there is some inhomogeneity in coverage with direction.

During a second half rotation of the scene, there is a display of the peculiar velocity vectors that are fundamental to the current analysis. Peculiar motions towards us are in blue and those away are in red. Measurements are dense within $7,500 \text{ km s}^{-1} = 100 \text{ Mpc}$ and fall off quickly beyond that distance.

The rotations continue, but now with the full velocity flows derived from the Wiener filter analysis. The immediate product of the filtering is a most probable three-dimensional peculiar velocity at each position within the volume of the study. An imaginary vehicle that starts from any seed location is passed to an adjacent position by the locally inferred vector of motion, then again passed onward, creating flows that are captured by the visualized stream lines. The dominant flow is towards the Shapley concentration of galaxies. A second feature involves an accumulation of flows in the Perseus–Pisces region that then streams down to the Lepus region. After a rotation, the implied large-scale potential field is superimposed. If the viewing box is extended, it is reasonably convincing that the flow to Lepus continues on to connect with Shapley. In other words, essentially the entire volume being displayed is involved in a flow towards Shapley. This proposition will be explored elsewhere.

The emphasis of the present discussion is on intermediate scales that are within the domain of most of the distance measurements. The contours that become superimposed on the flow lines represent the density field derived from the Wiener filter velocity field. The amplitudes of the density reconstructions are high near the centre of the cube where data constraints dominate errors, and taper to low values at the edges where uncertainties are large and the derived density field approaches the mean field. Extended Data Fig. 1 is extracted from this sequence.

After some more rotations, the flow lines are seen to break up into disjoint pieces. These disconnections occur because we have begun to use the trick of separating into local and tidal components. Local velocity flows are due to the mass distribution in a restricted region and tidal flows are due to mass outside the restricted region.

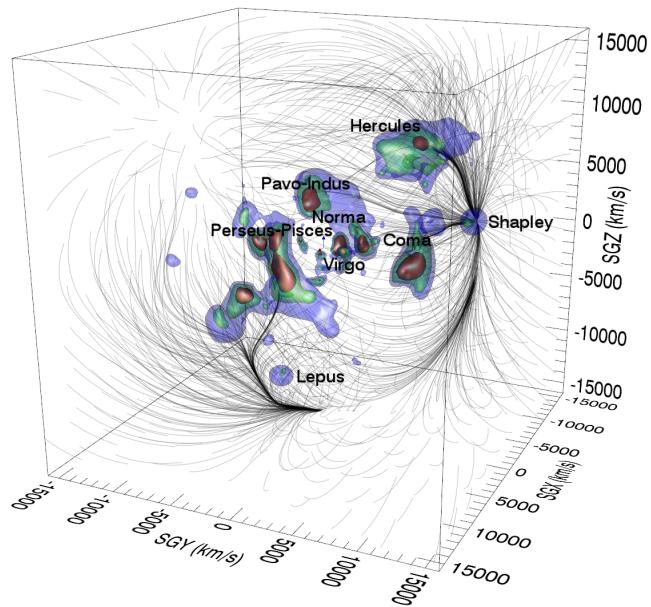
Next there is a transition to the V-web representation (see main text). The eigenvalues of the velocity shear can be evaluated at each local position. Surfaces are shown that enclose clusters, filaments, and sheets. Now in the video there is a zoom inward and flow lines are shown for local motions within a zone extending to $6,000 \text{ km s}^{-1} = 80 \text{ Mpc}$ from our central position. We are sampling two separate basins of attraction. The red flow lines are directed toward the Perseus–Pisces complex while the black flow lines are directed toward the Centaurus–Norma–Pavo–Indus structures. Our Earth lies within the domain of the black flow lines, although we are near the boundary where local flow directions flip. It is seen that the flows converge to follow along filaments. There are interesting places along the V-web filaments where local peculiar velocities flip in direction. After some time in this video sequence there is the superposition of individual galaxies from the redshift catalogue. Colours are given to members of different structures. The positions of the individual galaxies are not corrected for redshift distortion but reasonable agreement is seen between the locations of galaxies, the flow lines, and the inferred V-web structures. Extended Data Fig. 2 is extracted from the video as an illustration of this point.

In the concluding section of the Supplementary Video, the local/tidal decomposition is shifted to a centre near to the Norma cluster in order to isolate our basin of attraction. The flow lines now only include those that are moving inward. A surface is created at the boundary of this region. Figure 1 shows two projections from Supplementary Video frames and Fig. 2, the concluding frame of the Supplementary Video, is a projection onto the equatorial plane in supergalactic coordinates. This latter figure is augmented in Extended Data Fig. 3 to include flow lines away from, as well as towards, the local region of attraction. In addition there is an indication in Extended Data Fig. 3 of the swath of the obscuration in the plane of the Milky Way. Extended Data Figs 4 and 5 give two orthogonal views of the inward and outward local flows.

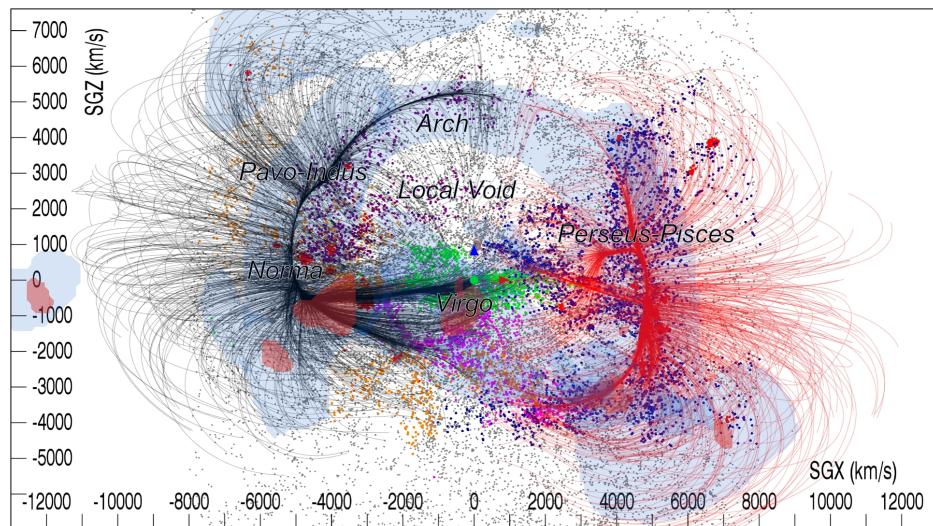
The orange outline in Fig. 2 and in Extended Data Figs 3–5 is the projected boundary of our local basin of attraction that we call the Laniakea supercluster. This domain

of diameter 160 Mpc, given the mean density of the Universe, encloses $10^{17} M_{\odot}$. We stress that in the fullness of distance measurements on a much larger scale it will almost certainly be demonstrated that Laniakea is not at rest with respect to the cosmic expansion and is only a part of something very large that is at rest in the cosmic reference frame. We remind those unfamiliar with the field that cosmic expansion velocities are removed in deriving peculiar motions. Infalling motions on large scales are only perturbations. All galaxies except those in the immediate vicinity of clusters and groups are flying apart.

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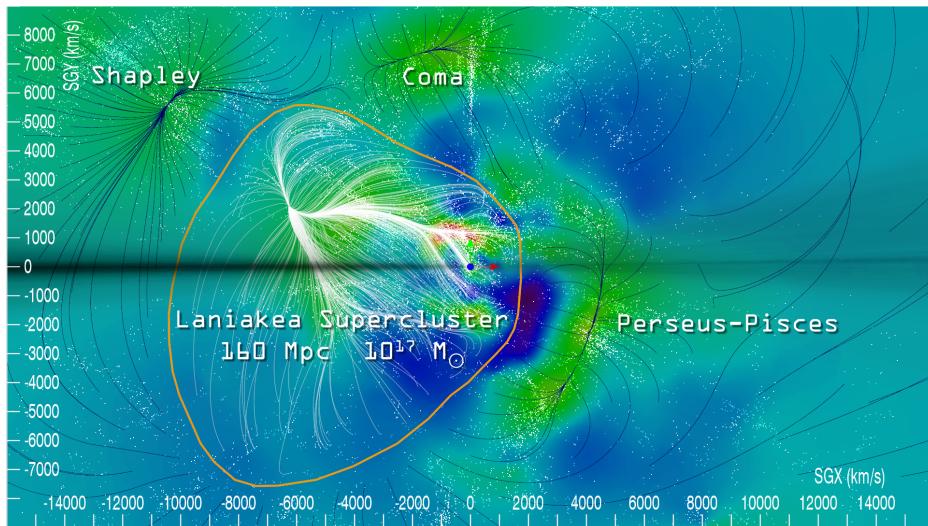


Extended Data Figure 1 | Structure within a cube extending $16,000 \text{ km s}^{-1}$ ($\sim 200 \text{ Mpc}$) on the cardinal axes from our position at the origin. Densities on a grid within the volume are determined from a Wiener filter reconstruction based on the observed velocity field. Three isodensity contours are shown. The density map is detailed near the centre of the box where observational constraints are dense and accurate, but tapers to the mean density as constraints weaken. Nevertheless, velocity flows illustrated by the black threads are defined on large scales. Ultimately all flows appear to drain towards Shapley, although flows through the Perseus–Pisces filament take a circuitous route through the poorly studied Lepus region.



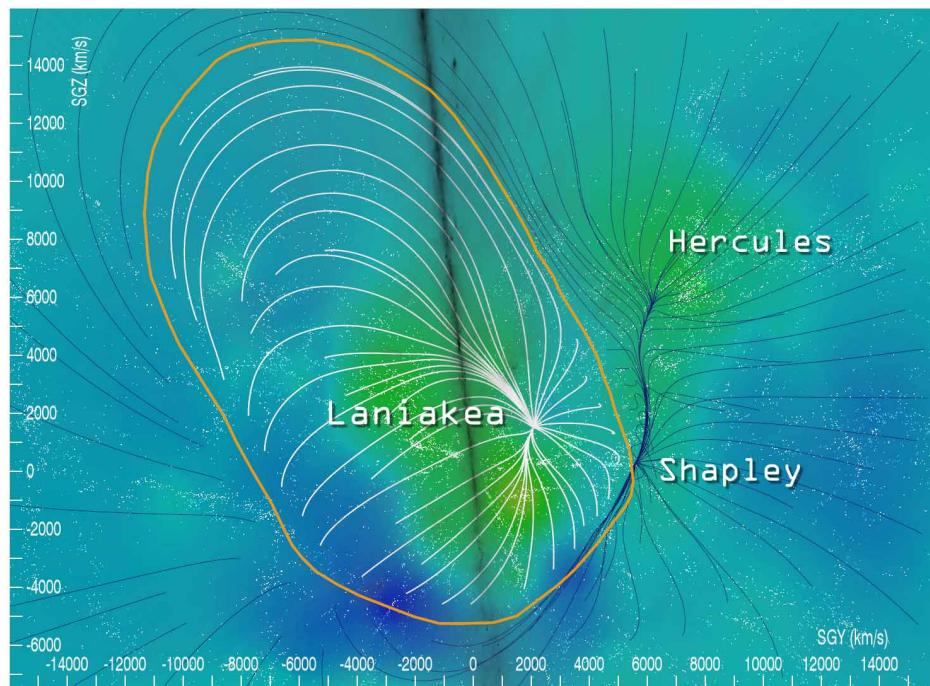
Extended Data Figure 2 | A representation of structure and flows due to mass within $6,000 \text{ km s}^{-1}$ ($\sim 80 \text{ Mpc}$). Surfaces of red and blue respectively represent outer contours of clusters and filaments as defined by the local eigenvalues of the velocity shear tensor determined from the Wiener filter analysis. Flow threads originating in our basin of attraction that terminate near

the Norma cluster are in black, and adjacent flow threads that terminate at the relative attractor near the Perseus cluster are in red. The Arch and extended Antlia Wall structures bridge between the two attraction basins. The arch is part of a wall surrounding the Local Void.



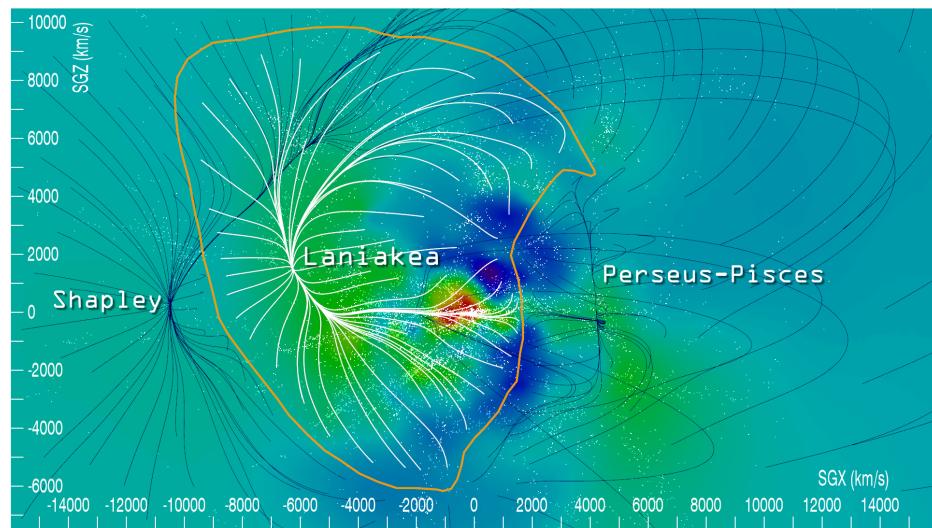
Extended Data Figure 3 | One of three orthogonal views that illustrate the limits of the Laniakea supercluster. This SGX-SGY view at SGZ = 0 extends the scene shown in Fig. 2 with the addition of dark blue flow lines away from the Laniakea local basin of attraction, and also includes a dark swath at

SGY = 0 showing the region obscured by the plane of the Milky Way. As in Fig 2 in the main text, the orange contour encloses the inflowing streams, hence, defines the limits of the Laniakea supercluster containing the mass of 10^{17} Suns and 100,000 large galaxies.



Extended Data Figure 4 | The second orthogonal view that illustrates the limits of the Laniakea supercluster. This is an SGY–SGZ slice at

$\text{SGX} = -4,750 \text{ km s}^{-1}$: the dark swath indicates the region obscured by the Milky Way.



Extended Data Figure 5 | The third orthogonal view that illustrates the limits of the Laniakea supercluster. Here we show an SGX–SGZ slice at

$\text{SGY} = +1,000 \text{ km s}^{-1}$. The Milky Way lies essentially in a plane parallel to the slice, at $\text{SGY} = 0$.