

Efficiency of transverse forces in dense systems

Federico Ghimenti¹, Ludovic Berthier², Grzegorz Szamel³ and Frédéric van Wijland¹

¹ Laboratoire Matière et Systèmes Complexes, Université Paris Cité and CNRS (UMR 7057), Paris, France

² Laboratoire Charles Coulomb (L2C), Université de Montpellier and CNRS (UMR 5221), Montpellier, France

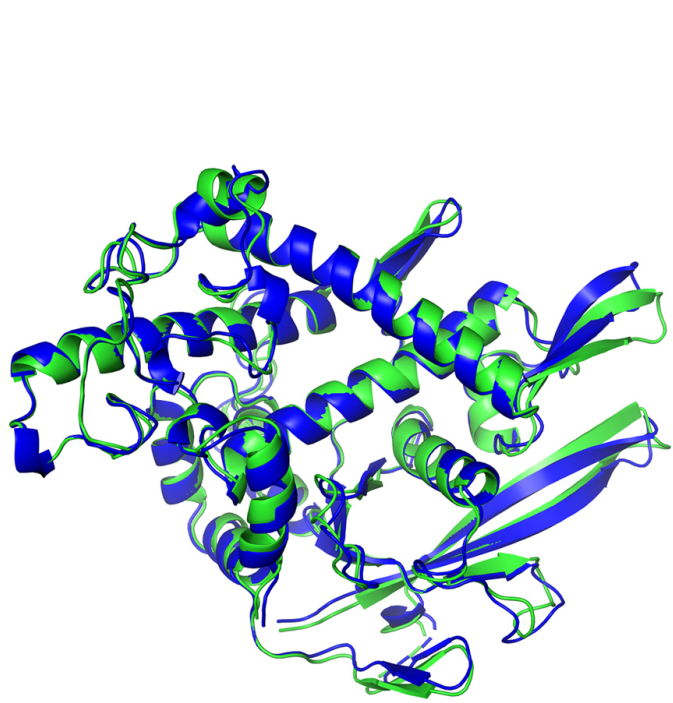
³ Department of Chemistry, Colorado State University, Fort Collins, Colorado, USA

federico.ghimenti@u-paris.fr

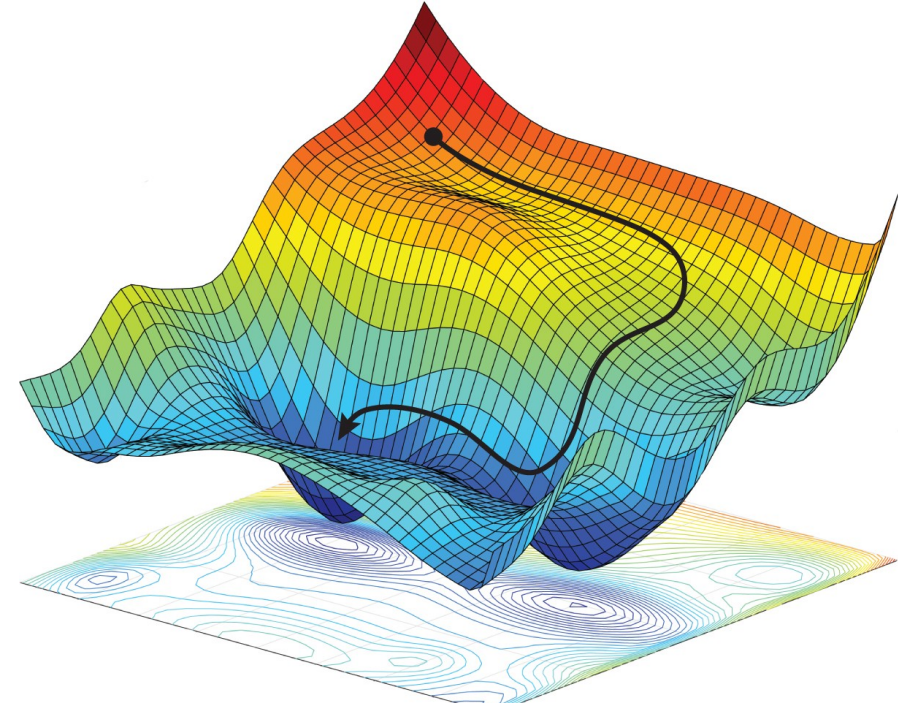


1. Sampling Boltzmann

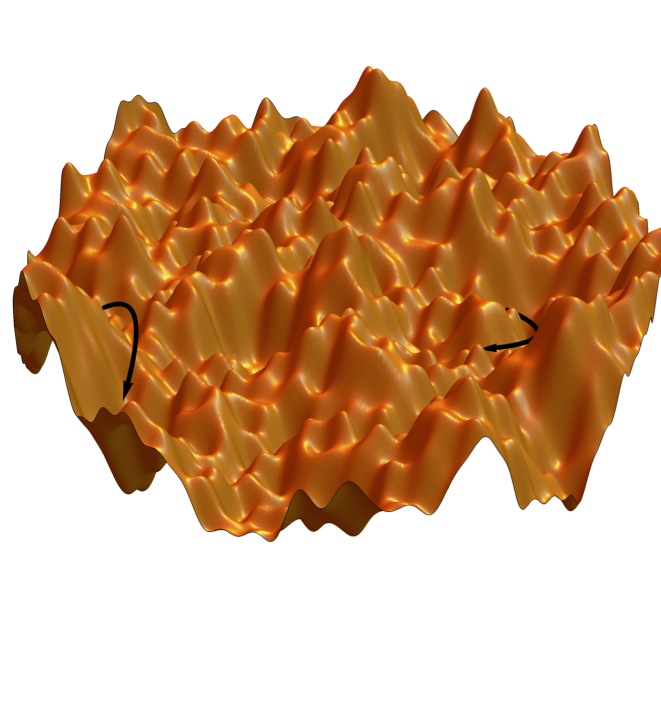
The goal: $\rho_B \propto e^{-\beta V}$
Hard in some relevant cases



DeepMind



Amini, NIPS 2017



C. Cammarota

Equilibrium sampling

$$\dot{\mathbf{r}} = -\nabla V(\mathbf{r}) + \sqrt{2T}\boldsymbol{\xi}$$

$$\langle \xi_i(t) \xi_j(t') \rangle = \delta_{ij} \delta(t - t')$$

$$\partial_t \rho(\mathbf{r}, t) = -\nabla \cdot \mathbf{j}(\mathbf{r}, t)$$

$$\rho(\mathbf{r}, t) \approx \rho_B(\mathbf{r}) + e^{-\frac{t}{\tau_R}} \phi_1(\mathbf{r})$$

τ_R can become very large ... new dynamical rules?

2. Transverse Forces

$$\dot{\mathbf{r}} = -(\mathbb{1} + \gamma \mathbf{A}) \nabla V(\mathbf{r}) + \sqrt{2T} \boldsymbol{\xi}$$

$$\mathbf{A}^T = -\mathbf{A}, \mathbf{A} \sim O(1)$$

Extra current $\partial_t \rho(\mathbf{r}, t) = -\nabla \cdot (\mathbf{j}(\mathbf{r}, t) + \mathbf{j}_\gamma(\mathbf{r}, t)).$

Sampling Boltzmann: $\nabla \cdot \mathbf{j}_\gamma(\mathbf{r}, t) = 0$ when $\rho(\mathbf{r}, t) = \rho_{ss}(\mathbf{r}) = \rho_B$

Out of equilibrium: Entropy production rate $\dot{\Sigma} = \frac{\gamma^2}{T} \langle (\nabla V)^2 \rangle > 0$

A theorem on dynamics' acceleration [1]

$$\tau_{R,irr} \leq \tau_{R,rev}$$

How big is the speedup in dense fluids?

3. Connection with other schemes

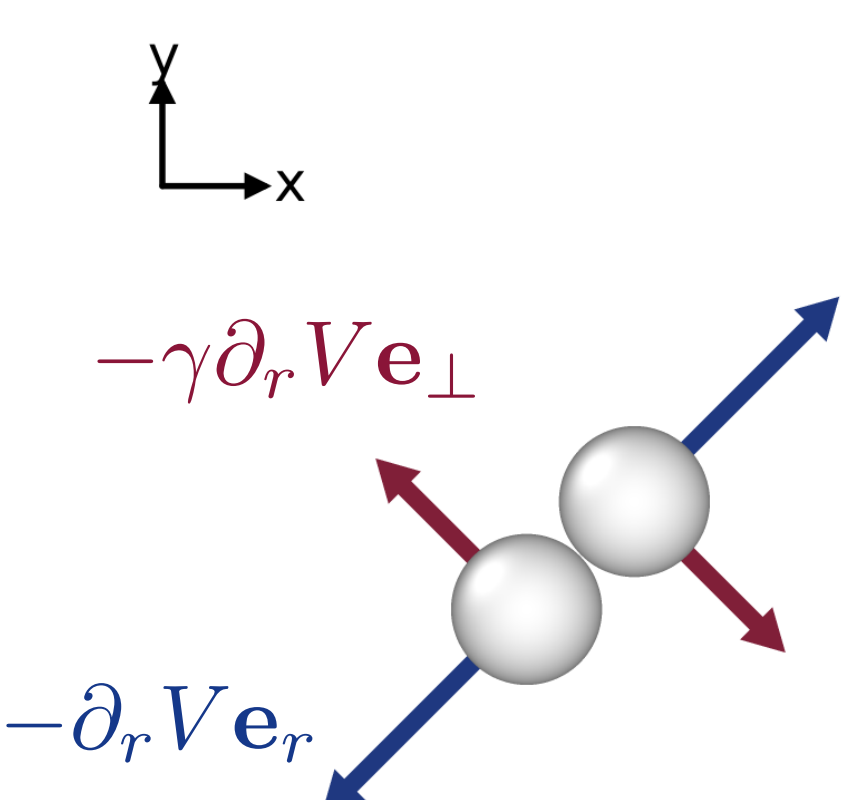
Ichiki-Ohzeki [2] Two copies of the same system

$$\mathbf{r} = (\mathbf{r}^{(1)}, \mathbf{r}^{(2)}), U(\mathbf{r}) = V(\mathbf{r}^{(1)}) + V(\mathbf{r}^{(2)})$$

$$\begin{bmatrix} \dot{\mathbf{r}}^{(1)} \\ \dot{\mathbf{r}}^{(2)} \end{bmatrix} = - \begin{bmatrix} \mathbb{1} & -\gamma \mathbb{1} \\ \gamma \mathbb{1} & \mathbb{1} \end{bmatrix} \cdot \begin{bmatrix} \nabla_{\mathbf{r}^{(1)}} U \\ \nabla_{\mathbf{r}^{(2)}} U \end{bmatrix} + \sqrt{2T} \begin{bmatrix} \boldsymbol{\xi}^{(1)} \\ \boldsymbol{\xi}^{(2)} \end{bmatrix}$$

- One can couple any two arbitrary systems
- An extended phase space with irreversible dynamics \rightarrow **lifting** [3]

4. Transverse forces in dense liquids



$$\frac{d\mathbf{r}_i}{dt} = \mathbf{F}_i + \sqrt{2T} \boldsymbol{\xi}_i$$

$$\mathbf{F}_i = -(\mathbf{1} + \gamma \mathbf{A}) \sum_{j \neq i} \nabla_{\mathbf{r}_i} V(|\mathbf{r}_i - \mathbf{r}_j|)$$

- Numerical simulations of structural glass-former
- Infinite dimensions
- Mode Coupling Theory

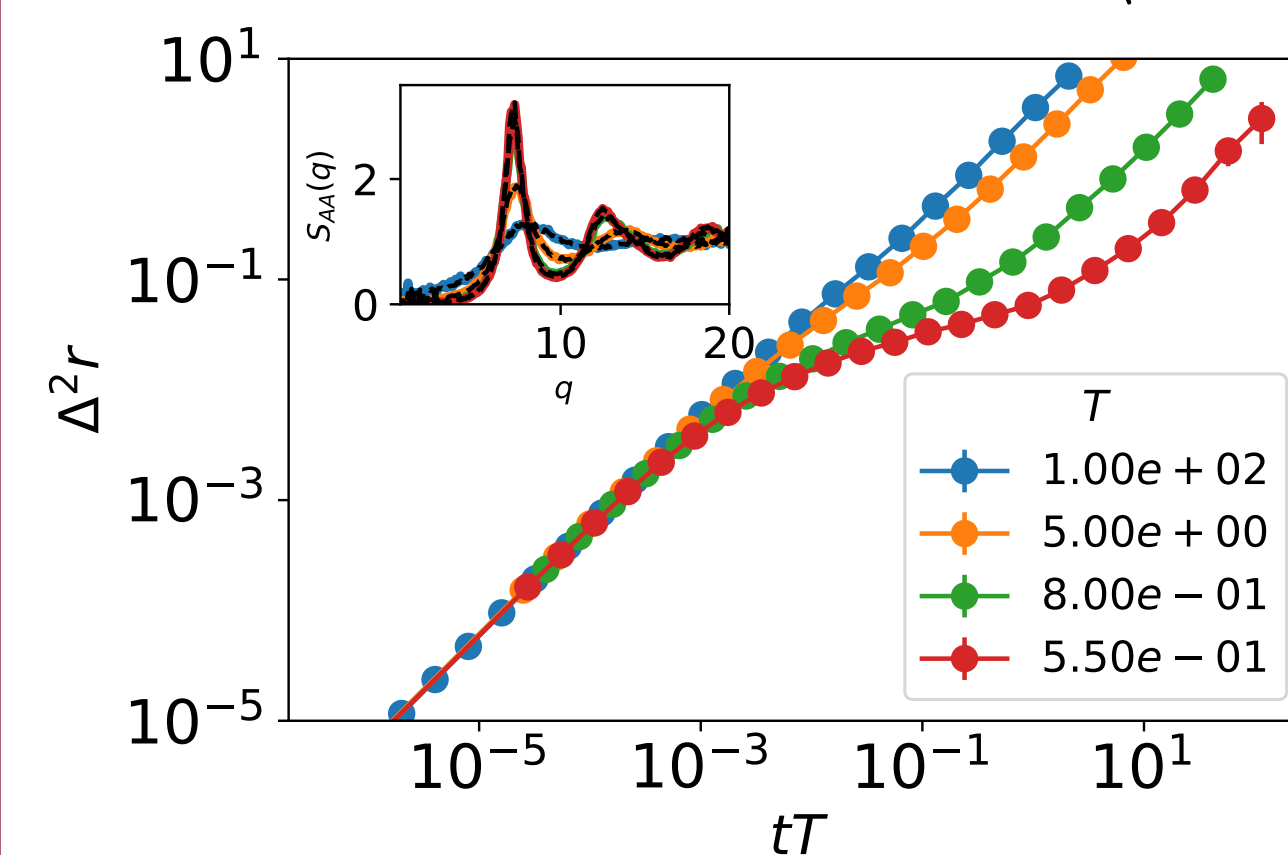
5. Numerics - Speedup

Brownian Dynamics simulations of a three dimensional 80:20 Kob-Andersen mixture [4, 5], $\mathbf{A} = \begin{bmatrix} 0 & -1 & 0 \\ 1 & 0 & 0 \\ 0 & 0 & 0 \end{bmatrix}$

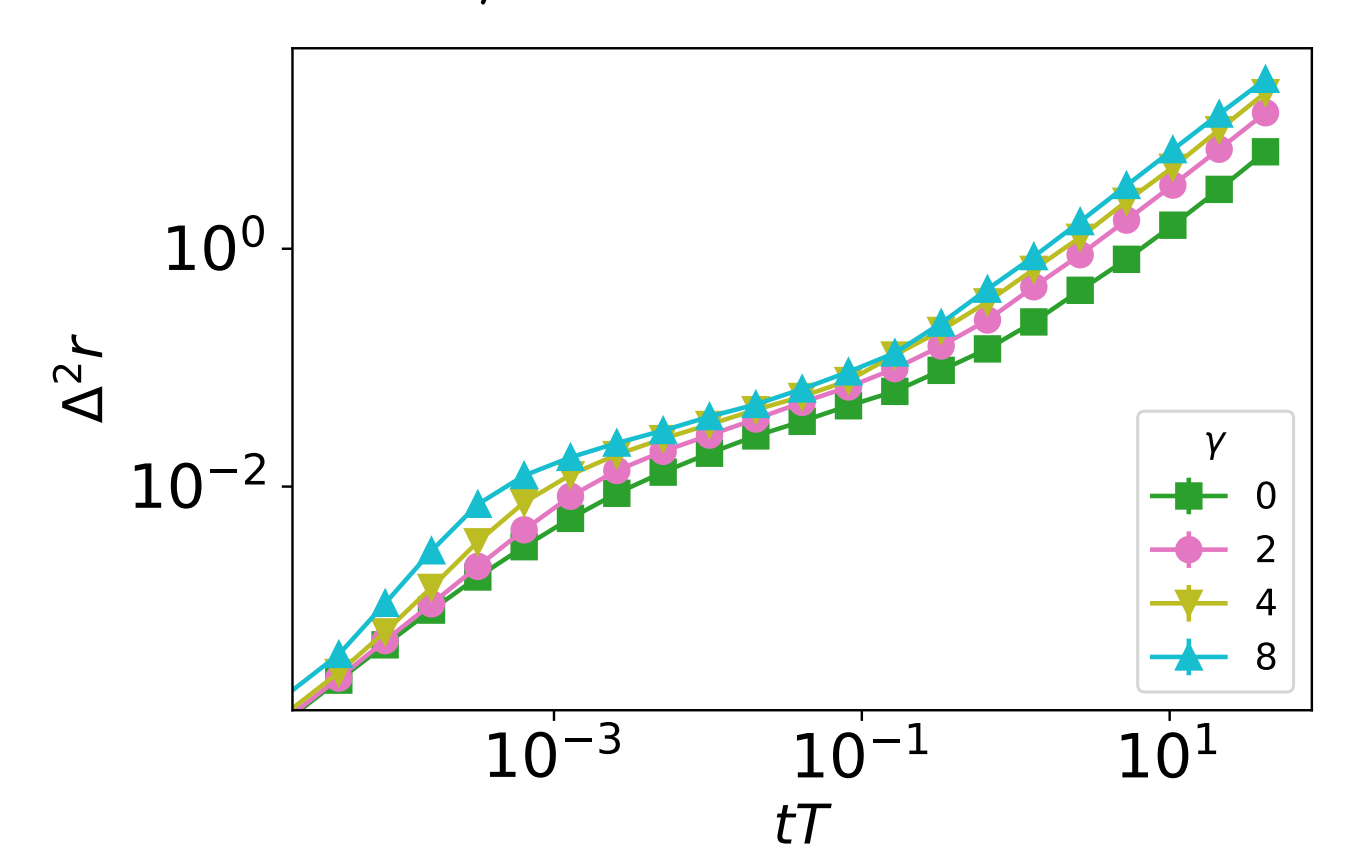
For $\gamma = 0$, $T \sim 1$ onset of slow dynamics

Structure: $S_{AA}(q) \equiv \frac{1}{N_A} \left\langle \sum_{i,j}^{N_A} e^{i\mathbf{q} \cdot (\mathbf{r}_i - \mathbf{r}_j)} \right\rangle$

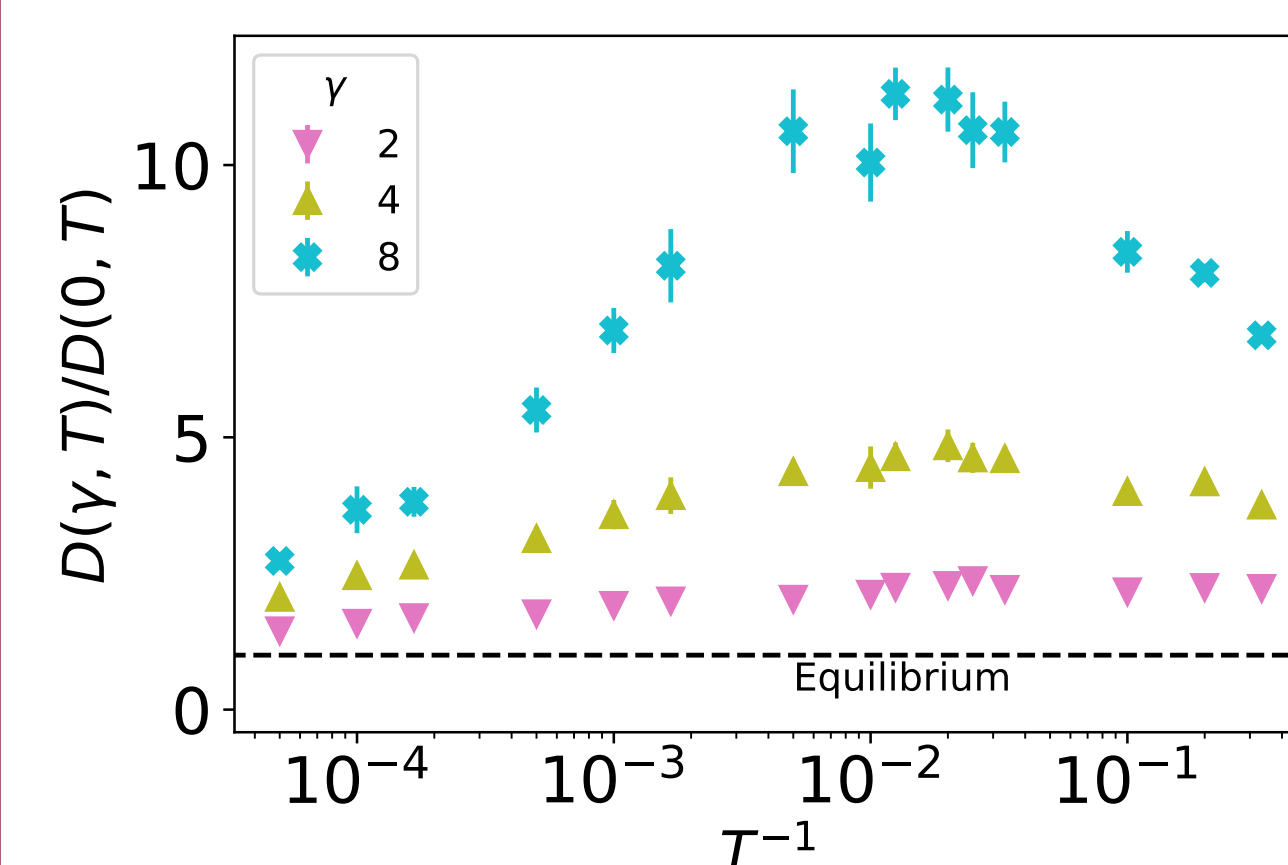
Dynamics: $\Delta^2 r(t) \equiv \frac{1}{N_A} \left\langle \sum_i^{N_A} [\mathbf{r}_i(t) - \mathbf{r}_i(0)]^2 \right\rangle \stackrel{t \rightarrow \infty}{\approx} 6D(T, \gamma)t$



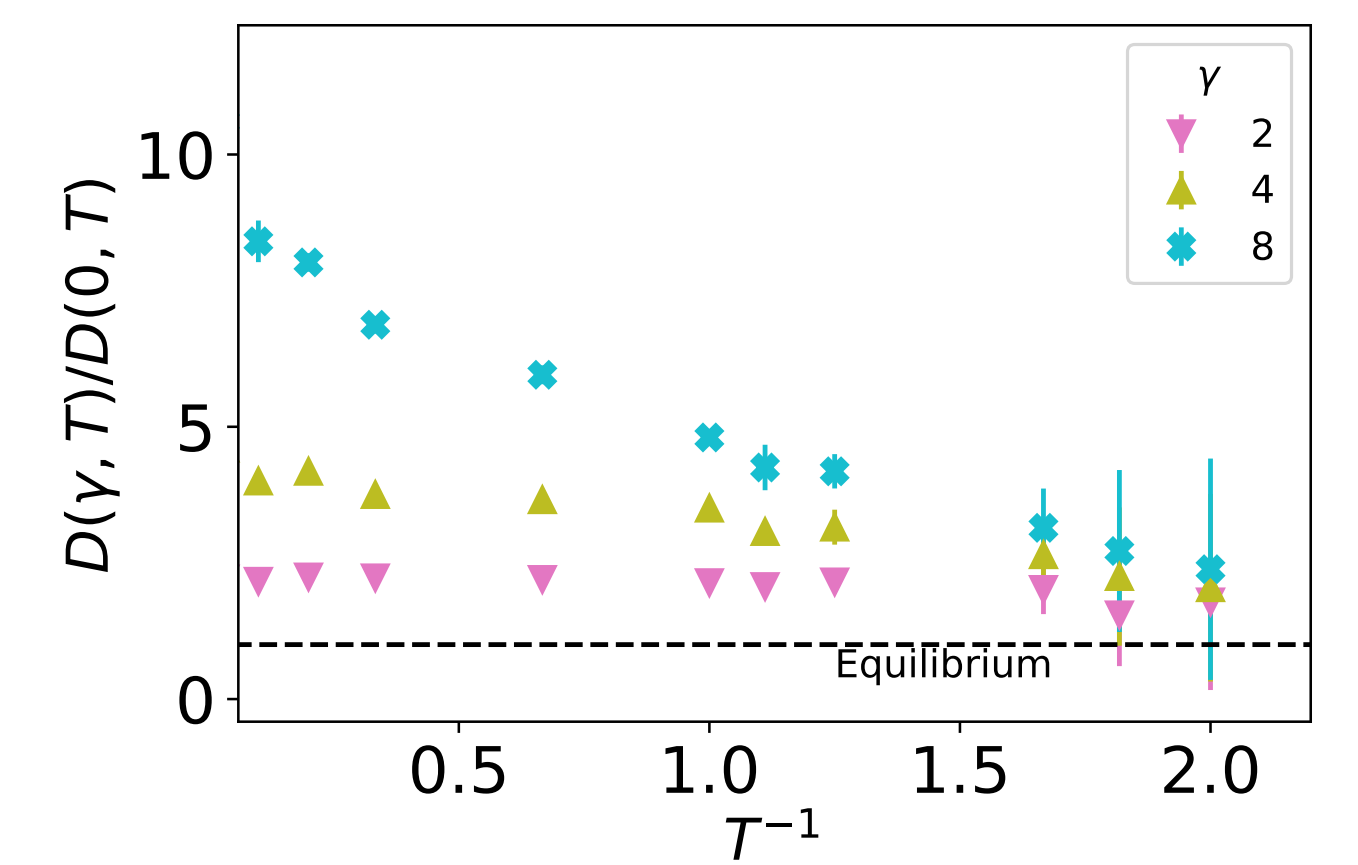
Calibration



A first glance at the speedup, T=0.8

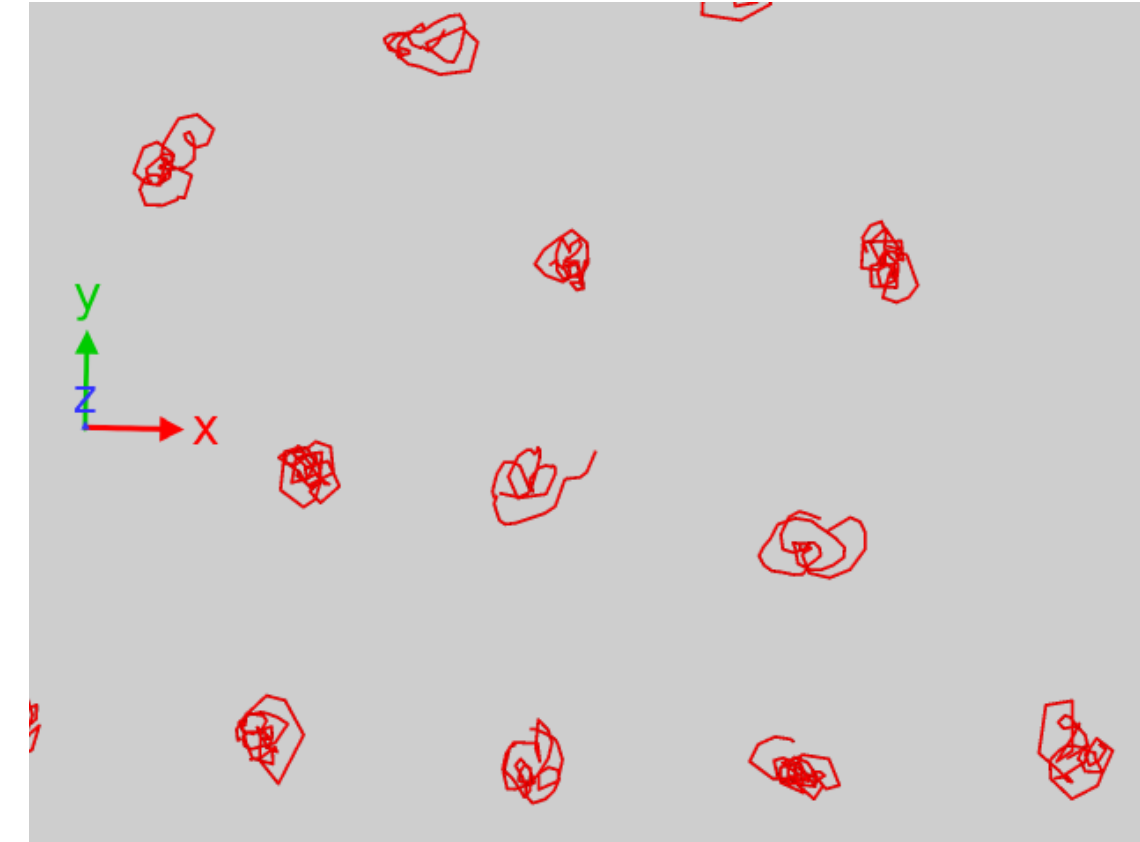


Maximal efficiency at **high T**



Efficiency **decreases** in the glass

6. Numerics - Odd transport



short time trajectories, $\gamma = 200$, $T = 0.8$

Motivation: odd diffusion helps relaxation [6]

$$D_{\perp} \equiv \frac{1}{2N_A} \sum_i^{N_A} \int_0^{+\infty} dt \langle \dot{\mathbf{r}}_i(t) \cdot \mathbf{A} \dot{\mathbf{r}}_i(0) \rangle$$

As T decreases, $D_{\perp} \rightarrow c(\gamma)$
Swirling without relaxing

7. Infinite dimensions

$$D(\gamma, T) = T \frac{1 + (1 + \gamma^2) \beta \widehat{M}}{(1 + \beta \widehat{M})^2 + (\gamma \beta \widehat{M})^2}$$

$$D_{\perp}(\gamma, T) = -\gamma T \frac{\beta \widehat{M} + (1 + \gamma^2)(\beta \widehat{M})^2}{(1 + \beta \widehat{M})^2 + (\gamma \beta \widehat{M})^2}$$

$$\widehat{M} = \frac{1}{d} \int_0^{+\infty} \langle \nabla V(t) \cdot \nabla V(0) \rangle, \widehat{M} \rightarrow \infty \text{ for } T \rightarrow T_d(\gamma) = T_d(0)$$

- $\mathbf{T} \rightarrow +\infty$: $D(\gamma, T)/D(0, T) = 1$, $D_{\perp}(\gamma, T) = 0$
- $\mathbf{T} \rightarrow \mathbf{T}_d$: $D(\gamma, T)/D(0, T) = a(\gamma)$, $D_{\perp}(\gamma, T) = -\gamma T$

Similar results for Mode Coupling Theory

8. Outlook

- Efficiency decreases when sampling gets harder
- Need for currents informed about the relevant degrees of freedom
- Monte Carlo schemes?

[1] C. Hwang, S. Hwang-Ma, and S. Sheu. *Ann. App. Prob.*, 1993.

[2] M. Ohzeki and A. Ichiki. *PRE*, 92(1):012105, 2015.

[3] M. Vucelja. *Am. J. Phys.*, 84(12), 2016.

[4] W. Kob and H. C Andersen. *PRE*, 51(5):4626, 1995.

[5] W. Kob and H. C Andersen. *PRE*, 52(4):4134, 1995.

[6] E. Kalz *et al.* *PRL*, 129(9):090601, 2022.