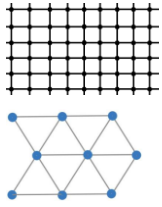


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## Model



Typical Hamiltonian:  $H = J \sum_{\langle i, j \rangle} (\sigma_i \sigma_j)$

where  $\langle i, j \rangle$  denotes neighboring sites

Scalar and mixed products of vectors of Pauli matrices:

$$(\sigma_1 \sigma_3) = \delta^{\alpha\beta} \cdot \sigma_1^\alpha \otimes 1_2 \otimes \sigma_3^\beta \otimes 1_4 \otimes 1_5$$

$$(\sigma_1 \sigma_3 \sigma_4) = \varepsilon^{\alpha\beta\gamma} \cdot \sigma_1^\alpha \otimes 1_2 \otimes \sigma_3^\beta \otimes \sigma_4^\gamma \otimes 1_5$$

$$\alpha, \beta, \gamma \in \{x, y, z\}$$

## Squaring parameterization of density matrix [1]

Common conditions:  $\rho^+ = \rho$ ;  $\text{Tr} \rho = 1$ ;  $\langle \rho \rangle \geq 0$

From  $\rho$  to  $\tau$ :  $\rho = \frac{\tau^2}{\text{Tr} \tau^2}$ ; for  $\tau$  only one condition:  $\tau^+ = \tau$

$$\rho = \frac{1}{2^n} a_i A_i; \quad \tau = b_i A_i$$

Basis:

$$\{A_i\} = \{1, (\sigma_j \sigma_k), (\sigma_j \sigma_k)(\sigma_l \sigma_m), \dots\}$$

Basis overflow:

$$+(\sigma_1 \sigma_2)(\sigma_3 \sigma_4 \sigma_5) - (\sigma_1 \sigma_3)(\sigma_2 \sigma_4 \sigma_5) + (\sigma_1 \sigma_4)(\sigma_2 \sigma_3 \sigma_5) - (\sigma_1 \sigma_5)(\sigma_2 \sigma_3 \sigma_4) = 0$$

## Symbolic relations and simplifications

• Algorithm which based on identity of Pauli ( $\sigma^\alpha \sigma^\beta = \delta^{\alpha\beta} + i \varepsilon^{\alpha\beta\gamma} \sigma^\gamma$ ) was implemented in Wolfram Mathematica

$$(\sigma_1 \sigma_2)^2 = 3 - 2(\sigma_1 \sigma_2)$$

$$(\sigma_1 \sigma_2)(\sigma_2 \sigma_3) = (\sigma_1 \sigma_3) - i(\sigma_1 \sigma_2 \sigma_3)$$

$$(\sigma_1 \sigma_2)(\sigma_1 \sigma_3 \sigma_4) = -(\sigma_1 \sigma_2 \sigma_3) - 2i(\sigma_1 \sigma_3) + 2i(\sigma_2 \sigma_3)$$

$$(\sigma_1 \sigma_2 \sigma_3)(\sigma_1 \sigma_2) = -(\sigma_1 \sigma_2 \sigma_3) + 2i(\sigma_1 \sigma_3) - 2i(\sigma_2 \sigma_3)$$

$$(\sigma_1 \sigma_2)(\sigma_2 \sigma_3 \sigma_4) = (\sigma_1 \sigma_3 \sigma_4) - i(\sigma_1 \sigma_3)(\sigma_2 \sigma_4) + i(\sigma_1 \sigma_4)(\sigma_2 \sigma_3)$$

$$(\sigma_2 \sigma_3 \sigma_4)(\sigma_1 \sigma_2) = (\sigma_1 \sigma_3 \sigma_4) + i(\sigma_1 \sigma_3)(\sigma_2 \sigma_4) - i(\sigma_1 \sigma_4)(\sigma_2 \sigma_3)$$

$$(\sigma_1 \sigma_2 \sigma_3)^2 = 6 - 2(\sigma_1 \sigma_2) - 2(\sigma_1 \sigma_3) - 2(\sigma_2 \sigma_3)$$

$$(\sigma_1 \sigma_2 \sigma_3)(\sigma_1 \sigma_2 \sigma_4) = -(\sigma_1 \sigma_3)(\sigma_2 \sigma_4) - (\sigma_1 \sigma_4)(\sigma_2 \sigma_3) + 2(\sigma_3 \sigma_4) + i(\sigma_1 \sigma_3 \sigma_4) + i(\sigma_2 \sigma_3 \sigma_4)$$

$$(\sigma_1 \sigma_2 \sigma_3)(\sigma_1 \sigma_4 \sigma_5) = +(\sigma_2 \sigma_4)(\sigma_3 \sigma_5) - (\sigma_2 \sigma_5)(\sigma_3 \sigma_4) - i(\sigma_1 \sigma_2)(\sigma_3 \sigma_4 \sigma_5) + i(\sigma_1 \sigma_3)(\sigma_2 \sigma_4 \sigma_5)$$

• Scalar product of matrices also was implemented in Wolfram Mathematica:

$$(A, B) = \text{Tr}(A^+ B) = \text{Tr}(AB); \quad A^+ = A$$

$$(A, B) = 0, \text{ if } A \text{ and } B \text{ contains different spins}$$

$$\text{Tr}((\sigma_i \sigma_j)(\sigma_j \sigma_k) \dots (\sigma_l \sigma_m)(\sigma_m \sigma_i)) = 3 \cdot 2^n$$

$$\text{Tr}((\sigma_i \sigma_j \sigma_k)(\sigma_i \sigma_j \sigma_n)(\sigma_k \sigma_l) \dots (\sigma_m \sigma_n)) =$$

$$= \text{Tr}((\sigma_i \sigma_j \sigma_k)(\sigma_i \sigma_j \sigma_k)) = 6 \cdot 2^n$$

$$(A, B) = 3^C \cdot 2^n, \text{ where } C - \text{number of cycles}$$

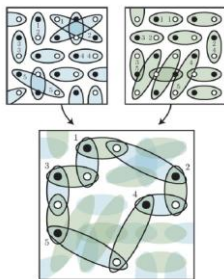

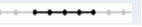





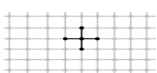
Illustration from [6]

## Results

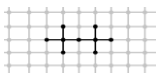
1d:

cluster	$E_{gs}/N$
	-2.1547
	-1.92789
	-1.99486
	-1.89083
	-1.92853

2d:



$$E_{gs}/N > 3$$



$$E_{gs}/N > -2,9685$$

1d exact theoretical [7]:  $E_{gs}/N = 1 - 4 \ln 2 = -1.7725$

## Challenge and main ideas

$$??? \leq E_{gs} \leq \langle \psi | H | \psi \rangle$$

Our challenge

well known variational method

Divide **full** crystal to **clusters**:

N – number of spins in crystal

M – number of clusters

d - number of bonds by 1 spin

m – number of bonds in cluster

$$M = \frac{Nd}{m}$$

$$\bullet \text{Schrödinger equation: } H_{full} = \sum_{i=1}^M H_{icl} \Rightarrow E_{gs, full} \geq \sum_{i=1}^M E_{i, gs, cl}$$

$$E_{gs, full} / N \geq \frac{d}{m} E_{gs, cl}$$

• Variational method:

$$E_{gs, full} = \min_{\psi} \langle \psi | H | \psi \rangle = \min_{\rho_{full} \in M_{full}^S} \text{Tr}_{full} H_{full} \rho_{full} = \min_{\rho_{full} \in M_{full}^S} \text{Tr}_{full} H_{full} \rho_{full} =$$

$$= M \min_{\rho_{full} \in M_{full}^S} \text{Tr}_{full} H_{cl} \rho_{full} = M \min_{\rho_{full} \in M_{full}^S} \text{Tr}_{cl} (H_{cl} \text{Tr}_{full-cl} \rho_{full}) =$$

$$= M \min_{\rho_{cl} \in M_{cl}^{trS}} \text{Tr}_{cl} H_{cl} \rho_{cl} \geq M \min_{\rho_{cl} \in M_{cl}'} \text{Tr}_{cl} H_{cl} \rho_{cl}$$

$$E_{gs, full} / N \geq \frac{d}{m} \min_{\rho_{cl} \in M_{cl}'} \text{Tr}_{cl} H_{cl} \rho_{cl};$$

## Schrödinger equation

$$H\rho = E\rho \Leftrightarrow H \frac{\tau^2}{\text{Tr} \tau^2} = E \frac{\tau^2}{\text{Tr} \tau^2} \Leftrightarrow H\tau = E\tau$$

$$HA_j a_j = A_i h_{ij} a_j + B_i d_{ij} a_j = EA_j a_j \Rightarrow \begin{cases} h_{ij} a_j = Ea_i \\ d_{ij} a_j = 0 \end{cases}$$

$$(A_i, A_j) = g_{ij}$$

$$\text{Tr}(A_i HA_j) = \eta_{ij} \quad h_{ij} = g_{ik}^{-1} \eta_{kj}$$

We take into account the symmetries of commutation of spins in cluster.

Example:

$$A_0 = 1$$

$$A_1 = (\sigma_1 \sigma_2) + (\sigma_2 \sigma_3) + (\sigma_3 \sigma_4) + (\sigma_4 \sigma_5)$$

$$A_2 = (\sigma_1 \sigma_2) + (\sigma_2 \sigma_3) + (\sigma_3 \sigma_4) + (\sigma_4 \sigma_5) + (\sigma_1 \sigma_3) + (\sigma_2 \sigma_4)$$

$$A_3 = (\sigma_1 \sigma_2)(\sigma_2 \sigma_3) + (\sigma_2 \sigma_3)(\sigma_3 \sigma_4) + (\sigma_3 \sigma_4)(\sigma_4 \sigma_5) + (\sigma_1 \sigma_3)(\sigma_3 \sigma_4)$$

$$+ (\sigma_2 \sigma_3)(\sigma_1 \sigma_3) + (\sigma_2 \sigma_3)(\sigma_1 \sigma_4) + (\sigma_2 \sigma_3)(\sigma_3 \sigma_4)$$

$$+ (\sigma_3 \sigma_4)(\sigma_1 \sigma_3) + (\sigma_3 \sigma_4)(\sigma_1 \sigma_4) + (\sigma_3 \sigma_4)(\sigma_3 \sigma_4)$$

$$+ (\sigma_4 \sigma_5)(\sigma_1 \sigma_3) + (\sigma_4 \sigma_5)(\sigma_1 \sigma_4) + (\sigma_4 \sigma_5)(\sigma_3 \sigma_4)$$

$$A_4 = (\sigma_1 \sigma_2)(\sigma_3 \sigma_4) + (\sigma_1 \sigma_3)(\sigma_2 \sigma_4) + (\sigma_1 \sigma_4)(\sigma_2 \sigma_3)$$

$$h_{ij} = \begin{pmatrix} 0 & 12 & 0 & 0 & 0 \\ 1 & -2 & 3 & 0 & 0 \\ 0 & 2 & 0 & 10 & 0 \\ 0 & 0 & 1 & 2 & 1 \\ 0 & 0 & 0 & 4 & 0 \end{pmatrix}$$

$$d_{ij} = 0$$

$$E_{gs} = -6, -4, 4, 2, 0$$

## Variational method

$$\text{Numeric method: } \min \text{Tr} H \rho = \min_{\{b_i\}} \frac{b_i b_j}{b_l b_m} \frac{\eta_{ij}}{g_{lm}}$$

Lagrange method:

$$\min \text{Tr} H \rho = \min_{\{b_i\}} b_i b_j \eta_{ij}$$

$$L(\{b_k\}, \lambda) = b_i b_j (\eta_{ij} - \lambda g_{ij}) + \lambda$$

$$\Rightarrow b_i (\eta_{ik} - \lambda g_{ik}) = 0$$

with condition  $b_l b_m g_{lm} = 1$

$$\Leftrightarrow b_i (g_{ij}^{-1} \eta_{jk}) = b_i h_{ik} = \lambda b_k$$

$$\min_{\{b_i\}} b_i b_j \eta_{ij} = \min \lambda$$

If we don't take into account the symmetries of commutation of spins

Variational method is identical to Schrodinger equation

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