1D Hubbard Model

BA-LDA

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Hubbard Model

Criation and Aniquilation Operators



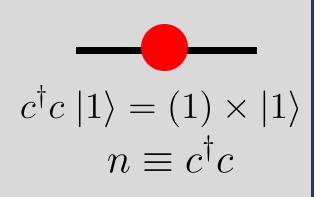
$$|1\rangle = c^{\dagger}|0\rangle$$

$$|0\rangle = c|1\rangle$$



$$0 = c^{\dagger} c^{\dagger} |0\rangle$$

$$0 = c|0\rangle$$



Let's consider ONLY ONE atom only one orbital for simplicity:



EMPTY

$$E = 0$$
$$|sitio\rangle = |0\rangle$$





SIMGLE OCCUPANCE

$$E = v_0$$
$$|sitio\rangle = c_{0,\sigma}^{\dagger}|0\rangle$$

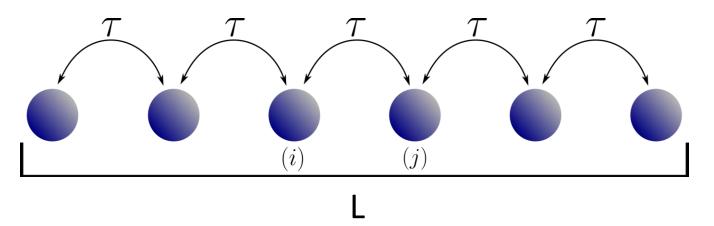


DUBLE OCCUPANCE

$$E = 2v_0 + \mathcal{U}$$
$$|sitio\rangle = c_{0,\uparrow}^{\dagger} c_{0,\downarrow}^{\dagger} |0\rangle$$

Hubbard Model

Now, we can set up a 1D chain with L those atoms. We may also allow the electrons to hop between them with a probability au.



The system's Hamiltonian is

$$\mathcal{H} = \left[-\tau \sum_{i,\sigma} \left(c_{i,\sigma}^{\dagger} c_{i+1,\sigma} + h.c. \right) + \left[\sum_{i,\sigma} v_{i,\sigma} n_{i,\sigma} \right] + \left[\mathcal{U} \sum_{i} n_{i,\uparrow} n_{i,\downarrow} \right] \right]$$

Kinetic Energy

Atomic Orbital Energy **Coulomb Repulsion**

Bethe Ansatz Solution*

The Bethe ansatz solution starts assuming a general form for the system's eigenstate

$$|\Psi\rangle = \sum_{1 < x_1 < \dots x_n < N} \sum_{p \in S_N} A_p e^{\sum_{j=1}^n i k_{p_j} x_j} |x_1, \dots, x_n\rangle$$

From this, we can obtain the ground-state energy in the half-filling case

$$E_0 = -4L \int_0^\infty \frac{J_0(x)J_1(x)}{x(1+e^{xU/2})} dx$$

While the energy per site is a function of the occupation number (density),

$$e^{BA}(n,U) = -\frac{2\beta(U)}{\pi} \sin\left(\frac{n\pi}{\beta(U)}\right) \qquad n \le 1$$

$$e^{BA}(n,U) = -\frac{2\beta(U)}{\pi} \sin\left(\frac{(2-n)\pi}{\beta(U)}\right) + U(n-1) \qquad 1 < n \le 2$$

Such that $\beta(U)$ can be obtained from, considering n=1,

$$-\frac{2\beta(U)}{\pi}\sin\left(\frac{\pi}{\beta(U)}\right) = -4\int_0^\infty \frac{J_0(x)J_1(x)}{x(1+e^{xU/2})}dx$$

DFT - Scheme*

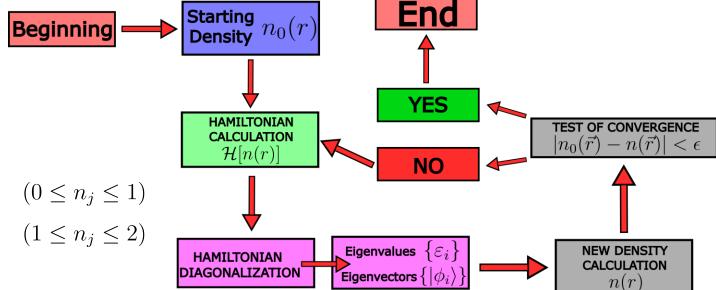
The Kohn-Sham Hamiltonian is

$$\mathcal{H}_{KS} = -\tau \sum_{j} (c_{j,\sigma}^{\dagger} c_{j+1,\sigma} + h.c.) + \sum_{j,\sigma} v_{j,\sigma}^{KS} n_{j,\sigma} - \mu \sum_{j,\sigma} n_{j,\sigma}$$

Such the external potentials are

$$v_{j,\sigma}^{KS} = v_{j,\sigma} + v_j^H[n_j] + v_j^{xc}[n_j]$$
$$v_j^H[n_j] = \frac{U}{2}n_j$$

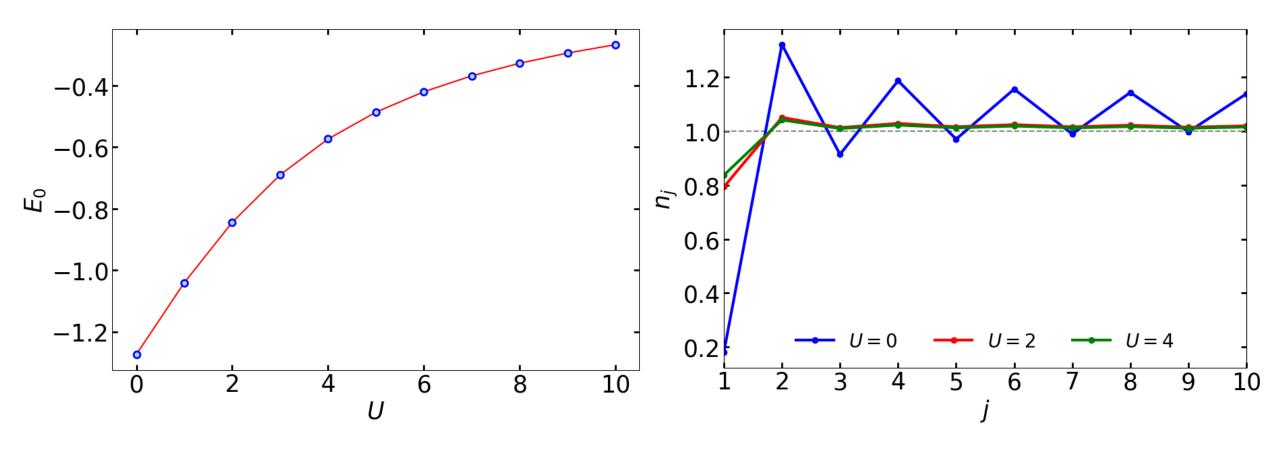
$$v_j^{xc}[n_j] = 2\sin\frac{\pi n_j}{2} - \frac{U}{2}n_j - 2\begin{cases} \cos\left(\frac{\pi n_j}{\beta(U)}\right), & (0 \le n_j \le 1) \\ -\cos\left(\frac{\pi(2-n_j)}{\beta(U)}\right) + U, & (1 \le n_j \le 2) \end{cases}$$



The total energy is

$$E[n] = -\frac{2\beta(U)}{\pi} \sum_{j} \begin{cases} \sin\left(\frac{\pi n_{j}}{\beta(U)}\right), & (0 \le n_{j} \le 1) \\ -\sin\left(\frac{\pi(2-n_{j})}{\beta(U)}\right) + U(n_{j} - 1), & (1 \le n_{j} \le 2) \end{cases}$$

DFT – Results



Ground state energy as a function of the Hubbard interaction in the homogenous case.

Occupation number for a chain of 10 sites considering different many-body interaction values. A positive potential is applied at the first site.

TD-DFT - Scheme*

We have to solve the Schrödinger equation

$$-i\frac{\partial |\Psi(t)\rangle}{\partial t} = \mathcal{H}|\Psi(t)\rangle$$

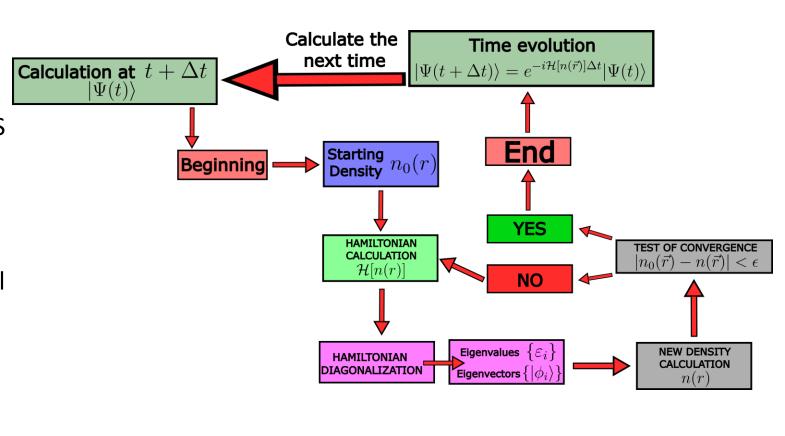
Whose Hamiltonian is calculated by the KS scheme.

Our initial state is the initial ground state.

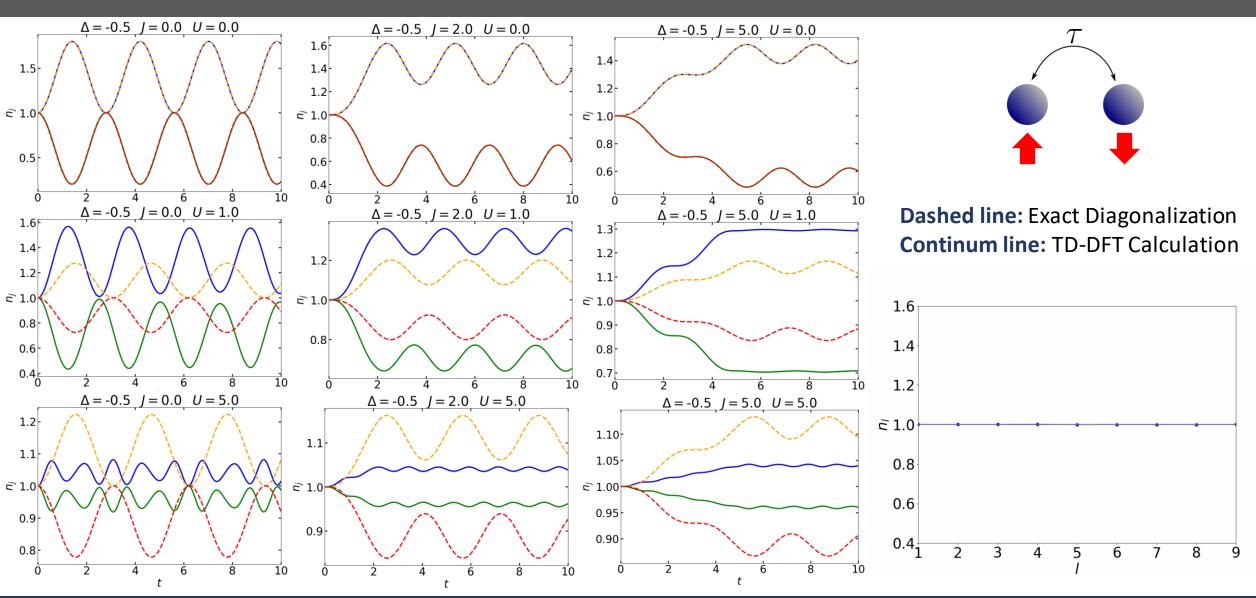
We have changed the external potential following the ramp up protocol,

$$v(t) = \frac{\Delta t}{J}$$

Where we apply positive (negative) potential in odd (even) sites.



TD-DFT – Results: Hubbard dimer



References:

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- Herrera, M., Serra, R.M. & D'Amico, I. Sci Rep 7, 4655 (2017).

THANKS!!

Code:

https://github.com/FelipePicoli/Hubbard-Model---DFT-and-TD-DFT