

## 18.303 Problem Set 3

Due Wednesday, 26 September 2012.

### Problem 1: Deriving diffusion

In this problem, you will be deriving the fact that a diffusion equation results from any microscopic picture in which the flow of solute is proportional to the gradient or difference in solute concentration (“Fick’s law”).

- (a) Consider the case of salt diffusing in water within a narrow straw of length  $L$ , which we will approximate as a one-dimensional system. Imagine dividing the straw into  $M$  small pieces of length  $\Delta x$ , and call the mass of salt in the  $m$ -th piece  $u_m \Delta x$  for  $m = 1, \dots, M$ , where  $u_m$  is a concentration (mass per  $\Delta x$ ). Suppose that the rate at which mass flows from one piece to the next is proportional to the difference in concentrations ( $u_m$ ), and is inversely proportional to the distance to the mass must travel ( $\Delta x$ ), so that the *mass per unit time* flowing from  $m$  to  $m + 1$  is given by:

$$D \frac{u_{m+1} - u_m}{\Delta x}$$

for some constant *diffusion coefficient*  $D$ . Derive an equation for the net rate of change  $du_m/dt$  of the concentration in piece  $m$ . Take the limit  $\Delta x \rightarrow 0$  and show that you obtain a diffusion equation for  $u(x)$ , where  $u(m\Delta x) = u_m$ . Don’t worry about the boundary conditions; just look at the interior  $1 < m < M$ .

- (b) Suppose that any salt that reaches the ends of the straw is immediately removed (there is a little salt-eating demon sitting at each end of the straw). What boundary condition on  $u_0$  and  $u_{M+1}$  does that imply in the discrete model? What boundary condition on  $u(x)$  in the continuum limit?
- (c) Suppose that we seal the ends of the straw, so that no salt can enter or leave through the ends. Write down a boundary condition on  $u_0$  and  $u_{M+1}$  that reflects this situation in the discrete model, write the corresponding matrix  $A$  in the discrete system ( $d\mathbf{u}/dt = A\mathbf{u}$ ), and write the corresponding boundary condition on  $u(x)$  in the continuum limit.

### Problem 2: Adjoint in higher dimensions

- (a) Suppose that we have an *anisotropic* wave equation in  $d$  dimensions,  $\frac{\partial^2 u}{\partial t^2} = \frac{1}{\rho(\mathbf{x})} \nabla \cdot C(\mathbf{x}) \nabla u$ , where  $\rho(\mathbf{x})$  is a scalar density (real,  $> 0$ ) and  $C(\mathbf{x})$  is a  $d \times d$  matrix: a “stiffness” that depends on direction. (e.g. think of a stretched drum in 2 dimensions, made of a material that is easier to stretch in one direction than in the other). Suppose we have Dirichlet boundary conditions  $u|_{\partial\Omega} = 0$ . If  $C$  is a real-symmetric positive-definite matrix at every  $\mathbf{x}$ , show that we still expect a superposition of “normal” (orthogonal) oscillating modes (real frequencies) by choosing the correct inner product and analyzing the operators, similar to class.
- (b) Consider 2-component vector fields  $\mathbf{u}(\mathbf{x})$  in two dimensions, under the inner product  $\langle \mathbf{u}, \mathbf{v} \rangle = \int_{\Omega} \mathbf{u}^* \mathbf{v}$  for some domain  $\Omega$  (where  $*$  applied to a vector is just the conjugate transpose). Give an example of boundary conditions under which the operator

$$\hat{A}\mathbf{u} = \nabla^2 \mathbf{u} + \nabla(\nabla \cdot \mathbf{u}) = \begin{pmatrix} \nabla^2 u_x + \frac{\partial}{\partial x} \nabla \cdot \mathbf{u} \\ \nabla^2 u_y + \frac{\partial}{\partial y} \nabla \cdot \mathbf{u} \end{pmatrix}$$

is self-adjoint under this inner product. (This operator shows up, for example, as part of the Lamé–Navier equations of elastic solids.)

### Problem 3: Symmetric matrices

In class, when looking at the balls-and-springs system for varying masses and spring constants, we showed that the  $M \times M$  matrix

$$A = -M^{-1}D^*KD\Delta x^2$$

is self-adjoint and negative-definite under the inner product  $\langle \mathbf{x}, \mathbf{y} \rangle_M = \mathbf{x}^* M \mathbf{y}$ , where  $M$  and  $K$  were positive diagonal matrices (in fact, any positive-definite matrices would do), and  $D\Delta x$  was our  $(M+1) \times M$  difference matrix.

- (a) Construct a  $5 \times 5$  example of such an  $A$  in Matlab, using the `diff1(5)` command from the last pset to construct  $D\Delta x$ , and pick some  $M$  and  $K$  with distinct diagonal elements. [e.g. the command `diag([1 2 3 4 5])` makes a  $5 \times 5$  diagonal matrix with the numbers 1 to 5 on its diagonal.] Using the `[V,S] = eig(A)` command, as in the last pset, to get the matrix  $V$  of eigenvectors and the diagonal matrix  $S$  of eigenvalues, verify that the eigenvalues are real and negative and that the eigenvectors are orthogonal under the  $M$  inner product above. e.g. `V(:,1)' * M * V(:,2)` computes the  $M$  inner product of the first two eigenvectors, and `V'*M*V` computes a matrix  $(V^*MV)$  of all the inner products of all the eigenvectors.
- (b) Show that *any* diagonalizable matrix  $A$  with real eigenvalues is self-adjoint under  $\langle \mathbf{x}, \mathbf{y} \rangle_W = \mathbf{x}^* W \mathbf{y}$  for  $W = (VV^*)^{-1}$ , in terms of the diagonalization  $A = V\Lambda V^{-1}$  of  $A$ . Compute this in Matlab for your  $V$  from (a); do you get  $W = M$ ? Explain why you can get many different  $W$  matrices from the same  $W = (VV^*)^{-1}$  formula.