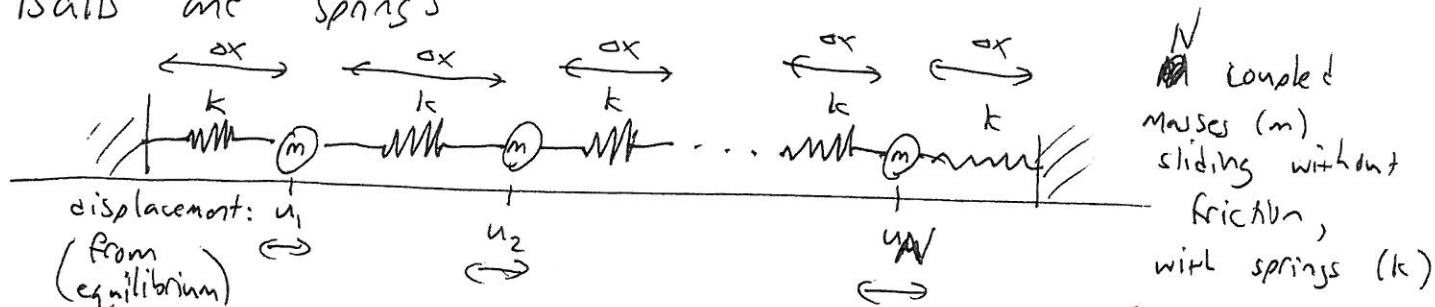


# Lecture 5.5 : from discrete to continuum

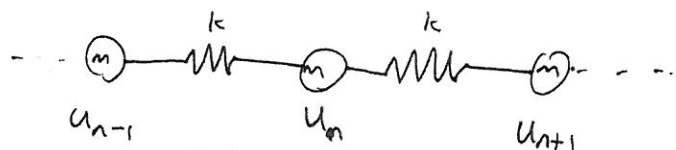
①

\* before, we viewed the discretized eq as an approx for the continuum equations — now we will do the reverse: start with the discrete problem, & derive continuum problem as a limit or approximation

## \* Balls and springs



( $u_n = 0$  at equilibrium)



net force on  $u_n$ :

$F_{n-1/2}$

$F_{n+1/2}$

$$k(u_{n+1} - u_n) - k(u_n - u_{n-1}) \quad (\text{Hooke's Law})$$

"  $F_{n+1/2}$  "      "  $F_{n-1/2}$  "

$$= k(u_{n+1} - 2u_n + u_{n-1})$$

(looks like  $\approx \frac{d^2}{dx^2}$  without the  $\Delta x^2$ !)

more systematically:

(i) get  $F_{n+1/2}$ 's from  $k \times$  (differences in  $u_n$ 's)

(ii) get net force from differences in  $F_{n+1/2}$ 's

(i) in matrix form:

$$\underbrace{\begin{pmatrix} F_{1/2} \\ F_{3/2} \\ \vdots \\ F_{N+1/2} \end{pmatrix}}_{\substack{\vec{F} \\ (N+1 \text{ components})}} = K \underbrace{\begin{pmatrix} 1 & & & \\ -1 & 1 & & \\ & -1 & 1 & \\ & & \ddots & \ddots \\ & & & -1 & 1 \end{pmatrix}}_{(N+1) \times N} \underbrace{\begin{pmatrix} u_1 \\ u_2 \\ \vdots \\ u_N \end{pmatrix}}_{\substack{\vec{u} \\ (N \text{ comp's})}}$$

$$\Rightarrow \vec{F} = K \underset{\substack{\uparrow \\ \text{same } D \text{ as for FD approx!}}}{D} \vec{u} \cdot \Delta x$$

~~with  $\Delta x$~~

(ii) in matrix form:

$$m \ddot{\vec{u}} = \underbrace{\text{net force}}_{N \text{ components}} = \underbrace{\begin{pmatrix} -1 & 1 & & \\ & -1 & 1 & \\ & & \ddots & \ddots \\ & & & -1 & 1 \end{pmatrix}}_{= -\Delta x D^T \text{ from before!}} \begin{pmatrix} F_{1/2} \\ F_{3/2} \\ \vdots \\ F_{N+1/2} \end{pmatrix}$$

$$\Rightarrow \ddot{\vec{u}} = \underbrace{-\frac{K}{m} \Delta x^2 D^T D}_{A} \vec{u} = \bullet \underset{\substack{\uparrow \\ \text{real-symmetric} \\ \text{negative definite}}}{A} \vec{u}$$

$$= \frac{K}{m} \Delta x^2 \underbrace{\begin{pmatrix} -2 & 1 & & \\ 1 & -2 & 1 & \\ & \ddots & \ddots & \ddots \\ & & 1 & -2 & 1 \\ & & & 1 & -2 \end{pmatrix}}_{\substack{\Delta x^2 \\ \text{same discrete Laplacian}}}$$

• Solution to  $\ddot{\vec{u}} = A \vec{u}$  :

$A$  diagonalizable :  $N$  eigenvectors  $\vec{u}_n$   
and eigenvalues  $\lambda_n$

orthonormal: choose  $\vec{u}_n^* \vec{u}_m = \begin{cases} 0 & n \neq m \\ 1 & n = m \end{cases}$   
real,  $< 0$

expand  $\vec{u}(t)$  in this basis:

$$\vec{u}(t) = \sum_{n=1}^N c_n(t) \vec{u}_n$$

some coefficients :  $c_n(t) = \vec{u}_n^* \vec{u}(t)$   
by orthonormality

plug in :  $\ddot{\vec{u}} = A \vec{u}$

$$\sum_n \ddot{c}_n \vec{u}_n = \sum_n c_n \lambda_n \vec{u}_n \Rightarrow \ddot{c}_n = \lambda_n c_n$$

$$\Rightarrow c_n(t) = \alpha_n \cos(\underbrace{\sqrt{-\lambda_n}}_{\omega_n} t) + \beta_n \sin(\underbrace{\sqrt{-\lambda_n}}_{\omega_n} t)$$

$\omega_n = \sqrt{-\lambda_n}$   
= "eigen frequency"  
(real since  $\lambda_n < 0$ )

where  $\alpha_n, \beta_n$  are some coefficients  
determined by initial conditions

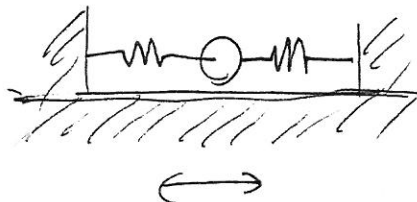
$$\Rightarrow \vec{u}(t) = \sum_{n=1}^N \left[ \alpha_n \cos(\omega_n t) + \beta_n \sin(\omega_n t) \right] \vec{u}_n$$

"normal modes"  
oscillating with frequencies  $\omega_n$

$$\begin{aligned} \vec{u}(0) &= \sum_n \alpha_n \vec{u}_n \Rightarrow \alpha_n = \vec{u}_n^* \vec{u}(0) \\ \dot{\vec{u}}(0) &= \sum_n \omega_n \beta_n \vec{u}_n \Rightarrow \beta_n = \frac{\vec{u}_n^* \dot{\vec{u}}(0)}{\omega_n} \end{aligned}$$

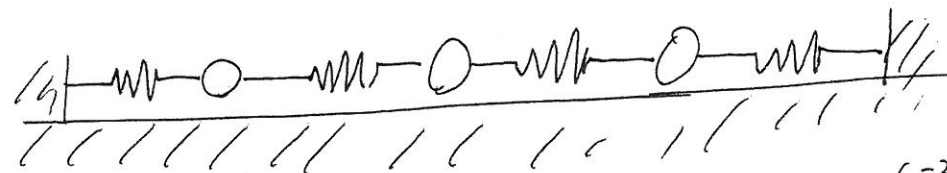
examples :

(4)

★  $N=1$  :   $A = \frac{k}{m} \begin{pmatrix} -2 \end{pmatrix}$   
 $\Rightarrow$  one mode, one  $\omega_n = \sqrt{\frac{2k}{m}}$

★  $N=2$  :   $A = \frac{k}{m} \begin{pmatrix} -2 & 1 \\ 1 & -2 \end{pmatrix}$

$\vec{u}_1 = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ 1 \end{pmatrix}$  :  $\rightarrow \rightarrow$  (moving together)  $\Rightarrow$  2 modes  $\vec{u}_1, \vec{u}_2$   
 $\perp \updownarrow$   
 $\vec{u}_2 = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ -1 \end{pmatrix}$  :  $\rightarrow \leftarrow$  (moving opposite)  
 $\omega_1 = \sqrt{\frac{k}{m}}$   
 $\omega_2 = \sqrt{\frac{3k}{m}}$

★  $N=3$  :   $A = \frac{k}{m} \begin{pmatrix} -2 & 1 & 0 \\ 1 & -2 & 1 \\ 0 & 1 & -2 \end{pmatrix}$

$\vec{u}_1 = \frac{1}{2} \begin{pmatrix} 1 \\ \sqrt{2} \\ 1 \end{pmatrix}$  :  $\rightarrow \rightarrow \rightarrow$   $\omega_1 \approx 0.765 \sqrt{\frac{k}{m}}$   
 $<$

(note  $\perp$ )  $\vec{u}_2 = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ 0 \\ -1 \end{pmatrix}$  :  $\rightarrow \cdot \leftarrow$   $\omega_2 \approx 1.414 \sqrt{\frac{k}{m}}$   
 $<$

$\vec{u}_3 = \frac{1}{2} \begin{pmatrix} 1 \\ -\sqrt{2} \\ 1 \end{pmatrix}$  :  $\rightarrow \leftarrow \rightarrow$   $\omega_3 \approx 1.848 \sqrt{\frac{k}{m}}$

# \* The Continuum Limit $N \rightarrow \infty$

- make  $\Delta x$  smaller & smaller

$\Rightarrow$  make mass  $m$  smaller : let  $m = \rho \Delta x$   
 $\uparrow$   
 density (per length)

- shortening springs increases  $k$  !

$$\left( \begin{array}{c} \text{---} \text{---} \text{---} \\ k_1 \quad k_2 \end{array} = \text{---} \text{---} \text{---} \right. \quad \left. \begin{array}{c} \text{---} \text{---} \text{---} \\ 2k \quad 2k \end{array} = \text{---} \text{---} \text{---} \right)$$

$$\left( \text{---} \text{---} \text{---} = \left( \frac{1}{k_1} + \frac{1}{k_2} \right)^{-1} \right. \quad \left. \text{---} \text{---} \text{---} = \text{---} \text{---} \text{---} \right)$$

$$\Rightarrow \text{let } k = c / \Delta x$$

$$\Rightarrow \ddot{\vec{u}} = - \frac{k}{m} \Delta x^2 \vec{D}^T \vec{D} \vec{u} = - \frac{c}{\rho} \vec{D}^T \vec{D} \vec{u}$$

$$\xrightarrow{\Delta x \rightarrow 0} \left[ \frac{\partial^2 u(x,t)}{\partial t^2} = + \frac{c}{\rho} \frac{\partial^2 u(x,t)}{\partial x^2} \right] \quad \left( - \vec{D}^T \vec{D} \rightarrow \frac{\partial^2}{\partial x^2} \right)$$

scalar wave equation !

$$\ddot{u} = \hat{A} u \quad , \quad \begin{array}{l} \text{fixed ends:} \\ u(0,t) \\ = u(L,t) = 0 \end{array}$$

$$\hat{A} = + \frac{c}{\rho} \frac{\partial^2}{\partial x^2} \quad \begin{array}{l} \text{negative} \\ \text{definite} \\ \text{self-adjoint} \end{array}$$

for usual  $\langle u, v \rangle = \int u v$

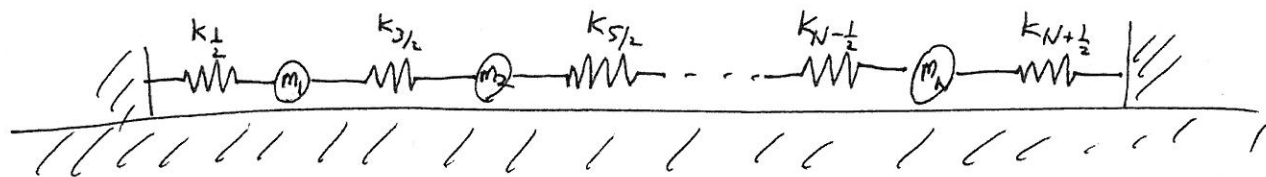
$$\Rightarrow \text{real, } \lambda_n < 0$$

$$\Rightarrow \text{oscillating sols with } \omega_n = \sqrt{-\lambda_n}$$

# \* Inhomogeneous materials :

(6)

- suppose each  $m, k$  is different :



$$\Rightarrow (i) \vec{F} = K D \vec{u} \Delta x$$

$$\begin{pmatrix} k_{1/2} & & \\ & k_{3/2} & \\ & & \ddots \\ & & & k_{N+1/2} \end{pmatrix} \quad (N+1) \times (N+1) \text{ diagonal matrix of } k\text{'s}$$

$$(ii) \ddot{\vec{u}} = - \underbrace{M^{-1}}_{\text{diagonal matrix of } 1/m \text{'s}} \Delta x^2 D^T K D \vec{u} = A \vec{u}$$

$$\begin{pmatrix} 1/m_1 & & \\ & 1/m_2 & \\ & & \ddots \\ & & & 1/m_N \end{pmatrix} = N \times N \text{ diagonal matrix of } 1/m \text{'s}$$

$$A = - \Delta x^2 M^{-1} D^T K D = A^*$$

under  $\langle \vec{u}, \vec{v} \rangle = \vec{u}^* M \vec{v}$   
 & negative-def for  $m, k > 0$

$$\xrightarrow[N \rightarrow \infty]{} \hat{A} = \frac{1}{\rho(x)} \frac{\partial}{\partial x} c(x) \frac{\partial}{\partial x}$$

$$= \hat{A}^* \text{ under } \langle u, v \rangle = \int \rho \bar{u} v$$

& negative-definite for  $\rho, c > 0$