

18.303 Midterm Solutions, Fall 2010

Problem 1: Finite differences (20 points)

- (a) At the right boundary it is the same as the Dirichlet matrix before. At the left boundary, evaluating at $\Delta x/2$, we have

$$u'_{1/2} \approx \frac{u_1 - u_0}{\Delta x} \approx -\frac{u_0 + u_1}{2} \approx -u_{1/2},$$

Solving for u_0 , we have:

$$\frac{u_1 - u_0}{\Delta x} = -\frac{u_0 + u_1}{2} \implies u_0 \left(\frac{1}{2} - \frac{1}{\Delta x} \right) = -u_1 \left(\frac{1}{2} + \frac{1}{\Delta x} \right) \implies u_0 = u_1 \frac{\frac{1}{\Delta x} + \frac{1}{2}}{\frac{1}{\Delta x} - \frac{1}{2}} = u_1 \frac{2 + \Delta x}{2 - \Delta x},$$

so that the first row of D , which computes $u'_{1/2}$, now computes $u'_{1/2} = \frac{u_1 - u_0}{\Delta x} = \frac{u_1}{\Delta x} \left(1 - \frac{2 + \Delta x}{2 - \Delta x} \right) = \frac{u_1}{\Delta x} \frac{-2\Delta x}{2 - \Delta x}$, and hence D could be written as the following $(M+1) \times M$ matrix that computes the derivatives from M unknowns $(u_1, u_2, \dots, u_M)^T$:

$$D = \begin{pmatrix} \frac{-2\Delta x}{2 - \Delta x} & & & & \\ -1 & 1 & & & \\ & \ddots & \ddots & & \\ & & -1 & 1 & \\ & & & & -1 \end{pmatrix}.$$

- (b) Since we applied the left boundary condition at the wrong place $\Delta x/2$, this will introduce an error $\sim \Delta x$.

[Like in the Neumann pset, however, we can correct this by changing Δx to “stretch” the grid so that the left boundary is at 0: $\Delta x = L/(M+0.5)$, in which case we would get errors $\sim \Delta x^2$ since the difference formulas are otherwise 2nd-order.]

Problem 2: Adjoint and stuff (20 points)

- (a) **Solution 1:** Writing out the integrals:

$$\begin{aligned} \langle u, \hat{B}v \rangle &= \int_0^{L_x} \int_0^{L_y} dx dy u \frac{\partial}{\partial x} \left[c(x, y) \frac{\partial v}{\partial y} \right] \\ &= - \int_0^{L_x} \int_0^{L_y} dx dy \frac{\partial u}{\partial x} \left[c(x, y) \frac{\partial v}{\partial y} \right] + \int_0^{L_y} dy u(x, y) \left[c(x, y) \frac{\partial v}{\partial y} \right] \Big|_{x=0}^{L_x} \\ &= + \int_0^{L_x} \int_0^{L_y} dx dy \frac{\partial}{\partial y} \left[c(x, y) \frac{\partial u}{\partial x} \right] v - \int_0^{L_x} dx \left[c(x, y) \frac{\partial u}{\partial x} \right] v(x, y) \Big|_{y=0}^{L_y} \\ &= \left\langle \frac{\partial}{\partial y} \left[c \frac{\partial u}{\partial x} \right], v \right\rangle, \end{aligned}$$

where we have integrated by parts twice (flipping the sign twice to get back to $+$) and the boundary terms vanish as usual from Dirichlet. Note that integrating by parts **switches the order** of the derivatives! Hence:

$$\hat{B}^* = \left(\frac{\partial}{\partial x} \left[c(x, y) \frac{\partial}{\partial y} \right] \right)^* = \frac{\partial}{\partial y} \left[c(x, y) \frac{\partial}{\partial x} \right],$$

and thus

$$\hat{A} = \nabla^2 + \hat{B} + \hat{B}^*.$$

It follows immediately that $\hat{A}^* = \nabla^2 + \hat{B}^* + \hat{B} = \hat{A}$ (since we already know ∇^2 is self-adjoint with this inner product and boundary condition).

Solution 2:

By now, you should be used to the fact that integrating by parts moves a first derivative from one side to

the other of an inner product with a sign flip, plus a boundary term which is zero for Dirichlet boundaries, hence the above solution could be written in short form as:

$$\begin{aligned}\langle u, \hat{B}v \rangle &= \left\langle u, \frac{\partial}{\partial x} \left[c \frac{\partial v}{\partial y} \right] \right\rangle \\ &= - \left\langle \frac{\partial u}{\partial x}, c \frac{\partial v}{\partial y} \right\rangle \\ &= - \left\langle c \frac{\partial u}{\partial x}, \frac{\partial v}{\partial y} \right\rangle \\ &= + \left\langle \frac{\partial}{\partial y} \left[c \frac{\partial u}{\partial x} \right], v \right\rangle = \langle \hat{B}^* u, v \rangle.\end{aligned}$$

Solution 3 (not expected): A more beautiful and general way to do this is to write

$$\hat{A}u = \nabla \cdot \left[\begin{pmatrix} 1 & c \\ c & 1 \end{pmatrix} \nabla u \right] = \nabla \cdot (C \nabla u)$$

with the real-symmetric matrix $C = C^T$, in which case we can integrate by parts using the divergence theorem for any shape Ω (and in any number of dimensions if we make C a bigger matrix):

$$\begin{aligned}\langle u, \nabla \cdot (C \nabla v) \rangle &= \langle -\nabla u, C \nabla v \rangle + \cancel{\oint_{\partial\Omega} u C \nabla v \cdot d\mathbf{A}} = - \int_{\Omega} \nabla u \cdot C \nabla v \\ &= - \int_{\Omega} (C^T \nabla u) \cdot \nabla v = - \langle C \nabla u, \nabla v \rangle \\ &= + \langle \nabla \cdot (C \nabla u), v \rangle - \cancel{\oint_{\partial\Omega} (C \nabla u) v \cdot d\mathbf{A}},\end{aligned}$$

quoting the integration-by-parts for ∇ as in class (where the boundary terms vanish by Dirichlet), hence $\hat{A} = \hat{A}^*$.

- (b) It will have (exponentially) decaying solutions if \hat{A} is negative-definite, i.e. $\langle u, \hat{A}u \rangle < 0$ for all $u \neq 0$ (hence negative λ 's). To check this, we integrate by parts *once*, to obtain:

$$\begin{aligned}\langle u, \hat{A}u \rangle &= - \langle \nabla u, \nabla u \rangle - \left\langle \frac{\partial u}{\partial x}, c \frac{\partial u}{\partial y} \right\rangle - \left\langle \frac{\partial u}{\partial y}, c \frac{\partial u}{\partial x} \right\rangle \\ &= - \int_{\Omega} \nabla u \cdot \begin{pmatrix} 1 & c \\ c & 1 \end{pmatrix} \nabla u,\end{aligned}$$

which is < 0 if $\nabla u \neq 0$ (true because of the boundary conditions: u cannot be a nonzero constant) and if $\begin{pmatrix} 1 & c \\ c & 1 \end{pmatrix}$ is positive-definite, i.e. its eigenvalues $1 \pm c$ must be > 0 , i.e. $|c| < 1$.

Problem 3: Thinking Green (20 points)

- (a) \hat{A} is self-adjoint under the inner product $\langle u, v \rangle = \int_{\Omega} \bar{u}v/c$, from class (and the 1d version in homework), and hence the eigenfunctions are orthogonal, hence $\alpha_n = \langle u_n, f \rangle / \langle u_n, u_n \rangle$. The expression $u(\mathbf{x}) = \int_{\Omega} G(\mathbf{x}, \mathbf{x}') f(\mathbf{x}') d^2 \mathbf{x}'$ is just the solution to $\hat{A}u = f$, i.e. it is $u = \hat{A}^{-1}f = \sum_n \frac{\alpha_n}{\lambda_n} u_n(\mathbf{x})$.
- (b) This is just $\langle u, \hat{A}^{-1}u \rangle / \langle u, u \rangle$, which is a Rayleigh quotient for the self-adjoint operator \hat{A}^{-1} , and hence (thanks to the min-max theorem) it is bounded above by the **largest** eigenvalue of \hat{A}^{-1} , which is $1/\lambda_1$ (the inverse of the **smallest** eigenvalue of \hat{A}). (Note that by positive-definiteness $0 < \lambda_1 < \lambda_2 < \dots$)