

## 18.303 Problem Set 3 Solutions

### Problem 1: (5+5+10+10+5)

- (a)  $\frac{d}{dt}u_m\Delta x$  is the net rate at which mass diffuses into piece  $m$  minuse the rate at which it dissociates:

$$\frac{du_m}{dt}\Delta x = D\frac{u_{m+1} - u_m}{\Delta x} + D\frac{u_{m-1} - u_m}{\Delta x} - Ru_m\Delta x$$

which can be rearranged into

$$\frac{du_m}{dt} = D\frac{u_{m+1} - 2u_m + u_{m-1}}{\Delta x^2} - Ru_m.$$

In the  $\Delta x \rightarrow 0$  limit, the first term is exactly the second derivative, from class, while the second term doesn't change, so we obtain the following PDE for  $u(x, t)$ :

$$\frac{\partial u}{\partial t} = D\frac{\partial^2 u}{\partial x^2} - Ru$$

- (b) This corresponds to  $u_0 = u_{M+1} = 0$ , or Dirichlet boundary conditions  $u(0) = u(L) = 0$ .  
(c) If there is no flow through the ends, this corresponds to setting

$$D\frac{u_{M+1} - u_M}{\Delta x} = 0 = D\frac{u_0 - u_1}{\Delta x},$$

or equivalently  $u_{M+1} = u_M$  and  $u_0 = u_1$ . This makes the  $u_1$  and  $u_M$  equations

$$\frac{du_1}{dt} = D\frac{u_2 - 2u_1 + u_1}{\Delta x^2} - Ru_1,$$

$$\frac{du_M}{dt} = D\frac{u_M - 2u_M + u_{M-1}}{\Delta x^2} - Ru_M.$$

If we write this in matrix form  $d\mathbf{u}/dt = \mathbf{A}\mathbf{u}$ , we get

$$\mathbf{A} = \frac{1}{\Delta x^2} \begin{pmatrix} -1 & 1 & & & \\ 1 & -2 & 1 & & \\ & \ddots & \ddots & \ddots & \\ & & 1 & -2 & 1 \\ & & & 1 & -1 \end{pmatrix} - R\mathbf{I}$$

where  $\mathbf{I}$  is the identity matrix, and the first matrix is the same as our second-derivative matrix from class except for the first and last rows. In the continuum limit, setting  $\frac{u_{m+1} - u_m}{\Delta x} = 0$  corresponds to setting the first derivative to zero at the ends, i.e. we have the (Neumann) boundary conditions  $\frac{\partial u}{\partial x} = 0$  at  $x = 0, L$ .

- (d) The operator in both cases is  $\hat{A} = \frac{\partial^2}{\partial x^2} - R$ . We will use the usual inner product  $\langle u, v \rangle = \int_0^L \bar{u}v$ . The  $R$  term is just a real number and hence is Hermitian ( $\langle u, Rv \rangle = \langle Ru, v \rangle$  by inspection). The  $\frac{\partial^2}{\partial x^2}$  term was already shown in class to be Hermitian for Dirichlet boundary conditions. For Neumann boundary conditions, we integrate by parts

$$\langle u, v'' \rangle = \int_0^L \bar{u}v'' = \bar{u}v'|_0^L - \int_0^L \bar{u}'v' = -\bar{u}'v|_0^L + \int_0^L \bar{u}''v = \langle u'', v \rangle,$$

in which the boundary terms again vanish so we obtain that  $\hat{A} = \hat{A}^*$ . To show definiteness, as in class we just integrate by parts once, which with either boundary condition gives

$$\langle u, \hat{A}u \rangle = - \int_0^L (|u'|^2 + R|u|^2) dx \leq 0,$$

which is  $< 0$  unless  $u = 0$  almost everywhere, due to the  $R$  term alone (since  $R > 0$  was given). Hence it is negative-definite. If we set  $R = 0$ , then with Dirichlet boundary conditions it is still negative-definite as proved in class. For  $R = 0$  and Neumann boundary conditions, however,  $\langle u, \hat{A}u \rangle = 0 \implies u' = 0 \implies u = \text{constant} \implies u = 0$ : nonzero constants are allowed by the boundary conditions, so, e.g.  $\langle 1, \hat{A}1 \rangle = 0$  is allowed. Hence  $\hat{A}$  is negative semidefinite for Neumann boundaries.

When it is negative-definite ( $R > 0$  and/or Dirichlet), that means that the solutions decay to zero as  $t \rightarrow \infty$ , as discussed in class (each eigenfunction decays as  $e^{\lambda t}$  and  $\lambda < 0$ ). When it is only negative semidefinite ( $R = 0$  with Neumann boundary conditions), then the solutions decay to constants (the eigenfunction  $u = 1$  with  $\lambda = 0$  does not decay, while all of the other eigenfunctions decay exponentially)

(e) In this case we get

$$\frac{du_m}{dt} = \frac{D_{m+0.5} \frac{u_{m+1} - u_m}{\Delta x} + D_{m-0.5} \frac{u_m - u_{m-1}}{\Delta x}}{\Delta x} - R_m u_m$$

which goes to the PDE

$$\frac{\partial u}{\partial t} = \frac{\partial}{\partial x} \left( D \frac{\partial u}{\partial x} \right) - Ru.$$

That is, our difference equation is a finite-difference approximation to taking the derivative, then multiplying by  $D(x)$ , then taking another derivative.

As shown in class, the operator on the right-hand-side is still Hermitian and negative definite (or semidefinite if  $R = 0$  and we have Neumann boundaries), so the qualitative behavior of the PDE does not change.