

# computational Many-Body Methods in Physics

Anton N. Torgersen

University of Oslo  
FYS4480

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# Outline

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- 3 Hartree-Fock Theory
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- 6 Coupled Cluster Theory
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# Introduction: The Many-Body Problem

**Goal:** Solve the Schrödinger equation for  $N$  interacting particles.

$$\hat{H} |\Psi\rangle = E |\Psi\rangle$$

**The Challenge:** The Hilbert space grows factorially with  $N$ .

⇒ Exact diagonalization (FCI) is impossible for large systems.

**The Model System: Lipkin-Meshkov-Glick (LMG)** We will use the following Hamiltonian, to analyse the HF method.

$$\hat{H} = \hat{H}_0 + \hat{H}_1 + \hat{H}_2$$

Defined using quasi-spin operators  $\hat{J}_+$ ,  $\hat{J}_-$ ,  $\hat{J}_z$ :

$$\hat{H} = \epsilon \hat{J}_z - \frac{1}{2} V (\hat{J}_+^2 + \hat{J}_-^2) - \frac{1}{2} W (\hat{J}_+ \hat{J}_- + \hat{J}_- \hat{J}_+ - \hat{N})$$

**The Model System: Pairing Model** This model will be used for a full comparison of FCI, HF and MBPT.

$$\hat{H} = \hat{H}_0 + \hat{H}_I$$

Where

$$\hat{H}_0 = \sum_{p,\omega} (p - 1) \hat{a}_{p\omega}^\dagger \hat{a}_{p\omega}$$

and

$$\hat{H}_I = -\frac{1}{2}g \sum_{p,q} \hat{P}_p^+ \hat{P}_q^-.$$

# Second Quantization

**Second Quantization** is a formalism which greatly simplifies many-body calculations, especially when combined with Wick's theorem. Where we know work in the Fock space of occupation numbers.

- Basis of creation ( $a_p^\dagger$ ) and annihilation ( $a_p$ ) operators.
- Anti-commutation relations for fermions:  $\{a_p, a_q^\dagger\} = \delta_{pq}$ . (Pauli exclusion principle built-in)

**Ground State Reference**  $|\Phi_0\rangle$  is defined as the Slater determinant where all the  $N$ -lowest energy states (states below the fermi level) are filled.

$$|\Phi_0\rangle = \prod_{i \leq F} a_i^\dagger |0\rangle$$

**Wick's Theorem** Allows us to evaluate vacuum expectation values of long operator strings.

$$\hat{A}\hat{B}\hat{C}\dots = \{\hat{A}\hat{B}\hat{C}\dots\} + \sum \text{all contractions}$$

- **Midterm 1 Application:** We used Wick's theorem to evaluate matrix elements like  $\langle \Phi_0 | \hat{H} | \Phi_0 \rangle$  and check commutation relations  $[\hat{H}, \hat{J}^2] = 0$ .

# Diagrammatic Representation



# Full Configuration Interaction (FCI)

## The Exact Solution

- Expands  $\Psi$  in the full basis of all possible Slater Determinants.
- $|\Psi_{FCI}\rangle = c_0 |\Phi_0\rangle + \sum_{ia} c_i^a |\Phi_i^a\rangle + \sum_{ijab} c_{ij}^{ab} |\Phi_{ij}^{ab}\rangle + \dots$

Diagonalization

# How it Scales

Is Exact but scales to the power of  $N!$

**Results from Midterm (Lipkin Model)** Exact diagonalization of the Hamiltonian matrix for  $N = 4$  particles.

Coupling $g$	$E_{FCI}$ (Ground State)	Character
-1.0	2.7799	Repulsive
0.0	2.0000	Non-interacting
0.5	1.4168	Attractive
1.0	0.6355	Strong Coupling

Table 1: Exact eigenvalues computed via `numpy.linalg.eigh` (from `Calculations.ipynb`).

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# Mean Field Ansatz

Explination of the mean field ansatz

# Hartree-Fock (HF)

Hartree Fock Equation.

**Stability Condition (Midterm 1)** We derived the stability of the HF solution against particle-hole mixing coefficients  $C_{\alpha\pm}$ .

$$\frac{\epsilon}{3} > -(V + W)$$

- If this holds: The non-interacting ground state is stable.
- If violated: A "mixed" state (deformed solution) lowers the energy.
- For parameter set (2) ( $V = -4/3$ ,  $W = -1$ ), we found a lower energy mixed state  $E \approx -7.57$ , significantly better than the standard reference.

# Hartree-Fock Results

Lower energy for mixed state compared to non-interacting.



# Many-Body Perturbation Theory (MBPT)

Formalism of MBPT.

# Diagrammatic Representation

# linked Diagram Theorem

# Third order Corrections

## Second Midterm

# Benchmarking MBPT With FCI

# MBPT vs FCI Results

## Convergence Analysis (from Calculations.ipynb)

Method	$g = -1.0$	$g = 0.5$	$g = 1.0$
$E_{FCI}$ (Exact)	<b>2.7799</b>	<b>1.4168</b>	<b>0.6355</b>
$E^{(0+1)}$ (HF)	3.0000	1.5000	1.0000
$E^{(2)}$	-0.4667	-0.0624	-0.2190
$E^{(3)}$	+0.5156	-0.0165	-0.1005
$E^{(4)}$	-0.7527	-0.0056	-0.0588
$E_{Total}$ (order 4)	2.2962	<b>1.4155</b>	<b>0.6217</b>

Table 3: Comparison of Perturbation theory orders vs Exact result.

### Observation:

- Excellent agreement for  $g = 0.5$  (weak coupling).
- Divergence/oscillations start appearing at strong coupling ( $g = -1.0$  or  $g = 1.0$ ).

# Motivation for DFT.

# The Hohenberg-Kohn Theorems



# The Kohn-Sham Equations

# The Exchange-Correlation Functional

# The exponential ansatz

# Coupled Cluster Theory (CC)

## Approximations

- **CCD (Doubles):**  $\hat{T} \approx \hat{T}_2 = \frac{1}{4} \sum t_{ij}^{ab} a_a^\dagger a_b^\dagger a_j a_i$ . Relevant for Lipkin model since Hamiltonian connects states by 2p-2h excitations.
- **CCSD (Singles + Doubles):**  $\hat{T} \approx \hat{T}_1 + \hat{T}_2$ . The "gold standard" for chemical accuracy ( $O(N^6)$  scaling).

## Energy Equation

$$E_{CC} = \langle \Phi_0 | \bar{H} | \Phi_0 \rangle = \langle \Phi_0 | e^{-\hat{T}} \hat{H} e^{\hat{T}} | \Phi_0 \rangle$$

**Amplitude Equations (for  $t$ -amplitudes)** We project onto excited determinants to solve for amplitudes:

$$\langle \Phi_{ij}^{ab} | \bar{H} | \Phi_0 \rangle = 0$$

*Comparison note:* While MBPT(4) failed for strong coupling in our previous slide, CCD typically remains robust and closer to FCI in the strong coupling regime because it sums the diagrams to infinite order.

# Summary and Comparison

Method	Pros	Cons
<b>FCI</b>	Exact.	Exponential cost ( $N!$ ). Tiny systems only.
<b>Hartree-Fock</b>	Simple, defines orbitals.	No correlation ( $E_{corr} = 0$ ).
<b>MBPT</b>	Systematically improvable.	Diverges for strong interaction ( $g$ ).
<b>Coupled Cluster</b>	High accuracy (Gold Standard).	High cost ( $O(N^6)$ ), complex to implement.
<b>DFT</b>	Efficient ( $O(N^3)$ ), includes correlation.	Approx. functional ( $V_{xc}$ ), no hierarchy.

**Conclusion from our work:** For the Lipkin model, MBPT works excellently at weak coupling ( $g = 0.5$ ) but breaks down at strong coupling ( $g = 1.0$ ), necessitating non-perturbative methods like Coupled Cluster or FCI.

# Thank You!

Questions?