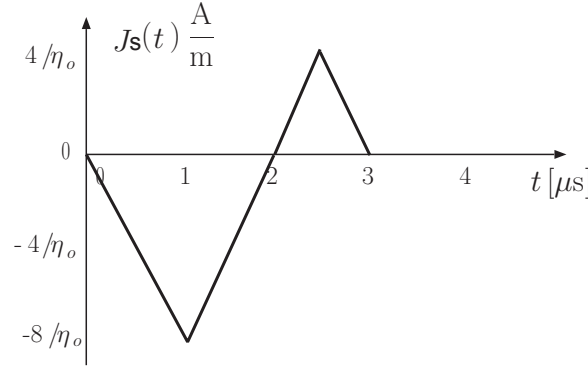


1. A z -polarized plane TEM wave propagates in $-\hat{y}$ direction. If the wave field varies at $y = 0$ according to $E_z(0, t) = 4\Delta\left(\frac{t-1}{\tau}\right) - 2\Delta\left(\frac{t-2.5}{\tau/2}\right)$ V/m, where $\Delta\left(\frac{t}{\tau}\right)$ is the unit triangle function with width $\tau = 2\mu\text{s}$, we can find that

- a) The surface current density can be expressed as:

$$\mathbf{J}_s(t) = -\frac{2}{\eta_o} E_z(t) \hat{z} = -\frac{2}{\eta_o} \left[4\Delta\left(\frac{t-1\mu\text{s}}{2\mu\text{s}}\right) - 2\Delta\left(\frac{t-2.5\mu\text{s}}{1\mu\text{s}}\right) \right] \hat{z} \frac{\text{A}}{\text{m}}$$

where $\eta_o \approx 120\pi\Omega$ is the intrinsic impedance of vacuum. And the plot is:



- b) The vector wave field $\vec{E}(y, t)$ is given by

$$\vec{E}(y, t) = \left[4\Delta\left(\frac{t-1+y/c}{\tau}\right) - 2\Delta\left(\frac{t-2.5+y/c}{\tau/2}\right) \right] \hat{z} \frac{\text{V}}{\text{m}},$$

where $c \approx 3 \times 10^8$ m/s is the wave propagation velocity in vacuum.

- c) The associated wave field $\vec{H}(y, t)$ is

$$\vec{H}(y, t) = \left[\frac{4}{\eta_o} \Delta\left(\frac{t-1+y/c}{\tau}\right) - \frac{2}{\eta_o} \Delta\left(\frac{t-2.5+y/c}{\tau/2}\right) \right] (-\hat{x}) \frac{\text{A}}{\text{m}},$$

- d) The Poynting vector is

$$\vec{S} = \vec{E} \times \vec{H} = -\frac{16}{\eta_o} \Delta^2\left(\frac{t-1+y/c}{\tau}\right) - \frac{4}{\eta_o} \Delta^2\left(\frac{t-2.5+y/c}{\tau/2}\right) \hat{y} \frac{\text{W}}{\text{m}^2},$$

The poynting vector includes 2 triangle function, so there are 2 peaks, and 1 peak will arrive a point later than the other.

So its maximum value is

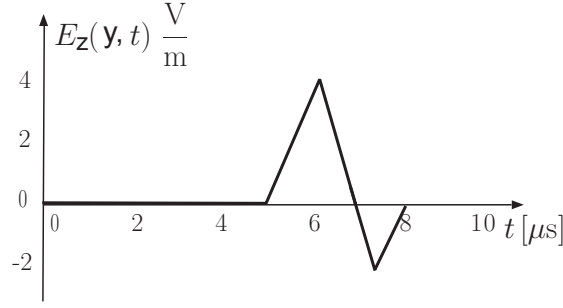
$$\max\left(\left|\vec{E} \times \vec{H}\right|\right) = \frac{16}{\eta_o} \approx \frac{2}{15\pi} \frac{\text{W}}{\text{m}^2}.$$

- e) The locations of the peak of $\vec{E} \times \vec{H}$ evolves according to

$$\frac{t-1+y/c}{\tau} = 0 \quad \rightarrow \quad y = (1-t)c.$$

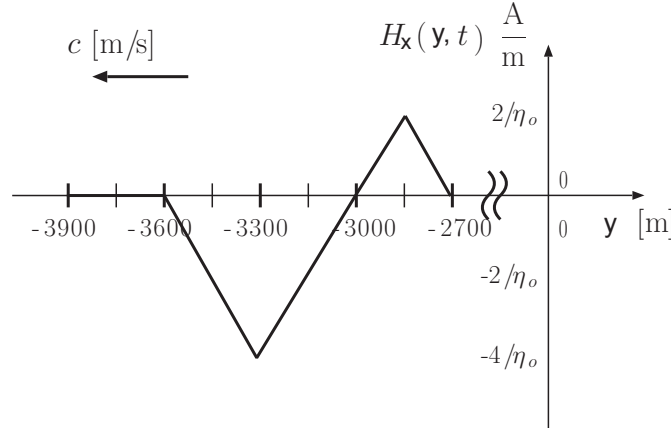
- f) Plotting E_z vs t at $y = -1500$ m.

$$\begin{aligned} E_z(y, t)|_{y=-1500} &= 4\Delta\left(\frac{t-1\mu\text{s} - \frac{15 \times 10^2}{3 \times 10^8} \text{s}}{2\mu\text{s}}\right) - 2\Delta\left(\frac{t-2.5\mu\text{s} - \frac{15 \times 10^2}{3 \times 10^8} \text{s}}{1\mu\text{s}}\right) \\ &= 4\Delta\left(\frac{t-6\mu\text{s}}{2\mu\text{s}}\right) - 2\Delta\left(\frac{t-7.5\mu\text{s}}{1\mu\text{s}}\right) \frac{\text{V}}{\text{m}} \end{aligned}$$



g) Plotting H_x vs y at $t = 12 \mu\text{s}$.

$$\begin{aligned} H_x(y, t)|_{t=12 \mu\text{s}} &= -\frac{4}{\eta_o} \Delta\left(\frac{12 - 1 \mu\text{s} + y \cdot \frac{1}{3} \times 10^{-8}}{2 \mu\text{s}}\right) + \frac{2}{\eta_o} \Delta\left(\frac{12 - 2.5 \mu\text{s} + y \cdot \frac{1}{3} \times 10^{-8}}{1 \mu\text{s}}\right) \\ &= -\frac{4}{\eta_o} \Delta\left(\frac{3300 \text{ m} + y}{600 \text{ m}}\right) + \frac{2}{\eta_o} \Delta\left(\frac{2850 \text{ m} + y}{300 \text{ m}}\right) \frac{\text{A}}{\text{m}} \end{aligned}$$



2. In a homogeneous lossless dielectric with $\epsilon = \epsilon_r \epsilon_o$ and $\mu = \mu_r \mu_o$ a plane TEM wave with the following components is observed,

$$\vec{E} = \hat{x} \frac{1}{2} u\left(t - \frac{z}{c/2}\right) + \hat{y} g\left(t - \frac{z}{c/2}\right) \frac{\text{V}}{\text{m}},$$

and

$$\vec{H} = \hat{x} \left(\frac{10z}{c} - 5t\right) + \hat{y} \frac{1}{120\pi} u\left(t - \frac{2z}{c}\right) \frac{\text{A}}{\text{m}}.$$

a) The intrinsic impedance of the medium is

$$\eta = \left| \frac{E_x}{H_y} \right| = 60\pi \Omega.$$

We have used the orthogonal pair E_x and H_y . The same relation should be valid for the orthogonal pair E_y and H_x .

b) The propagation velocity is

$$v = \frac{c}{2} = 1.5 \times 10^8 \frac{\text{m}}{\text{s}}.$$

c) If $\epsilon = \epsilon_r \epsilon_o$ and $\mu = \mu_r \mu_o$, then the intrinsic impedance η can take the following form

$$\eta = \sqrt{\frac{\mu}{\epsilon}} = \sqrt{\frac{\mu_r}{\epsilon_r}} \eta_o$$

while the wave propagation velocity v can be written as

$$v = \frac{1}{\sqrt{\epsilon\mu}} = \frac{c}{\sqrt{\epsilon_r\mu_r}}.$$

Using the results of parts (a) and (b) we find that

$$\epsilon_r = 4, \quad \mu_r = 1.$$

- d) Finally, since $E_y = -\eta H_x$ (recall that the propagation direction is $\hat{z} = -\hat{y} \times \hat{x}$, hence the minus sign), we have that

$$g\left(t - \frac{z}{c/2}\right) = -60\pi \times \left(\frac{10z}{c} - 5t\right) = 300\pi \times \left(t - \frac{z}{c/2}\right),$$

then

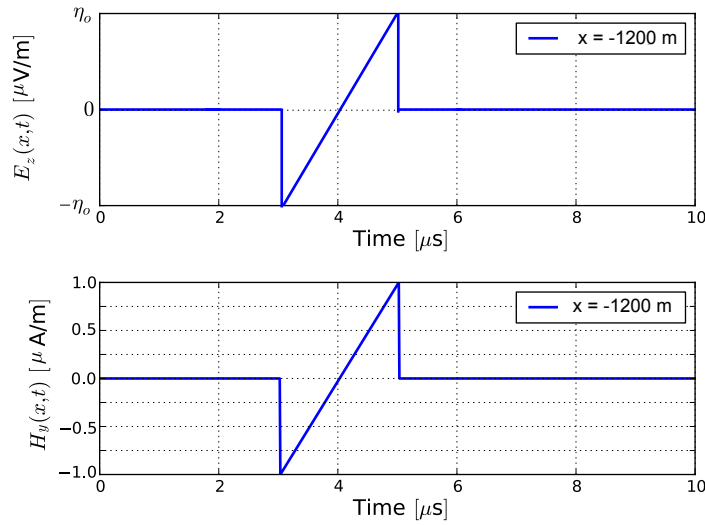
$$g(t) = 300\pi t.$$

3. Consider on $x = 0$ plane, a pulse of sheet current $\vec{J}_s(t) = -\hat{z} 2t \text{rect}\left(\frac{t}{\tau}\right)$ A/m, where $\tau = 2\mu\text{s}$. The corresponding magnetic field is: (it's direction can be verify using the right-hand-rule for Ampere's law $\oint_C \mathbf{H} \cdot d\mathbf{l} = I_C$)

$$\begin{aligned} \vec{H}^\pm(x, t) &= \mp \frac{1}{2} J_{so} \left(t \mp \frac{x}{c}\right) \hat{y} \frac{\text{A}}{\text{m}} \\ &= \mp \left(t \mp \frac{x}{c}\right) \text{rect}\left(\frac{t \mp \frac{x}{c}}{\tau}\right) \hat{y} \frac{\text{A}}{\text{m}} \quad \text{for } x \geq 0, \end{aligned}$$

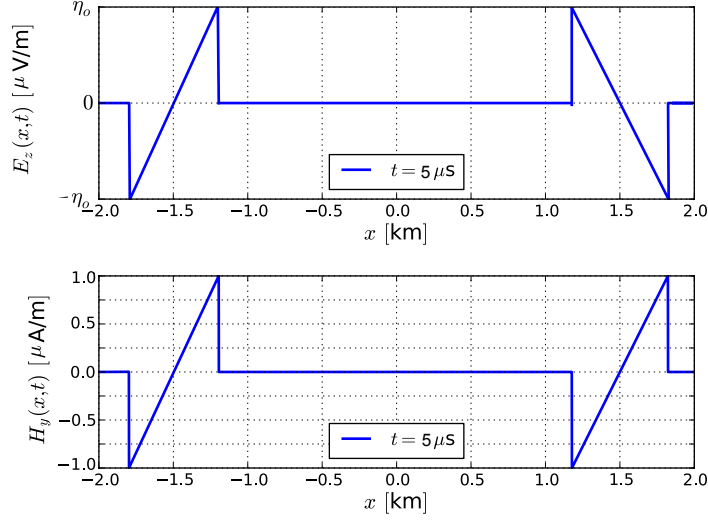
- (a) For $x = -1200\text{m}$, we have that

$$\begin{aligned} E_z^-(x, t) &= \eta_o (t - 4\mu\text{s}) \text{rect}\left(\frac{t - 4\mu\text{s}}{2\mu\text{s}}\right) \frac{\text{V}}{\text{m}} \\ H_y^-(x, t) &= (t - 4\mu\text{s}) \text{rect}\left(\frac{t - 4\mu\text{s}}{2\mu\text{s}}\right) \frac{\text{A}}{\text{m}}. \end{aligned}$$



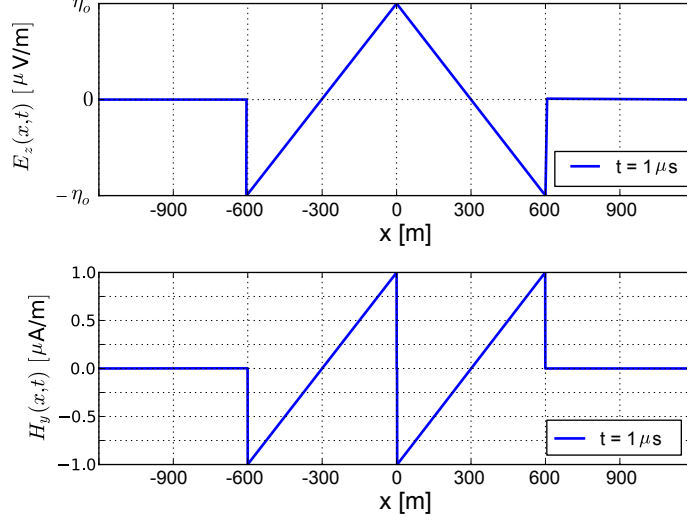
(b) For $t = 5 \mu s$, we have that

$$\begin{aligned}
E_z^\pm(x, t) &= \eta_o \left(5 \mu s \mp \frac{x}{c} \right) \text{rect} \left(\frac{5 \mu s \mp \frac{x}{c}}{2 \mu s} \right) \frac{V}{m} \quad \text{for } x \geq 0 \\
&= \eta_o \left(\frac{1.5 \text{ km} \mp x}{c} \right) \text{rect} \left(\frac{1.5 \text{ km} \mp x}{0.6 \text{ km}} \right) \frac{V}{m} \quad \text{for } x \geq 0, \\
H_y^\pm(x, t) &= \mp \left(5 \mu s \mp \frac{x}{c} \right) \text{rect} \left(\frac{5 \mu s \mp \frac{x}{c}}{2 \mu s} \right) \frac{A}{m} \quad \text{for } x \geq 0 \\
&= \mp \left(\frac{1.5 \text{ km} \mp x}{c} \right) \text{rect} \left(\frac{1.5 \text{ km} \mp x}{0.6 \text{ km}} \right) \frac{A}{m} \quad \text{for } x \geq 0.
\end{aligned}$$



(c) For $t = 1 \mu s$, we have that

$$\begin{aligned}
E_z^\pm(x, t) &= \eta_o \left(1 \mu s \mp \frac{x}{c} \right) \text{rect} \left(\frac{1 \mu s \mp \frac{x}{c}}{2 \mu s} \right) \frac{V}{m} \quad \text{for } x \geq 0 \\
&= \eta_o \left(\frac{300 \text{ m} \mp x}{c} \right) \text{rect} \left(\frac{300 \text{ m} \mp x}{600 \text{ m}} \right) \frac{V}{m} \quad \text{for } x \geq 0, \\
H_y^\pm(x, t) &= \mp \left(1 \mu s \mp \frac{x}{c} \right) \text{rect} \left(\frac{1 \mu s \mp \frac{x}{c}}{2 \mu s} \right) \frac{A}{m} \quad \text{for } x \geq 0 \\
&= \mp \left(\frac{300 \text{ m} \mp x}{c} \right) \text{rect} \left(\frac{300 \text{ m} \mp x}{600 \text{ m}} \right) \frac{A}{m} \quad \text{for } x \geq 0.
\end{aligned}$$



(d) The TEM wave power injected per unit volume is (see lecture 20). Recall J_s exists only at $x = 0$:

$$\begin{aligned} -\vec{J}_s \cdot \vec{E} &= -\left(-\hat{z} 2t \operatorname{rect}\left(\frac{t}{\tau}\right)\right) \cdot \left(\eta_o t \operatorname{rect}\left(\frac{t}{\tau}\right) \hat{z}\right) \\ &= 2\eta_o t^2 \operatorname{rect}^2\left(\frac{t}{\tau}\right) \frac{\text{W}}{\text{m}^2}. \end{aligned}$$

Then, the TEM wave density energy is

$$\begin{aligned} \int -\vec{J}_s \cdot \vec{E} dt &= \int 2\eta_o t^2 \operatorname{rect}^2\left(\frac{t}{\tau}\right) dt \\ &= \int_{-\tau/2}^{\tau/2} 2\eta_o t^2 dt = \frac{2}{3}\eta_o [t^3]_{-\tau/2}^{\tau/2} \\ &= \frac{1}{6}\eta_o \tau^3 = 160\pi \times 10^{-18} \frac{\text{J}}{\text{m}^2}. \end{aligned}$$

4. a) Wave: $\vec{E}_1 = -4 \cos(\omega t - \beta z) \hat{y} \frac{\text{V}}{\text{m}}$

- i. The electric and magnetic fields (\vec{E} and \vec{H}) of a uniform plane wave are orthogonal to each other and to their direction of propagation, thus, it can be verified that such fields satisfy the following relation

$$\vec{E} = \eta \vec{H} \times \hat{\beta},$$

where $\hat{\beta}$ is the unit vector parallel to the propagation direction and η is the intrinsic impedance. Using this relation and considering $\eta_o \approx 120\pi \Omega$ (free space), we can find the expressions for \vec{H} or \vec{E} that accompany the given wave fields,

$$\hat{\beta}_1 = \hat{z} \rightarrow \vec{H}_1 = \frac{4}{\eta_o} \cos(\omega t - \beta z) \hat{x} \frac{\text{A}}{\text{m}}$$

- ii. The instantaneous power flow density is given by the Poynting vector $\vec{S} = \vec{E} \times \vec{H}$. Therefore, the instantaneous power that crosses some surface A is given by $P = \int_A \vec{S} \cdot d\vec{A}$. In the case of uniform plane waves, this expression simplifies to

$$P = \vec{S} \cdot \hat{n} A,$$

where \hat{n} is the vector normal to the flat area A . Below, we are considering $A = 1 \text{ m}^2$ and $\hat{n} = \hat{z}$. For this case:

$$\vec{S}_1 = \vec{E}_1 \times \vec{H}_1 = \frac{16}{\eta_o} \cos^2(\omega t - \beta z) \hat{z} \frac{\text{W}}{\text{m}^2},$$

thus,

$$P_1 = \frac{16}{\eta_o} \cos^2(\omega t - \beta z) \text{ W}.$$

- iii. We can calculate the time-average of the Poynting vector using the trig. identity: $\cos^2 \theta = \frac{1}{2} (1 + \cos(2\theta))$, and the fact that the time average of the cosine wave is zero ($\frac{1}{T} \int_T \cos(\omega t) dt = 0$):

$$\langle \vec{S}_1 \rangle = \left\langle \frac{16}{\eta_o} \cos^2(\omega t - \beta z) \hat{z} \frac{\text{W}}{\text{m}^2} \right\rangle = \frac{8}{\eta_o} \hat{z} \frac{\text{W}}{\text{m}^2}$$

Therefore, the average power that crosses some surface A is given by $\langle P \rangle = \int_A \langle \vec{S} \rangle \cdot d\vec{A}$. In the case of uniform plane waves, this expression simplifies to

$$\langle P \rangle = \langle \vec{S} \rangle \cdot \hat{n} A,$$

where \hat{n} is the vector normal to the flat area A . Below, we are considering $A = 1 \text{ m}^2$ and $\hat{n} = \hat{z}$. For this case:

$$\langle P_1 \rangle = \frac{8}{\eta_o} \text{ W}.$$

- b) Wave: $\vec{E}_2 = E_o (\cos(\omega t - \beta x) \hat{y} - \sin(\omega t - \beta x) \hat{z}) \frac{\text{V}}{\text{m}}$
i. $\hat{\beta}_2 = \hat{x} \rightarrow \vec{H}_2 = \frac{E_o}{\eta_o} (\cos(\omega t - \beta x) \hat{z} + \sin(\omega t - \beta x) \hat{y}) \frac{\text{A}}{\text{m}}$
ii. The Poynting vector is

$$\begin{aligned} \vec{S}_2 &= \vec{E}_2 \times \vec{H}_2 \\ &= E_o (\cos(\omega t - \beta x) \hat{y} - \sin(\omega t - \beta x) \hat{z}) \times \frac{E_o}{\eta_o} (\cos(\omega t - \beta x) \hat{z} + \sin(\omega t - \beta x) \hat{y}) \\ &= \frac{E_o^2}{\eta_o} (\cos^2(\omega t - \beta x) \hat{z} + \sin^2(\omega t - \beta x) \hat{x}) \\ &= \frac{E_o^2}{\eta_o} \hat{x} \frac{\text{W}}{\text{m}^2}. \end{aligned}$$

The wave is propagating in the $+x$ direction, therefore there is no flux of energy flowing into the z direction. Therefore, the instantaneous power crossing a 1 m^2 area in the xy - plane from $-z$ to z is

$$P_2 = 0 \text{ W},$$

- iii. The time-average power crossing the xy - plane from $-z$ to z is

$$\langle P_2 \rangle = 0 \text{ W}.$$

- c) Wave: $\vec{H}_3 = \cos(\omega t + \beta z + \frac{\pi}{3}) \hat{x} + \sin(\omega t + \beta z - \frac{\pi}{6}) \hat{y} \frac{\text{A}}{\text{m}}$
i. $\hat{\beta}_3 = -\hat{z} \rightarrow \vec{E}_3 = \eta_o (\cos(\omega t + \beta z + \frac{\pi}{3}) \hat{y} - \sin(\omega t + \beta z - \frac{\pi}{6}) \hat{x}) \frac{\text{V}}{\text{m}}$

ii. The Poynting vector is

$$\begin{aligned}
\vec{S}_3 &= \vec{E}_3 \times \vec{H}_3 \\
&= \eta_o \left(\cos(\omega t + \beta z + \frac{\pi}{3})\hat{y} - \sin(\omega t + \beta z - \frac{\pi}{6})\hat{x} \right) \times \\
&\quad \left(\cos(\omega t + \beta z + \frac{\pi}{3})\hat{x} + \sin(\omega t + \beta z - \frac{\pi}{6})\hat{y} \right) \\
&= -\eta_o \left(\cos^2(\omega t + \beta z + \frac{\pi}{3})\hat{z} + \sin^2(\omega t + \beta z - \frac{\pi}{6})\hat{z} \right) \\
&= -\eta_o \left(\cos^2(\omega t + \beta z + \frac{\pi}{3})\hat{z} + \cos^2(\omega t + \beta z + \frac{\pi}{3})\hat{z} \right) \\
&= -2\eta_o \cos^2(\omega t + \beta z + \frac{\pi}{3})\hat{z} \frac{W}{m^2}.
\end{aligned}$$

Therefore, the instantaneous power crossing a 1 m^2 area in the xy - plane from $-z$ to z is

$$P_3 = -2\eta_o \cos^2(\omega t + \beta z + \frac{\pi}{3})\hat{z} \frac{W}{m^2} \cdot \hat{z} m^2 = -2\eta_o \cos^2(\omega t + \beta z + \frac{\pi}{3}) W.$$

iii. The time-average power crossing a 1 m^2 area in the xy - plane from $-z$ to z is

$$\langle P_3 \rangle = -\eta_o W.$$

d) Wave: $\vec{H}_4 = \cos(\omega t - \beta z - \frac{\pi}{2})\hat{x} + \sin(\omega t - \beta z)\hat{y} \frac{A}{m}.$

i. $\hat{\beta}_4 = \hat{z} \rightarrow \vec{E}_4 = \eta_o (-\cos(\omega t - \beta z - \frac{\pi}{2})\hat{y} + \sin(\omega t - \beta z)\hat{x}) = \eta_o \sin(\omega t - \beta z)(-\hat{y} + \hat{x}) \frac{V}{m}.$

ii. The Poynting vector is

$$\begin{aligned}
\vec{S}_4 &= \vec{E}_4 \times \vec{H}_4 \\
&= \eta_o \sin(\omega t - \beta z)(-\hat{y} + \hat{x}) \times [\cos(\omega t - \beta z - \frac{\pi}{2})\hat{x} + \sin(\omega t - \beta z)\hat{y}] \\
&= 2\eta_o \sin^2(\omega t - \beta z)\hat{z}
\end{aligned}$$

Therefore, the instantaneous power crossing the area $A = 1 \text{ m}^2$ is

$$P_4 = 2\eta_o \sin^2(\omega t - \beta z) W.$$

iii. The time-average power crossing a 1 m^2 area in the xy - plane from $-z$ to z is

$$\langle P_4 \rangle = \eta_o W.$$

5. (a) The propagation velocity is

$$v_p = \frac{1}{\sqrt{\epsilon\mu}} = \frac{c}{\sqrt{\epsilon_r\mu_r}} = \frac{c}{\sqrt{\frac{9}{4}}} = 2 \times 10^8 \frac{m}{s}.$$

(b) The wave number is

$$\beta = \frac{\omega}{v_p} = \frac{6\pi \times 10^8}{2 \times 10^8} = 3\pi \frac{\text{rad}}{m}.$$

The wavelength is

$$\lambda = \frac{2\pi}{\beta} = \frac{2}{3} m.$$

(c) The intrinsic impedance is

$$\eta = \sqrt{\frac{\mu}{\epsilon}} = \sqrt{\frac{\mu_r}{\epsilon_r}} \eta_o = \sqrt{\frac{4}{9}} \eta_o = 80\pi \Omega.$$

(d) The electric field is

$$\begin{aligned}\mathbf{E} &= -\eta \cos[\omega t \mp \beta(x-2)] \hat{y} \frac{V}{m}, \quad \text{for } x \geq 2, \\ \tilde{\mathbf{E}} &= -\eta e^{\mp j\beta(x-2)} \hat{y} \frac{V}{m}, \quad \text{for } x \geq 2,\end{aligned}$$

and the magnetic field is

$$\begin{aligned}\mathbf{H} &= \mp \cos[\omega t \mp \beta(x-2)] \hat{z} \frac{A}{m}, \quad \text{for } x \geq 2, \\ \tilde{\mathbf{H}} &= \mp e^{\mp j\beta(x-2)} \hat{z} \frac{A}{m}, \quad \text{for } x \geq 2,\end{aligned}$$

(e) At the location of $x = 4\text{m}$, the Poynting vector is

$$\mathbf{S} = \mathbf{E} \times \mathbf{H} = \eta \cos^2[\omega t \mp \beta(x-2)] \hat{x}$$

Therefore, the instantaneous power crossing the area $A = 2\text{ m}^2$ is

$$P = 2\eta \cos^2[\omega t \mp \beta(x-2)] W.$$

And the time-average power crossing the 2 m^2 square surface is

$$\langle P \rangle = \eta W.$$