Quantum harminic oscillator

N. Anders Petersson

Lawrence Livermore National Laboratory¹

April 17, 2019

¹LLNL-PRES-abcdef; This work was performed under the auspices of the U.S. Department of Energy by Lawrence Livermore National Laboratory under Contract DE-AC52-07NA27344. Lawrence Livermore National Security, ŁLC. ■

Notation and lowering operator

Ground state eigenfunction $\psi_0(x)$ and first excited state $\psi_1(x)$ using vector and bra-ket notation

$$\psi_0(x)=\ket{0}=egin{bmatrix}1\0\0\dots\ \end{bmatrix},\quad \psi_1(x)=\ket{1}=egin{bmatrix}0\1\0\dots\ \end{bmatrix},$$

The n^{th} excited state eigenfunction: $\psi_n(x) = |n\rangle$. The lowering operator satisfies $a\psi_n(x) = \sqrt{n}\psi_{n-1}(x)$. In matrix form,

$$a=egin{bmatrix} 0&1&&&&\ &0&\sqrt{2}&&&\ &&0&\sqrt{3}&&\ &&\ddots&\ddots \end{bmatrix},\quad a|n
angle=\sqrt{n}|n-1
angle.$$

Raising and number operators

The raising operator satisfies $a^{\dagger}\psi_n(x) = \sqrt{n+1}\psi_{n+1}(x)$,

$$a^\dagger = egin{bmatrix} 0 & & & & & \ 1 & 0 & & & & \ & \sqrt{2} & 0 & & & \ & & \sqrt{3} & 0 & & \ & & & \ddots & \ddots \ \end{pmatrix}, \quad a^\dagger | {\it n}
angle = \sqrt{n+1} | {\it n}+1
angle.$$

The number operator, N,

Hamiltonian and energy

Hamiltonian of a quantum harmonic oscillator in operator form,

$$H = \hbar\omega \left(a^{\dagger}a + \frac{1}{2} \right) = \hbar\omega \left(N + \frac{1}{2}I \right)$$

and in matrix form

$$H=rac{\hbar\omega}{2}egin{bmatrix}1&&&&&&\\&3&&&&\\&&&5&&&\\&&&&7&&\\&&&&\ddots\end{bmatrix}$$

The Hamiltonian of and eigenfunction $\psi_n(x)$,

$$H\psi_n = \hbar\omega \left(N + \frac{1}{2}I\right)\psi_n = \hbar\omega \left(n + \frac{1}{2}\right)\psi_n,$$

gives the energy level, $E_n = \hbar\omega \left(n + \frac{1}{2}\right)$.



Coupled oscillators

The system Hamiltonian for Q qudits is

$$H_0 = \sum_{j=0}^{Q-1} \left(\omega_j a_j^{\dagger} a_j - \chi_{jj} a_j^{\dagger} a_j^{\dagger} a_j a_j - \sum_{k \neq j} \chi_{jk} a_j^{\dagger} a_j a_k^{\dagger} a_k \right)$$

If we retain L levels in each qudit, the lowering operators are defined in terms of the L by L identity matrix I_L . If Q=2,

$$a_0 = a \otimes I_L, \quad a_1 = I_L \otimes a.$$

Thus,

$$\begin{split} a_0^\dagger a_0 &= (a^\dagger \otimes I_L)(a \otimes I_L) = (a^\dagger a) \otimes (I_L I_L) = N \otimes I_L, \\ a_1^\dagger a_1 &= (I_L \otimes a^\dagger)(I_L \otimes a) = (I_L I_L) \otimes (a^\dagger a) = I_L \otimes N. \end{split}$$

The matrices a_0^{\dagger} and a_1^{\dagger} have size L^2 by L^2 .



The system Hamiltonian

The matrices

$$a_0^{\dagger}a_0 = N \otimes I_L =: N_0,$$

 $a_1^{\dagger}a_1 = I_L \otimes N =: N_1,$

are both diagonal. Furthermore,

$$\begin{split} & a_0^{\dagger} a_0^{\dagger} a_0 a_0 = \textit{N}_0 \, \textit{N}_0 - \textit{N}_0, \\ & a_1^{\dagger} a_1^{\dagger} a_1 a_1 = \textit{N}_1 \, \textit{N}_1 - \textit{N}_1, \end{split}$$

are also diagonal. Thus all terms in the system Hamiltonian

$$H_0 = \sum_{j=0}^{Q-1} \left(\omega_j N_j - \chi_{jj} (N_j^2 - N_j) - \sum_{k \neq j} \chi_{jk} N_j N_k \right)$$

are diagonal.



The control Hamiltonian

The control Hamiltonian satisfies

$$H_c(t) = \sum_{k=0}^{Q-1} F(t) \left(a_k^{\dagger} + a_k \right) + i \sum_{k=0}^{Q-1} G(t) \left(a_k^{\dagger} - a_k \right),$$

where F and G are real-valued functions. These matrices are not diagonal, but block-diagonal. For example,

$$a_0^{\dagger} + a_0 = (a \otimes I_L) + (a^{\dagger} \otimes I_L) =$$

$$\begin{bmatrix} 0 & I_L \\ I_L & 0 & \sqrt{2}I_L \\ & \sqrt{2}I_L & 0 & \sqrt{3}I_L \\ & & & \ddots & \ddots \end{bmatrix}$$

The control Hamiltonian 2

Similarly,

$$a_1^{\dagger} + a_1 = (I_L \otimes a) + (I_L \otimes a^{\dagger}) =$$

$$\begin{bmatrix} a^{\dagger} + a & 0 & & & & \\ 0 & a^{\dagger} + a & 0 & & & \\ & 0 & a^{\dagger} + a & 0 & & \\ & & 0 & a^{\dagger} + a & \ddots & \\ & & & \ddots & \ddots & \ddots \end{bmatrix}$$

is a block-diagonal matrix.