

Group Meeting 07.05

Three-body reaction theory in a model space

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Overview

The Schrodinger equation of three-body System

$$\left[E - K - V(r) - U_1(r_1) - U_2(r_2) \right] \psi = 0,$$

The “Distorted-wave” Faddeev differential equations

$$\left[E - K - V(r) - \mathcal{P}_\lambda (U_1 + U_2) \mathcal{P}_\lambda \right] \psi_\lambda = V [\psi_{1\lambda} + \psi_{2\lambda}]$$

$$\left[E - K - U_1(r_1) \right] \psi_{1\lambda} = \mathcal{P}_\lambda \left[U_1 - \mathcal{P}_\lambda U_1 \mathcal{P}_\lambda \right] \psi_\lambda + \mathcal{P}_\lambda U_1 \psi_{2\lambda}$$

$$\left[E - K - U_2(r_2) \right] \psi_{2\lambda} = \mathcal{P}_\lambda \left[U_2 - \mathcal{P}_\lambda U_2 \mathcal{P}_\lambda \right] \psi_\lambda + \mathcal{P}_\lambda U_2 \psi_{1\lambda},$$

where $\psi = \psi_\lambda + \psi_{1\lambda} + \psi_{2\lambda}$

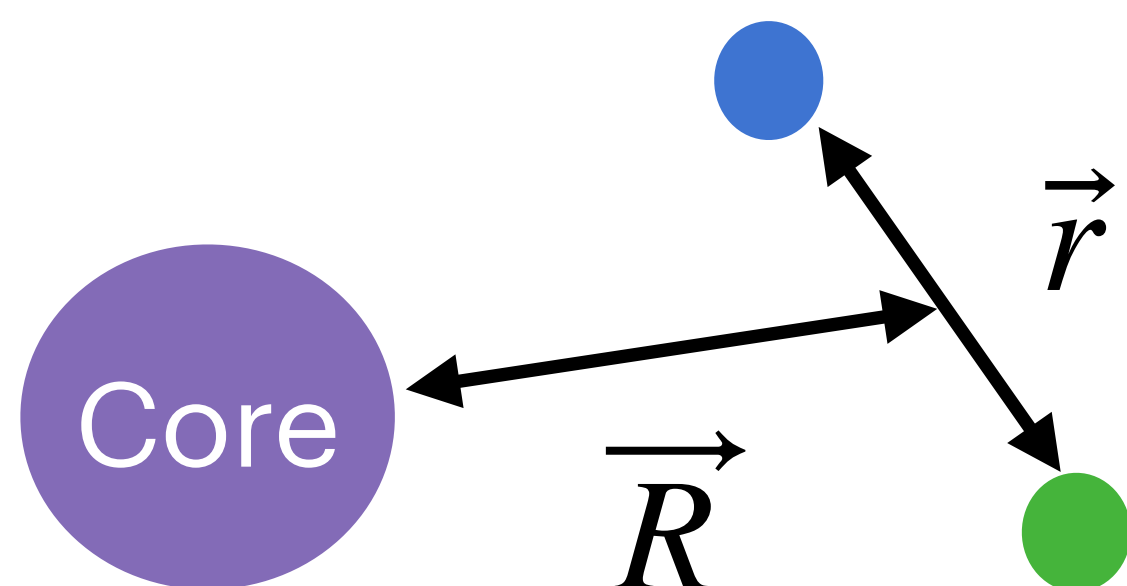
Overview

Expand it in the set of basis $|k[lL]JM) \equiv |\phi_l(k, r) [Y_l(\hat{r}), Y_L(\hat{R})]_{JM})$,

$$[E - K - V(r)] \psi_\lambda = \mathcal{P}_\lambda (U_1 + U_2) \mathcal{P}_\lambda \psi_\lambda + V [\psi_{1\lambda} + \psi_{2\lambda}]$$



$$[E - \epsilon(k) - K_R] g_{\lambda LL}^J[P(k), \vec{R}] = \int d^3r d\hat{R} \left\{ \phi_l(k', r) [Y_1(r), Y_L(R)]_{JM} \right\}^* \\ \left\{ \mathcal{P}_\lambda (U_1 + U_2) \mathcal{P}'_\lambda \psi' + V [\psi_{1\lambda} + \psi_{2\lambda}] \right\}$$



$$U_\lambda \equiv \mathcal{P}_\lambda (U_1 + U_2) \mathcal{P}'_\lambda$$

$$\xi_\lambda \equiv \psi_{1\lambda} + \psi_{2\lambda}$$

Overview

$$\left[E - K - U_1(r_1) \right] \psi_{1\lambda} = \mathcal{P}_\lambda \left[U_1 - \mathcal{P}_\lambda U_1 \mathcal{P}_\lambda \right] \psi_\lambda + \mathcal{P}_\lambda U_1 \psi_{2\lambda}$$

$$\left[E - K - U_2(r_2) \right] \psi_{2\lambda} = \mathcal{P}_\lambda \left[U_2 - \mathcal{P}_\lambda U_2 \mathcal{P}_\lambda \right] \psi_\lambda + \mathcal{P}_\lambda U_2 \psi_{1\lambda},$$

With “Distorted-wave” Faddeev differential equations, they get

$$\left[E - K - U_1 - U_2 \right] \xi_\lambda = \left[U_1 + U_2 - U_\lambda \right] \psi_\lambda,$$

$$\left[E - K - V - U_\lambda \right] \psi_\lambda = V \xi_\lambda,$$

$$U \equiv U_1 + U_2$$

$$\left[E - K - U \right] \xi_\lambda = \left[U - U_\lambda \right] \psi_\lambda,$$

$$\left[E - K - V - U_\lambda \right] \psi_\lambda = V \xi_\lambda,$$

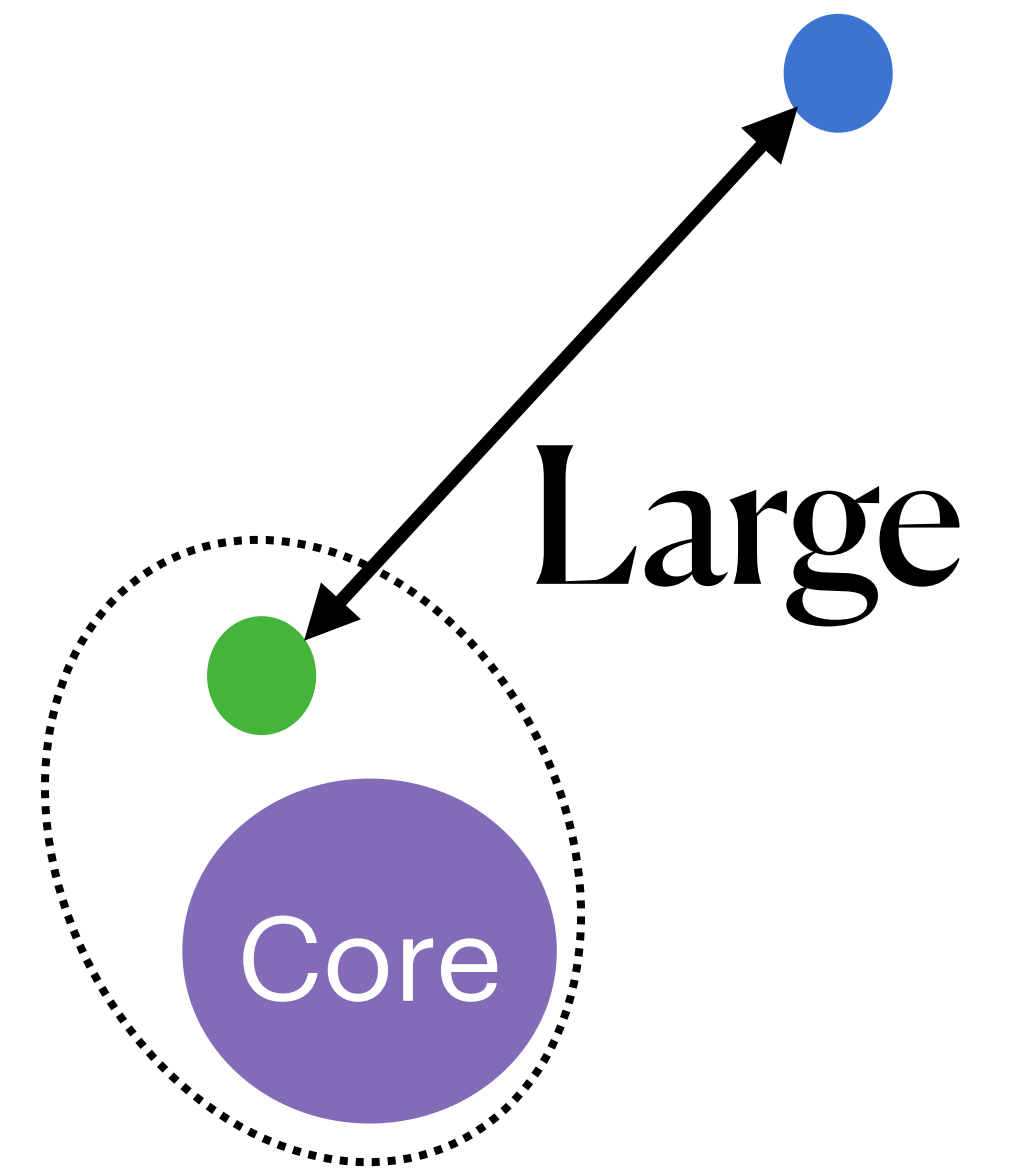
Overview

$\xi_\lambda \equiv \psi_{1\lambda} + \psi_{2\lambda}$ just has outgoing wave, Source term

$$\xi_\lambda = [E^+ - K - U]^{-1} [U - U_\lambda] \psi_\lambda,$$

This term will be small in the range of V .

Source term show there will be a 'hole' near the origin, where the low multipoles of U cancel with those of U_λ .



CDCC

$$(E - K - V - U)\psi = 0$$



$$(E - K - V - \mathcal{P}U)\mathcal{P}\psi = \mathcal{P}U\mathcal{Q}\psi,$$

$$(E - K - V - \mathcal{P}U)\mathcal{P}\psi = \mathcal{P}U(1 - \mathcal{P})\psi,$$

In the CDCC calculation, they assume

$$(E - K - V - \mathcal{P}U\mathcal{P})\psi^{CDCC} = 0,$$

CDCC

$$(E - K - V - \mathcal{P}U)\mathcal{P}\psi = \mathcal{P}U(1 - \mathcal{P})\psi,$$

$$(E - K - V - \mathcal{P}U\mathcal{P})\psi^{CDCC} = 0,$$

This term must be small.

$$[E - K - V - U_\lambda] \psi_\lambda = V\xi_\lambda,$$

$$\xi_\lambda = [E - K - U]^{-1} [U - U_\lambda] \psi_\lambda$$

Now, however, by the analysis given before, we recognize that $\mathcal{P}_\lambda U(1 - \mathcal{P}_\lambda)\psi$ must be negligible in a domain \mathcal{D} , defined by $R < R_c$. We assume that we can find R_c in such a way that it is small enough to ensure that $\mathcal{P}_\lambda U(1 - \mathcal{P}_\lambda)\psi$ is small in \mathcal{D} , so that partial waves with $l < \lambda$ are not much affected in \mathcal{D} by those with $l > \lambda$. At the same time, R_c is large enough to ensure that the truncated tail of $\mathcal{P}_\lambda U \mathcal{P}_\lambda$ beyond R_c is small and smooth, so that the reflection of $\mathcal{P}_\lambda \psi_\lambda$ from outside is negligible. Then $\mathcal{P}_\lambda \psi$ satisfies approximately the same equation in \mathcal{D} and approximately the same boundary conditions at $R = R_c$ as ψ_λ , so that $\psi_\lambda \approx \mathcal{P}_\lambda \psi$ in \mathcal{D} . The larger λ is, the larger R_c can be and the closer ψ_λ is to ψ .