

# In Search For The Optimal Phase Hologram



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## **Declaration**

I hereby declare that except where specific reference is made to the work of others, the contents of this dissertation are original and have not been submitted in whole or in part for consideration for any other degree or qualification in this, or any other university. This dissertation is my own work and contains nothing which is the outcome of work done in collaboration with others, except as specified in the text and Acknowledgements. This dissertation contains fewer than 65,000 words including appendices, bibliography, footnotes, tables and equations and has fewer than 150 figures.

Jinze Sha

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## **Acknowledgements**

And I would like to acknowledge ...



## **Abstract**

This is where you write your abstract ...



# List of Publications

- [1] Jana Skirnewskaja, Yunuen Montelongo, Jinze Sha, and Timothy D. Wilkinson. Holographic lidar projections with brightness control. In *Imaging and Applied Optics Congress 2022 (3D, AOA, COSI, ISA, pcAOP)*, page 3F2A.6. Optica Publishing Group, 2022
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- [7] Roubing Meng, Jinze Sha, Zhongling Huang, and Timothy D. Wilkinson. Extending FOV of holographic display with alternating lasers. In Peter Schelkens and Tomasz Kozacki,

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# **Chapter 1**

## **Introduction**

The pursuit of three-dimensional (3D) display has never stopped. Currently, most commercially available so-called ‘3D display’ products such as 3D cinema, 3D TV, handheld 3D devices (e.g. Nintendo 3DS, HTC Evo 3D) and Virtual Reality (VR) and Augmented Reality (AR) head sets are in fact stereoscopic displays where two different two-dimensional (2D) images are displayed to the left and right eye respectively, creating a 3D illusion in the brain. Despite its high image quality, the major issue with stereoscopic displays is that they cannot provide real defocusing effect in depth. Modern 3D cinema are able to provide good comfort because the polarisation glasses are as light as regular glasses, and the variable defocusing issue can be avoided by the combination of good design of point of interest in each scene and the according defocusing effect as captured by the camera, so most audience won’t experience much discomfort for around 2 to 3 hours. But the content, viewing angle and depth of focus are fixed at how they are captured. To provide an interactive and real-time rendered immersive experience, the VR/AR headset has frequently been advertised as the ‘gateway to metaverse’ in recent years. However, my personal experience with VR headset is far from comfortable, not only because of its heavy weight, but also because the display is physically at a very near distance, while my brain thinks the objects are at various distances and yet are still all in focus, which is very unnatural, because in real life, when the eye is focused on a near object, the far backgrounds would blur out. And also, the two displays in the VR headset needs to be rendered in real-time based on the location and angle of the user, which is nowhere near practical. Hence, the heavy weight, the lack of defocusing, the delay between the rendering and the change in my position are the three major factors causing my dizziness using VR headsets, either of which is quite impractical to solve, especially the

weight issue. Only if VR/AR headsets could be reduced to the weight of eyeglasses would I ever consider the possibility of those head-mounted devices leading us to the ‘metaverse’.

In comparison, the holography technique can produce the full 3D light field, which does not rely on any head mounted device, has the true depth of focus, and does not need to re-render according to change in viewer positions and viewing angle.



Fig. 1.1 A photo of the holographic portrait of Dennis Gabor [10]

Holography, taking its name from the Greek word *ολόσ* (holos), meaning *whole*, was first introduced in 1948 by Dennis Gabor [20], originally named as *wavefront reconstruction* [21]. It is a cool technology which generates 3D images via the diffraction of light. Similar to 2D photography, the earliest holography uses a piece of film to record the diffraction pattern, which can then reconstruct the 3D field, as shown in Fig. 1.1 which is a holographic recording of Dennis Gabor himself. After the invention of digital cameras, digital holography emerged. The limitation of both methods is that they require a physical object as a priori to record the hologram. In order to generate hologram for objects that do not physically exist, computer-generated holography (CGH) emerged where a hologram can be calculated through various algorithmic approaches and then displayed on a spatial light modulator (SLM) in order to create an image in the replay field through diffraction [15, 22, 23, 19]. Although being a fancy technology of true 3D display, CGH still has some fundamental issues, mainly its image quality and the heavy computation required, the solutions of which are my ultimate goals.

# **Chapter 2**

## **Literature Review**

This chapter lays down the fundamental theories of optoelectronics and search algorithms, which are essential to the research outlined in the later chapters of this thesis.

### **2.1 The Nature of Light**

#### **2.1.1 Wave-Particle Duality**

The problem of how light propagates has been troubling scientists for centuries. The journey began with Sir Isaac Newton, who advocated the particle theory of light in the 17th century. Newton proposed that light consists of particles, or ‘corpuscles’, which explained phenomena such as reflection and refraction [24]. In contrast, Christiaan Huygens, a contemporary of Newton, demonstrated that light behaves as a wave, as it is capable of diffraction and interference [25].



Fig. 2.1 An illustration of the Young's Double-slit experiment [11]

The wave theory gained significant support in the early 19th century through the experiments of Thomas Young. Young's double-slit experiment in 1801 provided clear evidence of the wave nature of light by showing that light passing through two slits creates an interference pattern on a screen [26], as illustrated in Fig. 2.1. Augustin-Jean Fresnel further advanced the wave theory by developing a comprehensive mathematical framework to describe light as a wave, explaining phenomena such as polarization and the diffraction of light [27].

This understanding was further reinforced by James Clerk Maxwell in the mid-19th century. In 1864, James Clerk Maxwell organised a set of four equations describing the space and time dependence of the electromagnetic field, which are:

$$\nabla \times \mathbb{E} = -\frac{\partial \mathbb{B}}{\partial t} \quad (2.1)$$

$$\nabla \times \mathbb{H} = \mathbb{J} + \frac{\partial \mathbb{D}}{\partial t} \quad (2.2)$$

$$\nabla \cdot \mathbb{D} = \rho \quad (2.3)$$

$$\nabla \cdot \mathbb{B} = 0 \quad (2.4)$$

where  $\mathbb{D}$  is the electric flux density,  $\mathbb{E}$  is the electric field intensity,  $\mathbb{B}$  is the magnetic flux density,  $\mathbb{H}$  is the magnetic field intensity,  $\rho$  is the volume charge density, and  $\mathbb{J}$  is the current density [28].

And the relation between  $\mathbb{D}$  and  $\mathbb{E}$  and between  $\mathbb{B}$  and  $\mathbb{H}$  for linear materials are:

$$\mathbb{B} = \mu \mathbb{H} \quad (2.5)$$

$$\mathbb{D} = \epsilon \mathbb{E} \quad (2.6)$$

where  $\mu$  is the magnetic permeability and  $\epsilon$  is the dielectric permittivity of the material [29].

Maxwell's equations unified electricity and magnetism into a single theory of electromagnetism, predicting that light is an electromagnetic wave that propagates through space [30].

Despite the success of the wave theory, it could not explain all light-related phenomena. The early 20th century brought a pivotal development with Albert Einstein's explanation of the photoelectric effect. In 1905, Einstein proposed that light also behaves as particles, or 'quanta' (later called photons), which could eject electrons from a metal surface when light is shone upon it [31]. This particle nature of light was critical in explaining observations that wave theory alone could not address and earned Einstein the Nobel Prize in Physics in 1921.

These discoveries collectively revealed that light exhibits both wave and particle properties, depending on the experimental context. This wave-particle duality became a cornerstone of quantum mechanics, fundamentally altering human's understanding of the nature of light. Although to date, it is still not yet known what light exactly is, it is now known how light behaves.

### 2.1.2 Wave Equation

To mathematically describe the propagation of light in free space (i.e. in absence of free charge), the Maxwell equations in Eq. (2.1) - Eq. (2.4) can be simplified as:

$$\nabla \times \mathbb{E} = -\mu \frac{\partial \mathbb{H}}{\partial t} \quad (2.7)$$

$$\nabla \times \mathbb{H} = \epsilon \frac{\partial \mathbb{E}}{\partial t} \quad (2.8)$$

$$\nabla \cdot \epsilon \mathbb{E} = 0 \quad (2.9)$$

$$\nabla \cdot \mu \mathbb{H} = 0 \quad (2.10)$$

Taking the curl on both the left and right hand sides of Eq. (2.7), and using the vector identity of  $\nabla \times (\nabla \times \mathbf{u}) = \nabla(\nabla \cdot \mathbf{u}) - \nabla^2 \mathbf{u}$ , we get:

$$\nabla \times (\nabla \times \mathbb{E}) = -\nabla \times (\mu \frac{\partial \mathbb{H}}{\partial t}) \quad (2.11)$$

$$\nabla(\nabla \cdot \mathbb{E}) - \nabla^2 \mathbb{E} = -\frac{\partial}{\partial t} \nabla \times (\mu \mathbb{H}) \quad (2.12)$$

Then, by substituting Eq. (2.8) and Eq. (2.9) in, Eq. (2.12) becomes:

$$-\nabla^2 \mathbb{E} = -\frac{\partial}{\partial t} (\mu \epsilon \frac{\partial \mathbb{E}}{\partial t}) \quad (2.13)$$

Hence, we have a generic form of wave equation, relating the space and time domain relation of electromagnetic waves propagating in free space:

$$\nabla^2 \mathbb{E} = \mu \epsilon \frac{\partial^2 \mathbb{E}}{\partial t^2} \quad (2.14)$$

A valid solution to Eq. (2.14) is:

$$\mathbb{E} = \mathbb{E}_0 e^{j(\omega t - kr)} \quad (2.15)$$

where  $\omega$  is the angular velocity of the wave,  $t$  is time,  $r$  is the propagation distance and  $k$  is called the wave number ( $k = \frac{2\pi}{\lambda}$ , where  $\lambda$  is the wavelength). From Eq. (2.15) we can see that the propagation of light in free space is essentially a phase shift. This suggests that, if we have a coherent light source and a device to manipulate light (called SLM, further explained in Section 2.2.3), we can produce an interference pattern reconstructing the target field we desire, and such method is called holographic projection.

## 2.2 Fundamentals of Holography

Holography is a technology that can fully reconstruct the wavefront of 3D objects, which is usually achieved by modulating a coherent light source. This section explains what a coherent light source is and how it is modulated and diffracted.

### 2.2.1 Light Source

The mechanism of holographic projection is to control the propagation of light in a way that, after diffraction, reconstructs a wavefront that matches the target field. We usually prefer to start from a coherent light source rather than a random one which will be a lot more difficult or even impossible to analyse and predict the interference pattern.

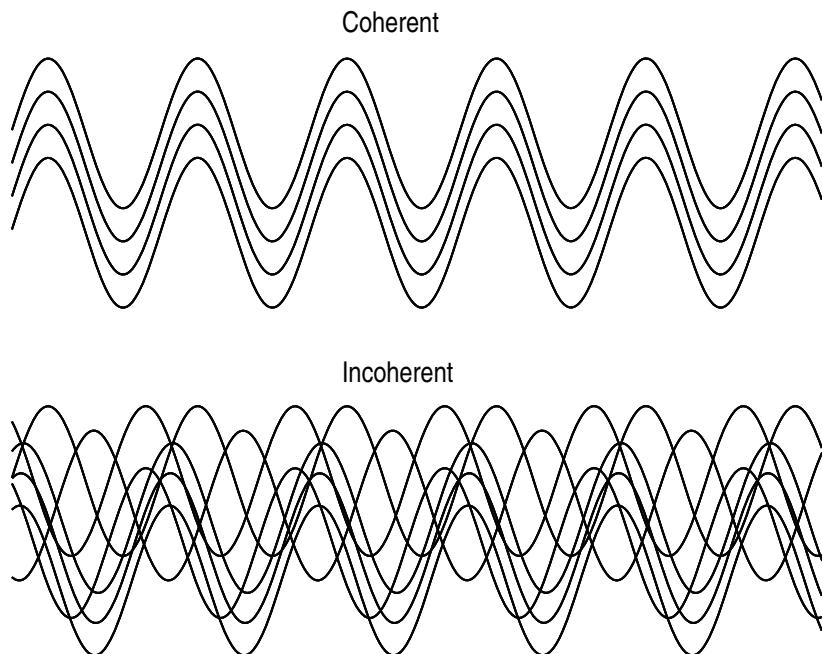


Fig. 2.2 Coherent v.s. incoherent light

The coherence of light refers to the property of light waves where the phase relationship between the waves is consistent over time and space, corresponding to temporal and spatial coherence:

- **Temporal coherence:** Temporal coherence describes the correlation between the phases of a light wave at different points along its propagation direction. It indicates how monochromatic (i.e. single-frequency) a light source is.

- **Spatial coherence:** Spatial coherence describes the correlation between the phases of a light wave at different points across the wavefront, perpendicular to the direction of propagation. It indicates the uniformity of the phase across the wavefront, as illustrated in Fig. 2.2. High spatial coherence means that the light waves across different points on the wavefront are in phase.



Fig. 2.3 Structure of the first laser [12]

The most common coherent light source is Laser, which stands for *Light Amplification by the Stimulated Emission of Radiation*. It was first invented by Theodore Maiman in 1959, with the structure shown in Fig. 2.3 [12, 32, 33]. It differs from other sources of light in that it emits coherent light, which is suitable for holographic projection. However, the coherent and monochromatic property of laser also has a side effect of speckle noise in the reconstructed image [34], which is one of the major problems affecting the image quality of holographic projections and has seen lots of efforts to cope with it in the literature [35–38].

## 2.2.2 Diffraction

This section delves into how light interacts with apertures, leading to diffractions. Understanding diffraction is essential for holography, as it explains how light can be manipulated to reconstruct three-dimensional light fields. The principles of diffraction and interference

underpin the essential process of holographic projection, making it possible to accurately recreate complex wavefronts and achieve true 3D visualization.



Fig. 2.4 Diffraction geometry

To model how light diffracts through a 2D aperture, we first set up a coordinate system as shown in Fig. 2.4, where the aperture is denoted by  $A(x,y)$  and the diffracted field is denoted by  $E(\alpha,\beta,z)$ .  $R$  defines the distance between point  $P$  and the origin of the aperture ( $(x,y) = (0,0)$ ),  $r$  defines the distance between point  $P$  and a point on the aperture, and  $\theta$  defines the angle  $r$  from the  $z$ -axis. Then by trigonometry we can have the following identities:

$$\cos(\theta) = \frac{z}{r} \quad (2.16)$$

$$R^2 = \alpha^2 + \beta^2 + z^2 \quad (2.17)$$

$$r^2 = (\alpha - x)^2 + (\beta - y)^2 + z^2 \quad (2.18)$$



Fig. 2.5 Huygens-Fresnel wavelet principle [13]

The Huygens-Fresnel principle states that every point on a wavefront is itself the source of outgoing secondary spherical wavelets, which can be expressed mathematically as follows when  $r \gg \lambda$  [39]:

$$E(\alpha, \beta, z) = \frac{1}{j\lambda} \iint A(x, y) \frac{e^{jkr}}{r} \cos(\theta) dx dy \quad (2.19)$$

Applying the identities in Eq. (2.16) - Eq. (2.18), Eq. (2.19) becomes:

$$E(\alpha, \beta, z) = \frac{z}{j\lambda} \iint A(x, y) \frac{e^{jkr}}{r^2} dx dy \quad (2.20)$$

$$= \frac{z}{j\lambda} \iint A(x, y) \frac{e^{jk\sqrt{(\alpha-x)^2 + (\beta-y)^2 + z^2}}}{(\alpha-x)^2 + (\beta-y)^2 + z^2} dx dy \quad (2.21)$$

Unfortunately, Eq. (2.21) cannot be solved analytically except for few specific aperture functions  $A(x, y)$ , so we have to make some approximations in order to solve for arbitrary  $A(x, y)$ , the common methods are *Fresnel* and *Fraunhofer* approximations for regions depicted in Fig. 2.6.

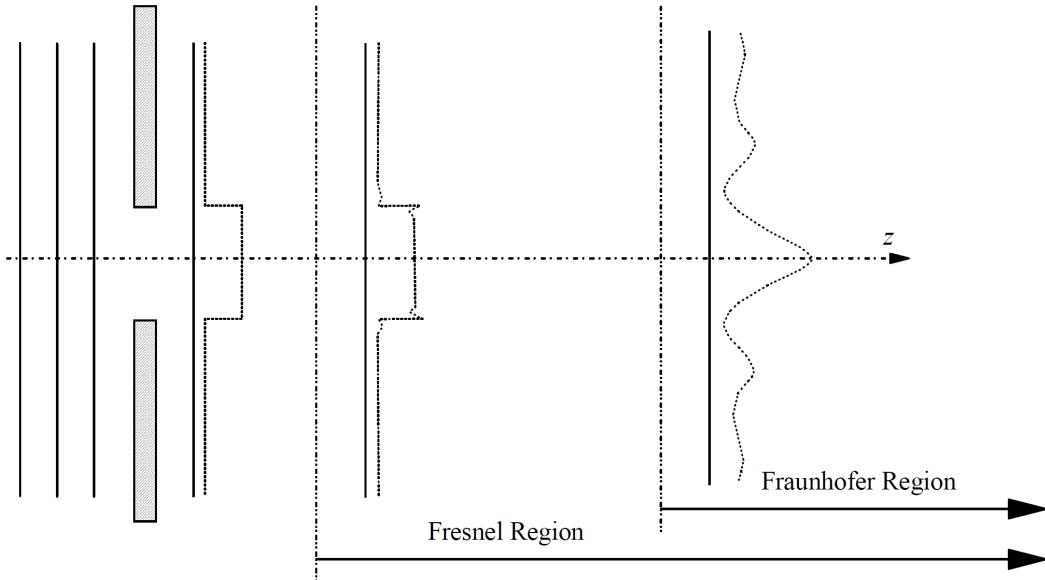


Fig. 2.6 Fresnel and Fraunhofer region [14]

### Fresnel Approximation

$$\sqrt{1+d} = 1 + \frac{1}{2}d - \frac{1}{8}d^2 + \dots \quad (2.22)$$

Fresnel approximation replaces expressions for spherical waves by quadratic-phase exponentials, using the binomial expansion of the square root (given in Eq. (2.22)) to approximate  $r$  in Eq. (2.20) [39].

Retaining only the first two terms of the expansion gives:

$$r = \sqrt{(\alpha-x)^2 + (\beta-y)^2 + z^2} \quad (2.23)$$

$$= z \sqrt{1 + \left(\frac{\alpha-x}{z}\right)^2 + \left(\frac{\beta-y}{z}\right)^2} \quad (2.24)$$

$$\approx z \left[ 1 + \frac{1}{2} \left( \frac{\alpha-x}{z} \right)^2 + \frac{1}{2} \left( \frac{\beta-y}{z} \right)^2 \right] \quad (2.25)$$

For the  $r^2$  in the denominator of Eq. (2.20), the error introduced by dropping all terms but  $z$  is generally acceptably small (i.e.  $r^2 \approx z^2$ ), and for the  $r$  appearing in the exponent in the numerator of Eq. (2.20), errors are much more critical [39]. So, by substituting Eq. (2.25) for

the  $r$  in the numerator of Eq. (2.20) and substituting  $z$  for the  $r$  in the denominator, we have:

$$E(\alpha, \beta, z) \approx \frac{z}{j\lambda} \iint A(x, y) \frac{e^{jkz \left[ 1 + \frac{1}{2} \left( \frac{\alpha-x}{z} \right)^2 + \frac{1}{2} \left( \frac{\beta-y}{z} \right)^2 \right]}}{z^2} dx dy \quad (2.26)$$

$$= \frac{e^{jkz}}{j\lambda z} e^{j\frac{k}{2z}(\alpha^2 + \beta^2)} \iint \left\{ A(x, y) e^{j\frac{k}{2z}(x^2 + y^2)} \right\} e^{-j\frac{2\pi}{\lambda z}(\alpha x + \beta y)} dx dy \quad (2.27)$$

$$= \frac{e^{jkz}}{j\lambda z} e^{j\frac{k}{2z}(\alpha^2 + \beta^2)} \mathcal{F} \left\{ A(x, y) e^{j\frac{k}{2z}(x^2 + y^2)} \right\} \quad (2.28)$$

where  $\mathcal{F}$  is the Fourier Transform. Such method of including Fourier Transform (FT) in the study of optics is also named ‘Fourier Optics’.

Now we have a more simple and solvable expression than Eq. (2.21). And also, as we are only interested in the scaling of relative points at  $P$  with respect to each other, so it is safe to normalise the multiplier term before the Fourier Transform to 1 [14]. So we can express the diffraction pattern in Fresnel region as:

$$E_{\text{Fresnel region}}(\alpha, \beta, z) = \mathcal{F} \left\{ A(x, y) e^{j\frac{k}{2z}(x^2 + y^2)} \right\} \quad (2.29)$$

### Fraunhofer Approximation

Fraunhofer diffraction is a form of diffraction in which the distance between the light source and the receiving screen are in effect at infinite, so that the wave fronts can be treated as planar rather than spherical [28]. Fraunhofer approximation is very stringent, it assumes that the distance between the light source and the receiving screen are in effect at infinite:

$$z \gg \frac{k(x^2 + y^2)_{\max}}{2} \quad (2.30)$$

so that the wave fronts can be treated as planar rather than spherical [28], then the  $e^{j\frac{k}{2z}(x^2 + y^2)}$  term tends to 1, and Eq. (7.1) becomes:

$$E_{\text{Fraunhofer region}}(\alpha, \beta) = \mathcal{F} \{ A(x, y) \} \quad (2.31)$$

which suggests that the far field pattern is simply the Fourier Transform of the aperture function.

### 2.2.3 Spatial Light Modulator (SLM)

SLMs are critical components in computer-generated holography (CGH). SLM is a device used to control the amplitude or phase of light waves in a spatially varying manner. SLMs typically consist of a two-dimensional(2D) array of pixels, each of which can modulate the light either passing through or reflected from it. These pixels are usually addressed by electronic signals, allowing precise manipulation of the light wavefront.

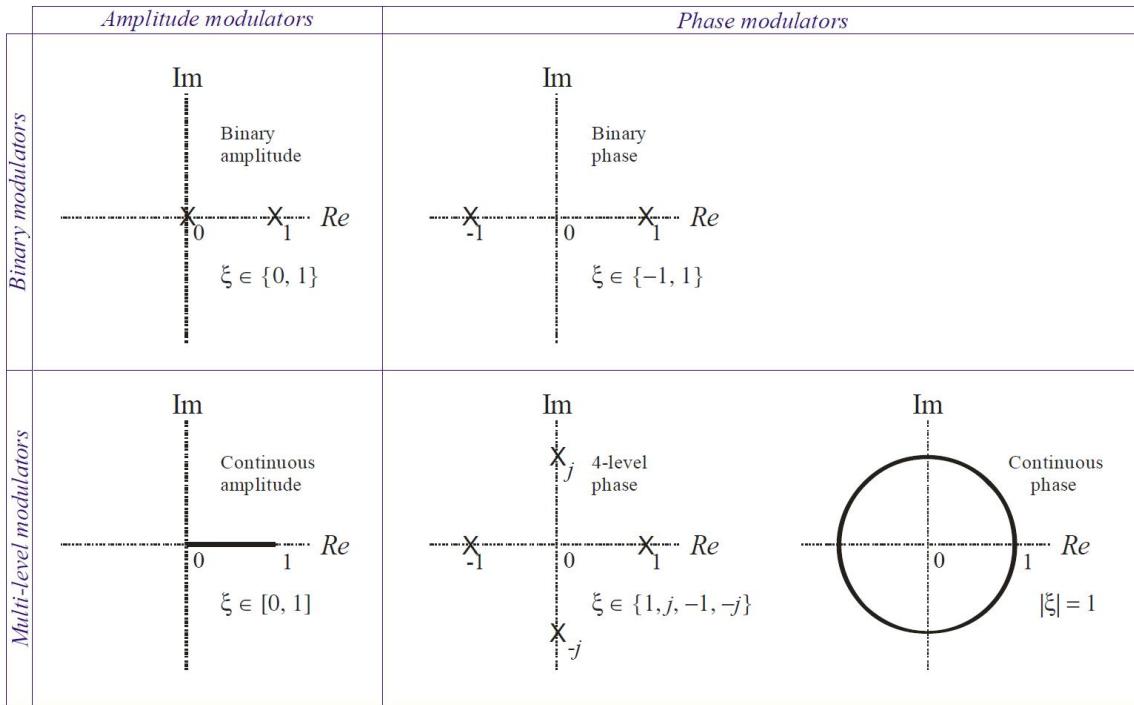


Fig. 2.7 Modulation loci in the complex plane [15]

The modulation can be achieved through various mechanisms, such as liquid crystal SLMs(LC-SLM), magneto-optic SLMs, deformable mirror SLMs, multiple-quantum-well SLMs, or acousto-optic Bragg cells [39], which all fall into the four modulation categories, as illustrated in Fig. 2.7 [15]. The four modulation schemes are:

- **Multi-level Amplitude** modulators can modulate each pixel from zero transmission (0) to full transmission (1), either continuously or in discrete steps. (e.g. nematic liquid crystal display [40], found for example in laptops and many conventional video projectors)

- **Binary Amplitude** modulators can switch each pixel to zero transmission (0) or full transmission (1), but nothing in between. (e.g. deformable mirror device [41], ferroelectric liquid crystal display [42], both used in high-end video projectors)
- **Multi-level Phase** modulators can modulate the phase shift imparted by each pixel from 0 to  $2\pi$  radians, either continuously or in discrete steps. (e.g. Nematic liquid crystal devices [43])
- **Binary Phase** modulators can switch each pixel for a phase shift of either 0 or  $\pi$  radians. (e.g. Ferroelectric liquid crystal displays [44])

Among the four modulation schemes, phase modulations are of higher interests for the purpose of holography, because amplitude modulators, either multi-level or binary, blocks light at the SLM, causing waste of energy, leading to poorer energy efficiency. And also, amplitude modulations always have a zero-order (forming a central bright spot), because the average amplitude is always between 0 and 1; on contrary, phase modulation can suppress the zero-order by designing the hologram to have zero average.

As there is no complex modulator available yet, we need algorithms to generate phase-only holograms, such process is called phase retrieval. There are currently many algorithms for such purpose, which will be discussed in Section 2.3.

### **Rotational symmetries in the binary phase modulation**

The spatial light modulator (SLM) used in this thesis is a binary phase modulator. As the binary phase modulation is purely real (i.e. it's only switching between  $0^\circ$  and  $180^\circ$ , corresponding to 1 and -1 values of  $A^*(x,y)$ ), the complex conjugate  $A^*(x,y)$  is the same as  $A(x,y)$ :

$$A^*(x,y) = A(x,y) \quad (2.32)$$

because the Fourier transform of  $A^*(x,y)$  is the same as the Fourier transform of  $A(x,y)$

$$E(-\alpha, -\beta) = \mathcal{F}[A^*(x,y)] = \mathcal{F}[A(x,y)] = E(\alpha, \beta) \quad (2.33)$$

So, at the Fraunhofer region, there is no distinction between the desired image and its  $180^\circ$  rotation in the replay field, causing a symmetrical conjugate image rotated  $180^\circ$  from the target image.

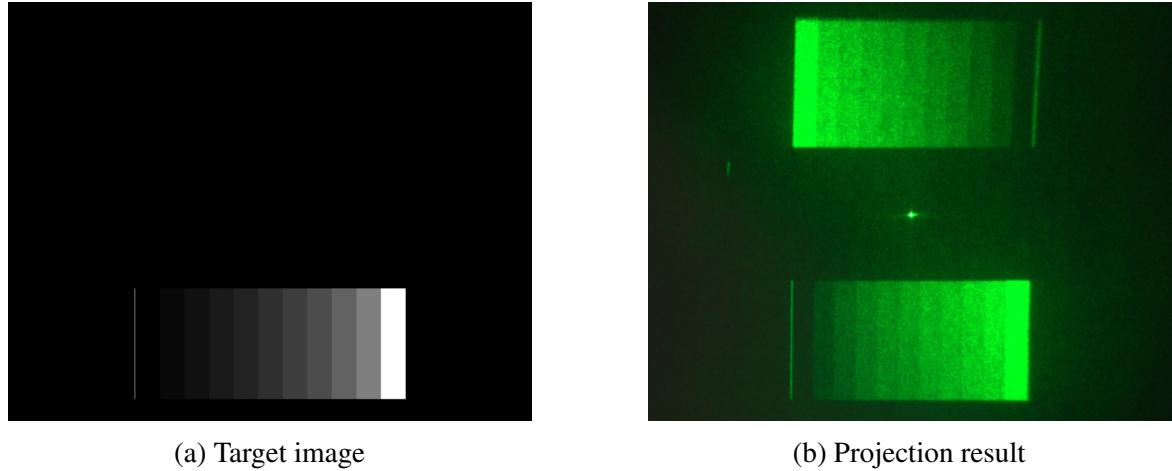


Fig. 2.8 Rotational symmetry in the projection result using the binary phase SLM

To demonstrate such phenomena, the example target image shown in Fig. 2.8a ran through the binary-phase hologram generation algorithm called one-step phase-retrieval (OSPR), which will be explained in detail in Section 2.3.5, and the projection output is shown in Fig. 2.8b. The simplest workaround for this issue is to use only half of the reconstruction field, like the target image in Fig. 2.8a. However, the side effect of this workaround is that half of the energy will be wasted, leading to higher power consumption and heat dissipation.

## 2.3 Computer-Generated Holography (CGH)

Different from the traditional analogue optical holography and digital holography which both need a physically existing object to record the hologram, CGH is a method to use computer algorithms to generate holograms without the need for the physical target object. Ideally, holograms can be simply computed using the inverse Fourier Transform, using the inverse functions of Eq. (7.1) and Eq. (7.2) derived in Section 2.2.2. The inverse Fourier Transforms on most targets will end up with results in complex numbers; however, as mentioned in Section 2.2.3, currently available SLMs cannot achieve complex modulation yet. Therefore, computer algorithms are needed to compute either phase-only or amplitude-only holograms, among which the former is preferred, as explained in Section 2.2.3.

This section reviews the existing methods in the literature for calculating phase-only holograms. The phase-only hologram is labelled as  $H$ , so that it differs from the previous notation of  $A$  for complex-valued hologram apertures. And the propagation function is unified as  $\mathcal{P}$ , which can be either the Fraunhofer propagation equation in Eq. (7.1) or the Fresnel propagation equation in Eq. (7.2) depending on the distance of the target field. Lastly,  $R$  denotes the reconstruction from the hologram, which is the amplitude of the result from  $\mathcal{P}$ .



Fig. 2.9 Sample target image of a mandrill ( $T$ ) [16]

A sample image of a mandrill (shown in Fig. 2.9) is chosen from the University of Southern California Signal and Image Processing Institute (USC-SIPI) Image Database [16] to test and compare the classical phase retrieval algorithms in the literature. As the square of the amplitude, also known as the ‘intensity’ of light, is the only visible component with human eyes [45], the diffracted electric field’s amplitude ( $|E|$ ) will be targeted to match the square root of  $T$ .

### 2.3.1 Naive Method

The simplest method to get a phase-only hologram  $H$  is by directly extracting the phase of the reverse propagation from the target field, discarding the amplitude component (e.g. for Fraunhofer propagation,  $H$  will simply be the phase of the inverse Fourier transform  $\mathcal{F}^{-1}$  of the target field). This method is named as ‘Naive’ method in this thesis for its simplicity. The pseudocode of the Naive method is shown in Algorithm 1 below:

---

**Algorithm 1** Naive method

---

**Input:** Target image  $T$ , Propagation function  $\mathcal{P}$  (e.g. Fresnel or Fraunhofer propagation)

**Output:** Phase hologram  $H$  and its reconstruction intensity  $R$

$$\begin{aligned} E &\leftarrow \sqrt{T} \\ A &\leftarrow \mathcal{P}^{-1}[E] \\ H &\leftarrow \angle A \\ R &\leftarrow |\mathcal{P}[e^{jH}]|^2 \end{aligned}$$


---

where  $j = \sqrt{-1}$ , the  $\angle$  sign means phase extraction (i.e. arguments of complex numbers element-wise in the matrix), and  $e^{jH}$  converts the angles back to complex numbers. All exponentials, modulus and square-root operators are carried out in an element-wise manner, so that the dimensions of  $T, A, H, R$  and  $E$  are all the same.

The Naive method (described in Algorithm 1) was then implemented in MATLAB and ran on the sample target image shown in Fig. 2.9, with the distance set to infinity (i.e. in the Fraunhofer region using the propagation formula Eq. (7.2)). The results are shown in Fig. 2.10 below:

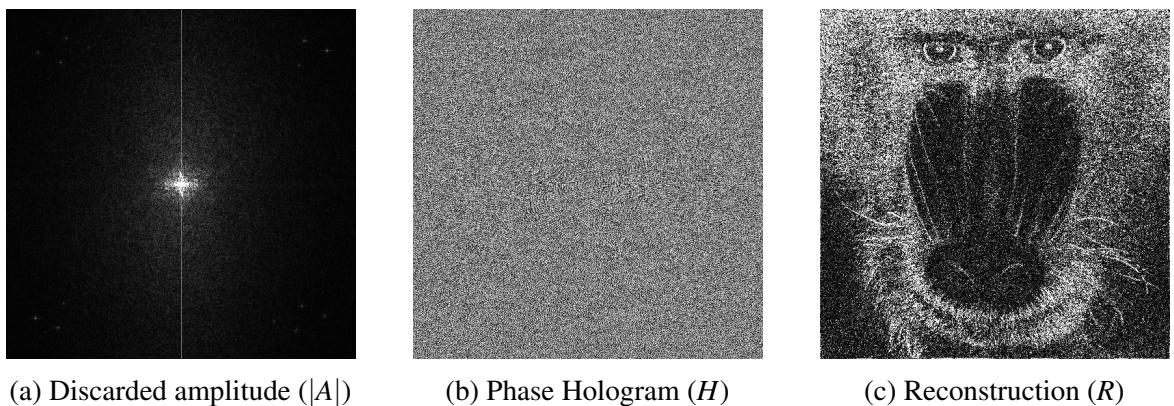


Fig. 2.10 Naive method output

After taking the inverse Fourier Transform, the amplitude and the phase of  $A$  are shown in Fig. 2.10a and Fig. 2.10b respectively. The next step then discards the amplitude component (Fig. 2.10a) and uses the phase component (in Fig. 2.10b) as the phase hologram  $H$ . Then the phase hologram went through a forward propagation and the resulting reconstruction intensity is shown in Fig. 2.10.

The reconstruction intensity is very far from the desired target image in Fig. 2.9. It shows that, discarding amplitude has introduced a significant loss of information. From a signal processing point of view, the peak around the centre in Fig. 2.10a corresponds to low spatial frequency signals, and discarding them causes the reconstruction in Fig. 2.10c to lose low frequency components and effectively becomes an edge detector. Another explanation of the poor reconstruction quality is that, this method is assuming a uniform phase profile for the target image, which is physically difficult to achieve. A simple improvement can be made by adding a random phase to the target, as shown in the pseudocode below:

---

**Algorithm 2** Improved Naive method with random phase added to the target field

---

**Input:** Target image  $T$ , Propagation function  $\mathcal{P}$  (e.g. Fresnel or Fraunhofer propagation)

**Output:** Phase hologram  $H$  and its reconstruction intensity  $R$

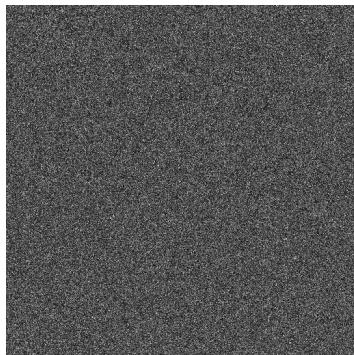
```

 $E \leftarrow \sqrt{T} \times \text{RandomPhase}()$ 
 $A \leftarrow \mathcal{P}^{-1}[E]$ 
 $H \leftarrow \angle A$ 
 $R \leftarrow |\mathcal{P}[e^{jH}]|^2$ 

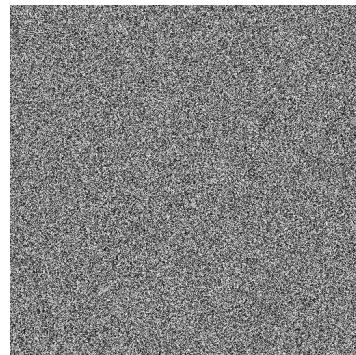
```

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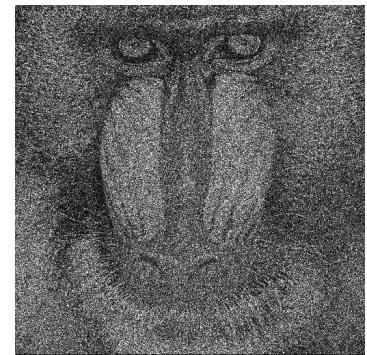
The improved Naive method was implemented in MATLAB and produced the results in Fig. 2.11 below:



(a) Discarded amplitude ( $|A|$ )



(b) Phase Hologram ( $H$ )



(c) Reconstruction ( $R$ )

Fig. 2.11 Output of the improved Naive method

As shown in Fig. 2.11c, the reconstruction quality has been greatly improved, although still quite noisy. The amplitude of the hologram being discarded (shown in Fig. 2.11a) is a lot more uniformly distributed than the one in Fig. 2.10a, so that the loss of information evenly spread across all spatial frequencies, leading to the much better reconstruction quality. However, the reconstruction quality is still quite noisy. To quantify the noise, the normalised mean squared error (NMSE) is calculated using equation  $NMSE = \frac{\frac{1}{n} \sum (x - \hat{x})^2}{\sum (x)^2}$ , where  $x$  is the target,  $\hat{x}$  is the measured output and  $n$  is the number of pixels in  $x$  and  $\hat{x}$ , together with the structural similarity index (SSIM) [46]. The NMSE is the lower the better while the SSIM is the higher the better. The reconstruction in Fig. 2.11c has an NMSE of  $1.0228 \times 10^{-6}$  and an SSIM of 0.1603.

Moreover, additional error will be introduced during the quantisation step, which is necessary for the phase hologram to be displayed on SLMs with limited bit depth. As the SLM used in this thesis is a binary phase SLM, which has a rotational symmetry property as explained in Section 2.2.3, a target image is specifically designed as shown in Fig. 2.12a which is rotational symmetrical by itself.

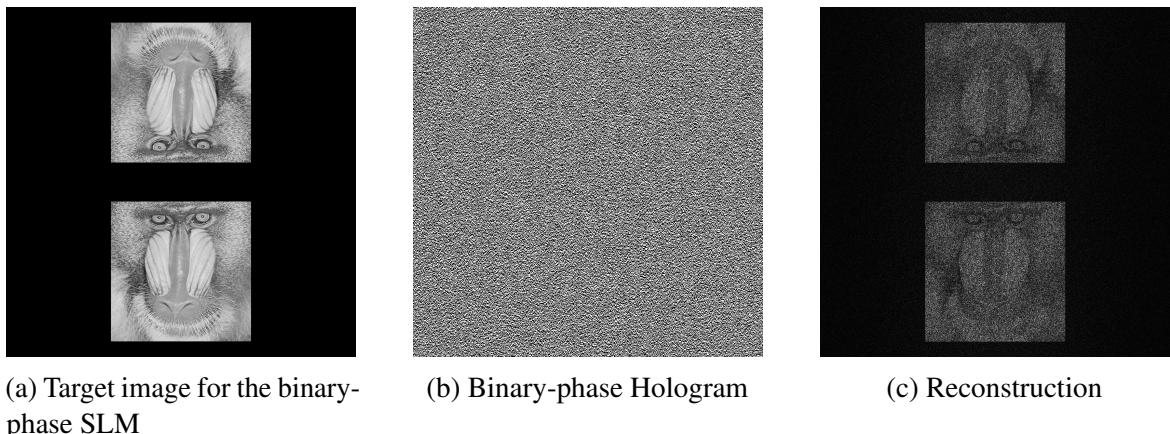


Fig. 2.12 Output of the improved Naive method with binary-phase quantisation

The binary phase hologram in Fig. 2.12b is generated by adding an additional binary quantisation step ( $\mathcal{Q}$ ) on the phase hologram computed using Algorithm 2, which is simply rounding all phases to 0 rad and  $\pi$  rad. The reconstruction of the binary-phase hologram is shown in Fig. 2.12c, which is of poor quality with NMSE being  $4.5452 \times 10^{-7}$  and SSIM being 0.0603. The only advantage of this method is its speed, as it only requires one inverse Fourier Transform calculation. To improve the reconstruction quality, better algorithms are needed. The following sections explores predecessors' efforts in quality improvement.

### 2.3.2 Direct Binary Search (DBS) Algorithm

Direct Binary Search (DBS) algorithm [22] is an algorithm that generates the hologram by randomly flipping each pixel in the SLM between binary states (0 and  $\pi$ ), one by one for many times in order to minimise the difference between its reconstruction intensity  $R$  and the target image  $T$ . The detailed algorithm is described in Algorithm 3 below:

---

**Algorithm 3** Direct Binary Search (DBS) algorithm

---

**Input:** Target image  $T$ , Propagation function  $\mathcal{P}$ , Loss function  $\mathcal{L}$  (e.g. mean-squared error), Number of iterations  $N$

**Output:** Phase hologram  $H$  and its reconstruction intensity  $R$

// Start with a random hologram with a size matching  $T$

$H \leftarrow \text{Rand}(\text{Size}(T))$

$R \leftarrow |\mathcal{P}[e^{jH}]|^2$

$L \leftarrow \mathcal{L}[R, T]$

**for**  $n = 1$  to  $N$  **do**

// Flip a random pixel in the hologram

$H_n \leftarrow \text{FlipRandomPixel}(H)$

// Calculate the loss function for the new hologram

$R_n \leftarrow |\mathcal{P}[e^{jH_n}]|^2$

$L_n \leftarrow \mathcal{L}[R_n, T]$

// Compare the new loss with the old one

**if**  $L_n < L$  **then**

// Accept the new hologram if loss is lower

$H \leftarrow H_n$

$R \leftarrow R_n$

$L \leftarrow L_n$

**end if**

**end for**

---

Although the DBS algorithm is specifically suited for generating binary phase holograms, it can also be adapted for generating multi-level phase holograms, by representing each level as binary numbers, at the cost of more computation.



Fig. 2.13 DBS algorithm running on the rotationally symmetrical mandrill target

DBS algorithm can sometimes find very accurate hologram if the run is lucky; however, it is extremely slow, because it takes numerous iterations (as shown in Fig. 2.13d, even  $10^5$  iterations has not reached a good convergence) and each iteration requires a Fourier Transform which is computationally expensive, the example run on the target image of resolution  $1024px \times 1024px$  in Fig. 2.13a took more than one hour to run the  $10^5$  iterations, which is still nowhere near convergence. The binary-phase hologram produced is shown in Fig. 2.13b and its corresponding reconstruction in Fig. 2.13c is of very poor quality.

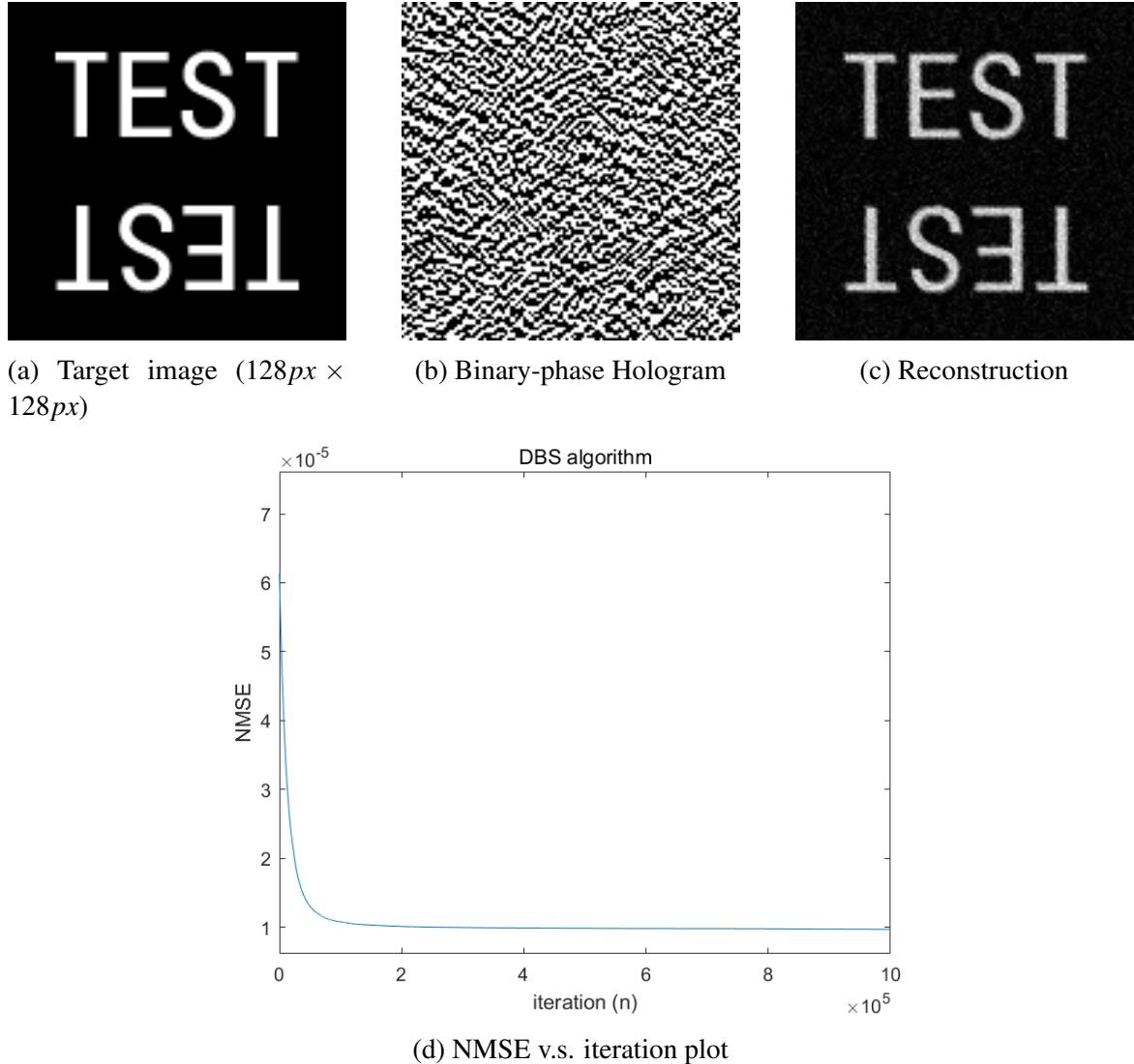


Fig. 2.14 DBS algorithm running on the low resolution target

As the DBS algorithm only flips one pixel per iteration, it naturally takes significantly longer to generate holograms with higher resolution. To test the programme on a smaller image for much quicker convergence, another target image has been designed as shown in Fig. 2.14a, which is also rotationally symmetrical as it is used for the binary-phase hologram generation. After  $10^6$  iterations, which took 10 minutes, the hologram generated is shown in Fig. 2.14b and its resulting reconstruction is shown in Fig. 2.14c, which has an NMSE of  $9.6881 \times 10^{-6}$  and an SSIM of 0.2871. The NMSE v.s. iteration plot in Fig. 2.14d shows that it reaches a good convergence at around  $2 \times 10^5$  iterations, corresponding to around 2 minutes. And the

curve of NMSE is monotonically decreasing with iteration number, as only holograms with better results are accepted during the iterations.

In summary, the DBS algorithm is a slow but working algorithm for binary phase hologram generation. The programme running time scales up significantly when the target image's resolution gets higher. And also, as it only cares about local optimality at each iteration, it is a greedy algorithm that only follows the steepest descent route, which could easily get trapped in a local minimum where flipping any bit is not getting better reconstruction. Another consequence of the random nature is that the generated hologram will be different at each run, so the quality of the resulting reconstruction ( $R$ ) will depend on how 'lucky' each run is. The Simulated Annealing (SA) algorithm [47] in the next session aims to resolve this issue.

### 2.3.3 Simulated Annealing (SA) Algorithm

Simulated Annealing (SA) algorithm [47] is a variant of the DBS algorithm. It adopts a probabilistic approach to avoid the steepest gradient descent. Its name derives from the fact that it approximates the recrystallisation process during metal annealing and is particularly well-suited to avoiding the trap of local minima [23]. To implement this idea, we then need a function ( $\mathcal{L}$ ) to calculate the probability of the hologram ( $H$ ), and a threshold  $p_t$  to decide whether the probability is high enough for the according hologram to be accepted. In this thesis, the probability is selected to be a random function and the threshold is chosen to be 0.9. The pseudocode for this algorithm is listed in Algorithm 4.

**Algorithm 4** Simulated Annealing (SA) algorithm

---

**Input:** Target image  $T$ , Propagation function  $\mathcal{P}$ , Loss function  $\mathcal{L}$ , Number of iterations  $N$ , Probability function  $\mathcal{Z}$ , Probability threshold  $p_t$

**Output:** Phase hologram  $H$  and its reconstruction intensity  $R$

```

// Start with a random hologram with a size matching  $T$ 
 $H \leftarrow \text{Rand}(\text{Size}(T))$ 
 $R \leftarrow |\mathcal{P}[e^{jH}]|^2$ 
 $L \leftarrow \mathcal{L}[R, T]$ 
for  $n = 1$  to  $N$  do
    // Flip a random pixel in the hologram
     $H_n \leftarrow \text{FlipRandomPixel}(H)$ 

    // Calculate the loss function for the new hologram
     $R_n \leftarrow |\mathcal{P}[e^{jH_n}]|^2$ 
     $L_n \leftarrow \mathcal{L}[R_n, T]$ 

    // Compare the new loss with the old one
    if  $L_n < L$  then
        // Accept the new hologram if loss is lower
         $H \leftarrow H_n$ 
         $R \leftarrow R_n$ 
         $L \leftarrow L_n$ 
    else
        // Calculate the probability of the hologram
         $p_n \leftarrow \mathcal{Z}[H_n]$ 
        if  $p_n > p_t$  then
            // Accept the new hologram if the probability exceeds the threshold
             $H \leftarrow H_n$ 
             $R \leftarrow R_n$ 
             $L \leftarrow L_n$ 
        end if
    end if
end for

```

---



Fig. 2.15 SA algorithm running on the low resolution target

An implementation of SA algorithm with  $p_t = 0.9$  was carried out on the low resolution target in Fig. 2.15a, and the resulting binary-phase hologram and its reconstruction are shown in Fig. 2.15b and Fig. 2.15c respectively. From the NMSE v.s. iteration plot in Fig. 2.15d, it can be seen that, instead of the monotonic decrease observed in Fig. 2.14d for DBS algorithm, the SA algorithm has occasional rises in NMSE, which happens when the probability  $p_n$  exceeds the threshold  $p_t$ . The final NMSE was recorded to be  $1.0542 \times 10^{-5}$  and the final SSIM was 0.2750, which are both slightly worse than the DBS algorithm in this case. Due to the probabilistic nature of the SA algorithm, although it can avoid being trapped in local optimal points, the ‘jump backs’ can also cause delays in convergence.

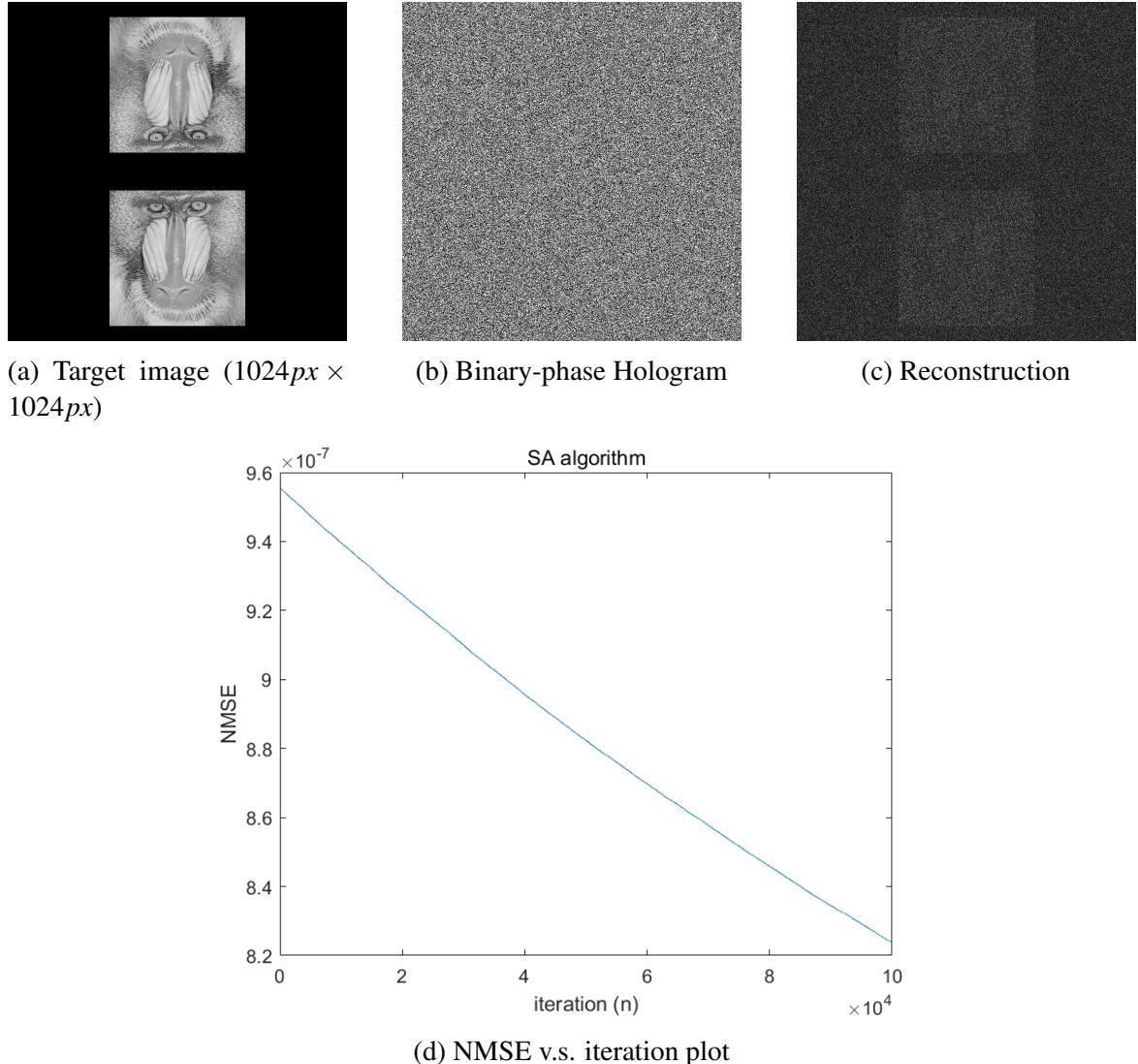


Fig. 2.16 SA algorithm running on the rotationally symmetrical mandrill target

Then the SA algorithm was run for the mandrill target in Fig. 2.16a. The convergence plot in Fig. 2.16d shows that it did not converge within  $10^5$  iterations, which took around one hour. The binary phase hologram generated is shown in Fig. 2.16b and its corresponding reconstruction intensity is shown in Fig. 2.16c, which is of very poor quality. Both the DBS and the SA algorithms rely on flipping only a single pixel per iteration, which is very inefficient. A better algorithm should change the value of multiple pixels at every iteration for better efficiency.

### 2.3.4 Gerchberg-Saxton (GS) Algorithm

The Gerchberg-Saxton (GS) algorithm [19] is a revolutionary algorithm and is much better and more robust than the algorithms introduced in the previous sections. Although being more than 50 years old, the GS algorithm is still frequently used and has lots of variants [48–50]. It functions by iteratively determining the phase profile of the hologram required to reconstruct a target image, looping between the hologram and the reconstruction plane, and applying constraints to each plane accordingly during each iteration. GS algorithm is very easy to implement, its pseudocode is shown in Algorithm 5.

---

**Algorithm 5** Gerchberg-Saxton (GS) Algorithm
 

---

**Input:** Target image  $T$ , Propagation function  $\mathcal{P}$ , Number of iterations  $N$ , Initial phase  $\Phi$  (e.g. random, zeros, or other patterns)

**Output:** Phase hologram  $H$  and its reconstruction intensity  $R$

```

// Initiate  $E$  with amplitude  $\sqrt{T}$  and initial phase  $\Phi$ 
 $E \leftarrow \sqrt{T} \times e^{j\Phi}$ 
for  $n = 1$  to  $N$  do
    // Compute the hologram plane
     $A \leftarrow \mathcal{P}^{-1}[E]$ 
    // Apply the phase-only constraint at the hologram plane
     $A \leftarrow e^{j\angle A}$ 

    // Compute the propagation for the new hologram
     $E \leftarrow \mathcal{P}[A]$ 
    // Apply the target field amplitude constraint at the reconstruction plane
     $E \leftarrow \sqrt{T} \times e^{j\angle E}$ 
end for
 $H \leftarrow \angle A$ 
 $R \leftarrow |\mathcal{P}[A]|^2$ 

```

---

The GS algorithm described in Algorithm 5 was implemented in MATLAB and was first run on the mandrill target image in Fig. 2.9, and the output results are shown in Fig. 2.17. It can be seen from Fig. 2.17c that, the reconstruction after 30 iterations of GS algorithm reached a very good result, having an NMSE of  $2.6612 \times 10^{-8}$  and an SSIM of 0.7940, which are both much better than the single-iteration Naive method's result in Fig. 2.11c.

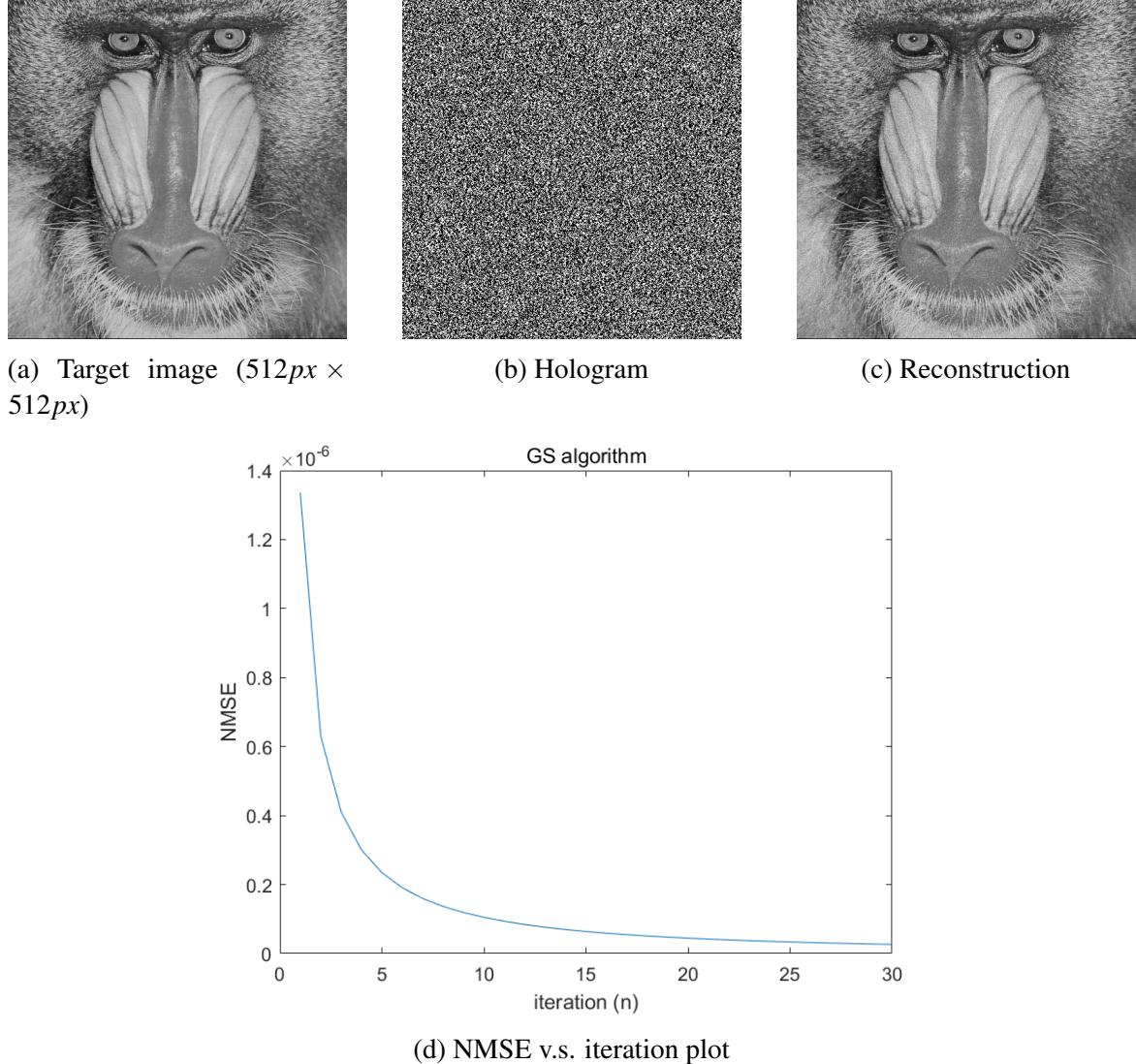


Fig. 2.17 GS algorithm output on the mandrill target

The NMSE v.s. iteration plot in Fig. 2.17d shows that the GS algorithm converged quickly, providing very good result in tens of iterations, which is much fewer than the DBS and SA algorithms. Although the GS algorithm is more computationally expensive at each iteration, as it needs to compute both a forward and an backward propagation, leading to two Fourier transforms every iteration, the GS algorithm is still much faster and provides much better reconstruction quality than the DBS and SA algorithms.

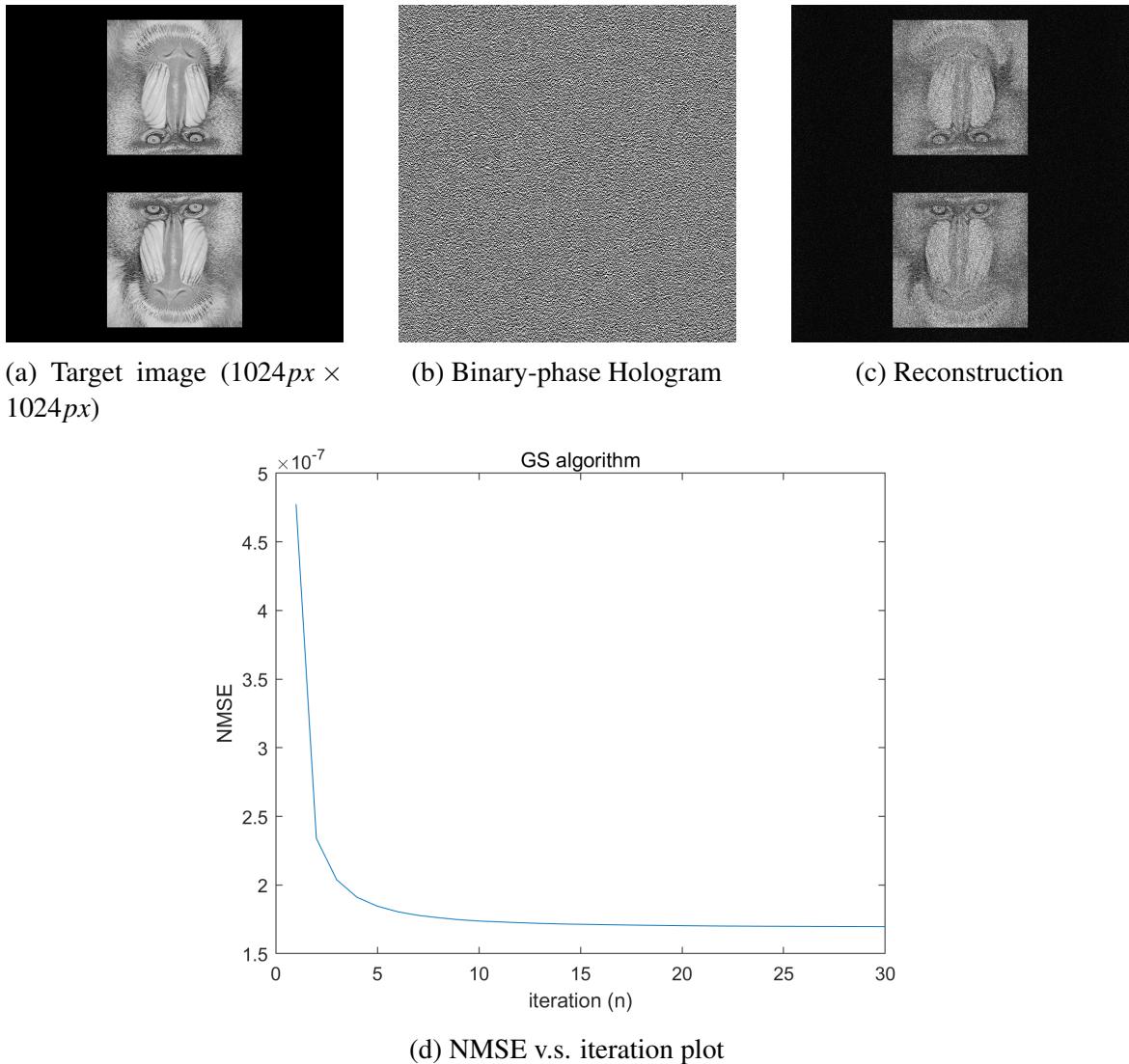


Fig. 2.18 GS algorithm running on the rotationally symmetrical mandrill target

Then the GS algorithm was adapted to generate binary-phase holograms, for use on the binary-phase SLM in this thesis. The change was simply implemented by adding a quantisation function ( $\mathcal{Q}$ ) when applying the phase-only constraint at the hologram plane (i.e. the line ' $A \leftarrow e^{j\angle A}$ ', in Algorithm 5 is changed to ' $A \leftarrow e^{j\mathcal{Q}[\angle A]}$ '). The results are shown in Fig. 2.18. Fig. 2.18d shows good convergence within 20 iterations. The resulting binary-phase hologram is shown in Fig. 2.18b and its corresponding reconstruction in Fig. 2.18c has an NMSE of  $1.6968e-07$  and an SSIM of 0.0619, which are both better than the Naive method's result in Fig. 2.12. In summary, the GS algorithm is quick and robust. On my laptop computer of model ASUS ROG Zephyrus M16, which has a CPU of model i7-11800H and a GPU of

model RTX3060, the 30 iterations took 1.5 seconds to complete. It reached convergence in tens of iterations. However, as it is still iterative, generating holograms in real-time is still a challenge, and the reconstruction still suffers from noise.

### 2.3.5 One-Step Phase Retrieval (OSPR) Algorithm

OSPR algorithm was first demonstrated by Buckley [51]. It is a solution to high-quality hologram reconstruction that relies on time multiplexing of holograms, exploiting the response time of eye in order to reduce noise in the replay field [15]. The random noises are averaged by the eye, while the target image stays, so that the average noise can be reduced. The perceived noise is lessened by the temporal average detected by the eye, rather than computational optimisation of the hologram [15]. The pseudocode for OSPR is shown in Algorithm 6 below.

---

**Algorithm 6** One-Step Phase Retrieval (OSPR) algorithm
 

---

**Input:** Target image  $T$ , Propagation function  $\mathcal{P}$ , Number of sub-frames  $S$ , Quantisation function  $\mathcal{Q}$

**Output:** List of phase holograms  $H[1 \dots S]$

```

// Compute a list of hologram sub-frames
for  $i = 1$  to  $S$  do
     $E \leftarrow \sqrt{T} \times \text{RandomPhase}()$ 
     $A \leftarrow \mathcal{P}^{-1}[E]$ 
     $H[i] \leftarrow \mathcal{Q}[\angle A]$ 
end for

// Then display the sub-frames on the phase modulator sequentially
 $i \leftarrow 1$ 
while True do
    Display( $H[i]$ )
     $i \leftarrow i + 1$ 
    if  $i > S$  then
         $i \leftarrow 1$ 
    end if
end while
  
```

---

When generating the list of holograms, it repetitively computes the inverse propagation of the target amplitude (which is the square-root of the target intensity  $T$ ) multiplied by different random phases, for a total of  $S$  times to generate  $S$  hologram sub-frames ( $H[1 \dots S]$ ). The computation of each hologram sub-frame is the same as the Naive method in Algorithm 2 discussed in Section 2.3.1. Then the  $S$  hologram sub-frames are displayed sequentially on a SLM having a refresh rate being so fast that the average reconstruction intensity is perceived by the human eyes. As currently available fast SLMs are binary-phase modulators, an example run on the rotationally symmetrical target previously used in Fig. 2.13a was carried out for 24 sub-frames ( $S = 24$ ). The results are summarised in Fig. 2.19.

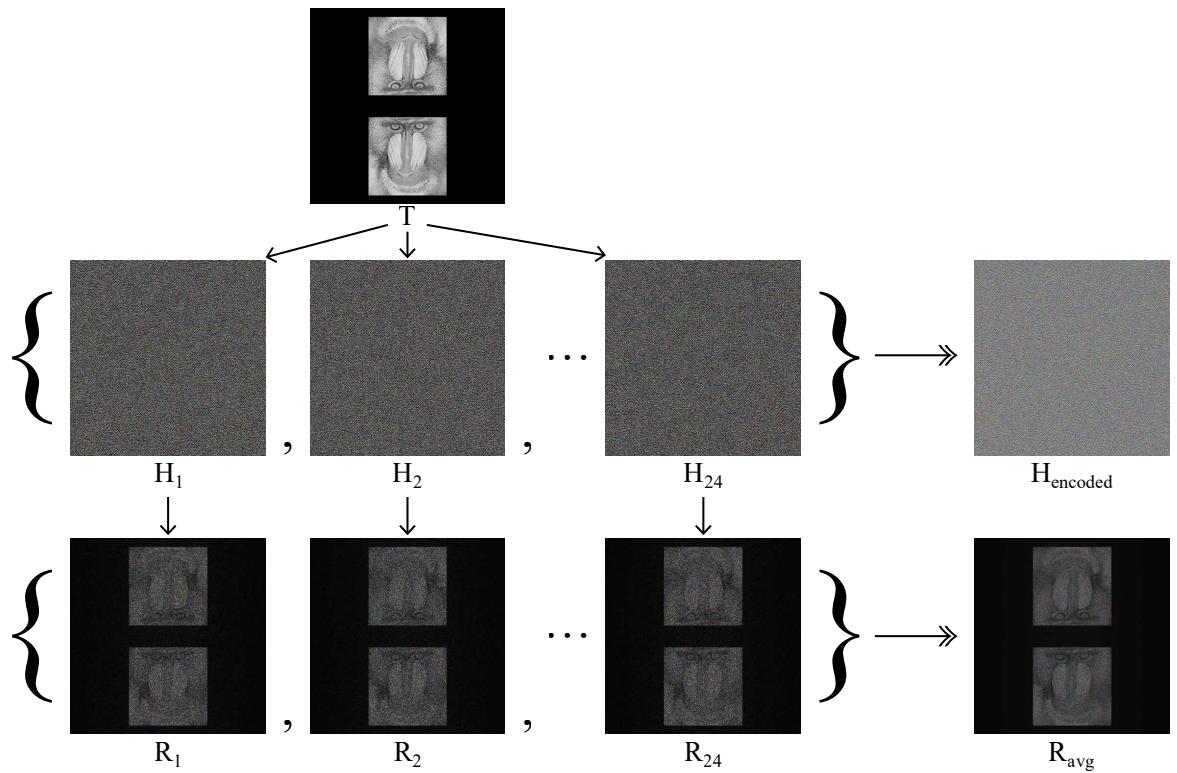


Fig. 2.19 OSPR algorithm running on the rotationally symmetrical mandrill target

In Fig. 2.19, a total of 24 binary-phase hologram sub-frames ( $H_1, H_2, \dots, H_{24}$ ) were generated for the target image  $T$ . For easier data transfer and to fit with the common display signal formats, the 24 binary-phase hologram sub-frames are encoded into a single file with 8 bit depth and RGB (red-green-blue) channels, so that each of the  $(8 \times 3 =) 24$  bit-planes corresponds to a single binary-phase hologram sub-frame. For a quantitative analysis, the reconstruction intensities of the hologram sub-frames are computed in  $R_1, R_2, \dots, R_{24}$  respectively, whose average is  $R_{\text{avg}}$  in Fig. 2.19. The average reconstruction intensity

$R_{avg}$  has an NMSE of  $9.8632 \times 10^{-8}$  and an SSIM of 0.1321, which are both significantly better than the GS algorithm (NMSE= $1.6968 \times 10^{-7}$ , SSIM=0.0619) and the Naive method (NMSE= $4.5452 \times 10^{-7}$ , SSIM=0.0603).

The major advantage of the OSPR algorithm is that it is superfast. It is non-iterative and requires only one Fourier Transform per frame, taking less than a second to generate a 24-frame hologram set. Its non-iterative nature also allows it to be parallelised to further improve computation speed, which is crucial for the Light Blue Optics who made the real-time holographic laser projector commercially available in 2010 [52], although the product was later discontinued for financial reasons. The downside of this algorithm is that the sub-frames are independent from each other, and the final reconstruction output is still subject to some noise, as they are only more randomly distributed instead of being reduced. There is a variant of slight improvement on the OSPR algorithm, called Adaptive OSPR (AD-OSPR), to be introduced in Section 2.3.6.

### 2.3.6 Adaptive One-Step Phase Retrieval (AD-OSPR) Algorithm

The AD-OSPR algorithm [53] is a variant of the OSPR algorithm. It aims to improve the reconstruction quality without introducing a significant amount of additional computational cost. It functions in such a way that, when computing the second sub-frame onwards, it subtracts the average reconstruction from the target image to get the error, so that it can be compensated in the next iteration. To help explain the process in detail, a pseudocode is written for the AD-OSPR algorithm, as shown in Algorithm 7.

**Algorithm 7** Adaptive One-Step Phase Retrieval (AD-OSPR) algorithm

---

**Input:** Target image  $T$ , Propagation function  $\mathcal{P}$ , Number of sub-frames  $S$ , Quantisation function  $\mathcal{Q}$

**Output:** List of phase holograms  $H[1 \dots S]$

```

// Compute a list of hologram sub-frames
 $T[1] \leftarrow T$ 
 $R_{avg}[0] \leftarrow 0$ 
for  $i = 1$  to  $S$  do
     $E \leftarrow \sqrt{T[s]} \times \text{RandomPhase}()$ 
     $A \leftarrow \mathcal{P}^{-1}[E]$ 
     $H[i] \leftarrow \mathcal{Q}[\angle A]$ 

    // Compute average reconstruction
     $R \leftarrow |\mathcal{P}[e^{jH[i]}]|^2$ 
     $R_{avg}[i] \leftarrow (R_{avg}[i-1] \times (i-1) + R)/i$ 

    // Update the target for the next iteration
     $T[i+1] \leftarrow T + T - R_{avg}[i]$ 
end for

// Then display the sub-frames on the phase modulator sequentially
 $i \leftarrow 1$ 
while True do
    Display( $H[i]$ )
     $i \leftarrow i + 1$ 
    if  $i > S$  then
         $i \leftarrow 1$ 
    end if
end while

```

---

In the pseudocode in Algorithm 7, the target intensity for the first iteration ( $T[1]$ ) is initialised as the input target image  $T$ , and the average reconstruction intensity for the holograms generated up to each iteration ( $R_{avg}[i]$ ) is initialised as 0. Then the **for** loop starts from the same routine as the OSPR algorithm, multiplying a random phase to the square root of target intensity, taking a inverse propagation, and applying the binary phase quantisation. Then, the

average reconstruction is computed by propagating the quantised hologram of each iteration and take the weighted average between the current reconstruction intensity and the historical average reconstruction value. The average reconstruction is then used to update the target intensity for the next iteration, by subtracting the average from the target and adding the difference to the input target image. To provide a visual illustration, an example run of the AD-OSPR algorithm was carried out on the mandrill target as shown in Fig. 2.20.

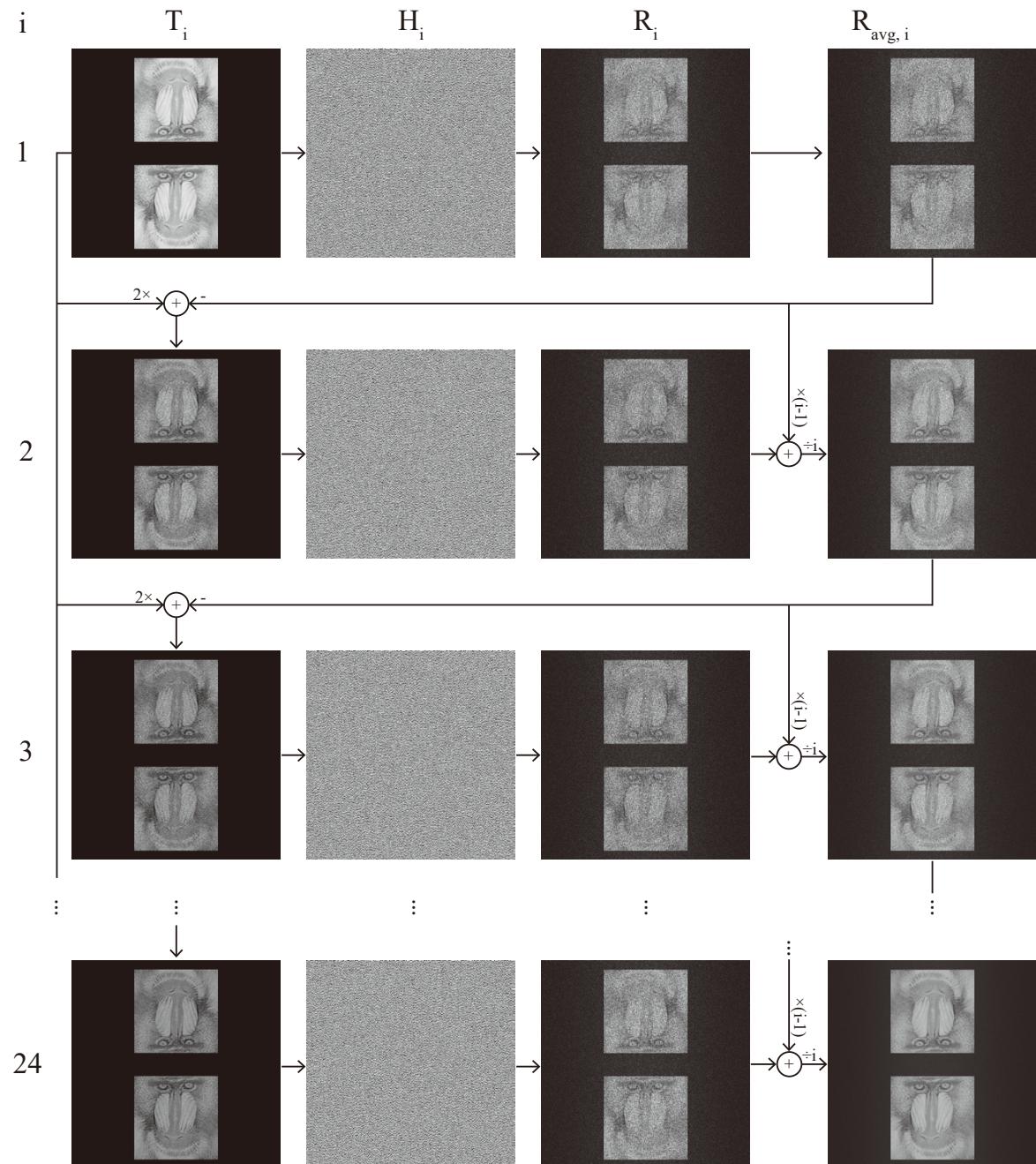


Fig. 2.20 AD-OSPR algorithm running on the rotationally symmetrical mandrill target

In Fig. 2.20, each row of images correspond to one iteration that is indexed in the ‘i’ column, where only iterations 1, 2, 3, 24 are shown here, while iterations 4-23 are omitted to save space. At each iteration, the binary-phase hologram  $H_i$  is computed from its target  $T_i$ , and the corresponding reconstruction is shown in  $R_i$ . For the first iteration, the target image of the same one as in Fig. 2.19 is used, and the first average reconstruction  $R_{avg,1}$  is just the first reconstruction  $R_1$ . Then, from the second iteration, the target intensity  $T_i$  is updated as the target image ( $T$ ) plus the difference between the target image and the average reconstruction so far ( $T - R_{avg,i}$ ). After computing new reconstructions  $R_i$ , the average reconstruction  $R_{avg,i}$  is updated using the weighted average function  $R_{avg}[i] \leftarrow (R_{avg}[i-1] \times (i-1) + R)/i$ . After 24 iterations, the 24 hologram sub-frames are computed and their average reconstruction ( $R_{avg,24}$ ) had an NMSE of  $9.1161 \times 10^{-8}$  and a SSIM of 0.1992, which are both better than the OSPR algorithm in Section 2.3.5 who achieved an NMSE of  $9.8632 \times 10^{-8}$  and an SSIM of 0.1321.

The programme running time of the AD-OSPR algorithm is nearly the same as the OSPR algorithm, both being around 0.7 second. Therefore the AD-OSPR algorithm is quite an effective improvement on the original OSPR algorithm, achieving a 7.6% reduction in NMSE and 50.8% improvement in SSIM without adding much computation. The disadvantage is however, that it cannot be parallelised, unlike the OSPR algorithm. As it needs to have the average reconstruction intensity result from the previous iteration for use in the next iteration.

### 2.3.7 3D CGH

Section 2.3.1 - Section 2.3.6 described several algorithms to generate a phase hologram for a single target image. However, the major benefit of holography is that it can produce true 3D light field, much more than just a single slice 2D image. Then the problem arises as how to generate a hologram for 3D targets to make full use of the major benefit of holography. The simplest method is to slice the 3D target into a set of 2D layers, like computed tomography (CT) scanning, and then generate a hologram so that its Fresnel propagation (in Eq. (7.1)) at each depth ( $z$ ) matches the according layer. This subsection therefore reviews how the current methods in Section 2.3.1 - Section 2.3.6 can be adapted to multi-depth hologram generation.

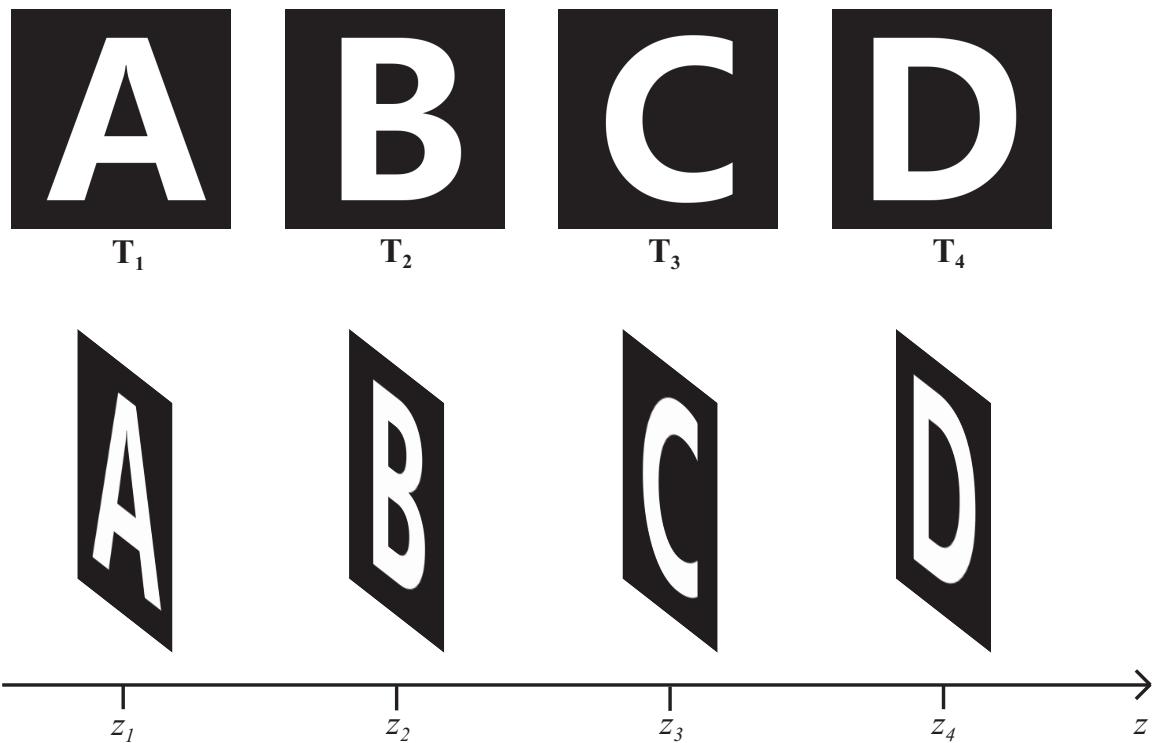


Fig. 2.21 Multi-slice target consisted of 4 different characters at different distances

For fair tests on algorithms, an example 3D target field consisted of 4 slices has been created using alphabets, as shown in Fig. 2.21.  $T_1$  to  $T_4$  have the same resolution of  $512px \times 512px$ , and the distances  $z_1$  to  $z_4$  are set to  $1, 2, 3, 4cm$  respectively.

### Naive phase addition

To adapt the naive method in Section 2.3.1 to compute multi-depth hologram, firstly, a set of holograms are generated corresponding to each layer of the target field respectively, using the inverse of the Fresnel propagation function in Eq. (7.1). Then, based on the principle of superposition, the set of holograms are added up to form the final hologram, whose phase is directly extracted to be the phase-only hologram while the amplitude component is discarded.

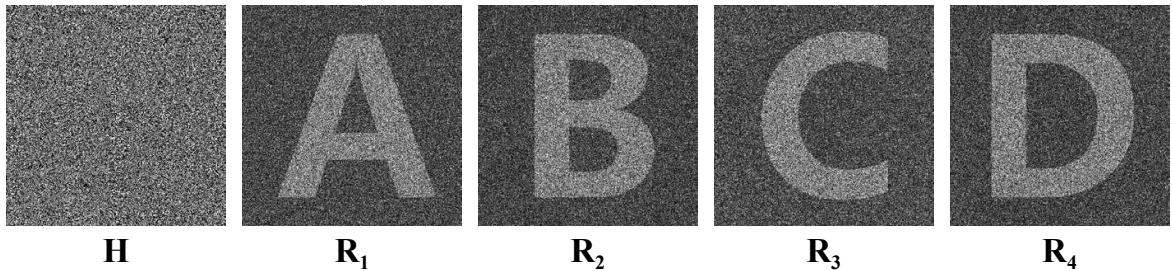


Fig. 2.22 Naive phase addition method's result on the 4-slice target

The simulation results shown in Fig. 2.22 demonstrates the effectiveness of the naive phase addition method. The reconstructions at each depth are legible, despite the presence of some noise. However, after applying a binary-phase quantisation on the hologram, the reconstruction quality deteriorates significantly, as shown in Fig. 2.23.

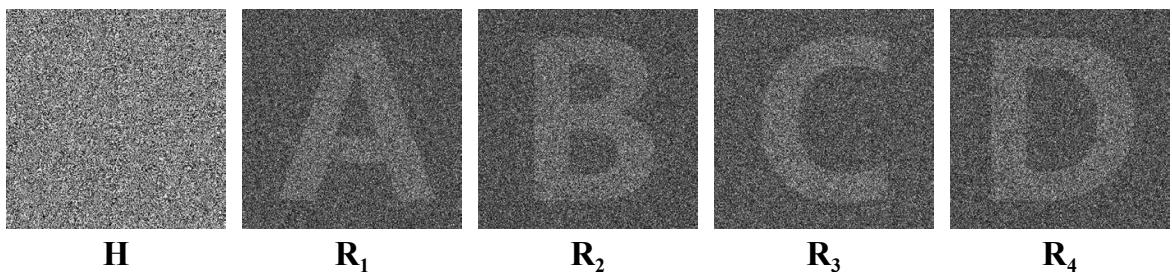


Fig. 2.23 Naive phase addition method's result on the 4-slice target after binary quantisation

To improve the reconstruction quality, the time-multiplexed OSPr algorithm (in Section 2.3.5) is adapted to produce multi-frame holograms for 3D targets.

### OSPR algorithm adaptation

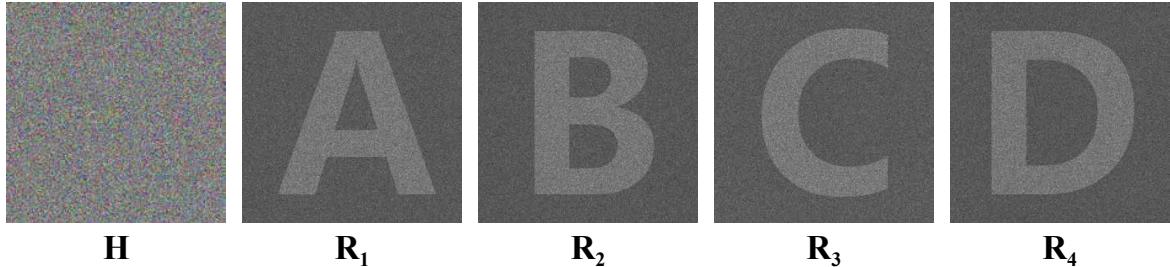


Fig. 2.24 OSPR algorithm's result on the 4-slice target

The OSPR algorithm in Section 2.3.5 was adapted to generate multi-depth targets by doing the Naive phase addition method  $S$  times, where  $S$  is the total number of sub-frames. The results of the example run with  $S = 24$  is shown in Fig. 2.24. The average reconstructions has much less noise than the single frame one in Fig. 2.23. However, the contrast is rather low. Therefore, one of the major objective of this thesis is to search for time-multiplexed binary-phase holograms with better reconstruction quality than the existing methods.

### GS algorithm adaptation

The GS algorithm in Section 2.3.4 can also be adapted to generate phase-only holograms for multi-depth 3D targets. Instead of propagating to a fixed target image at each iteration, it is modified to propagate to a different distance at each iteration, and the target amplitude constraint of the according distance is applied (i.e. line  $E \leftarrow \sqrt{T} \times e^{j\angle E}$  in Algorithm 5 becomes  $E \leftarrow \sqrt{T_{i\%n}} \times e^{j\angle E}$ , where  $i$  is the iteration number and  $n$  is the total number of slices). Such method is named sequential slicing, where the algorithm sweeps through the slices one by one sequentially during its iterations.

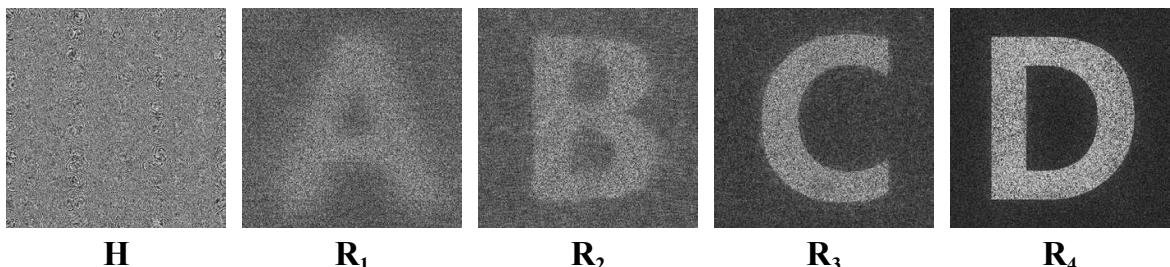


Fig. 2.25 GS with sequential slicing algorithm's result on the 4-slice target

The GS with sequential slicing (SS) algorithm was run on the target field in Fig. 2.21 for 100 iterations, and the result is shown in Fig. 2.25. The interesting phenomena is that, as the 100 iterations terminated at the 4<sup>th</sup> layer, the 4<sup>th</sup> reconstruction ( $R_4$ ) has the best quality among all slices, and the reconstruction quality deteriorates as the slice index reduces.

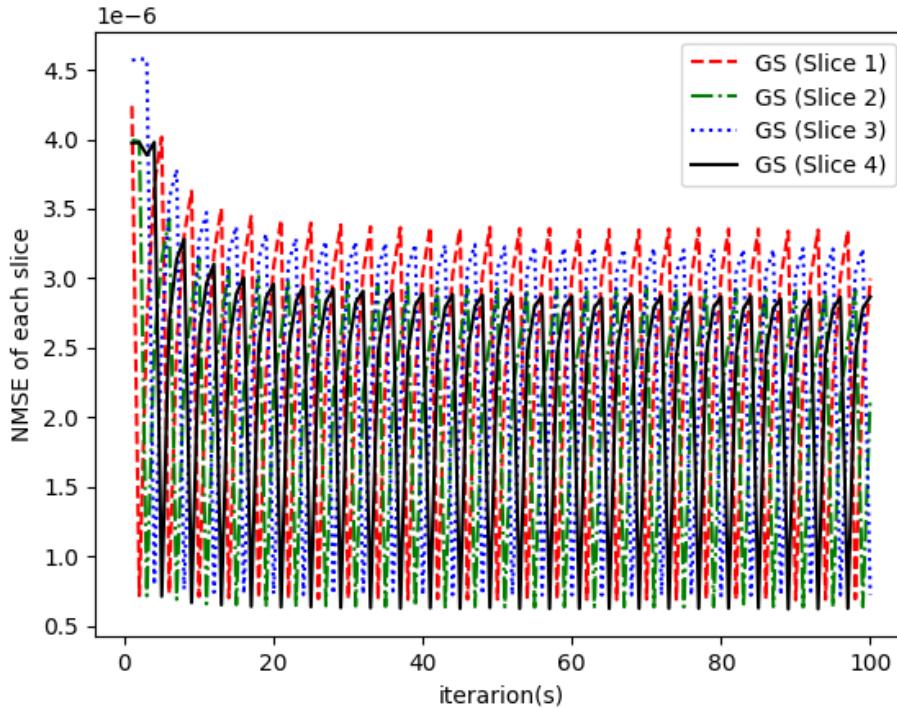


Fig. 2.26 GS with SS algorithm's NMSE v.s. iteration number plot

As a further investigation, the NMSE of each slice is plotted against iteration number in Fig. 2.26. It shows that the iterations did not converge. The NMSE curves are oscillating severely. When the NMSE of one slice decreases, the NMSE of all other slices increase. Applying the amplitude constraint at a single depth worsen all other slices. There is a solution in the literature, called dynamic compensatory GS (DCGS) [50].

The DCGS algorithm ‘softens’ the target field amplitude constraint. Instead of forcing the amplitude of the reconstructed field to the target amplitude directly, it allows a fraction  $\alpha$  of the original amplitude to retain (i.e.  $E \leftarrow \sqrt{T_{i\%n}} \times e^{j\angle E}$  becomes  $E \leftarrow [\alpha \times |E| + (1 - \alpha) \times \sqrt{T_{i\%n}}] \times e^{j\angle E}$ ), where  $\alpha$  is adjusted dynamically at each iteration. The DCGS algorithm was implemented and run on the 4-slice target in Fig. 2.21. The result is shown in Fig. 2.27.

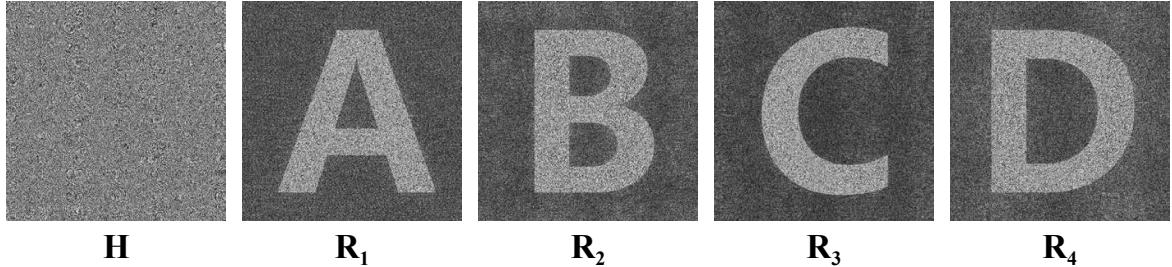


Fig. 2.27 DCGS algorithm's result on the 4-slice target

Fig. 2.27 shows the effectiveness of the DCGS algorithm, where all slices have good qualities instead of the huge inter-slice quality imbalance observed in Fig. 2.25. The NMSE of each slice is plotted against the iteration number in Fig. 2.28.

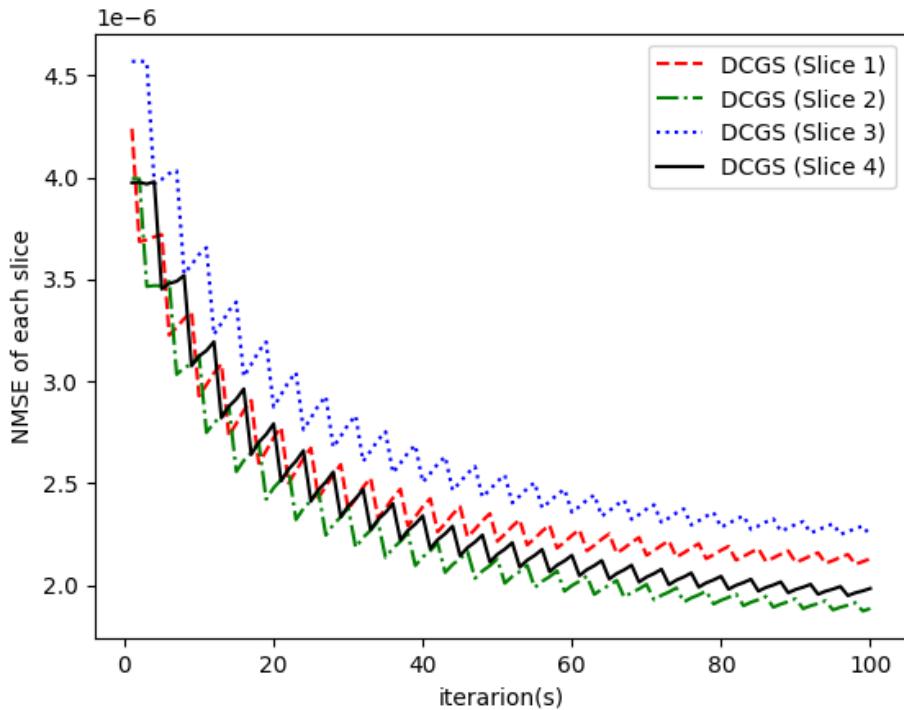


Fig. 2.28 DCGS algorithm's NMSE v.s. iteration number plot

Fig. 2.28 shows a much better convergence than Fig. 2.26. Although the bounce-backs still present (i.e. when the NMSE of one slice decreases, those of other slices increase), the overall trend of NMSE is decreasing as the iterations continue.

# Chapter 3

## Digital Pre-Distorted One-Step Phase Retrieval (DPD-OSPR) Algorithm

*Note: The work in this chapter has been published in Ref. [5]*

In a computer-generated holographic projection system, the image is reconstructed via the diffraction of light from a spatial light modulator. In this process, several factors could contribute to non-linearities between the reconstruction and the target image. This chapter evaluates the non-linearity of the overall holographic projection system experimentally, using binary phase holograms computed by the OSPR algorithm introduced in Section 2.3.5, and then applies a digital pre-distortion (DPD) method to correct for the non-linearity. Both a notable increase in reconstruction quality and a significant reduction in mean squared error were observed, proving the effectiveness of the proposed DPD-OSPR algorithm.

### 3.1 Introduction

Chapter 2 reviewed the real-time CGH method called OSPR (in Section 2.3.5), which was fast enough for real-time holography, but the reconstruction quality has still got potential for improvement. Hence, a computationally inexpensive method is needed to improve the reconstruction quality whilst maintaining the real-time property of the OSPR algorithm.

There are several factors contributing to the non-linearities between the reconstruction of hologram and the target image, including the calculation and quantization of the hologram, the modulation of the light and the imperfections in the optical setup. This chapter proposes the digital pre-distorted one-step phase retrieval (DPD-OSPR) algorithm. The digital pre-distortion (DPD) is carried out on the holographic projection system by measuring the non-linearity experimentally and applying the according pre-distortion curve on target images. DPD can be done via a one-to-one correction curve or a look-up table (LUT) which allows the relationship between the input and output to be adjusted without any heavy computation.

The intuition of the proposed DPD-OSPR algorithm for CGH comes from the gamma correction method for conventional displays, such as cathode-ray tube (CRT) monitor [54], plasma display panel television (PDP-TV) [55] and thin-film-transistor liquid-crystal display (TFT LCD) [56, 57]. Gamma correction for conventional displays were originally developed to mimic the perceptual response of human vision [58]. The work presented here is a logical continuation of this approach applied to holographic displays.

## 3.2 Experimental setup

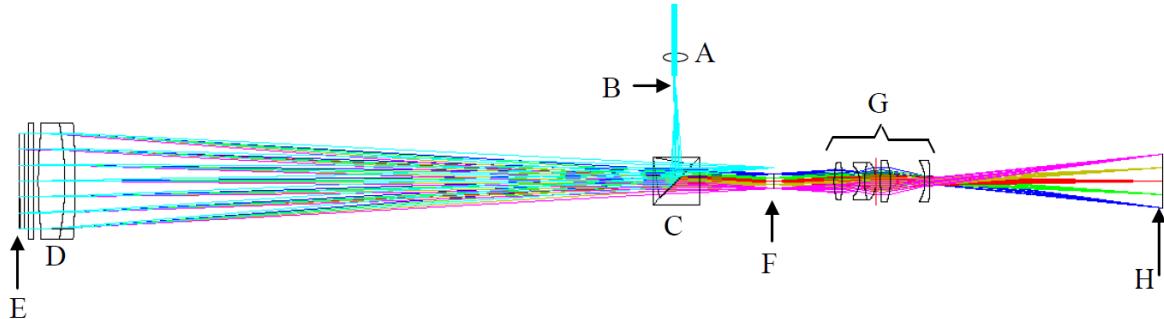


Fig. 3.1 Optical setup [17]

The holographic projector used in this experiment is a Fourier projection system developed by Freeman [17], as shown in Fig. 3.1. The design is consisted of a diode-pumped solid-state (DPSS) 532 nm 50mW laser source, focussed down by an aspheric singlet (A), the focus of which becomes the diffraction limited point source (B) for the projector. The beam then passes through a polarizing beam splitter cube (C) to a collimating lens (D), which illuminates the SLM (E). The SLM is a binary phase SXGA-R2 Ferroelectric Liquid crystal on silicon (LCOS) micro-display with a refresh rate of 1440Hz, a pixel pitch of 13.6  $\mu\text{m}$  and a resolution of 1280px  $\times$  1024px. An aperture at point (F) spatially filters out the other orders,

leaving only one first order, which is then magnified up by a finite conjugate lens group (G) to produce an image, of the required size, on the screen (H). [17]

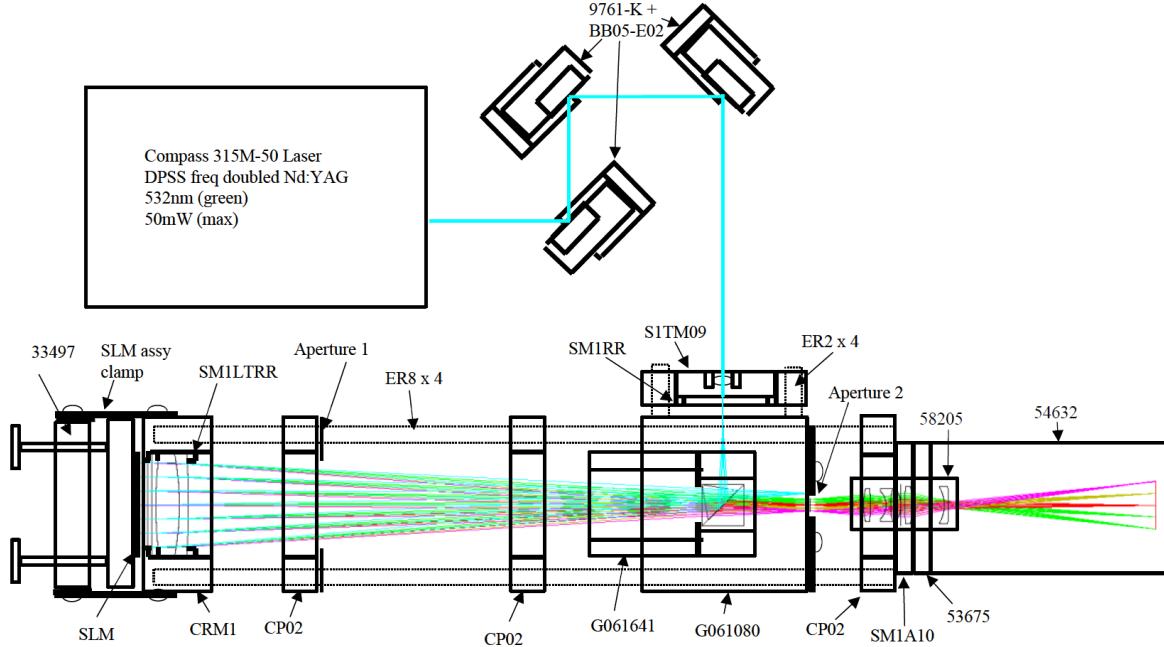


Fig. 3.2 Mechanical components [17]

The mechanical components are listed in Fig. 3.2 with parts number. The holograms displayed on the SLM are generated using the OSPr algorithm [35], and as the SLM is a binary-phase modulator, each hologram sub-frame needs to be binary quantized. Then each group of the 24 binary-phase hologram sub-frames are encoded as the 8-bit red, green, blue (RGB) channels of a 24-bit image to interface with the SLM driver electronics. The SLM displays each bit plane sequentially, with ones and zeros mapping to opposing phase modulations at each pixel. The reconstructions were captured using a Canon 550D camera with an EFS 18-55 mm lens. To ensure fair comparison, the camera was set to the same manual setting when comparing each pair of replay fields before and after DPD. It takes  $24/1440 = 1/60$ s to display all 24 sub-frames on a 1440Hz SLM, so the camera shutter speed was set to 1/30s to capture all frames twice. The images captured are raw format in RGB colour, which are subsequently converted to grey-scale (using the `rgb2gray` function in Matlab[59]) when calculating NMSE against the target images, using the equation  $NMSE = \frac{\frac{1}{n} * \sum(x - \hat{x})^2}{\sum(x)^2}$ , where  $x$  is the target,  $\hat{x}$  is the measured output and  $n$  is the dimension of  $x$  and  $\hat{x}$ .

### 3.3 Determining the DPD curve

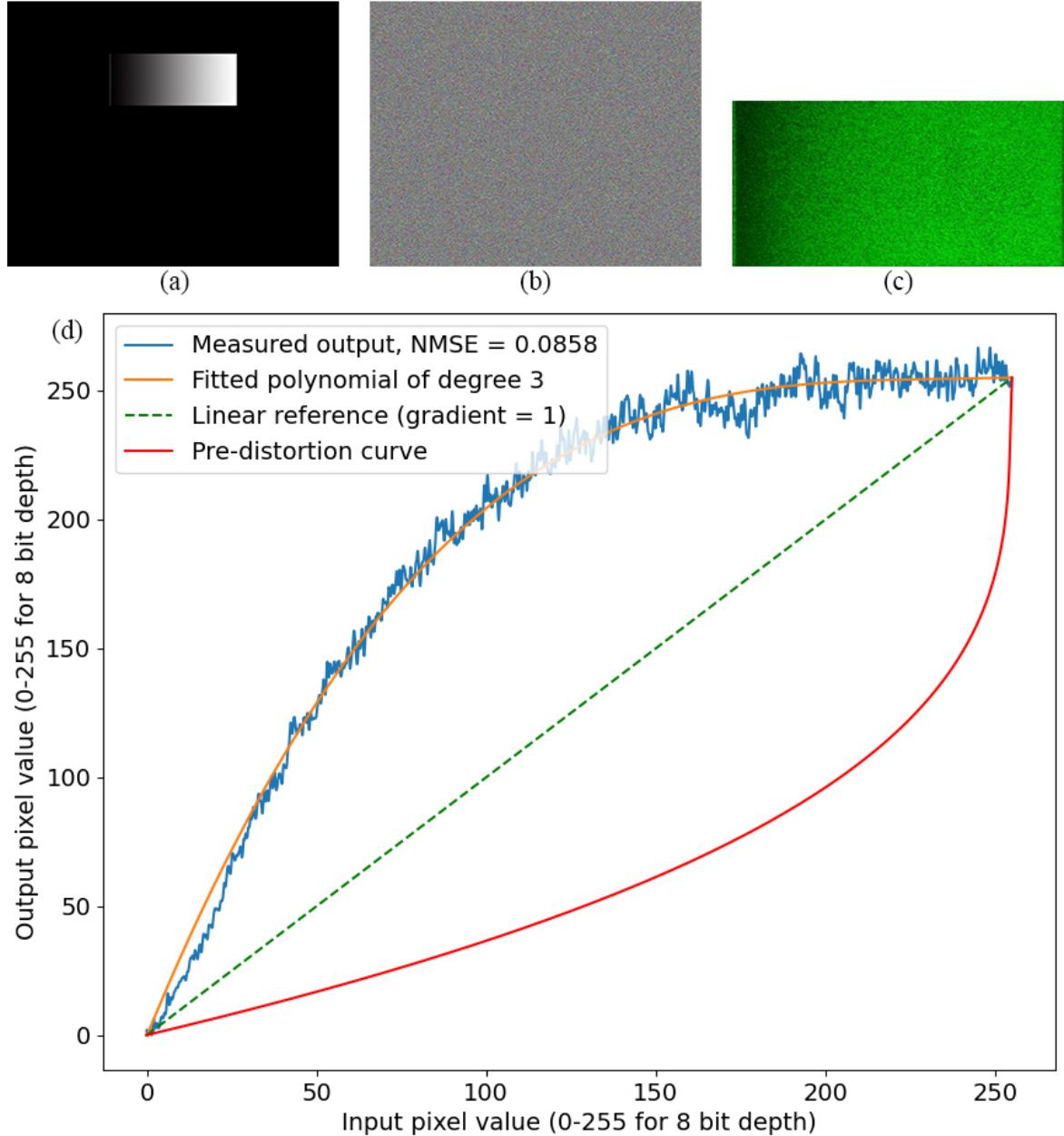


Fig. 3.3 Determining the DPD curve. (a) Input linear grey-scale ramp. (b) Corresponding CGH of (a) with 24-subframe binary phase encoding. (c) Holographic projection replay field of (b). (d) Plot of non-linearity measurement and according pre-distortion curve.

To determine the DPD curve of the holographic projection system, the non-linearity needs to be measured first. The hologram in Fig. 3.3(b) was first generated using OSPR algorithm for the linear grey-scale ramp of pixel value increasing linearly from 0 to 255, as shown in

Fig. 3.3(a), along with a single pixel white (255) strip at the left end as a fiducial marker to demonstrate the beginning of the grey-scale region [15].

The projection output of the linear grey-scale ramp was then captured and cropped as shown in Fig. 3.3(c), from which the non-linearity curve was determined, by averaging each column of pixels in the image and discarding the fiducial marker, forming the blue line in Fig. 3.3(d). A third-order polynomial fit was applied, generating a smoothed non-linearity curve (yellow line in Fig. 3.3(d)).

There exhibits a high degree of non-linearity. By taking the mean of the square of the error between the measured output (blue line) and the linear reference (green dashed line), the normalized mean squared error (NMSE) of the measured output was calculated to be 0.0858. To correct for the non-linearity, the DPD curve (red line) was formed by inverting the smoothed non-linearity curve (yellow line) in Fig. 3.3(d).

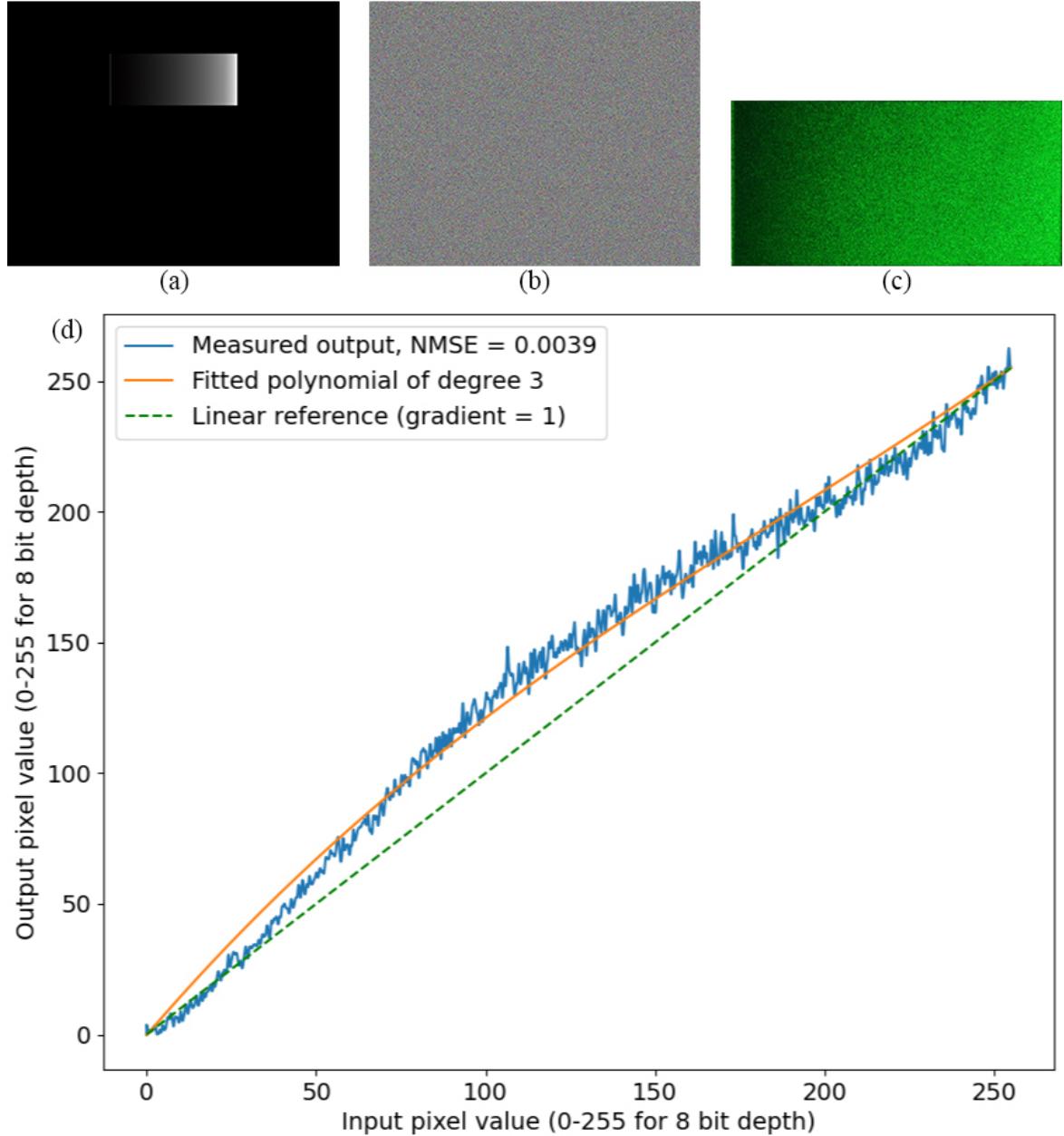


Fig. 3.4 Validation of DPD curve on the grey-scale ramp. (a) Pre-distorted ramp. (b) Corresponding CGH of (a) with 24-subframe binary phase encoding. (c) Holographic projection replay field of (b). (d) Non-linearity measurement after DPD.

Subsequently, the DPD curve (red line in Fig. 3.3(d)) was used to adjust the grey-scale ramp, achieving the pre-distorted grey-scale ramp as shown in Fig. 3.4(a). The according projection output was then captured as shown in Fig. 3.4(c). By using the same method of averaging columns of pixels, the measured output was plotted in Fig. 3.4(d). It can be seen that the

corrected non-linearity was much closer to linear comparing to the original non-linearity, and the NMSE was calculated to be 0.0039.

	NMSE	Percentage
Before DPD	0.0858	100%
After DPD	0.0039	4.55%

Table 3.1 Non-linearity results before and after DPD

Hence, as demonstrated in Table 3.1, DPD achieved a 95.45% reduction in NMSE, which was a significant improvement in non-linearity, therefore the DPD curve measured is validated.

### 3.4 Applying the DPD Curve

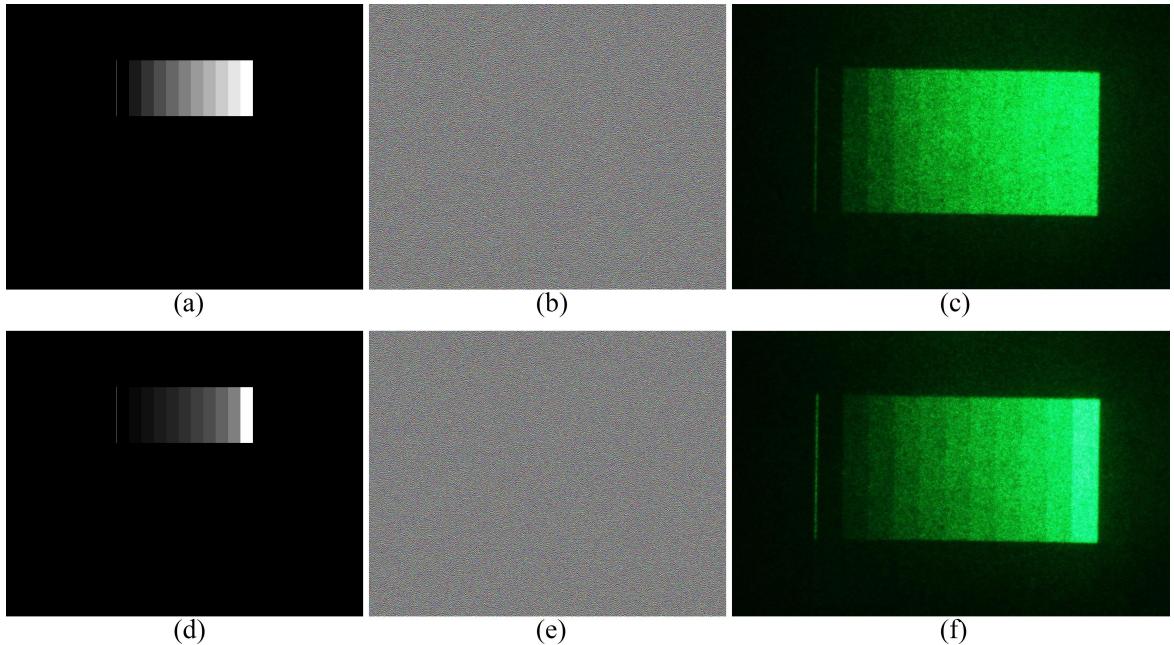


Fig. 3.5 Application of DPD on the 10-step strips. (a) 10 strips with equal step of pixel value. (b) CGH of (a). (c) Holographic projection replay field of (b). (d) After DPD of (a). (e) CGH of (d). (f) Holographic projection replay field of (e).

To qualitatively demonstrate the effectiveness of our approach, we project a simple test pattern of a graduated ramp test pattern consisted of 10-step strips in Fig. 3.5(a), which is commonly employed in gamma-correction calibration of many display systems. As shown

in the projection replay field captured in Fig. 3.5(c), before DPD, the right few strips are barely distinguishable. In comparison, after carrying out DPD, it can be seen that each pair of adjacent strips in Fig. 3.5(f) are much more distinguishable, qualitatively showing the effectiveness of the DPD method.

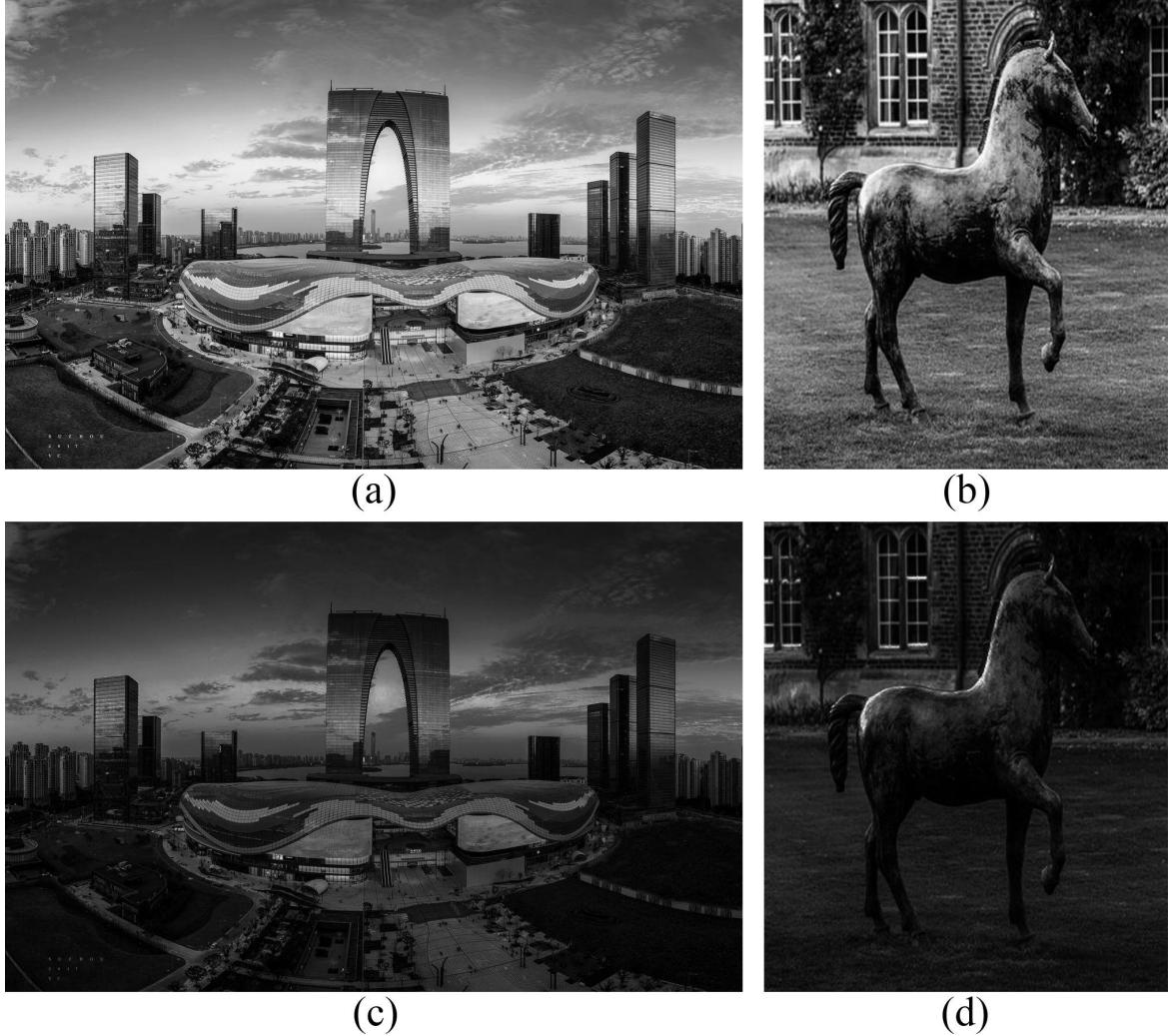


Fig. 3.6 Application of DPD on two sample real-word image. (a) Sample image 1: City Scene [18]. (b) Sample image 2: Horse. (c) Sample image 1 after DPD. (d) Sample image 2 after DPD.

Then the DPD curve was applied to the two sample images as shown in Fig. 3.6 (a) and (b), producing pre-distorted images in Fig. 3.6 (c) and (d). Holograms were generated for each image using the OSPR algorithm and loaded onto the SLM respectively. The according replay fields were captured as shown in Fig. 3.7.

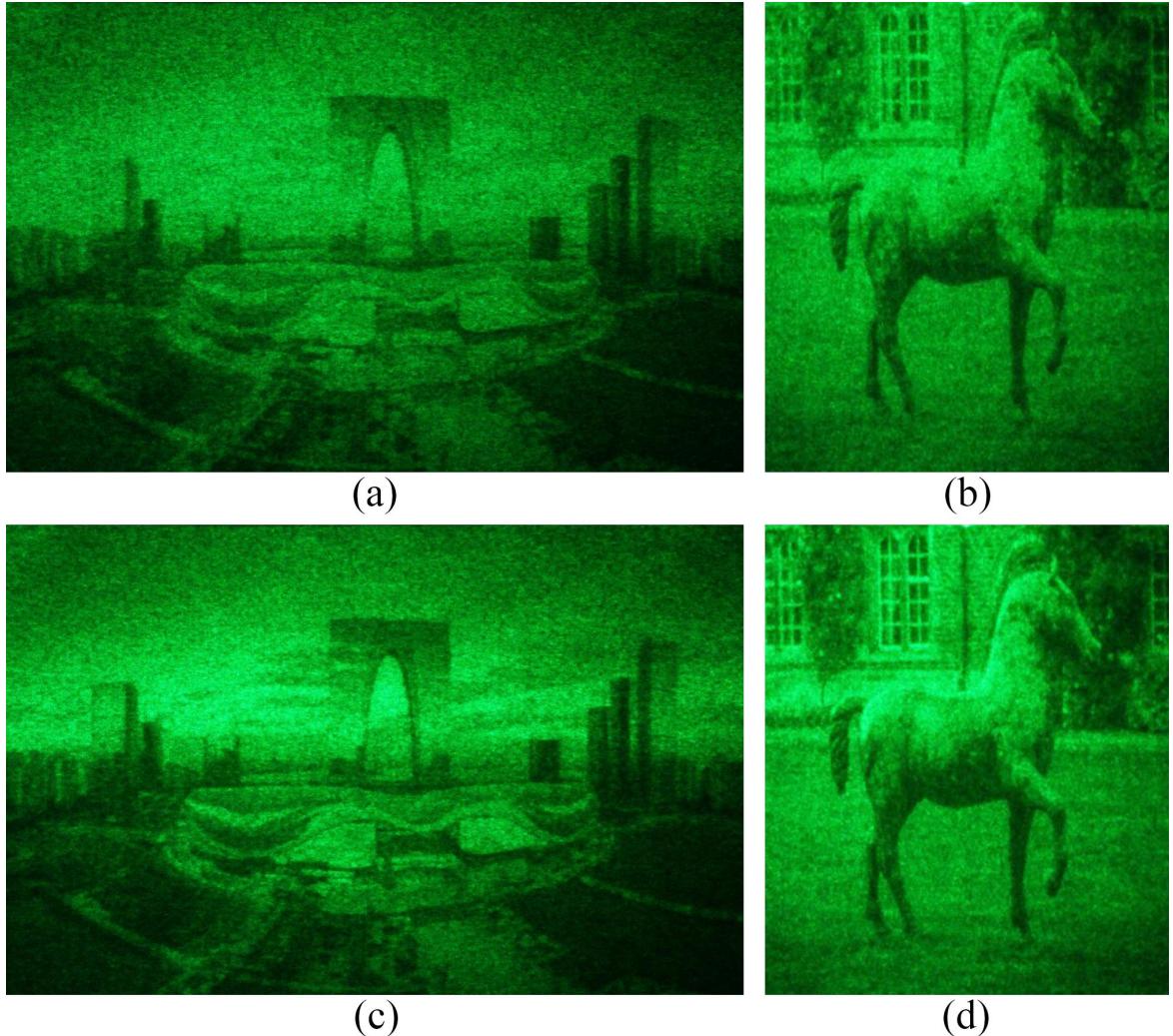


Fig. 3.7 Projection output of the two sample images before and after DPD. (a) Replay field of Sample image 1 before DPD ( $\text{NMSE}=0.06139$ ). (b) Replay field of Sample image 2 before DPD ( $\text{NMSE}=0.04309$ ). (c) Replay field of Sample image 1 after DPD ( $\text{NMSE}=0.04920$ ). (d) Replay field of Sample image 2 after DPD ( $\text{NMSE}=0.03635$ ).

The replay fields of the holographic projection of original images are shown in Fig. 3.7 (a) and (b), and the replay fields of the holographic projection of images after DPD are shown in Fig. 3.7 (c) and (d).

As shown in Fig. 3.7(a), it can be seen that, before DPD, the edges between the buildings and the sky were quite ambiguous, with most detail of the sky being lost. In comparison, after DPD, the replay field in Fig. 3.7(c) provided not only sharper edges between buildings and

the sky, but also more detail of clouds in the sky. The NMSE of the replay field for sample image 1 decreased from 0.06139 to 0.04920, which was a 19.86% reduction.

In Fig. 3.7(b), before DPD, the horse was difficult to distinguish from the background, especially around the horse's back area. But after DPD, as shown in Fig. 3.7(d), contrast has been significantly boosted and the fine detail around this part of the horse is more evident. The NMSE of the replay field for sample image 2 decreased from 0.04309 to 0.03635, which was a 15.64% reduction.

Sample image 1	NMSE	Percentage
Before DPD	0.06139	100%
After DPD	0.04920	80.15%
Sample image 2	NMSE	Percentage
Before DPD	0.04309	100%
After DPD	0.03635	84.36%

Table 3.2 DPD results for sample images

Hence, as summarised in Table 3.2, DPD achieved a 19.86% reduction in NMSE for sample image 1 and a 15.64% reduction in NMSE for sample image 2, quantitatively proving the effectiveness of DPD method for CGH of real-world test images using OSPR algorithm.

Lastly, as the DPD is a one-to-one mapping, the computation time is negligible. In practice, the computational overhead is too small to be measured against randomness between subsequent runs. DPD can also be further accelerated in hardware using a hardware LUT, so that the DPD can be carried out instantly. This approach is widely adopted in gamma correction for displays.

### 3.5 Summary

The non-linearity between target image and reconstructed image was measured for the overall holographic projection system by projecting a linear grey-scale ramp. Then the DPD curve was applied to the grey-scale ramp and successfully reduced the NMSE by 95.45%. To examine its effectiveness on real world images, the DPD method was applied on two sample images, it was observed that more details were shown in the replay field after DPD, and the NMSE's of the two example images were reduced by 19.86% and 15.64%. As

the DPD is a one-to-one mapping, the extra computation required is negligible. Hence, we have demonstrated the effectiveness of the proposed DPD-OSPR method to improve reconstruction quality on the existing OSPR algorithm while still keeping its ability for real-time holography.



# Chapter 4

## L-BFGS Optimisation of Phase-Only Hologram

As previously introduced in Chapter 2, currently available spatial light modulators (SLMs) can only modulate either phase or amplitude, so algorithms are needed to compute amplitude-only or phase-only holograms, among which the phase-only holograms are usually preferred due to their higher energy efficiency, leading to the emergence of the classical phase retrieval algorithms reviewed in Section 2.3. With the developments in modern numerical optimization methods and computational power, advances in CGH algorithms can be made. A search in the literature has found some recent work generating CGH using numerical optimisation methods [60–63]. This chapter therefore implements the use of numerical optimisation methods for hologram generation, including the novel method using the L-BFGS optimiser to compute phase-only holograms. The following sections start from the background knowledge of numerical optimization including L-BFGS algorithm, which also serves as the theoretical background for Chapter 5 and Chapter 6, then introduces and carries out the optimization of phase hologram.

### 4.1 Numerical Optimisation Methods

#### 4.1.1 Optimisation framework

Numerical optimisation methods aim to find an optimal solution which minimise an objective function numerically. They begin with an initial guess of the optimal solution ( $\mathbf{x}_0$ ) and

then, after iterations, generate a sequence of gradually improved estimates until they reach a solution [64]. If we have  $\mathbf{x}$  as the vector of variables, and denote  $f(\mathbf{x})$  as the objective function, which is a function of  $x$  we want to minimise, any unconstrained optimisation problem can be written as:

$$\underset{\mathbf{x} \in R^n}{\text{minimise}} \quad f(\mathbf{x}) \quad (4.1)$$

Numerical optimisation then calculate the optimal solution  $\mathbf{x}^*$  iteratively, the iteration is given by:

$$\mathbf{x}_{k+1} = \mathbf{x}_k + \alpha_k \mathbf{p}_k \quad (4.2)$$

where the positive scalar  $\alpha_k$  is called step length, or sometimes may be referred as ‘learning rate’ in some context, and the vector  $\mathbf{p}_k$  is the search direction, which usually takes the form of

$$\mathbf{p}_k = -\mathbf{B}_k^{-1} \nabla f_k \quad (4.3)$$

where  $\mathbf{B}_k$  is a nonsingular matrix that varies for different optimisation methods, and  $\nabla f_k$  is the gradient, which, if unable to evaluate directly, can be approximated by:

$$\nabla f_k \approx \frac{f_{k+1} - f_k}{\mathbf{x}_{k+1} - \mathbf{x}_k}$$

where  $f_k$  denotes  $f(\mathbf{x}_k)$  (4.4)

The strategy used to determine  $\mathbf{p}_k$  distinguishes one algorithm from another. Most methods make use of the values of  $f$ ,  $\nabla f$  and  $\nabla^2 f$ , and some methods even make use of the accumulated historical values of those derivatives, which are further discussed in Section 4.1.2 - Section 4.1.5.

### 4.1.2 Gradient Descent

Gradient descent (GD) is a first-order optimisation method, it finds a local minimum by following the negative of the gradient (i.e. the steepest descent direction). The  $\mathbf{B}_k$  (in Eq. (4.3)) for gradient descent simply takes the value of  $\mathbf{I}$ , which is the identity matrix. And the search direction becomes:

$$\mathbf{p}_k = -\nabla f_k \quad (4.5)$$

The steepest descent method is very intuitive: among all possible directions to move away from  $\mathbf{x}_k$ , the steepest gradient direction is the one which  $f$  decreases most rapidly. The

advantage of this method is that it requires few computation and memory resource, because it only requires a computation of the first derivative, and it does not require any accumulation of historical gradients. However, it is a greedy method that only considers the current iteration without any global consideration, so it can be extremely slow on complicated problems. [64]

To work around the disadvantage, a few variants have emerged, such as AdaGrad [65], RMSProp [66] and Adam [67] which combines the advantages of AdaGrad and RMSProp. The Adam method is an iconic variant of GD, often referred to as ‘gradient descent with momentum’, as the name Adam is derived from ‘adaptive moment estimation’. Adam algorithm is based on adaptive estimates of lower-order moments [67]. It computes individual adaptive learning rates for different parameters from estimates of first and second moments of the gradients [67].

### 4.1.3 Newton’s Method

Newton’s method is a second-order optimisation method. Its search direction is derived from the second-order Taylor series approximation to  $f(\mathbf{x}_k + \mathbf{p})$ , which is

$$f(\mathbf{x}_k + \mathbf{p}) \approx f_k + \mathbf{p}^T \nabla f_k + \frac{1}{2} \mathbf{p}^T \nabla^2 f_k \mathbf{p} \stackrel{\text{def}}{=} m_k(\mathbf{p}) \quad (4.6)$$

The Newton direction can then be obtained by finding the vector  $\mathbf{p}$  that minimises  $m_k(\mathbf{p})$ . By setting the derivative of  $m_k(\mathbf{p})$  to zero,  $\mathbf{p}$  can be obtained as:

$$\mathbf{p}_k = -\nabla^2 f_k^{-1} \nabla f_k \quad (4.7)$$

By comparing Eq. (4.7) to Eq. (4.3), it can be seen that the Newton’s method has a  $\mathbf{B}_k$  of  $\nabla^2 f_k$ . Unlike the gradient descent method, there is a “natural” step length of 1 associated with the Newton direction, so  $\alpha_k = 1$  by default and is only adjusted when it does not produce a satisfactory reduction in the value of  $f$ .

The Newton direction is reliable when the difference between the true function  $f(\mathbf{x}_k + \mathbf{p})$  and its quadratic model  $m_k(\mathbf{p})$  is not too large. Methods that use the Newton direction have a fast rate of local convergence, typically quadratic. After a neighbourhood of the solution is reached, convergence to high accuracy often occurs in just a few iterations. The main drawback of the Newton direction is the need for the Hessian  $\nabla^2 f_k$ . Explicit computation of this matrix of second derivatives can sometimes be a cumbersome, error-prone, and expensive process. [64]

#### 4.1.4 Quasi-Newton Method

Quasi-Newton method provides an attractive alternative to Newton's method, in that they do not require computation of the Hessian and yet still attain a super linear rate of convergence. In place of the true Hessian  $\nabla^2 f_k$ , they use an approximation  $\mathcal{H}_k \stackrel{\text{def}}{=} \mathbf{B}_k^{-1}$ , which is updated after each step to take account of the additional knowledge gained during the step. The updates make use of the fact that changes in the gradient provide information about the second derivative of  $f$  along the search direction. The most popular quasi-Newton algorithm is the Broyden-Fletcher-Goldfarb-Shanno (BFGS) method, named for its discoverers Broyden, Fletcher, Goldfarb, and Shanno. [64]

The process of the BFGS method is shown below:

$$\text{denote } \begin{cases} \mathcal{H}_k &= \mathbf{B}_k^{-1} \\ \mathbf{p}_k &= -\mathcal{H}_k \nabla f_k \end{cases} \quad (4.8)$$

$$\text{Initiate } \mathcal{H}_0 \leftarrow \frac{\mathbf{y}_k^T \mathbf{s}_k}{\mathbf{y}_k^T \mathbf{y}_k} \mathbf{I} \quad (4.9)$$

$$\text{update } \mathcal{H}_{k+1} = (\mathbf{I} - \rho_k \mathbf{s}_k \mathbf{y}_k^T) \mathcal{H}_k (\mathbf{I} - \rho_k \mathbf{y}_k \mathbf{s}_k^T) + \rho_k \mathbf{s}_k \mathbf{s}_k^T \quad (4.10)$$

$$\text{where } \begin{cases} \mathbf{s}_k &= \mathbf{x}_{k+1} - \mathbf{x}_k \\ \mathbf{y}_k &= \nabla f_{k+1} - \nabla f_k \\ \rho_k &= \frac{1}{\mathbf{y}_k^T \mathbf{s}_k} \end{cases} \quad (4.11)$$

The algorithm is robust, and its rate of convergence is super linear, which is fast enough for most practical purposes. Even though Newton's method converges more rapidly (that is, quadratically), its cost per iteration usually is higher, because of its need for second derivatives and solution of a linear system. The drawback is that, it is not directly applicable to large optimisation problems because  $\mathcal{H}_k$ 's are usually dense, requiring large storage and computational requirements. [64]

#### 4.1.5 Large Scale Quasi-Newton Method: Limited Memory BFGS (L-BFGS)

L-BFGS algorithm [68] modifies the technique described in Section 4.1.4 to obtain Hessian approximations that can be stored compactly in just a few vectors of length  $n$ , where  $n$  is the number of unknowns in the problem. The main idea of this method is to use curvature information from only the most recent iterations to construct the Hessian approximation.

Curvature information from earlier iterations, which is less likely to be relevant to the actual behaviour of the Hessian at the current iteration, is discarded in the interest of saving storage. [64]

Denoting  $\mathbf{V}_k = \mathbf{I} - \rho_k \mathbf{y}_k \mathbf{s}_k^T$ , Eq. (4.10) can be written as:

$$\mathcal{H}_{k+1} = \mathbf{V}_k^T \mathcal{H}_k \mathbf{V}_k + \rho_k \mathbf{s}_k \mathbf{s}_k^T \quad (4.12)$$

The inverse Hessian approximation  $\mathcal{H}_k$  will generally be dense, so that the cost of storing and manipulating it is prohibitive when the number of variables is large. To circumvent this problem, we store a modified version of  $\mathcal{H}_k$  implicitly, by storing a certain number (say,  $m$ ) of the vector pairs  $\{\mathbf{s}_i, \mathbf{y}_i\}$  used in the Eq. (4.10) and Eq. (4.11). The product  $\mathcal{H}_k \nabla f_k$  can be obtained by performing a sequence of inner products and vector summations involving  $\nabla f_k$  and the pairs  $\{\mathbf{s}_i, \mathbf{y}_i\}$ . After the new iterate is computed, the oldest vector pair in the set of pairs  $\{\mathbf{s}_i, \mathbf{y}_i\}$  is replaced by the new pair  $\{\mathbf{s}_k, \mathbf{y}_k\}$  obtained from the current step (Eq. (4.11)). In this way, the set of vector pairs includes curvature information from the  $m$  most recent iterations. Practical experience has shown that modest values of  $m$  (between 3 and 20, say) often produce satisfactory results. We now describe the updating process in a little more detail. At iteration  $k$ , the current iterate is  $\mathbf{x}_k$  and the set of vector pairs is given by  $\{\mathbf{s}_i, \mathbf{y}_i\}$  for  $i = k-m, \dots, k-1$ . We first choose some initial Hessian approximation  $\mathcal{H}_k^0$  (in contrast to the standard BFGS iteration, this initial approximation is allowed to vary from iteration to iteration) and find by repeated application of Eq. (4.10) that the L-BFGS approximation  $\mathcal{H}_k$  satisfies the following formula: [64]

$$\begin{aligned} \mathcal{H}_k = & (\mathbf{V}_{k-1}^T \cdots \mathbf{V}_{k-m}^T) \mathcal{H}_k^0 (\mathbf{V}_{k-m} \cdots \mathbf{V}_{k-1}) \\ & + \rho_{k-m} (\mathbf{V}_{k-1}^T \cdots \mathbf{V}_{k-m+1}^T) \mathbf{s}_{k-m} \mathbf{s}_{k-m}^T (\mathbf{V}_{k-m+1} \cdots \mathbf{V}_{k-1}) \\ & + \rho_{k-m+1} (\mathbf{V}_{k-1}^T \cdots \mathbf{V}_{k-m+2}^T) \mathbf{s}_{k-m+1} \mathbf{s}_{k-m+1}^T (\mathbf{V}_{k-m+2} \cdots \mathbf{V}_{k-1}) \\ & + \cdots \\ & + \rho_{k-1} \mathbf{s}_{k-1} \mathbf{s}_{k-1}^T \end{aligned} \quad (4.13)$$

From this expression we can derive a recursive procedure (Algorithm 8) to compute the product  $\mathcal{H}_k \nabla f_k$  efficiently.

**Algorithm 8** L-BFGS two-loop recursion [64]

---

```

 $\mathbf{q} \leftarrow \nabla f_k$ 
for  $i = k - 1, k - 2, \dots, k - m$  do
     $\alpha_i \leftarrow \rho_i \mathbf{s}_i^T \mathbf{q}$ 
     $\mathbf{q} \leftarrow \mathbf{q} - \alpha_i \mathbf{y}_i$ 
end for
 $\mathbf{r} \leftarrow \mathcal{H}_k^0 \mathbf{q}$ 
for  $i = k - m, k - m + 1, \dots, k - 1$  do
     $\beta \leftarrow \rho_i \mathbf{y}_i^T \mathbf{r}$ 
     $\mathbf{r} \leftarrow \mathbf{r} + \mathbf{s}_i (\alpha_i - \beta)$ 
end for
Step with  $\mathbf{p}_k \leftarrow -\mathcal{H}_k \nabla f_k = -\mathbf{r}$ 

```

---

Apart from being inexpensive, L-BFGS has the advantage that the multiplication by the initial matrix  $\mathcal{H}_k^0$  is isolated from the rest of the computations, allowing this matrix to be chosen freely and to vary between iterations. A method for choosing  $\mathcal{H}_k^0$  that has proved effective in practice is to use the same as BFGS as stated in Eq. (4.9). [64]

## 4.2 Optimisation of Phase Hologram for a Target Image

To implement the optimisation methods listed in Section 4.1 on CGH, firstly the optimisation framework needs to be adapted. The objective of CGH is to find the phase hologram ( $H$ ) that has the optimal reconstruction ( $R$ ) matching the target image ( $T$ ), which can be formulated as minimising the difference between  $R$  and  $T$  by varying  $H$ , leading to the mathematical expression below:

$$\arg \min_{\mathbf{H}} Loss(\mathbf{T}, \mathbf{R}) \quad (4.14)$$

where ‘*Loss*’ denotes an error function quantifying the difference between  $T$  and  $R$ , and ‘*arg*’ returns the argument ( $H$ ) instead of the error value. For the *Loss* function, the classic error function mean-squared error (MSE) [69] is selected, with its definition as shown below:

$$MSE(\mathbf{T}, \mathbf{R}) = \frac{1}{X \times Y} \sum_{x=1}^X \sum_{y=1}^Y (\mathbf{T}_{x,y} - \mathbf{R}_{x,y})^2 \quad (4.15)$$

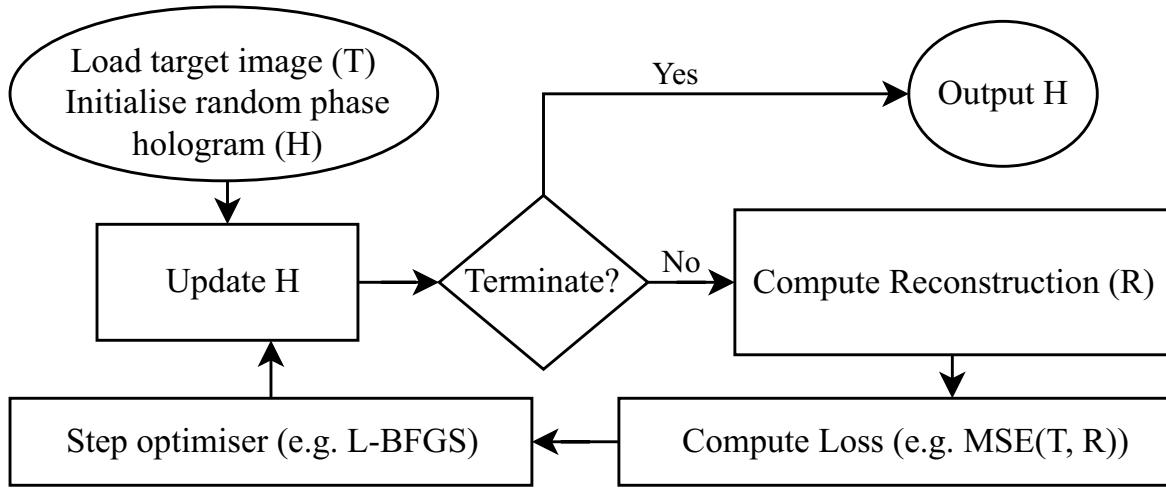


Fig. 4.1 Flowchart of the optimisation process

The optimisation process for generating a phase hologram to reconstruct a target image is depicted in Fig. 4.1. The process begins by loading the target image ( $T$ ) and initializing a random phase hologram ( $H$ ). Using this initial phase hologram, a reconstruction ( $R$ ) of the target image is computed via either the Fraunhofer or Fresnel propagation equation introduced in Section 2.2.2 based on the distance needed. The difference between the target image ( $T$ ) and the reconstruction ( $R$ ) is quantified by computing the loss, such as the MSE in Eq. (4.15). An optimizer (such as the GD, Adam or L-BFGS algorithms mentioned in Section 4.1) then updates the phase hologram based on the computed loss. This iterative process of reconstruction, loss calculation, and hologram update continues until a termination condition is met, usually a fixed total number of iterations, or when optimization has reached convergence. The final optimised phase hologram ( $H$ ) is then outputted.

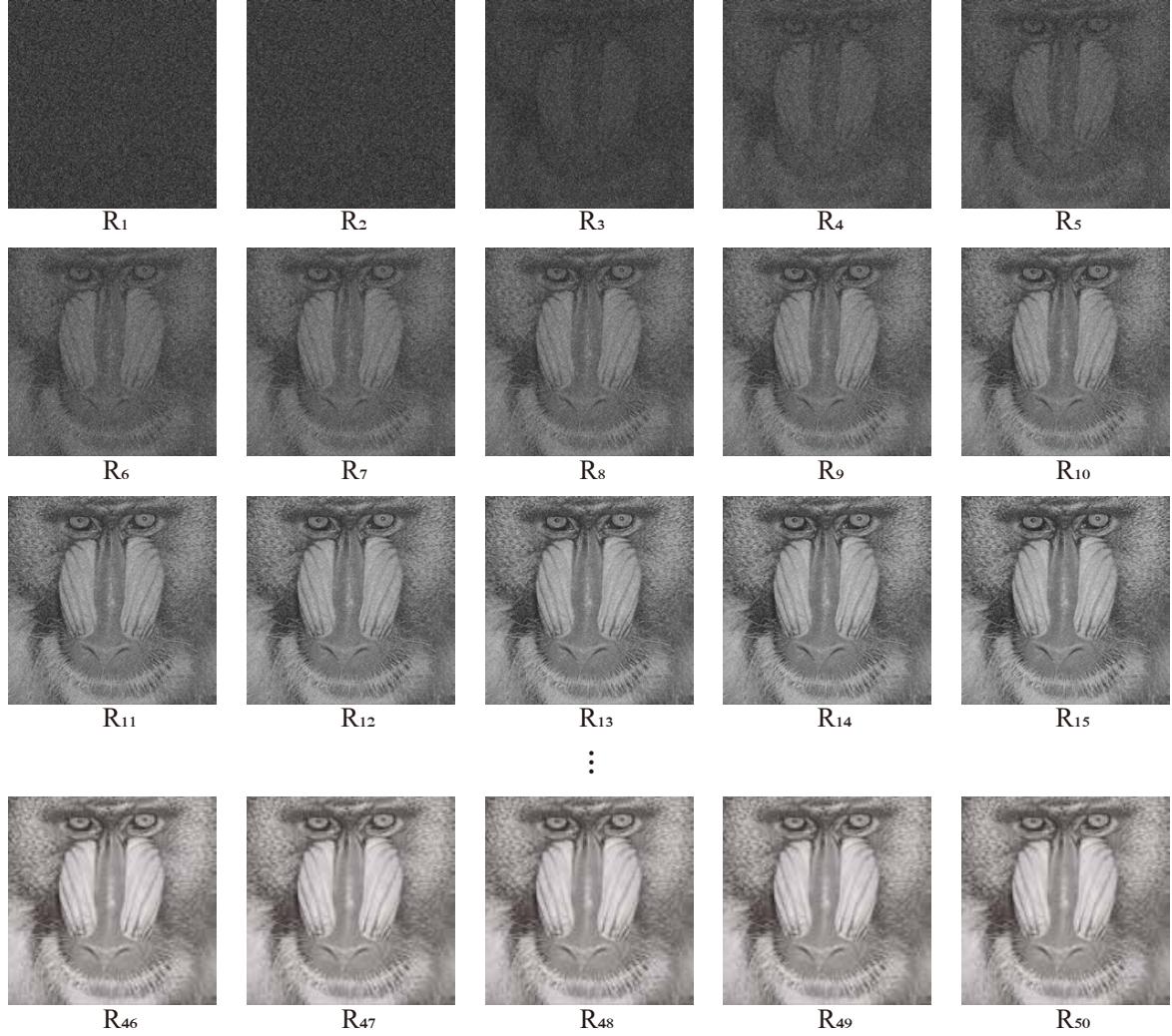


Fig. 4.2 Reconstructions at each iteration of the L-BFGS optimisation

The optimisation flowchart in Fig. 4.1 is run on the example target image in Fig. 2.9 for a total of 50 iterations. The reconstruction (R) at each iteration is listed in Fig. 4.2, with iterations 16 to 45 omitted to save space. The list of reconstructions demonstrate visually how the optimisation converges to the resulting reconstruction ( $R_{50}$ ) which has a very good quality. It shows that the proposed method using the L-BFGS algorithm to generate phase hologram is effective.

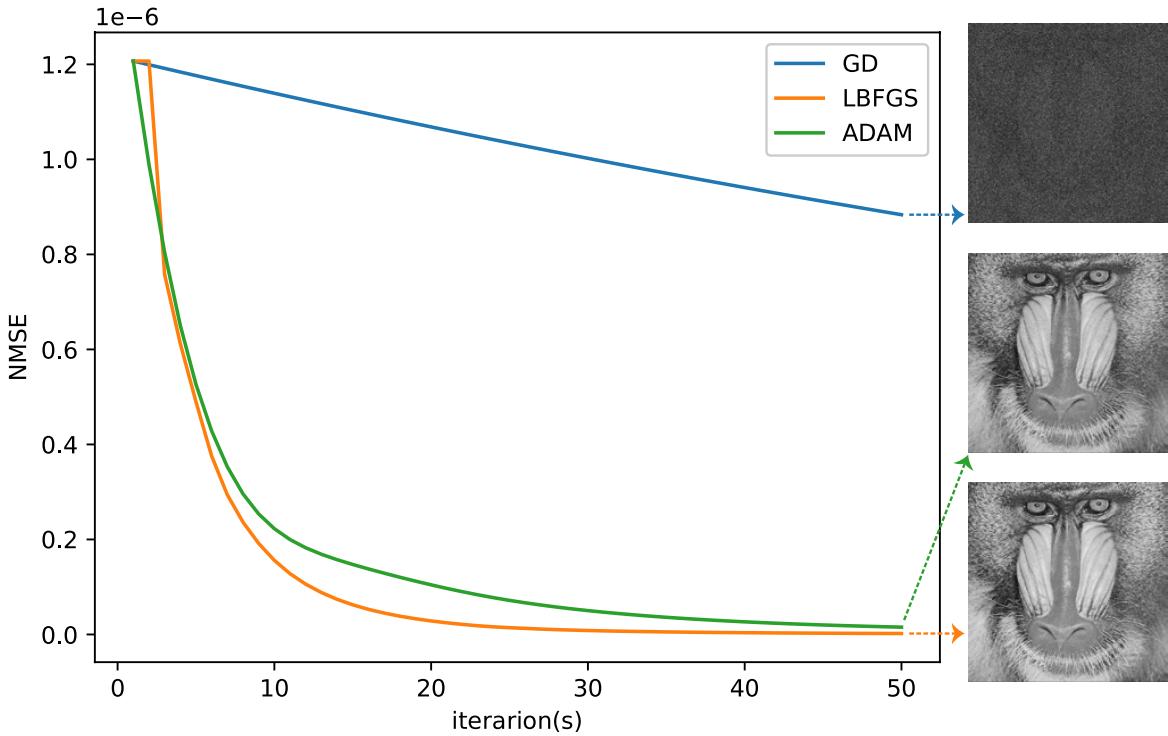


Fig. 4.3 Convergence plot for comparison between the GD, Adam and L-BFGS optimisations

Then the proposed method of using L-BFGS algorithm to generate phase hologram is quantitatively compared against the existing methods using GD and Adam algorithms, by plotting the normalised mean-squared error (NMSE) which is simply the MSE in Eq. (4.15) divided by the sum of squares of the target image value, so that pixel values are normalised, avoiding the huge difference in MSE value which is based on how the value range is defined (typically 0-1 or 0-255).

The proposed method using L-BFGS algorithm (the yellow line in Fig. 4.3) stagnates for the first two iterations as it needs to estimate the hessian as explained in Section 4.1.5. After the brief stagnation, the L-BFGS converges quickly, surpassing the existing methods in the literature using GD and Adam optimiser, corresponding to the blue line and the green line in Fig. 4.3 respectively. The final reconstructions are shown aside the NMSE plots, from which it can be seen that the L-BFGS algorithm produces a reconstruction of similar quality to the Adam algorithm, and both of them are much better than the GD algorithm. Such observation matches the final NMSE values. In summary, the L-BFGS algorithm was proven to be able to optimise a phase-only hologram for a target image, and is better than the Adam and GS optimisation algorithms.

### 4.3 Target Image Phase Optimisation

This section proposes a novel method of optimising the phase of the target image instead of optimising the phase of the hologram as previously introduced in Section 4.2, under the same ‘phase only’ constraint of the hologram. The objective is to find an optimal phase profile to be attached to the target image, so that its inverse propagation to the hologram plane produces a hologram whose phase has a reconstruction that best matches the target image. As it optimises the phase of the target image instead of the traditional way optimising the phase of the hologram, this method is named as Target Image Phase Optimisation (TIPO). A flowchart is drawn in Fig. 4.4 to clarify the detailed operation of the proposed TIPO algorithm.

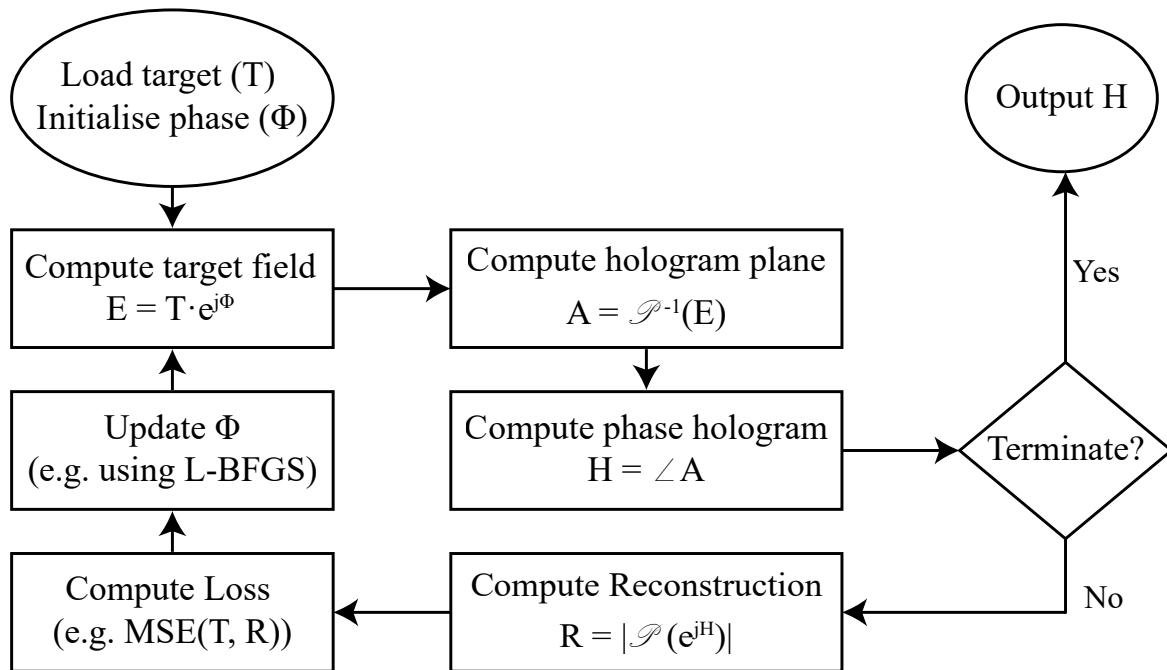


Fig. 4.4 TIPO flowchart

The flowchart in Fig. 4.4 outlines the TIPO algorithm process for generating a phase-only hologram ( $H$ ) to reconstruct a target image ( $T$ ). The procedure starts with loading the target image and initializing the phase ( $\Phi$ ). The target field ( $E$ ) is then computed from the target image and the target image phase ( $E = T \cdot e^{j\Phi}$ ). Next, the hologram aperture ( $A$ ) is computed by applying the inverse propagation ( $\mathcal{P}^{-1}$ ) to the target field, where  $\mathcal{P}$  is chosen from the Fraunhofer and Fresnel diffraction equations in Section 2.2.2. The phase-only hologram ( $H$ ) is then derived from the angle of the complex hologram aperture ( $H = \angle A$ ). Subsequently, a reconstruction ( $R$ ) is computed using the forward propagation ( $R = \mathcal{P}(e^{jH})$ ). The loss,

such as the Mean Squared Error (MSE), is calculated between the target image ( $T$ ) and the reconstruction ( $R$ ). An optimizer (e.g. SGD or L-BFGS) then updates the phase ( $\Phi$ ) based on the computed loss. This iterative process of computing the target field, hologram plane, phase hologram, reconstruction, and loss calculation continues, followed by phase updates, until a termination condition is met, which is usually set for a fixed number of iterations. Upon convergence, the final optimized phase hologram ( $H$ ) is produced.

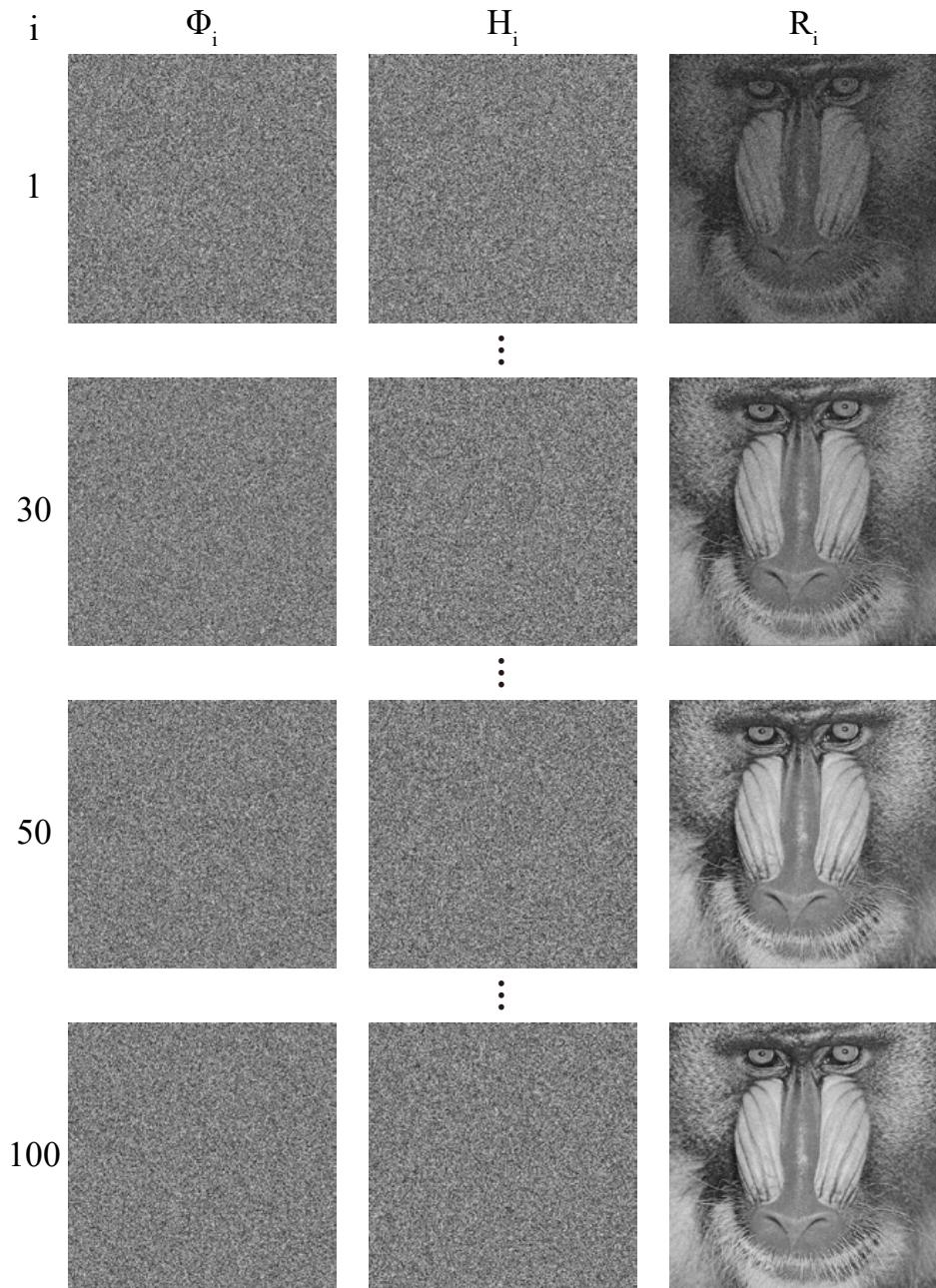


Fig. 4.5 TIPO iterations on the mandrill target

The results for an example run of the TIPO algorithm on the mandrill target (in Fig. 2.9) with the total number of iterations set to 100 are shown in Fig. 4.5 (only iterations number 1, 30, 50 and 100 are shown for space saving). As the first hologram is computed by the inverse propagation ( $\mathcal{P}^{-1}$ ) of the target image with a random phase, the reconstruction ( $R_1$ ) is already showing the target image. Then as the iterations continue, the reconstruction quality gets better, proving the effectiveness of the proposed TIPO algorithm.

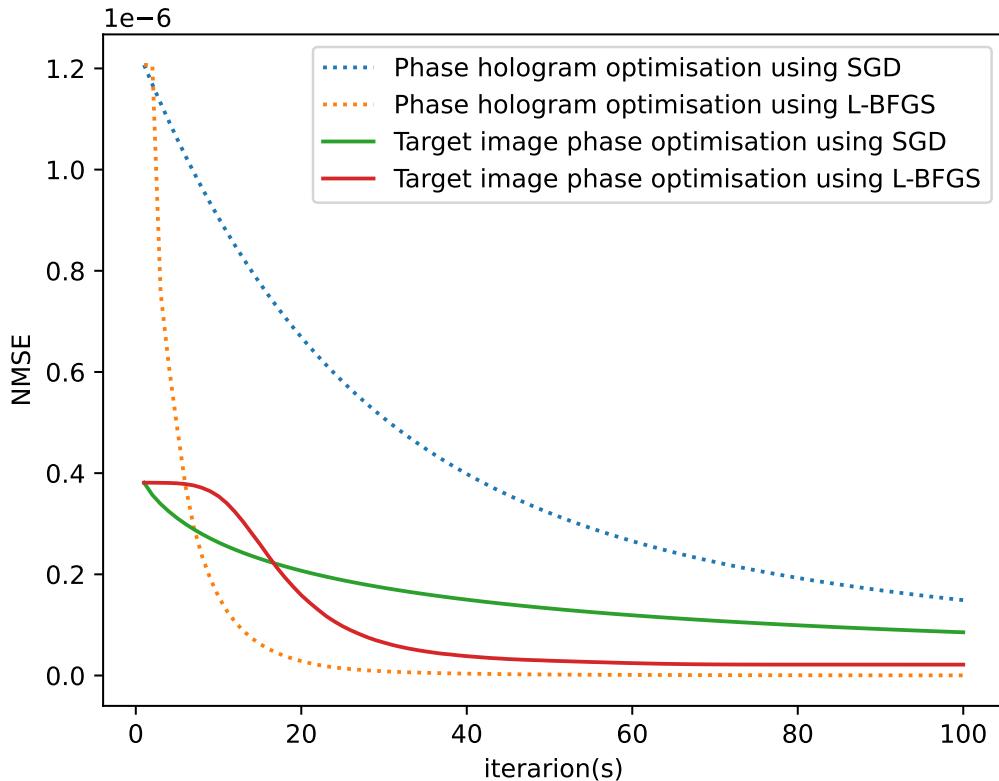


Fig. 4.6 TIPO convergence plot

To quantitatively analyse the results, the convergence of the TIPO algorithm is plotted in Fig. 4.6, where the NMSE between the reconstruction and the target image are plotted against the number of iterations. Both SGD and L-BFGS optimisers are run for the TIPO algorithm (corresponding to the green and red line in Fig. 4.6 accordingly), and are compared against the regular phase hologram optimisation algorithm in Section 4.2 (corresponding to the dotted blue line and the dotted yellow line in Fig. 4.6 respectively). The TIPO algorithms using both SGD and L-BFGS optimisers start with a lower NMSE of the reconstruction as their holograms in the first iteration are extracted from the inverse propagation of the target

image, instead of pure random holograms as done in the regular phase hologram optimisation. However, the regular phase hologram optimisation using L-BFGS optimiser quickly surpassed the TIPO algorithm within 10 iterations, and reached the lowest reconstruction error. For the SGD optimiser, the TIPO method (green line) has an significant improvement than the regular phase hologram optimisation method (dotted blue line). The two TIPO methods (in solid lines) lie between the two regular phase hologram optimisation method (in dotted lines), proving that although not being the best, the TIPO method is still an effective method for phase hologram generation.



# Chapter 5

## Multi-Depth Phase-Only Hologram Optimization with Sequential Slicing

*Note: The work in this chapter has been published in Ref. [4]*

This chapter implements the limited-memory Broyden-Fletcher-Goldfarb-Shanno (L-BFGS) optimization method introduced in Section 4.1.5 to generate phase-only computer-generated hologram (CGH) for both 2D and 3D targets. For 3D targets, instead of computing the full 3D reconstruction of the hologram at every iteration, a novel method called sequential slicing (SS) is proposed for partial evaluation of the hologram during optimization.

### 5.1 Introduction

The most common multi-depth CGH optimization methods either evaluate their loss against the entire multi-depth 3D target, which is time-consuming, or evaluate the hologram for each plane and then sum the holograms, which introduces quality degradation.

A search in the literature has found some recent work that compute CGH using numerical optimization methods [60–63], but speed and quality are still the major challenges.

proposes the sequential slicing (SS) technique for the optimization of CGH for multi-depth 3D target, which evaluates the loss for a single slice at each iteration, aiming for quicker

hologram generation with proper overall quality and low quality imbalance across the multiple depths enabled by the second-order nature of the L-BFGS optimization algorithm.

## 5.2 Method

The optimization algorithms used for CGH in this paper are GD and L-BFGS. The phase-only constraint of CGH can be easily applied by fixing a constant amplitude of the hologram, while keeping its phase varying and being the argument of optimization ( $\mathbf{x}$ ).

### 5.2.1 Hologram Optimization for Multi-Depth Targets

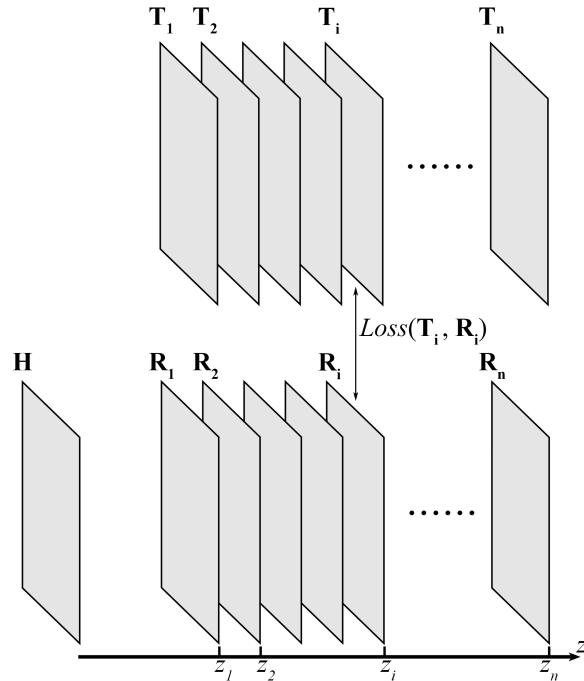


Fig. 5.1 Loss between the multi-depth targets ( $\mathbf{T}_1$  to  $\mathbf{T}_n$ ) and the reconstructions ( $\mathbf{R}_1$  to  $\mathbf{R}_n$ ) of hologram  $\mathbf{H}$

As shown in Fig. 5.1, the multi-depth target is set up as a collection of  $n$  slices ( $\mathbf{T}_1$  to  $\mathbf{T}_n$ ), each slice  $\mathbf{T}_i$  is at a distance  $z_i$  to the hologram plane. And for the hologram  $\mathbf{H}$ , its reconstruction at each distance  $z_i$  is computed using Fresnel diffraction formula in Eq. (7.1), which is labelled as the propagation function  $\mathcal{P}(\mathbf{H}, z_i)$ .

Then, to create an objective function ( $f(\mathbf{x})$  in Eq. (4.1)) to optimize, we need to quantify the difference between each target slice ( $\mathbf{T}_i$ ) and the respective reconstruction ( $\mathbf{R}_i$ ) numerically, and minimize such difference. Among the lots of options available, the loss functions selected

are mean squared error (MSE) [69], cross entropy (CE) [70] and relative entropy (RE) [71]. To adapt the loss functions for two-dimensional (2D) target image  $\mathbf{T}_i$  and reconstructed image  $\mathbf{R}_i$  of dimension  $X \times Y$ , the loss functions are adapted as shown in Eq. (5.1) to Eq. (6.2).

$$MSE(\mathbf{T}_i, \mathbf{R}_i) = \frac{1}{X \times Y} \sum_{x=1}^X \sum_{y=1}^Y (\mathbf{T}_{i;x,y} - \mathbf{R}_{i;x,y})^2 \quad (5.1)$$

$$CE(\mathbf{T}_i, \mathbf{R}_i) = - \sum_{x=1}^X \sum_{y=1}^Y \mathbf{T}_{i;x,y} \log(\mathbf{R}_{i;x,y}) \quad (5.2)$$

$$RE(\mathbf{T}_i, \mathbf{R}_i) = - \sum_{x=1}^X \sum_{y=1}^Y \mathbf{T}_{i;x,y} \log \left( \frac{\mathbf{R}_{i;x,y}}{\mathbf{T}_{i;x,y}} \right) \quad (5.3)$$

Mean squared error (MSE) is adapted as shown in Eq. (5.1). MSE is a traditional metric averaging the squared error between the target and observed values. Cross entropy (CE) is adapted as shown in Eq. (5.2). CE is often used in classification problems, such as language modelling [72]. Relative entropy (RE), also called Kullback-Leibler divergence (usually denoted as  $D_{KL}(P||Q)$ ), is adapted as shown in Eq. (6.2). For the uniformity of symbols in this paper, relative entropy is denoted as RE. It is a measure of how much a probability distribution  $P$  is different from another probability distribution  $Q$ . Both CE and RE are usually computed between the true probabilistic distribution and the predicted probabilistic distribution, while the images are not probability distributions, the pixel values are normalized to decimal numbers in the range of 0 to 1 so that CE and RE can be applied.

The effectiveness of L-BFGS optimization of phase-only CGH for a single slice target image ( $n = 1$ ) has been demonstrated in the previous research [2]. However, for a 3D target consisted of multiple slices at different depths, the optimization of CGH becomes challenging.

The typical technique is to sum the losses computed for each slice for each iteration during optimization, which is called the Sum-of-Loss (SoL) method in this paper. At every iteration, it computes the full 3D reconstructions  $(\mathbf{R}_1, \dots, \mathbf{R}_n)$  of the hologram  $\mathbf{H}$  at every distance  $z_i$ , requiring a total of  $n$  Fourier Transforms, and then computes the total loss between all the target slices and the reconstructed slices, therefore fully evaluating the hologram at each step. Although the total number of iterations does not scale with  $n$ , the number of Fourier Transforms performed at each iteration scales up with  $n$ , as it needs to compute the full

multi-depth reconstructions at every iteration.

$$\text{SoL: } \arg \min_{\mathbf{H}} \sum_{i=1}^n Loss(\mathbf{T}_i, \mathbf{R}_i) \quad (5.4)$$

Another universal technique for 3D CGH is to sum the sub-holograms  $\mathbf{H}_i$  generated for each target slice  $\mathbf{T}_i$  to form a total hologram based on the principle of superposition, which is called the Sum-of-Hologram (SoH) method in this paper. For a fixed number of iterations for each sub-hologram, the total computation scales up linearly with the number of slices  $n$ . SoH method's advantage is its ease of implementation, that it can compute multi-depth 3D CGH based on any existing single slice CGH algorithm. Its major disadvantage for phase-only hologram generation is that, the final summing of sub-holograms will result in a non-uniform amplitude hologram, and taking the phase of which will result in discarding the amplitude information of the summed hologram, leading to worse reconstructions quality. And also, SoH method suffers from the defocusing effect from one slice to another, causing additional noise.

$$\text{SoH: } \sum_{i=1}^n \arg \min_{\mathbf{H}_i} Loss(\mathbf{T}_i, \mathbf{R}_i) \quad (5.5)$$

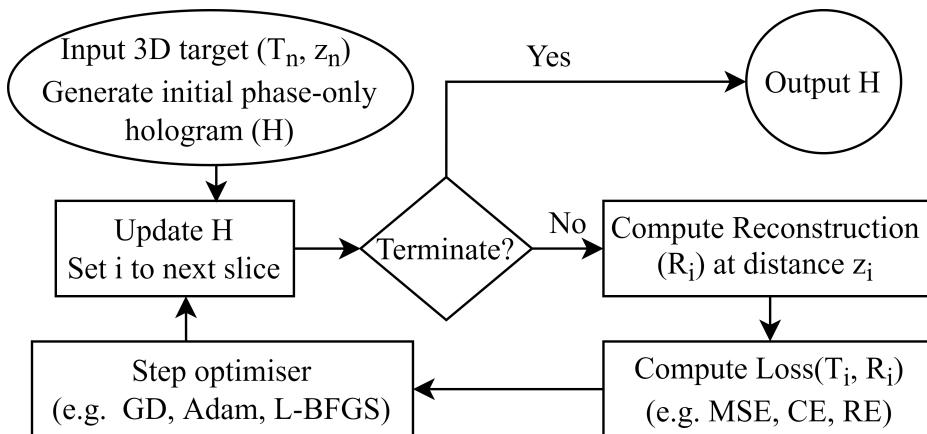


Fig. 5.2 Optimization of CGH with sequential slicing (SS) flowchart

This paper proposes a novel CGH optimization with sequential slicing (SS) technique, as shown in the flowchart in Fig. 5.2, that only computes the loss for a single slice at each iteration (between a reconstruction  $\mathbf{R}_i$  at a single distance  $z_i$  and its according target slice  $\mathbf{T}_i$ ), where  $i$  sweeps through the multi-layer 3D target sequentially when the algorithm iterates. When the final slice is reached ( $i = n$ ), it goes back to the first slice ( $i \leftarrow 1$ ). The proposed method only needs to carry out one Fourier Transform at each iteration, and the number of

iterations does not need to scale up with  $n$ . So it is expected to be much quicker than SoL and SoH techniques while producing proper resulting hologram.

## 5.3 Results

### 5.3.1 CGH for 4-slice targets

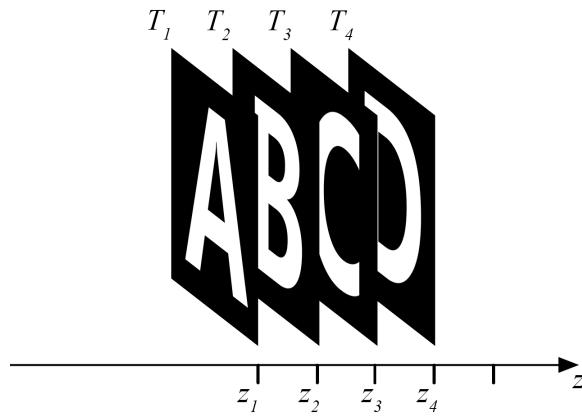


Fig. 5.3 Layout of the 4-slice target ( $z_1 = 1\text{ cm}$ ,  $z_2 = 2\text{ cm}$ ,  $z_3 = 3\text{ cm}$ ,  $z_4 = 4\text{ cm}$ )

The first example 3D target used is consisted of 4 slices made from alphabets ‘A’, ‘B’, ‘C’, ‘D’, each has  $512 \times 512$  pixels. The positions of the four slices range from  $1\text{ cm}$  to  $4\text{ cm}$  with  $1\text{ cm}$  gap between each other (i.e.  $z_i = i\text{ cm}$ ). The overall layout is shown in Fig. 5.3.

As there are two optimization algorithms (GD and L-BFGS), three techniques (SoL, SoH and SS) and three loss functions (MSE, CE, RE) in consideration, they can form a total of 18 combinations. In order to control the number of variables, all 18 combinations are set to start from the same initial random hologram and run for the same amount of 100 iterations for the optimization of each hologram, on the same laptop of model ASUS ROG Zephyrus M16, which has a CPU of model i7-11800H and a GPU of model RTX3060. For the L-BFGS algorithm, the gradient history of size 10 ( $m = 10$  in Algorithm 8) is used for all techniques and loss functions. And to ensure a sensible comparison, although three different loss functions are used for optimization of CGH, the metric used to assess the quality of multi-depth reconstructions of the final hologram is the normalized mean squared error (NMSE). As there are a total of 4 slices in this example, the final NMSE of each are computed separately and the total optimization run time is recorded. The final results are gathered in Fig. 5.4. As each slice has a different final error, their mean and standard

deviation (SD) are computed for investigations. Three columns are colour coded where green indicates better result while red indicates worse result.

Loss		Final NMSE ( $\times 10^{-6}$ )						Time (s)
		Slice 1	Slice 2	Slice 3	Slice 4	Mean	SD	
L-BFGS	MSE	1.62	1.23	1.43	1.19	1.37	0.17	1.04
	SoL	CE	1.76	1.29	1.52	1.24	1.45	0.21
	RE	1.62	1.23	1.43	1.19	1.37	0.17	1.02
	SoH	MSE	3.09	2.87	3.31	2.91	3.05	0.17
	RE	3.05	2.83	3.39	2.87	3.03	0.22	2.25
	SS	MSE	2.57	2.52	2.50	2.48	2.52	0.03
	CE	2.76	2.63	2.63	2.64	2.67	0.05	0.74
	RE	2.57	2.52	2.50	2.48	2.52	0.03	0.47
GD	MSE	1.94	1.65	1.96	1.66	1.80	0.14	0.86
	SoL	CE	2.20	1.94	2.30	1.97	2.10	0.15
	RE	2.23	1.98	2.35	2.01	2.14	0.15	0.85
	SoH	MSE	3.33	3.06	3.67	3.09	3.29	0.25
	RE	3.58	3.29	3.96	3.30	3.53	0.27	1.14
	SS	MSE	2.53	2.06	2.99	2.15	2.43	0.37
	CE	2.99	2.57	3.46	2.64	2.92	0.35	0.54
	RE	2.99	2.58	3.46	2.64	2.92	0.35	0.28

Fig. 5.4 Final NMSE and run time comparison across the three techniques

Comparing the mean of final NMSE and the run time of the proposed sequential slicing (SS) technique to those of the sum-of-loss (SoL) and sum-of-hologram (SoH) techniques in Fig. 5.4, it can be concluded that, for all combinations of optimizers and loss functions attempted, the proposed SS technique runs much quicker than the existing SoL and SoH techniques, while still providing a proper result, sitting between the SoL and SoH techniques. So the SS technique has not demonstrated absolute advantage to the SoL technique yet on this 4-slice example, therefore further investigation is done on a 30-slice example in Section 5.3.2. However, the SS technique is both quicker and has better reconstructions quality than the SoH technique, demonstrating an absolute advantage. Meanwhile, the combinations with CE as loss function are much slower and has not demonstrated advantage in average NMSE, demonstrating an absolute disadvantage, so the results with CE loss or SoH method will not be shown in the per-iteration plots.

For comparison among the sequential slicing (SS) techniques, the NMSE for each slice and their average and maximum difference values are plotted for GD and LBFGS algorithms with MSE and RE loss functions in Fig. 5.5, and sequential GS and DCGS [50] are also implemented and plotted for reference purpose. Apart from the L-BFGS algorithm, all the other algorithms are showing a staircase-like plot, where a decrease in error on one slice results in an increase in error on all other slices, and so the final NMSE of each slice distinguishes a lot from each other. The sequential GS algorithm suffers the most from the quality imbalance between each slice, and the sequential GD algorithm follows. The DCGS

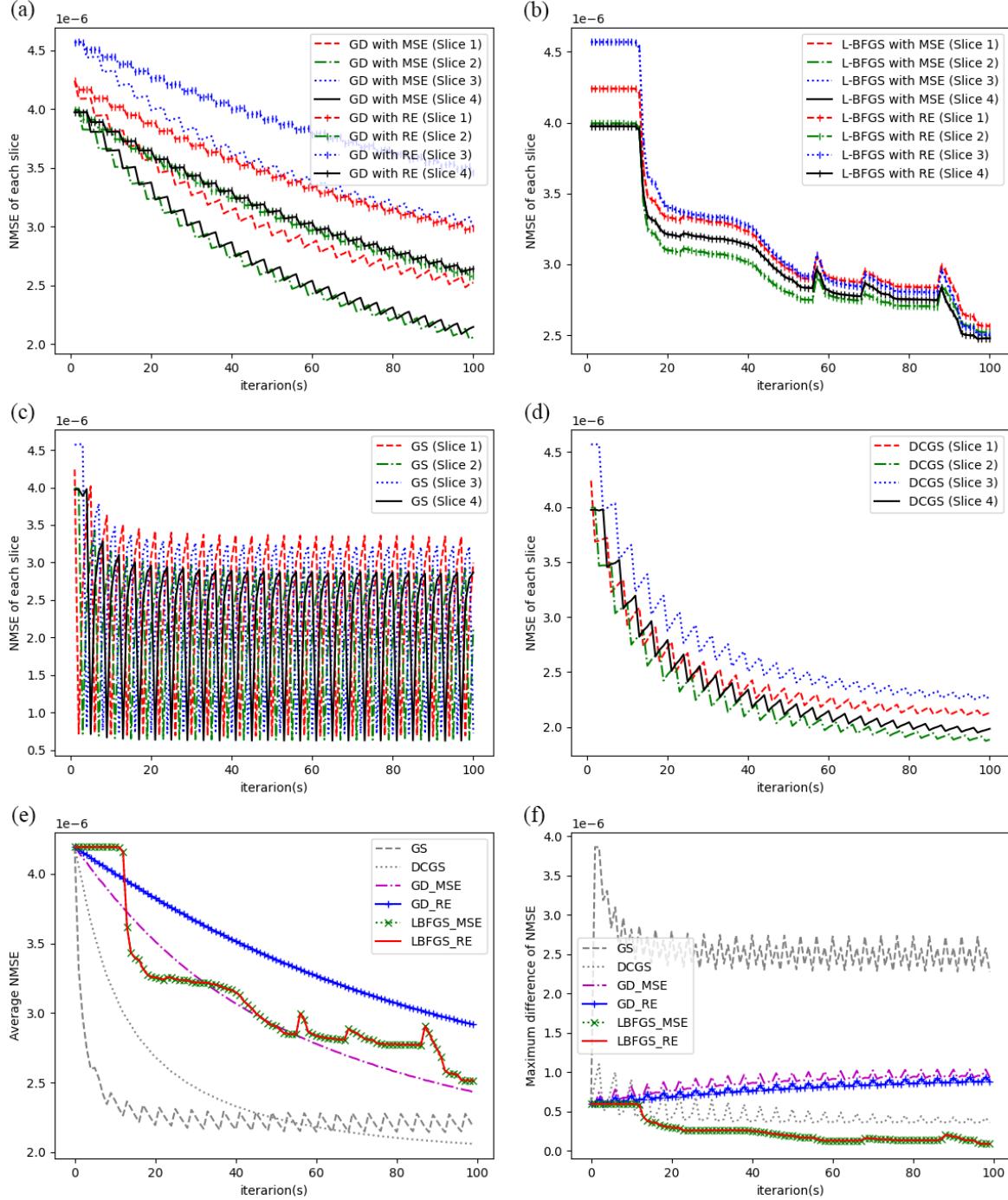


Fig. 5.5 Comparison among SS techniques for the 4-slice target using (a) GD algorithm, (b) L-BFGS algorithm, (c) GS algorithm, (d) DCGS algorithm. (e) Average NMSE, and (f) difference between the maximum and minimum NMSE across all slices.

algorithm benefits from its modification of the inclusion of a weighting factor consisting historical amplitude, therefore managed to converge. From the average NMSE plot (Fig. 5.5

(e)), the proposed sequential L-BFGS method does not appear to have the lowest average NMSE, but it has the lowest quality imbalance across the slices as shown in the maximum difference plot (Fig. 5.5 (f)). The L-BFGS algorithm mainly benefits from its inclusion of curvature information during optimization, so that the update of hologram  $\mathbf{H}$  at each iteration takes into account not only the loss for current slice, but also all historical iterations up to the set history size ( $m$  in Algorithm 8). So the NMSE of each slice stays close for L-BFGS and at each iteration, the NMSE of all slices behave in the same way, ensuring each slice to have the similar quality.

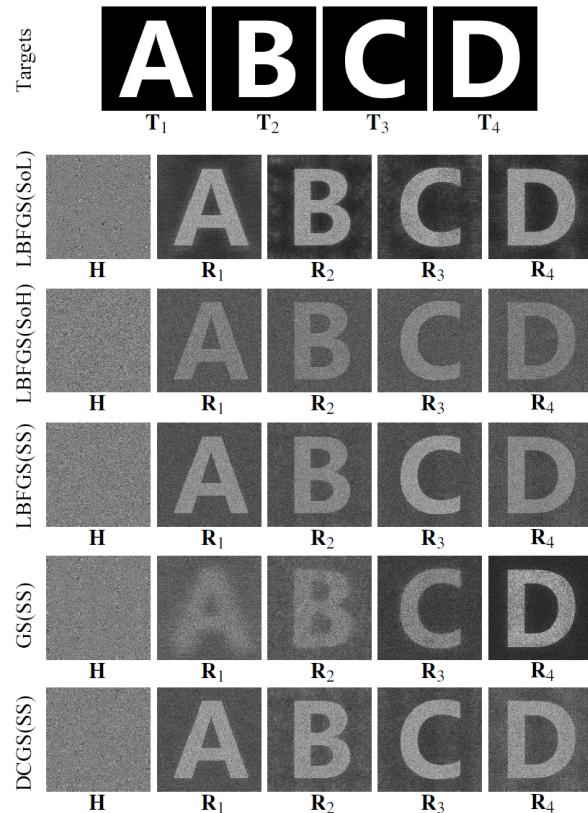


Fig. 5.6 Comparison of final holograms and reconstructions

The final holograms and their reconstructions for L-BFGS algorithm with SoL, SoH and SS techniques are shown in Fig. 5.6, with GS with SS technique and DCGS also shown as reference. The reconstructed images confirm the SS technique having a quality between SoH and SoL method (for the same amount of iterations), and has a much better quality imbalance than sequential GS, which has a very clear reconstruction at the fourth slice (letter 'D') because the iteration stopped at the fourth slice but much worse reconstructions at other slices. Admittedly, the proposed L-BFGS with SS method cannot surpass the GS-based DCGS algorithm yet.

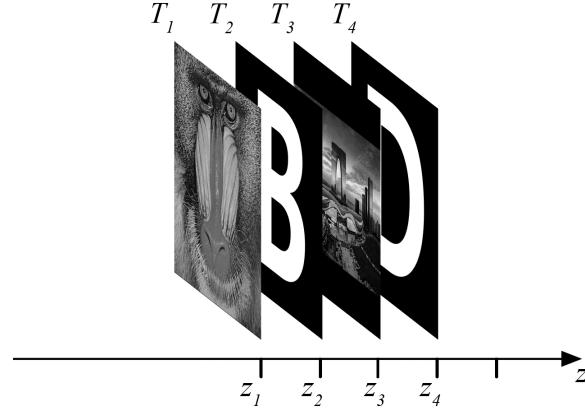


Fig. 5.7 Layout of the non-binary 4-slice target

To prove that the proposed method also works for non-binary targets, another example of a 4-slice 3D target was attempted, as shown in Fig. 5.7, where two of the slices are replaced by an image of the mandrill [16] and an image of the city scene [60] respectively.

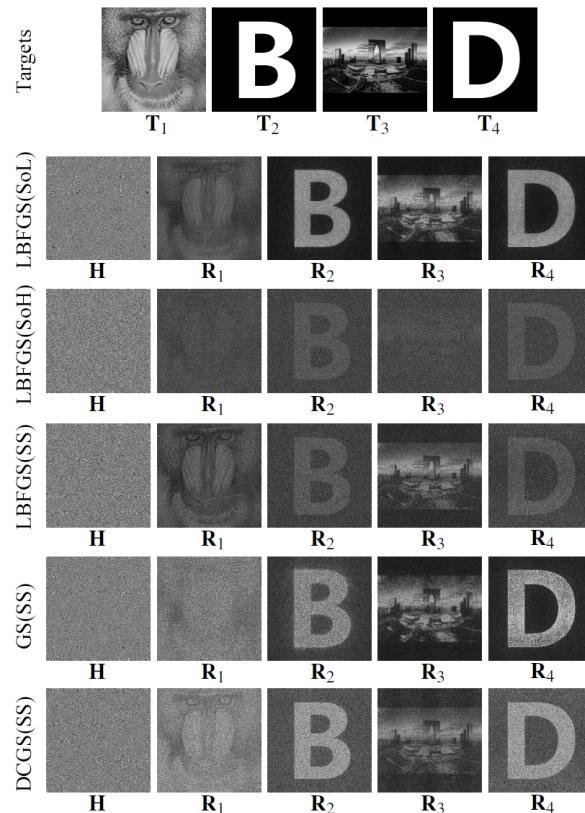


Fig. 5.8 Comparison of final holograms and reconstructions for non-binary target

As shown in the final holograms and reconstructions in Fig. 5.8, the proposed LBFGS with SS technique still managed to converge, with final reconstructions quality sitting between the SoL and SoH method, and also having a good quality balance across all slices.

### 5.3.2 CGH for a 30-Slice Target

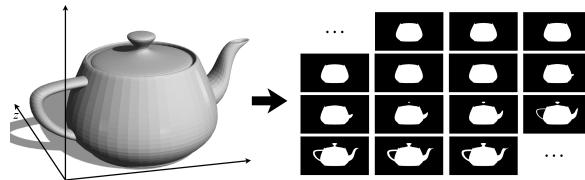


Fig. 5.9 30-slice target sliced from a 3D Teapot mesh

Another example of a 30-slice target sliced from a 3D teapot is attempted for speed analysis when the number of slices go higher. As shown in Fig. 5.9, the Utah teapot [73] is sliced into 30 planes, each of 720p resolution ( $1280 \times 720$ ). Each combination in Fig. 5.4 were run on a laptop of model ASUS ROG Zephyrus M16 with a CPU of model i7-11800H and a GPU of model RTX3060. The average and maximum difference of NMSE across all slices plotted against time (in Fig. 5.10 and Fig. 5.11 respectively) for all combinations except those with SoH technique or with CE as loss function as they have absolute disadvantage, for clearer plots.

Fig. 5.10 shows that, for both GD and L-BFGS optimization algorithms, the SS technique is faster than the SoL technique. Among the SS techniques, GD, GS and DCGS algorithms achieved less average NMSE than the proposed L-BFGS with SS method, but from the maximum difference of NMSE plot in Fig. 5.11, it can be shown that the L-BFGS algorithm has less quality imbalance than the GD algorithm and the GS algorithm, although admittedly it is not as good as the DCGS algorithm. Nevertheless, the proposed method has achieved an improvement of phase retrieval using optimization algorithms, in speed and quality imbalance suppression.

## 5.4 Conclusion

This paper has proposed the method of using L-BFGS optimization algorithm to generate phase-only hologram for multi-depth target, and discussed its suitability with the sequential slicing (SS) technique. The L-BFGS with SS method has demonstrated a good suppression on the quality imbalance across the multi-depth slices, benefiting from the nature of being a second order optimization algorithm, which implicitly records the historical gradients by other slices for the determination of the descent direction. For both GD and L-BFGS optimization algorithms, the SS technique ran faster and produced better reconstruction quality than the simple SoH technique, and it is quicker than the SoH technique when the

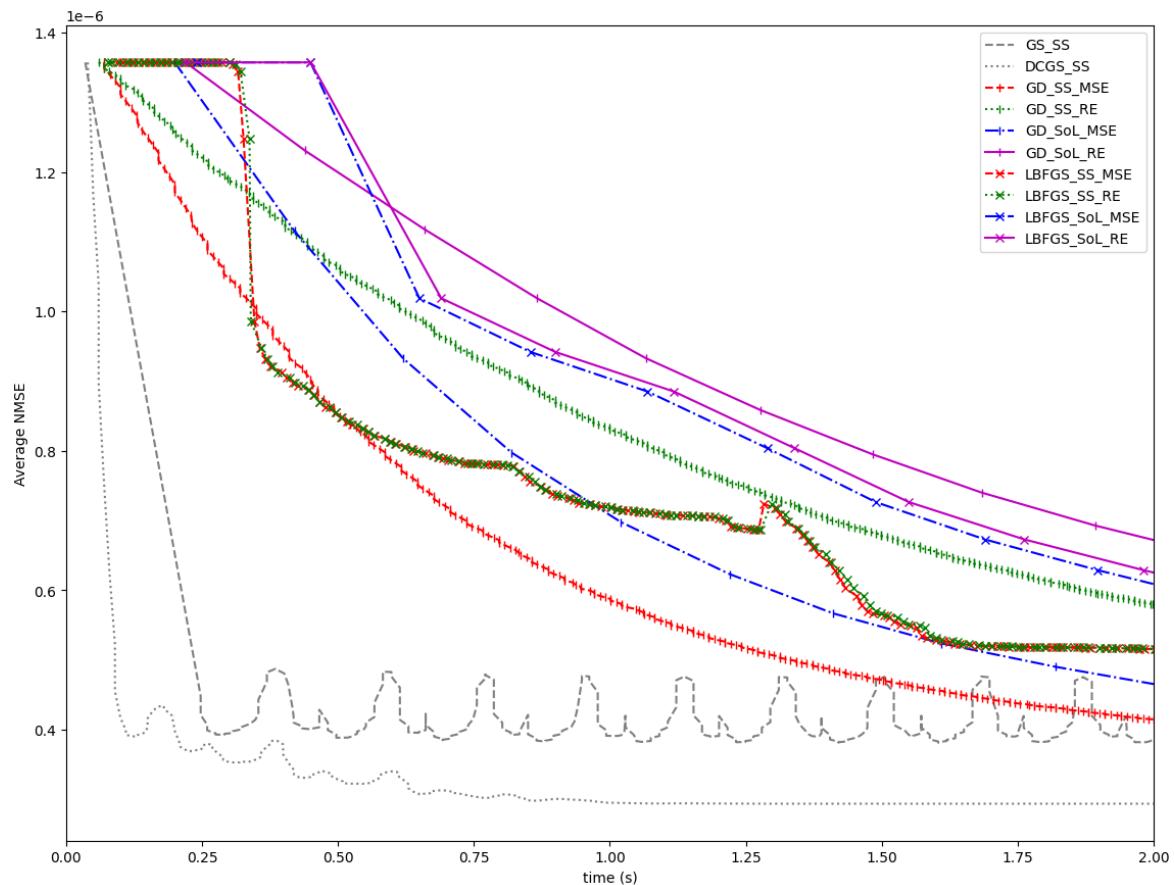


Fig. 5.10 Average NMSE v.s. time plot for the 30-slice target

number of depths get large. The proposed method has demonstrated great potential of time constrained optimization of multi-depth CGH.

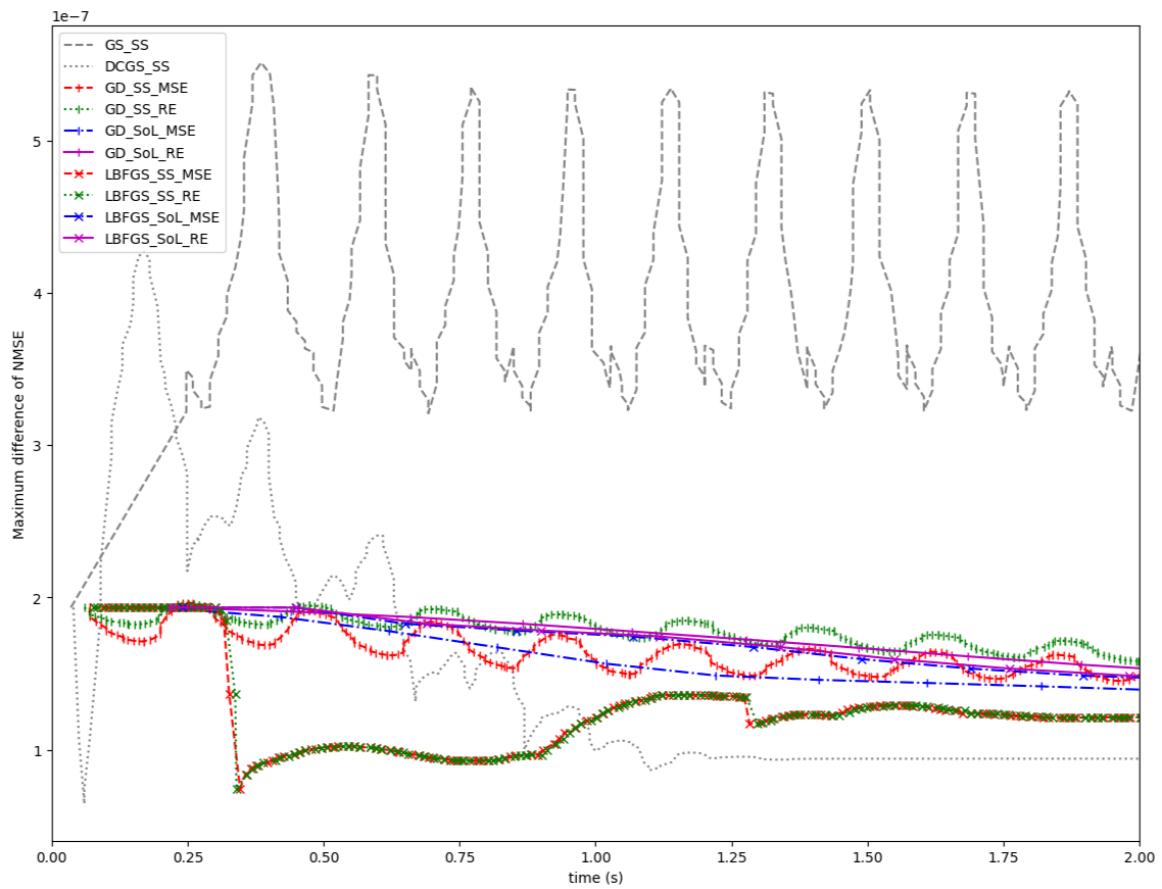


Fig. 5.11 Difference between the maximum and minimum NMSE across all slices v.s. time plot for the 30-slice target



# **Chapter 6**

## **Multi Frame Holograms Batched Optimization**

### **6.1 abstract**

Phase retrieval methods used in computer generated holograms such as Gerchberg-Saxton and gradient descent give results which are prone to noise and other defects. This work builds up on the idea of time-averaging multiple hologram frames, first introduced in methods like One-Step Phase-Retrieval and Adaptive One-Step Phase-Retrieval. The proposed technique called Multi-Frame Holograms Batched Optimization uses the L-BFGS optimization algorithm to simultaneously generate a batch of binary phase holograms which result in an average reconstructed image of improved fidelity and fast algorithmic convergence, both in the Fraunhofer and the Fresnel regimes. The results are compared to One-Step Phase-Retrieval and Adaptive One-Step Phase-Retrieval in simulation and experimentally, proving the superiority of the proposed approach. This technique can be easily extended to other spatial modulation methods.

### **6.2 Introduction**

Computer-generated hologram (CGH) enables producing three-dimensional (3D) multi-depth image reconstruction via modulating the wavefront of a coherent light source. Currently available spatial light modulators (SLM) can only modulate either phase or amplitude, so algorithms are needed to compute amplitude-only or phase-only holograms. The classic

phase-retrieval algorithms include direct binary search [22], simulated annealing [47] and Gerchberg-Saxton [19]. With the developments in modern numerical optimization methods and increase in computational power, phase retrieval with new numerical optimization methods has also been found in the literature such as: gradient descent [60, 61], its stochastic variations [63, 62, 3], and its derivative L-BFGS [2, 4]. However, all of these are single-frame hologram generation methods. In contrast, time multiplexing multi-frame holograms seeks to improve a time-averaged response by displaying different hologram sub-frames at a high refresh rate [74]. Such approach can exploit the finite response time of human vision, where human eyes average out the unwanted noise while the wanted signal remains. Similarly, such method could be used in holographic systems which require high precision without any hard restrictions on the projection refresh rate, such as holographic photo-lithography [75]. A few time-multiplexed multi-frame holograms generation methods have been explored in the literature, including the One-Step Phase-Retrieval (OSPR) algorithm [35] and the Adaptive One-Step Phase-Retrieval (AD-OSPR) algorithm [53]; however, both OSPR and AD-OSPR are still subject to defects in reconstruction quality.

This paper therefore extends on the previous research using the L-BFGS optimization algorithm for single-frame phase-only hologram generation [2, 4], and proposes a novel time-multiplexed multi-frame holograms generation method using L-BFGS optimization, called Multi-Frame Holograms Batched Optimization (MFHBO), to produce better reconstruction quality than the existing OSPR and AD-OSPR methods.

### 6.3 Method

This paper proposes a novel method of using numerical optimization algorithm L-BFGS [68] to generate multi-frame binary-phase holograms. The L-BFGS algorithm had previously been used for single-frame hologram optimization [2, 4]. To implement it onto multi-frame holograms generation, the argument to vary becomes the set of holograms with  $n$  sub-frames ( $\{\mathbf{H}_1, \mathbf{H}_2, \dots, \mathbf{H}_n\}$ ), each having a resolution of  $X \times Y$  pixels matching the resolution of the target image, and the objective function to minimise is therefore the difference between the average reconstruction amplitude ( $\mathbf{R}_{avg} = \frac{1}{n} \sum_{i=1}^n \mathbf{R}_i$ ) and the target image ( $\mathbf{T}$ ), which is denoted as  $Loss(\mathbf{T}, \mathbf{R}_{avg})$ , where  $n$  is the total number of frames,  $\mathbf{R}_i$ 's are reconstructions from individual hologram sub-frames  $\mathbf{H}_i$ 's for  $i \in [1, n]$ . To compute each  $\mathbf{R}_i$  from the corresponding  $\mathbf{H}_i$ , we start from the Fresnel diffraction formula given in Eq. (7.1) [39]

$$\mathbf{E}_{Fresnel\ region}(\alpha, \beta, z) = \mathcal{F} \left\{ \mathbf{A}(x, y) \cdot e^{j \frac{k}{2z} (x^2 + y^2)} \right\} \quad (6.1)$$

where  $\mathbf{E}$  is the reconstructed electric field, in complex form,  $\mathbf{A}$  is the hologram aperture, also in complex form, and  $\mathcal{F}$  denotes the Fourier Transform, implemented on computers using the Fast Fourier Transform (FFT) function. As eyes cannot perceive phase, the reconstruction amplitude is therefore the absolute value, giving  $\mathbf{R} = |\mathbf{E}|$ . And as we are generating holograms for phase-only SLM's,  $\mathbf{A}$  is then comprised of a uniform amplitude with phase  $\mathbf{H}$ , giving  $\mathbf{A} = e^{j\mathbf{H}}$ , where the exponential is taken element-wise.

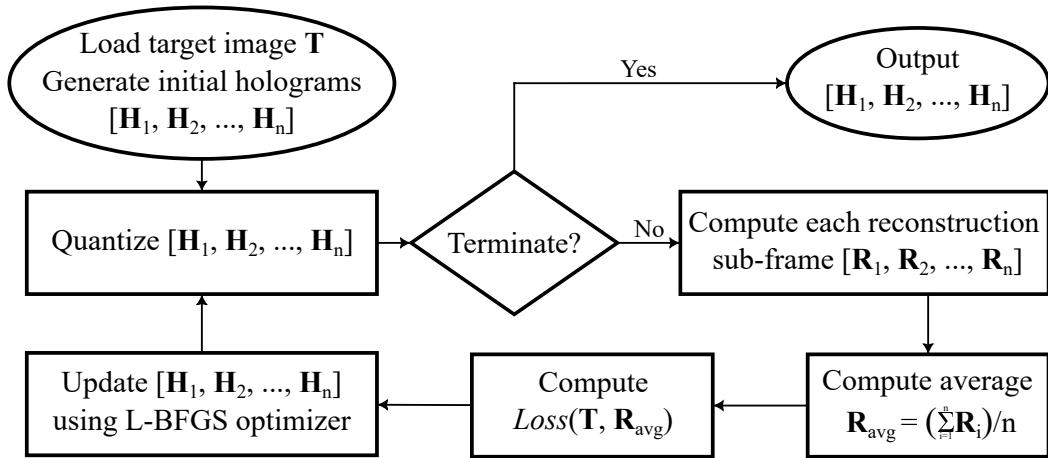


Fig. 6.1 MFHBO flowchart

To help explain the optimization process, a flow chart is drawn in Fig. 6.1. As shown in the flowchart, the target image  $\mathbf{T}$  is first loaded, with a set of  $n$  hologram sub-frames ( $\{\mathbf{H}_1, \mathbf{H}_2, \dots, \mathbf{H}_n\}$ ) generated randomly. Then at every iteration, each hologram sub-frame  $\mathbf{H}_i$  is quantized to the bit-depth constraint of the SLM, and propagated to the reconstruction plane  $\mathbf{R}_i$ , and the average of the amplitudes of all reconstructions  $\mathbf{R}_{avg}$  is computed and compared against the target image  $\mathbf{T}$  using a loss function  $Loss(\mathbf{T}, \mathbf{R}_{avg})$ , after which the search direction is computed using the L-BFGS optimizer and the hologram sub-frames are updated accordingly. Here the loss function selected is the relative entropy[71] given in Eq. (6.2).

$$Loss(\mathbf{T}, \mathbf{R}_{avg}) = - \sum_{x=1}^X \sum_{y=1}^Y \mathbf{T}_{(x,y)} \log \left( \frac{\mathbf{R}_{avg(x,y)}}{\mathbf{T}_{(x,y)}} \right) \quad (6.2)$$

Since fast SLM's available in the lab are binary-phase devices, the quantization step in the flowchart in Fig. 6.1 is carried out with bit-depth limit of 1, hence producing binary-phase holograms. However, the optimization algorithm does not converge with a straight binary quantization as integers are discrete, therefore a Sigmoid function [76] is used for a smoother and differentiable quantization, as defined in Eq. (6.3). The output of the Sigmoid function is then scaled by  $\pi$  so that the binary phase levels are 0 and  $\pi$ .

$$\text{sigmoid}(x) = \frac{1}{1 + e^{-x}} \quad (6.3)$$

And finally, when displaying the multi-frame holograms, each of the  $n$  frames generated are then rounded to binary phase values and displayed on the binary phase SLM sequentially. And when the first round finished, the second round starts with the first frame again (i.e. after frame  $n$ , the next frame displayed is frame 1), and such infinite loop doesn't stop until another set of holograms are uploaded.

## 6.4 Results

### 6.4.1 Simulation results

To test the proposed MFHBO method, a target image  $\mathbf{T}$  as shown in Fig. 6.2 was used. It was designed from the widely used mandrill image [16]. A rotational symmetry was introduced to match the rotational symmetric property of the far field projections from binary phase holograms. It was then zero padded to a resolution of  $1024px \times 1024px$  and subsequently interpolated to a resolution of  $1280px \times 1024px$  to match the resolution of the SLM in our lab. Note that the target image was zero padded to a square aspect ratio and then stretched to the non-square aspect ratio because more pixels in the horizontal axis only means higher sampling rate as part of the features of the FFT, the replay field is continuous and is not pixelated and the simulated reconstruction of  $1280px \times 1024px$  resolution is the sampled results, which will be illustrated visually in Fig. 6.4 later.

To further explain the optimization process described in Fig. 6.1, an example iteration with  $n = 24$  is shown in Fig. 6.2. At each iteration, every hologram is quantized and propagated to the reconstruction plane, forming  $\{\mathbf{R}_1, \mathbf{R}_2, \dots, \mathbf{R}_{24}\}$ . The average reconstruction amplitude  $\mathbf{R}_{avg}$  is then compared against the target image  $\mathbf{T}$ , using the loss function in Eq. (6.2). The holograms  $\{\mathbf{H}_1, \mathbf{H}_2, \dots, \mathbf{H}_{24}\}$  are then updated according to the search direction calculated

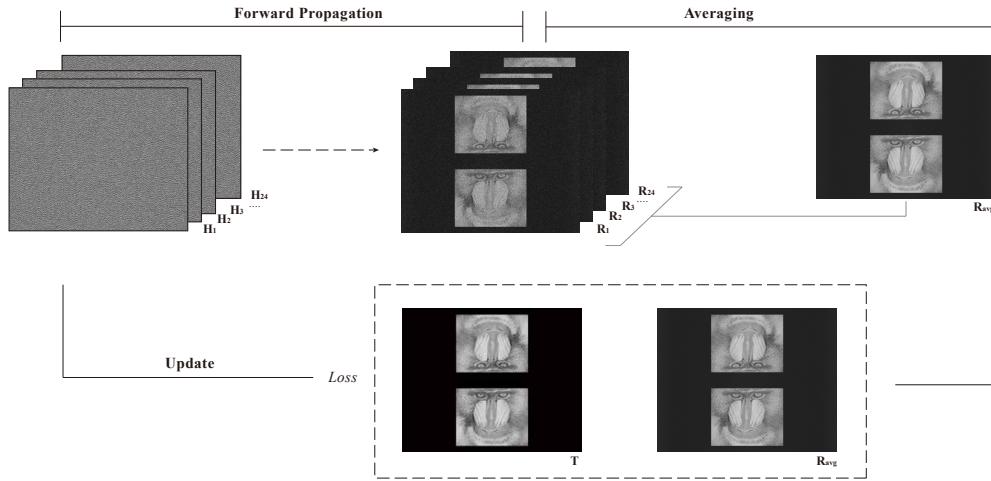


Fig. 6.2 An example iteration in the optimization process

using the L-BFGS optimizer. After setting the optimization to terminate when the number of iterations reach 1000, the same algorithm was run on the same target for different number of frames ( $n$ ), the normalised mean squared error (NMSE) and the peak signal-to-noise ratio (PSNR) between the average reconstructions  $\mathbf{R}_{avg}$  and the target image  $\mathbf{T}$  were calculated at every iteration and plotted in Fig. 6.3a and Fig. 6.3b respectively.

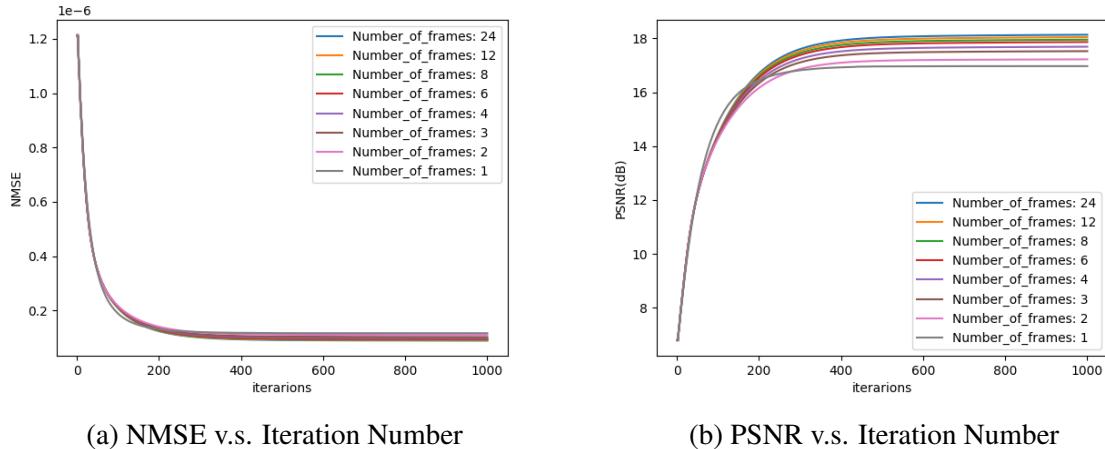


Fig. 6.3 Convergence of optimization

The plots in Fig. 6.3 show that the proposed MFHBO method has achieved good convergence within 400 iterations, for the various number of frame settings  $n$  in  $\{1, 2, 3, 4, 6, 8, 12, 24\}$ . The final NMSE values in Fig. 6.3a are difficult to distinguish in the plot, therefore it will be further compared in the bar chart in Fig. 6.4. The number of frames are chosen to be integer

factors of 24, which is determined by our experimental setup, further explained in the next subsection.

<b>Number of frames</b>	<b>200 iterations</b>	<b>400 iterations</b>	<b>600 iterations</b>	<b>800 iterations</b>	<b>1000 iterations</b>
1	1.79	3.59	5.31	7.00	8.72
2	2.82	5.59	8.34	11.11	13.88
3	3.84	7.67	11.45	15.21	19.00
4	4.93	9.83	14.70	19.58	24.47
6	6.95	13.87	20.76	27.58	34.50
8	8.87	17.67	26.54	35.60	44.56
12	12.95	25.79	38.63	51.47	64.30
24	51.81	101.43	151.09	201.08	251.15

Table 6.1 MFHBO runtime (s)

The programme runtime of the proposed MFHBO method has been measured on a laptop computer of model ASUS ROG Zephyrus M16 (GU603H) with a CPU of model i7-11800H and a GPU of model NVIDIA RTX3060 and the results for different combinations of number of frames and number of iterations are listed in Table 6.1. It can be concluded that the application of the proposed method is for pre-computed high-quality holograms, instead of real-time holographic projections.

#### 6.4.2 Optical Experiment results

The holographic projection system used in this experiment is the same as the one used in Chapter 3, with the optical setup shown in Fig. 3.2. Since the SLM has a refresh rate of 1440Hz and modern computer monitors have refresh rate of at least 60Hz, the maximum number of frames was chosen to be  $1440/60 = 24$ , so that each set of 24 frames will take a total of 1/60 seconds to display, therefore giving an equivalent refresh rate of 60Hz. Then the integer factors of 24 were chosen so that the equivalent refresh rate becomes integer multiples of 60Hz. The number of frames starts from 1 to help illustrate how the increase in number of frames positively affect the reconstruction quality.

The results in Fig. 6.4 further compares the final results for different number of frames. The histogram in Fig. 6.4 shows that, as the number of frames increases, the NMSE between the average reconstructions  $\mathbf{R}_{avg}$  and the target image  $\mathbf{T}$  decreases and the structural similarity

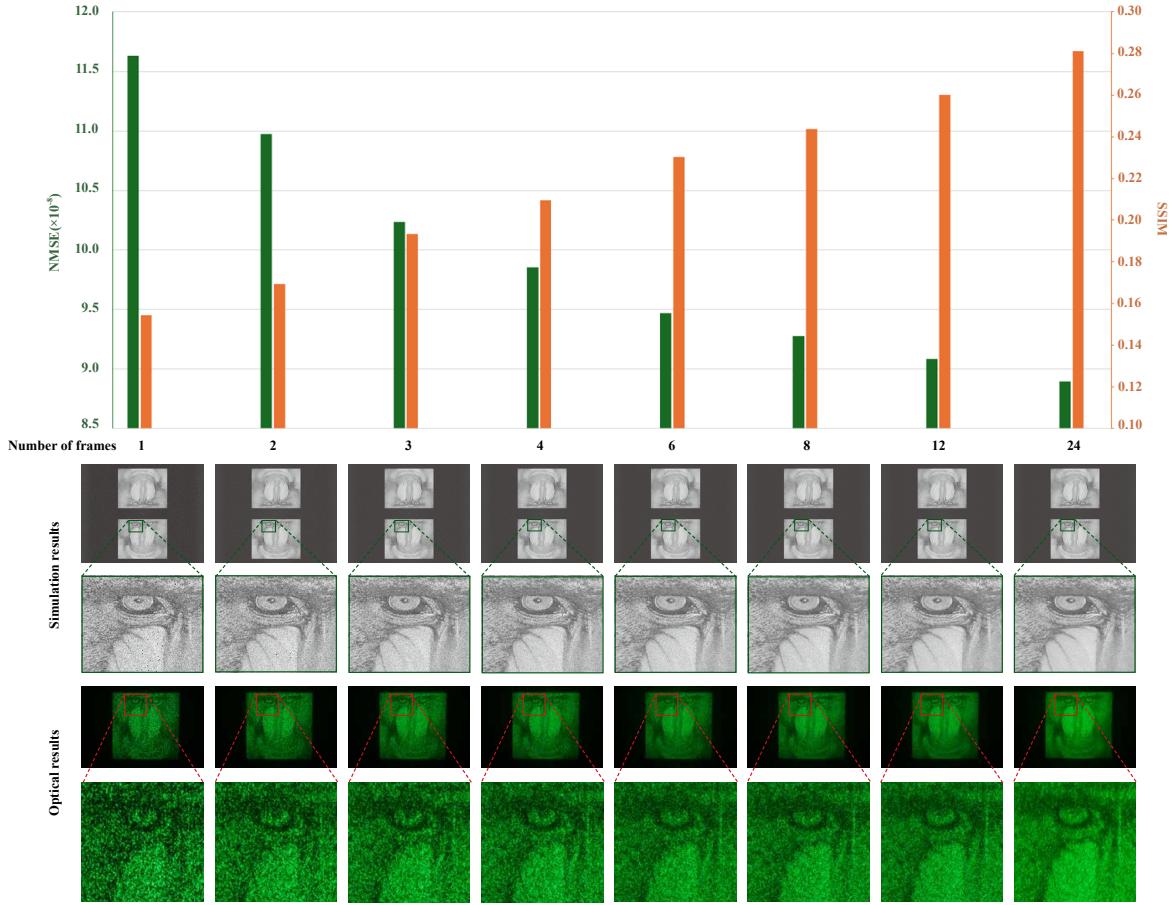


Fig. 6.4 Simulation and optical reconstruction results for different number of frames

index (SSIM)[46] increases, showing a trend of better reconstruction quality with higher number of frames. Such trend is expected as more frames provide higher information capacity, which agrees with the previous research where holograms with higher bit depth were found to achieve better reconstruction quality [8]. The trend is also shown visually via the simulation results and their detail enlargements. The corresponding multi-frame holograms are then loaded onto the SLM, and the reconstructed field is captured using a camera of model Cannon EOS 1000D. Only the bottom halves of the reconstructed field were captured as the symmetrical conjugates were unwanted feature of far field projections from binary-phase SLM's. The raw data including multi-frame binary-phase holograms, simulated reconstructions and optical results captured are accessible in the database [77].

Then another target image was tested, which is the holography ambigram shown as shown in Fig. 6.5.<sup>1</sup> The term ambigram is used to refer to (often typographical) designs that are

<sup>1</sup>Adapted, with colours reversed, from *holography* - Benjamin Wetherfield, 2022.

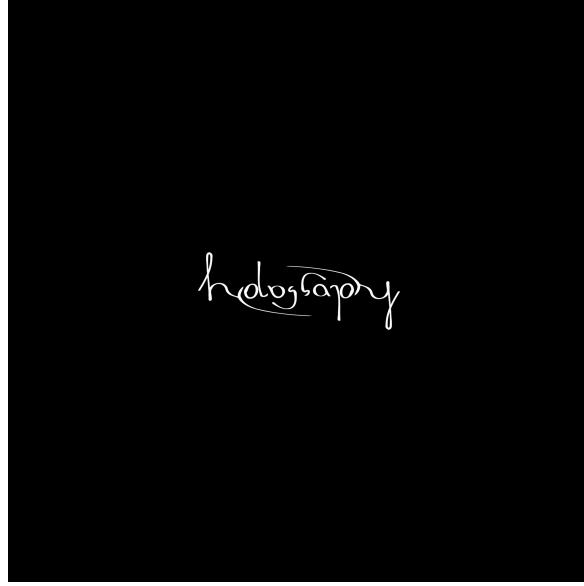


Fig. 6.5 Sample target image - ‘holography’ ambigram

invariant under a reflection, rotation or other symmetry. The ‘holography’ design contains 180-degree rotational symmetry, which makes it especially well suited to binary Fourier-holographic projection, where this symmetry is unavoidable. Multi-frame holograms were then generated using the proposed MFHBO method and the existing OSPR and AD-OSPR methods, for the same number of frames  $n = 24$ . And the optical results are shown in Fig. 6.6.

As shown in Fig. 6.6, for the Mandrill target image, it can be seen that the proposed MFHBO method achieved a much better optical reconstruction quality than the existing OSPR and AD-OSPR methods, with clearer details and better contrasts; for the ‘holography’ ambigram target image, the proposed MFHBO method is shown to have a much lower background noise around the centre, than the existing OSPR and AD-OSPR methods. The intended black regions are represented much more cleanly, with an elimination of speckle-like artefacts in the zero-valued space around the lettering, and an overall increase in discernible contrast.

<b>Image</b>	<b>Metric</b>	<b>MFHO</b>	<b>OSPR</b>	<b>AD-OSPR</b>
Mandrill	NMSE ( $\times 10^{-4}$ )	0.84	1.00	1.12
	SSIM	0.124	0.076	0.078
Ambigram	NMSE ( $\times 10^{-5}$ )	2.29	3.31	3.23
	SSIM	0.795	0.826	0.827

Table 6.2 Quantitative analysis of the optical results in Fig. 6.6

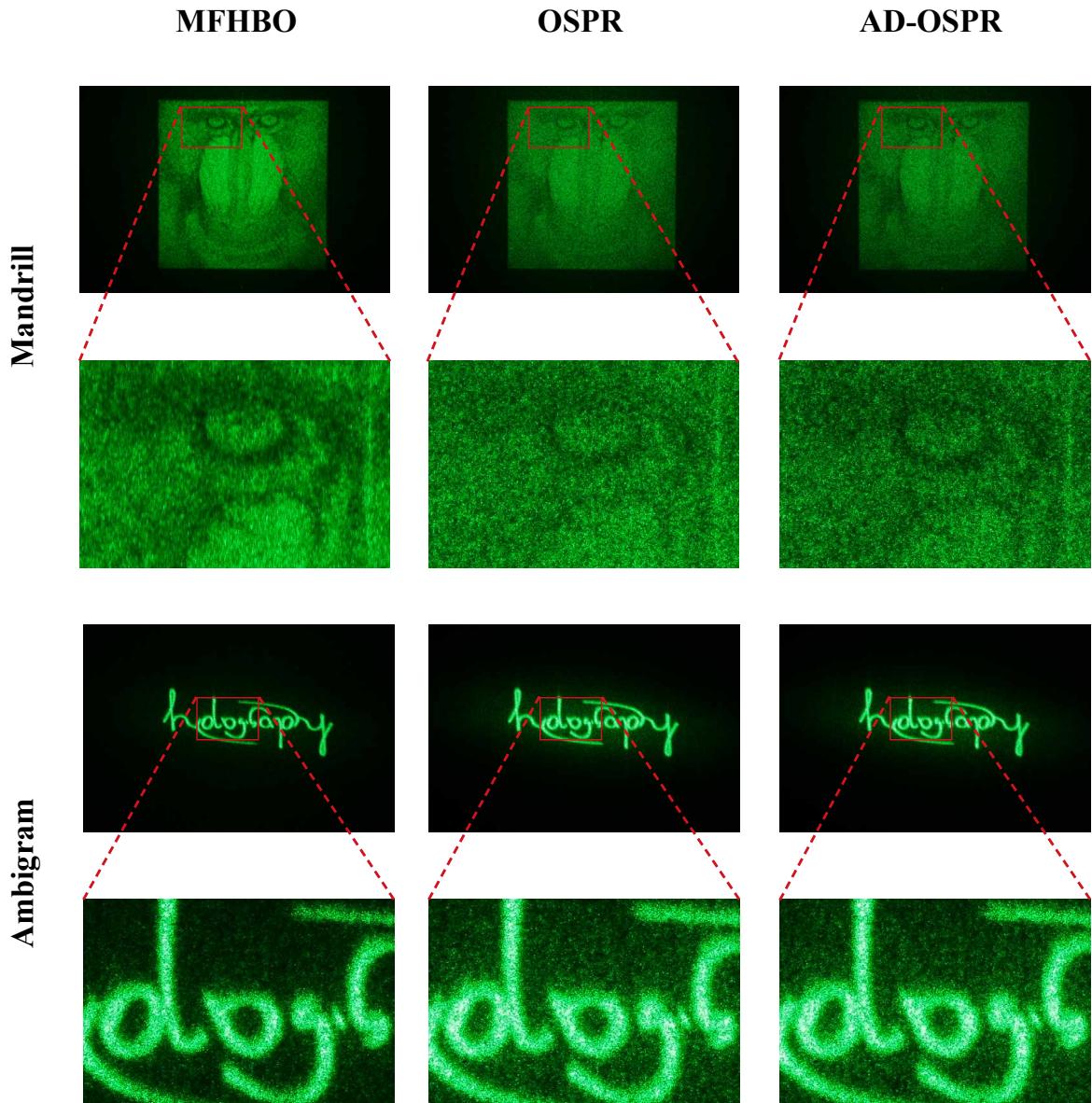


Fig. 6.6 Optical results comparison of the proposed MFHBO method against the existing OSPR and AD-OSPR methods

A quantitative analysis was then conducted on the optical results in Fig. 6.6, the NMSE and SSIM between the captured reconstructions and their corresponding targets are computed and listed in Table 6.2. The NMSE results of the proposed MFHBO method are lower than those of the existing OSPR and AD-OSPR methods, with a 25% reduction on average among both target images. On the other hand, the SSIM results have shown a 62% increase using MFHBO than OSPR and AD-OSPR for the mandrill target image, but a slight decrease of

3.7% for the ‘holography’ ambigram target image, which is negligible as it is less than 5% and the SSIM metric is not originally designed for binary-valued non-grayscale images.

### 3D Holography

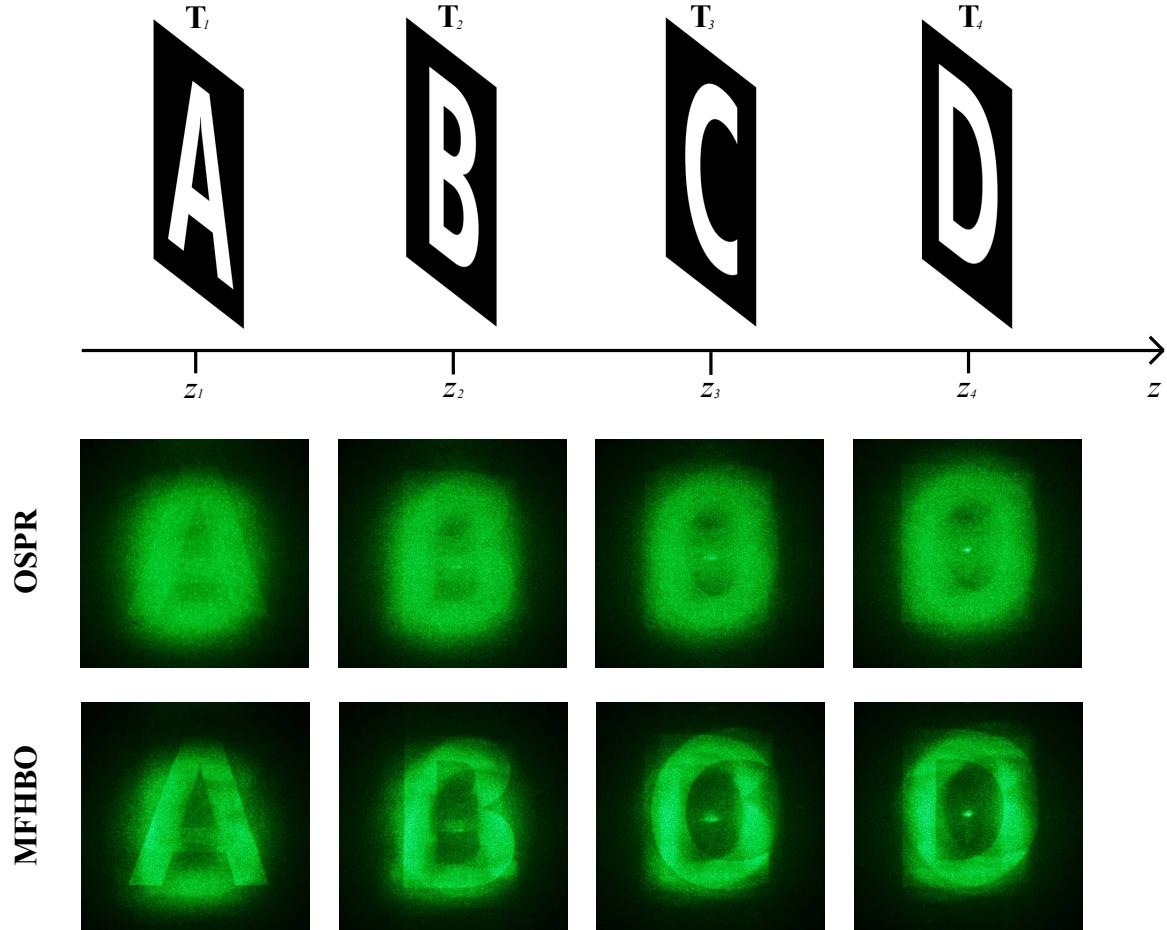


Fig. 6.7 4-slice target and according reconstruction results

The proposed MFHBO method was extended to multi-slice targets, by computing the loss between all 4 slices of reconstructions and target images (the Sum-of-Loss method in [4]). An example 4-slice target made from alphabets ‘A, B, C, D’ is shown in Fig. 6.7. The  $z$  values, corresponded to the  $z$  variable in Eq. (7.1), were chosen to be 1.1, 1.9, 3.5, 7.7 for the 4 slices respectively (as there’s no correlation between each slice, larger separation was chosen for fewer cross-talks across different planes). It can be seen that the proposed MFHBO method has produced sharper edges in reconstructions than the existing OSPR method. (The AD-OSPR method was not attempted here as its application to multi-slice targets was not defined).

<b>Method</b>	<b>Metric</b>	<b>Slice 1</b>	<b>Slice 2</b>	<b>Slice 3</b>	<b>Slice 4</b>	<b>Average</b>
<b>OSPR</b>	NMSE( $\times 10^{-4}$ )	4.980	4.484	5.644	4.846	4.988
	SSIM	0.072	0.061	0.048	0.060	0.060
<b>MFHO</b>	NMSE( $\times 10^{-4}$ )	4.484	4.230	4.990	4.289	4.498
	SSIM	0.063	0.067	0.058	0.072	0.065

Table 6.3 Quantitative analysis of the optical results in Fig. 6.7

Then a quantitative analysis was carried out, with NMSE and SSIM values measured and shown in Table 6.3. The proposed MFHBO method has shown a 10% reduction in NMSE and a 8% improvement in SSIM on average than the existing OSPR method, demonstrating the effectiveness of the proposed method.

Lastly, a set of real-life scene was captured in the lab using near, middle and far focus, as shown in  $\mathbf{T}_1, \mathbf{T}_2, \mathbf{T}_3$  in Fig. 6.8 respectively. The  $z$  values were set to 1.1, 1.2, 1.3 for hologram generation, and the reconstruction results of the existing OSPR and the proposed MFHBO methods are compared in Fig. 6.8. The proposed MFHBO method is shown to have achieved much better reconstruction quality than the existing OSPR method.

<b>Method</b>	<b>Metric</b>	<b>Slice 1</b>	<b>Slice 2</b>	<b>Slice 3</b>	<b>Average</b>
<b>OSPR</b>	NMSE( $\times 10^{-6}$ )	3.70	3.69	3.47	3.62
	SSIM	0.37	0.28	0.34	0.33
<b>MFHO</b>	NMSE( $\times 10^{-6}$ )	3.20	2.78	3.06	3.01
	SSIM	0.42	0.32	0.32	0.35

Table 6.4 Quantitative analysis of the optical results in Fig. 6.8

A quantitative analysis was conducted again, with NMSE and SSIM values measured and listed in Table 6.4. The proposed MFHBO method has shown a 17% reduction in NMSE and a 7% improvement in SSIM on average than the existing OSPR method, proving the effectiveness of the proposed method.

## 6.5 Conclusion

This paper proposed the MFHBO method to generate multi-frame binary-phase holograms to be displayed on high refresh rate binary-phase SLM. The proposed MFHBO method was

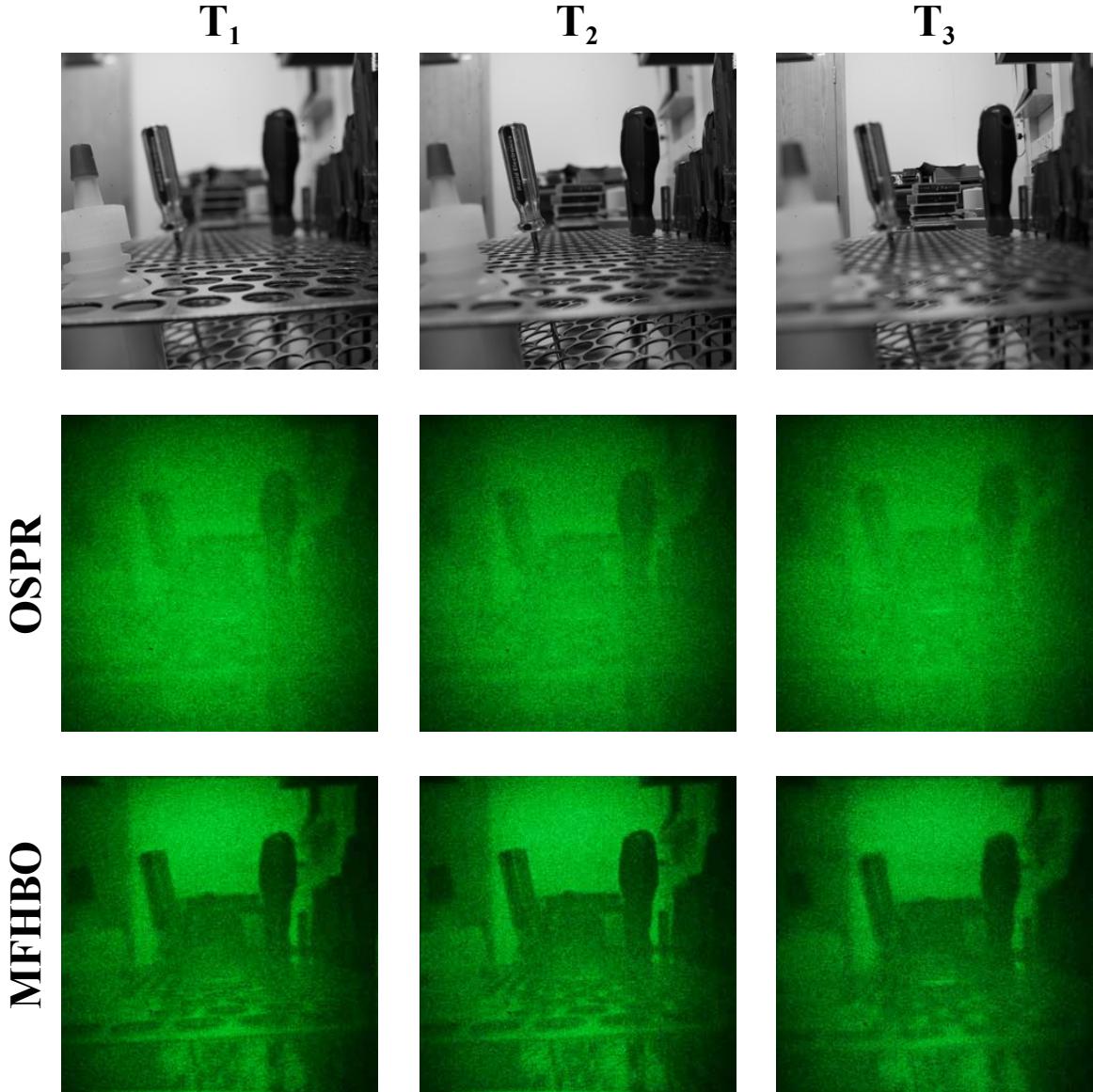


Fig. 6.8 Real-life captured image as target field and their reconstruction results

shown to achieve much better reconstruction quality and higher contrast than the existing multi-frame binary-phase holograms generation methods OSPR [35] and AD-OSPR [53] on the holographic projector with binary-phase SLM, for all the single-slice far-field targets and the multi-slice near-field targets tested. Although the propose MFHBO method is slower than the existing OSPR and AD-OSPR methods, its much better reconstruction quality makes it suitable for pre-computed high-quality hologram applications. Its strong advantage for high contrast target such as the ‘holography’ ambigram, with much suppressed speckle noise in the background, makes it well-suited for photo-lithography applications. The proposed

method can also be adapted for multi-level SLM's by simply removing the quantization step (in Fig. 6.1). This could be the case for applications such as photo-lithography, where the time response of the system is much longer than it is for human vision, and the high refresh rates of the SLM are not necessary.



# **Chapter 7**

## **Information capacity of phase-only computer-generated holograms for holographic displays**

### **7.1 abstract**

Despite many years of development in computer-generated holography, perfect phase-only holograms for most target images are still yet possible to compute. All computational phase retrieval algorithms end up with some error between the target image and the reconstruction of the computer-generated hologram (CGH), except for specific targets. This research focuses on the fundamental limits of phase-only CGH quantized to limited bit-depth levels, from the information theory point of view, revealing the information capacity of CGH and its effect on reconstruction quality, with an attempt to quantify how hard a target image is for phase-only hologram computation.

### **7.2 Introduction**

Holography is a technology that can fully reconstruct the wavefront of three-dimensional (3D) objects. Computer-generated holography (CGH) is a technique for converting a 3D object scene into a two-dimensional (2D), complex-valued hologram [78], without the need for the 3D object to physically exist. However, despite many advancements in liquid crystals and micro mirrors technologies, complex-valued spatial light modulators (SLM) are still

not yet available, and the currently available SLM can only achieve either amplitude or phase modulation, among which the phase modulation is usually preferred in holographic projections for its lower zero order and higher energy efficiency due to lower blockage of light. There are many algorithms available to compute good quality phase-only holograms, such as direct binary search [22], simulated annealing [47], Gerchberg-Saxton [19] and other optimization based methods [60–63, 3, 4]; however, none of them can guarantee to compute a perfect phase hologram for a 3D or even a 2D image, where they always end up with some error between the reconstruction of the hologram and the target scene, especially when the phase holograms are quantized to be able to display on SLM with limited bit depth. An intuition therefore arose that the bit depth of the phase hologram is limiting the reconstruction quality, and that the target scene's entropy seems to denote how difficult it is for phase hologram generation. It is obvious that the entropy of the target can certainly never exceed the bit depth limit of its corresponded perfect hologram, otherwise a lossless compressor breaking the Shannon's information theory [79] would be invented; however, such statement is not quite useful in practice as perfect holograms are generally not even possible to compute.

A literature review had found no related work on the information entropy of computer-generated holograms, with the closest match being the research on hologram compression using optical method to achieve lossy compression [80], which cannot be integrated in CGH processes. Therefore, this paper aims to investigate how much information content can a bit-depth-limited phase hologram contain, taking from an information theory point of view. Previously, work had been done to investigate the effect of Bit-depth in Stochastic Gradient Descent (SGD) performance for phase-only computer-generated holography displays [3]. This paper extends on previous research onto the Gerchberg-Saxton (GS) [19] algorithms for hologram generation, and investigates the correlation between the quality of the reconstructed image from the hologram and the information entropy of the target image, with an aim to reduce the hologram's entropy during the CGH process, for smaller size holograms.

## 7.3 Methods

### 7.3.1 Computation of Diffraction

To compute the spatial diffraction of light, the Fresnel diffraction formula is used, as shown in Equation (7.1) [39].

$$E_{Fresnel}(\alpha, \beta, z) = \mathcal{F} \left\{ A(x, y) e^{j \frac{k}{2z} (x^2 + y^2)} \right\} \quad (7.1)$$

where the hologram aperture is denoted by  $A(x, y)$  and the diffracted field is denoted by  $E(\alpha, \beta, z)$ , with  $z$  being the distance from the hologram aperture  $A(x, y)$ , and  $j$  is  $\sqrt{-1}$ ,  $k$  is the wave number ( $k = \frac{2\pi}{\lambda}$ , where  $\lambda$  is the wavelength), and  $\mathcal{F}$  denotes the Fourier Transform. The amplitude of  $E(\alpha, \beta, z)$ , which is  $|E(\alpha, \beta, z)|$ , is referred to as the reconstruction of the hologram throughout this paper.

When  $z \gg \frac{k(x^2+y^2)_{max}}{2}$  (i.e. in the far field), Equation (7.1) can be further approximated to give Equation (7.2), which is the Fraunhofer diffraction formula [39].

$$E_{Fraunhofer}(\alpha, \beta) = \mathcal{F} \{ A(x, y) \} \quad (7.2)$$

To implement the formula on computers, both  $A(x, y)$  and  $E(\alpha, \beta)$  are discretized, just like digital images. The Fourier Transform  $\mathcal{F}$  is discretized using the Fast Fourier Transform (FFT); similarly the inverse Fourier Transform  $\mathcal{F}^{-1}$  is discretized using the Inverse Fast Fourier Transform (IFFT) [81].

### 7.3.2 Computer-Generated Hologram (CGH) Algorithm

To compute phase-only holograms, the Gerchberg-Saxton (GS) [19] algorithm, which is the classic and robust phase retrieval algorithm, is selected. Taking the Fraunhofer propagation in Equation (7.2) as an example, the operation flowchart of GS algorithm is illustrated in Fig. 7.1. GS algorithm functions that it iteratively determines the phase profile ( $\angle A$ ) of the hologram ( $A$ ) required to reconstruct a target image ( $T$ ); it loops between the hologram ( $A$ ) and the diffracted plane ( $E$ ), and applying constraints to each plane accordingly during each iteration [19], with the total number of iterations denoted by  $N$ .

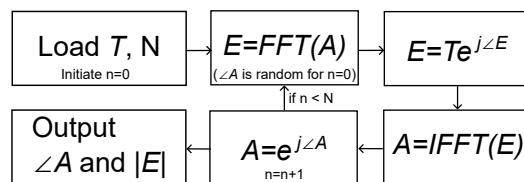


Fig. 7.1 Gerchberg-Saxton (GS) [19] algorithm flowchart

Then, as illustrated in Fig. 7.2, a quantization function ( $Q$ ) can be defined by finding the closest point from one of the  $2^d$  quantization levels, given the phase ( $\angle A$ ) and the quantization bit depth  $d$  as input.

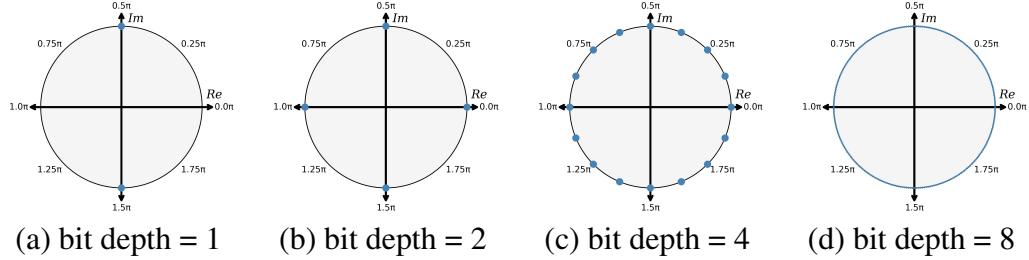


Fig. 7.2 Quantization of phase holograms

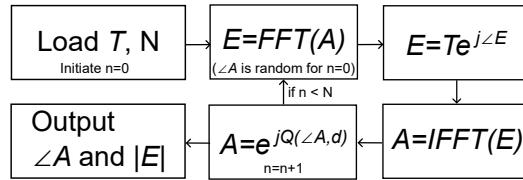


Fig. 7.3 Quantized Gerchberg-Saxton (GS) algorithm flowchart

To compute phase holograms quantized to certain bit depth ( $d$ ), the GS algorithm is modified to include an additional quantization operation ( $Q$ ) when applying the 'phase-only' constraint as shown in Fig. 7.3. Such method is better than applying the quantization at the end of the loop, as it includes the quantization constraint throughout the iterations, instead of introducing significant quantization noise in the end.

### 7.3.3 Measurement of Information

#### Shannon entropy

To quantify the information content, the classical one-dimensional (1D) Shannon entropy [79] with equation shown in Equation (7.3) is selected.

$$H(X) = - \sum_{x \in X} p(x) \log_2 p(x) \quad (7.3)$$

Although the Shannon entropy was designed for 1D data, it can also be implemented to two-dimensional (2D) data by ignoring the 2D spatial correlations and summing  $p(x) \log_2 p(x)$  over the histogram of the 2D data. As only discrete data can have a meaningful Shannon

entropy, the entropy can only be calculated for quantized holograms and target images, which are usually quantized to less than 8 bit depth.

### Delentropy

To account for 2D spatial correlation, the delentropy [82] is also used. Delentropy is an extension of the 1D Shannon entropy that it first computes the gradient (del) vector field image, whose entropy is then named as the delentropy, so that the spatial image structure and pixel co-occurrence can be captured [82].

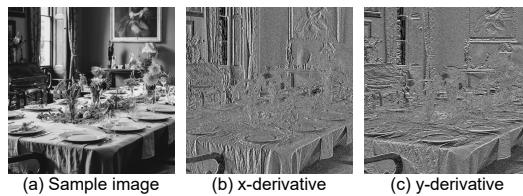


Fig. 7.4 Del operation on a sample image

As an example, the sample image in Fig. 7.4 (a) is the file ‘0500.png’ under the ‘DIV2K\_train\_HR’ folder sourced from the DIV2K dataset [83]. The sample image is calculated to have a Shannon entropy of 7.502 bits/pixel. By taking the  $x$ -derivative ( $f_x$ ) and  $y$ -derivative ( $f_y$ ) as shown in Fig. 7.4 (b) and (c), and using the Papoulis Generalized Sampling (PGS) [84] theory, the delentropy is calculated using Equation (7.4)[82] to be 5.867 bits/pixel.

$$H_{PGS}(f_x, f_y) \leq \frac{H(f_x) + H(f_y)}{2} \quad (7.4)$$

## 7.4 Results

### 7.4.1 Targets at far field (Fraunhofer region)

The target images used in this paper are sourced from the DIV2K dataset [83], among which, all of the 800 images in the ‘DIV2K\_train\_HR’ folder are selected. The quantized GS algorithm in Fig. 7.3 was run on each of the 800 target images set at far field (using the Fraunhofer diffraction formula in Equation (7.2)), for hologram bit depth set to integers ranging from 1 to 8, with the total number of iterations ( $N$ ) set to 100.

An example of quantized phase hologram generation for the sample image in Fig. 7.4 (a) is shown in Fig. 7.5, which demonstrates qualitatively how the reconstruction quality improves

## 10Information capacity of phase-only computer-generated holograms for holographic displays

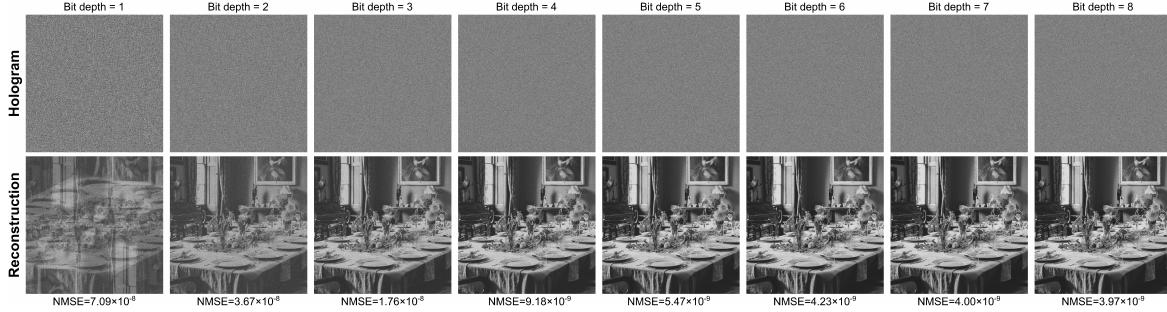


Fig. 7.5 Hologram generated at certain bit depths and their according reconstructions at far field

with the increase in the bit depth of the hologram, and also quantitatively, the normalized mean squared error (NMSE) between the reconstruction and target image has shown a decreasing trend as the bit depth of hologram increases. The rotational symmetry in the reconstruction of the hologram with bit depth 1 can be explained by the conjugate properties of Fourier Transforms, because the binary phase holograms have pixel values whose complex conjugates are the same as themselves.

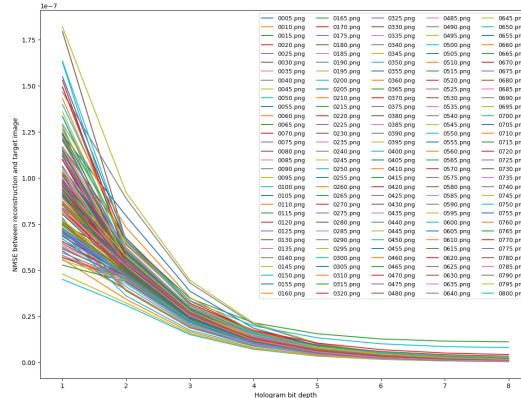


Fig. 7.6 NMSE v.s. Hologram bit depth for target images set at far field

Fig. 7.6 plots the NMSE between the reconstruction of the hologram and the target image against the hologram bit depth for every 5 images (as there are a total of 800 images, only every 5 images are shown in the plot to avoid overcrowding, where the full data can be accessed from the published research data [85]). It can be observed that, for each target image, the NMSE between the resulting reconstructions and their according target images decreases as the hologram bit depth increases. It infers that holograms with higher bit depth carries more sufficient information in order to better reconstruct the target images.

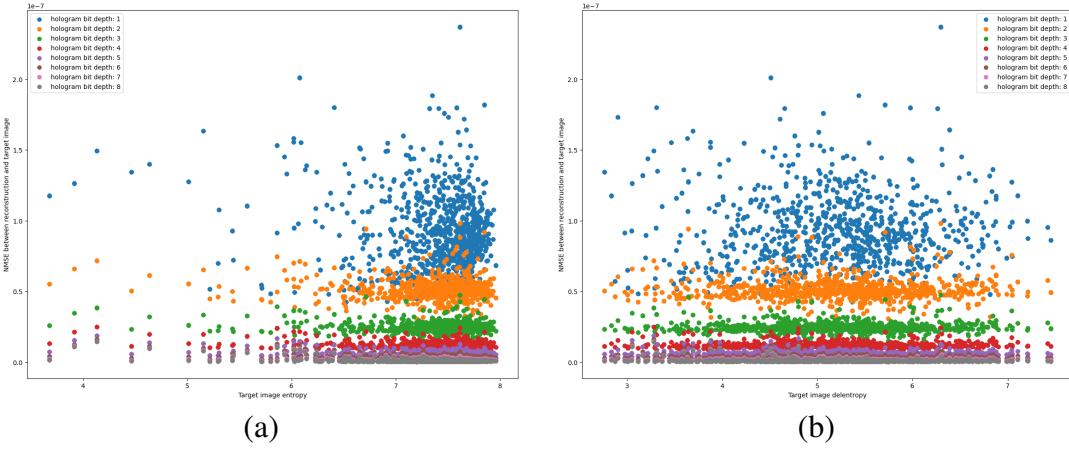


Fig. 7.7 Scatter plot for target images set at far field. (a) NMSE v.s. Entropy. (b) NMSE v.s. Delentropy.

Then the entropy of each of the 800 target images are calculated using Equation (7.3), and a scatter plot of the NMSE between the reconstruction and target image against the target image entropy is plotted for all 800 target images as shown in Fig. 7.7 (a), with the difference in hologram bit depth distinguished by different colours. To avoid the effect of initial random phase on the final result, 5 different randomly generated initial phases are used for each run. However, the result [85] had rarely shown any difference between each run, leading to 5 dots overlapping at the same spot in the scatter plot. Unfortunately, no correlation has been found between the NMSE and the target image entropy, inferring that the Shannon entropy cannot be used to quantify the difficulty of CGH for a target image.

Lastly, the delentropy of each of the 800 target images are calculated following the method in Section 7.3.3, and a scatter plot of the NMSE between the reconstruction and target image against the target image delentropy is plotted in Fig. 7.7 (b) for the 800 target images. And again, no correlation has been found between the NMSE and the target image delentropy either, inferring that the 2D delentropy also fails to quantify the difficulty of CGH for a target image.

#### 7.4.2 Targets at near field (Fresnel region)

The target images are now set to near field, where the Fresnel diffraction formula in Equation (7.1) applies; therefore the FFT and IFFT stages in Fig. 7.3 needs to be modified to include the phase term in Equation (7.1), and for experimental purpose, the distance ( $z$ ) is set

at  $10\text{cm}$ , the hologram's pixel pitch (sampling resolution of  $x$  and  $y$ ) has a size of  $13.62\mu\text{m}$  and the incident light's wavelength ( $\lambda$ ) is  $532\text{nm}$ .

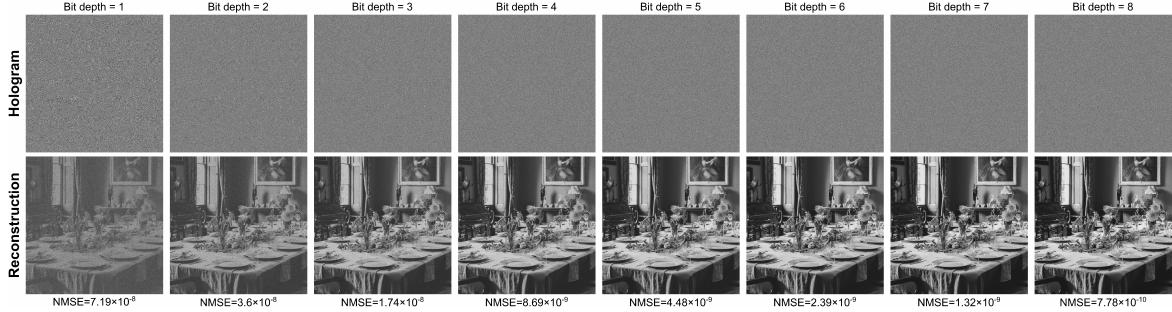


Fig. 7.8 Hologram generated at certain bit depths and their according reconstructions at near field

Fig. 7.8 shows both qualitatively and quantitatively how the reconstruction quality improves (i.e. NMSE decreases) with the increase in the bit depth of the hologram. Such trend is the same for target images places at near field as those placed at far field in Fig. 7.5. The rotational symmetry is gone for the binary phase hologram (bit depth = 1) due to the extra phase term in the Fresnel diffraction formula making the product of binary phase hologram and the phase term to be complex-valued whose complex conjugate does not equal to itself; however, the conjugate wouldn't disappear, but to appear at a different distance to where the target image is set at, leading to extra defocused noise onto the reconstruction plane. Nevertheless, the trend infers that holograms with higher bit depth produces better quality in the reconstruction plane.

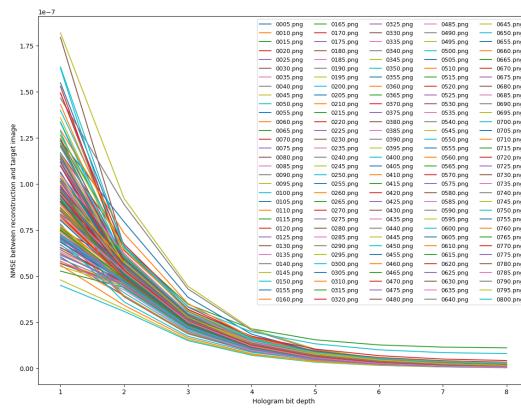


Fig. 7.9 NMSE v.s. Hologram bit depth for target images set at near field

Fig. 7.9 plots the trend of NMSE against increasing hologram bit depth, for every 5 target images to avoid overcrowding the plot, with all raw data accessible at the published research

data [85] (as the results for Fraunhofer has shown that the result from different initial random phases are the same, each run for Fresnel propagation only starts with one initial random phase as opposed to 5 different random phases for Fraunhofer propagation). In Fig. 7.9, the same trend as the one for far field in Fig. 7.6 can be observed, where NMSE decreases as hologram bit depth increases, for every single target image.

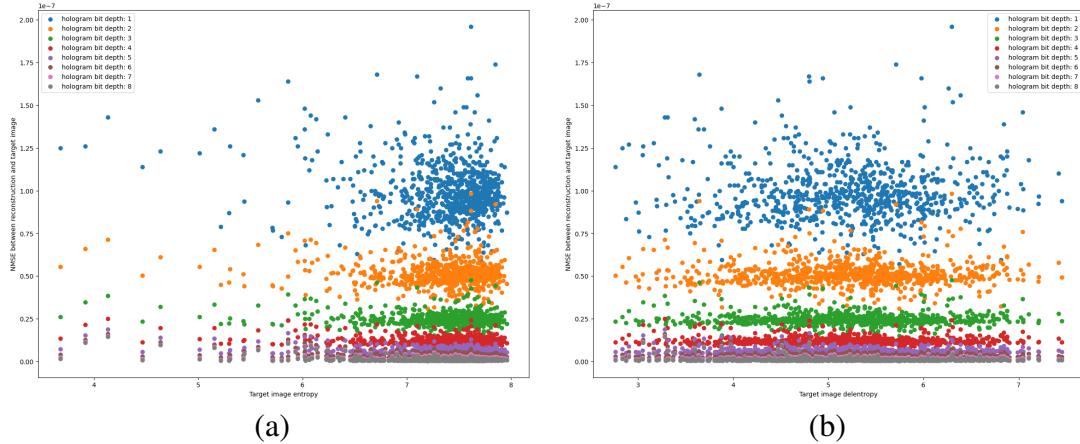


Fig. 7.10 Scatter plot for target images set at near field. (a) NMSE v.s. Entropy. (b) NMSE v.s. Delentropy.

The scatter plot of NMSE against target image's entropy and delentropy are plotted, in Fig. 7.10 (a) and (b) respectively. No correlation has been observed between the NMSE and either the entropy or delentropy of target image. Such 'no correlation observed' result is the same as the results for the targets in far field in Section 7.4.1, confirming that neither entropy nor delentropy is suitable for quantifying how difficult a target image is for phase hologram computation.

The research moves on to investigate the entropy of holograms. In an example run for a quantized hologram generation in the near field (Fresnel propagation in Equation (7.1), both the hologram entropy and the NMSE between reconstruction and target image are recorded and plotted in Fig. 7.11 (a) and (b) for the first 20 iterations. As randomly generated holograms will by definition create holograms with high entropy, additional experiment has been carried out with initial phase set to zeros (i.e.  $\angle A$  is set to zeros at  $n = 0$  in Fig. 7.3). Therefore, two diagrams are plotted in Fig. 7.11, with Fig. 7.11 (a) having initial phase of zeros and Fig. 7.11 (b) having random initial phase. Both Fig. 7.11 (a) and (b) have the horizontal axis to be the iteration number  $n$ , and the vertical axis in the left is the entropy of hologram (corresponding to solid lines in the plot), while the vertical axis in the right is

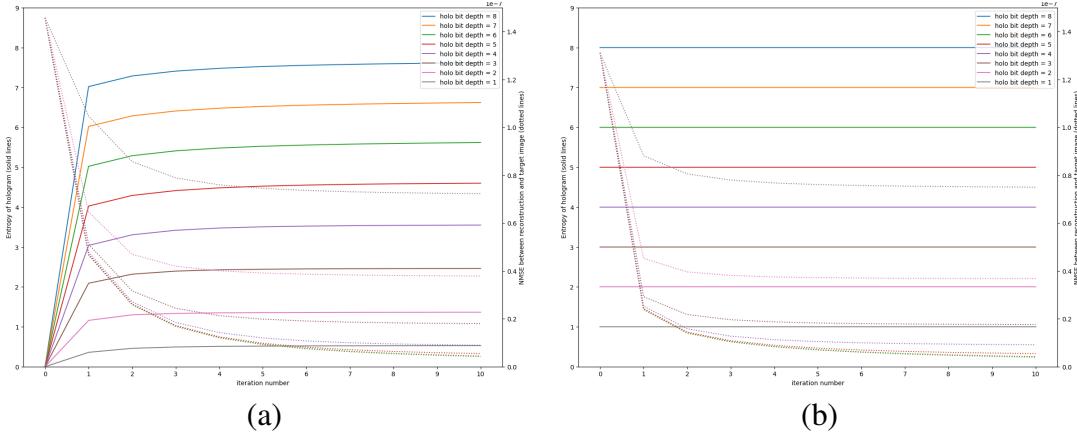


Fig. 7.11 Hologram entropy and NMSE v.s. iteration number for target images set at near field, with the initial phase ( $\angle A$ ) being (a) zeros (b) random

the NMSE (corresponding to dotted lines in the plot). Colour coding is used to distinguish between the 8 different runs where hologram bit depth is set from 1 to 8.

The solid lines in Fig. 7.11 (a) show that the entropy of hologram keeps increasing towards a value lower than the bit depth, with their corresponding NMSE between reconstruction and target image (dotted lines) decreasing. Such trend can be explained qualitatively that, as the iteration goes on, the hologram is attempting to contain more information to sustain a better reconstruction, while the entropy cannot exceed or even reach the bit depth level. On the other hand, the random initial phase plotted in Fig. 7.11 (b) has a constant entropy approximately equal to the bit depth, which infers that the iterations are improving the reconstruction quality without reducing the information entropy of the hologram. In both cases, the entropy of the hologram does not decrease at any iteration, and the final NMSE does not have significant reduction when the hologram's bit depth exceeds 5.

The entropy of final hologram in (a) is lower than that in (b), although the final NMSE are similar. If the computer-generated holograms were to undergo a lossless compression, the hologram in (a) would be better compressed than the hologram in (b) as the entropy denotes the compression limit, assuming that the holograms are treated as 1D arrays when compressing. Therefore if hologram compression is of a concern, then starting with a low entropy initial hologram in the CGH process is recommended. In case if the reconstruction quality is not as high of a priority than making holograms to occupy less storage space, this paper also recommends quantized GS algorithm with 5 bit depth instead of 6-7 bit depth as the final reconstruction quality will not degrade significantly.

## 7.5 Conclusion

By carrying out the quantized GS algorithm on 800 sample target images placed at both far field (Fraunhofer diffraction) and near field (Fresnel diffraction), and computing the entropy and delentropy of the target images and holograms, this paper reaches the conclusion that, holograms with higher bit depth can sustain more information therefore producing better quality reconstructions. However, the quality of the reconstruction is not correlated to either the entropy or the delentropy of the target image, so neither entropy nor delentropy can quantify how difficult an image is for phase-only hologram generation. Additionally, the entropy of the hologram generated using quantized GS algorithm is not only bounded by the hologram bit depth, but also affected by the entropy of the initial phase. For applications where holograms file size is a high priority, if the quantized GS method is used, this paper advises a low entropy initial phase (e.g. all zeros) rather than a random initial phase and recommends to reduce the hologram bit depth limit, for lower entropy hologram generation.



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