Primer on Semiconductors

Unit 3: Equilibrium Carrier Concentrations

Lecture 3.2: Fermi-Dirac integrals

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The Fermi function

The Fermi function gives the probability that a state (if it exists) is occupied in equilibrium.

$$f_0(E) = \frac{1}{1 + e^{(E - E_F)/k_B T}}$$
(Fermi function)

The two key parameters in the Fermi function are the Fermi level and the temperature.

Equilibrium carrier densities

$$f_{0}(E) = \frac{1}{1 + e^{(E - E_{F})/k_{B}T}} - - E_{F} - - - - -$$

 $n \propto e^{(E_F - E_C)/k_B T}$

$$E_{V}$$

 $p \propto e^{(E_V - E_F)/k_BT}$

nondegenerate semiconductor

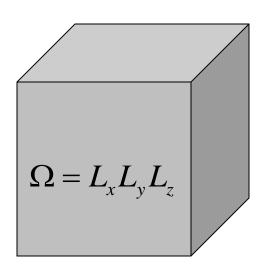
Two options

Our goal is to relate the carrier density to the Fermi level and to the properties of the semiconductor.

We can do the calculation two ways:

- 1) In k-space
- 2) In energy space

3D bulk semiconductor: k-space

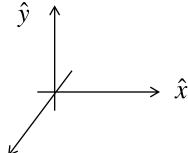


$$N = \sum_{\vec{k}} f_0(E_k)$$

$$n_0 = \frac{1}{\Omega} \sum_{\vec{k}} f_0(E_k)$$

$$n_0 = \frac{1}{\Omega} \int_{\vec{k}} f_0(E_k) N_k d^3 k$$

$$\sum_{\vec{k}} \bullet \to \frac{\Omega}{4\pi^3} \int_{BZ} \bullet \ d^3k$$



$$N_k = 2 \times \left(\frac{\Omega}{8\pi^3}\right) = \frac{\Omega}{4\pi^3}$$

Density-of-states in k-space

1D:
$$N_k = 2 \times \left(\frac{L}{2\pi}\right) = \frac{L}{\pi}$$

dk

$$N_k = 2 \times \left(\frac{A}{4\pi^2}\right) = \frac{A}{2\pi^2} \qquad dk_x dk_y$$

$$N_k = 2 \times \left(\frac{\Omega}{8\pi^3}\right) = \frac{\Omega}{4\pi^3} \qquad dk_x dk_y dk_z$$

independent of E(k)

$$\sum_{\vec{k}} \bullet \to N_k \int_{BZ} \bullet \ dk$$

(Should include valley degeneracy factor, g_{V} .)

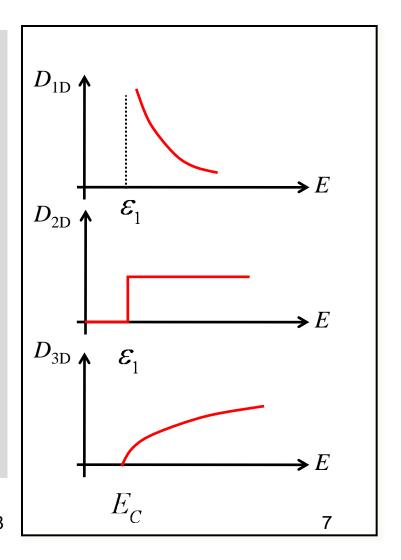
Energy space (parabolic bands)

$$D_{1D}(E) = g_V \frac{1}{\pi \hbar} \sqrt{\frac{2m^*}{(E - \varepsilon_1)}} \Theta(E - \varepsilon_1)$$

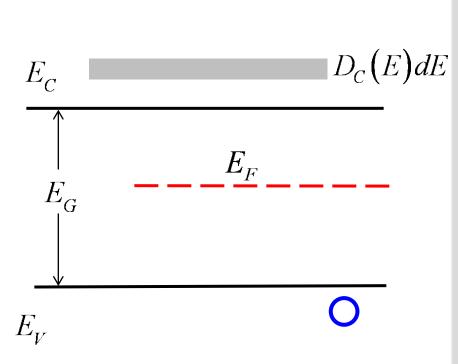
$$D_{2D}(E) = g_V \frac{m^*}{\pi \hbar^2} \Theta(E - \varepsilon_1)$$

$$D_{3D}(E) = g_V \frac{m^* \sqrt{2m^*(E - E_C)}}{\pi^2 \hbar^3} \Theta(E - E_C)$$

$$(E(k) = E_C + \hbar^2 k^2 / 2m^*)_{\text{Lundstrom: 2018}}$$



Distribution of electrons in the conduction band



$$n_0(E)dE = f_0(E)D_C(E)dE$$

$$f_0(E) = \frac{1}{1 + e^{(E - E_F)/k_B T}}$$

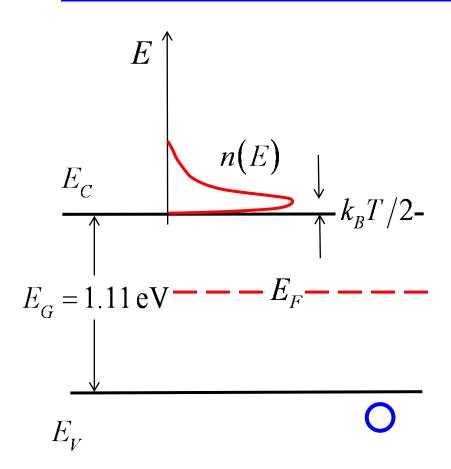
$$f_0(E) \approx e^{-(E-E_F)/k_BT}$$

$$D_C(E) \propto \sqrt{E - E_C}$$

$$n_0(E)dE \propto \sqrt{E-E_C}e^{-(E-E_F)/k_BT}dE$$

Lundstrom: 2018

Distribution of electrons in the conduction band



Electrons and holes are very near the band edges.

$$n_0(E) \propto \sqrt{E - E_C} e^{-(E - E_F)/k_BT}$$

$$\frac{dn_0(E)}{dE} = 0$$

$$E = k_{\rm\scriptscriptstyle R} T/2$$

$$E = 0.013 \text{ eV}$$
 $T = 300 \text{ K}$

Compute the electron density

$$n_0 = \int_{E_C}^{\infty} n_0(E) dE = \int_{E_C}^{\infty} f_0(E) D_C(E) dE$$

$$f_0(E) = \frac{1}{1 + e^{(E - E_F)/k_B T}}$$

$$D_{C}(E) = g_{V} \frac{m_{n}^{*} \sqrt{2m_{n}^{*}(E - E_{C})}}{\pi^{2}\hbar^{3}}$$

Fermi level and temperature

effective mass and valley degeneracy

Energy space (3D)

$$n_0 = \int_{E_C}^{\infty} f(E) D_C(E) dE$$

$$n_0 = \frac{\left(2m^*\right)^{3/2}}{2\pi^2\hbar^3} \int_{E_C}^{\infty} \frac{\left(E - E_C\right)^{1/2}}{1 + e^{(E - E_F)/k_B T}} dE$$

$$n_0 = \frac{\left(2m^* k_B T\right)^{3/2}}{2\pi^2 \hbar^3} \int_0^\infty \frac{\eta^{1/2} d\eta}{1 + e^{\eta - \eta_F}}$$

$$n_0 = N_C \mathcal{F}_{1/2}(\eta_F) \,\mathrm{cm}^{-3}$$

$$f_0(E) = \frac{1}{1 + e^{(E - E_F)/k_B T}}$$

$$D_{C}(E) = \frac{(2m^{*})^{3/2}}{2\pi^{2}\hbar^{3}} (E - E_{C})^{1/2}$$

$$\eta_F = (E_F - E_C)/k_B T$$

$$\eta = (E - E_C)/k_B T$$

$$N_{C} = \frac{1}{4} \left(\frac{2m^{*}k_{B}T}{\pi\hbar^{2}} \right)^{3/2}$$

Fermi-Dirac integrals

$$\mathcal{F}_{j}(\eta_{F}) \equiv \frac{1}{\Gamma(j+1)} \int_{0}^{\infty} \frac{\eta^{j} d\eta}{1 + e^{\eta - \eta_{F}}}$$

$$\Gamma(n) = (n-1)!$$
 (*n* integer)

$$\Gamma(1/2) = \sqrt{\pi}$$

$$\Gamma(p+1) = p\Gamma(p)$$

$$\mathcal{F}_{j}(\eta_{F}) \rightarrow e^{\eta_{F}} \quad \eta << 1$$

$$(E_{F} - E_{C})/k_{B}T << 1$$

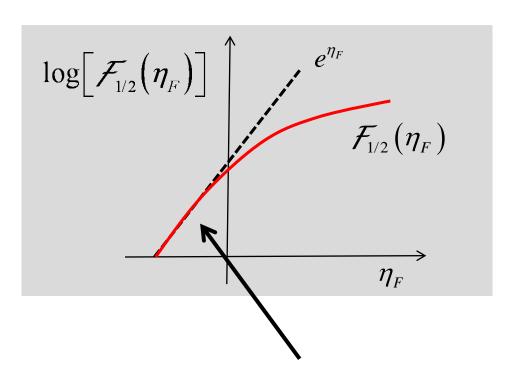
nondegenerate semiconductor

$$\frac{d\mathcal{F}_{j}}{d\eta_{F}} = \mathcal{F}_{j-1}$$

don't confuse with....
$$F_j(\eta) = \int_0^{+\infty} \frac{x^j dx}{1 + e^{x - \eta}}$$

For an introduction to Fermi-Dirac integrals, see: "Notes on Fermi-Dirac Integrals," 3rd Ed., by R. Kim and M. Lundstrom) https://www.nanohub.org/resources/5475

Fermi-Dirac integral of order 1/2



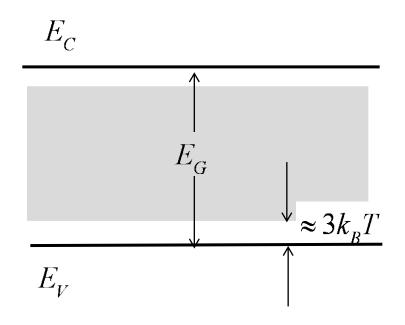
$$\mathcal{F}_{1/2}(\eta_F) \equiv \frac{2}{\sqrt{\pi}} \int_0^\infty \frac{\eta^{1/2} d\eta}{1 + e^{\eta - \eta_F}}$$

$$\eta_F = (E_F - E_C)/k_B T$$

$$\eta_F \ll 0$$
 $E_F \ll E_C$ $\mathcal{F}_{1/2}(\eta_F) \rightarrow e^{\eta_F}$ $n_0 = N_C e^{\eta_F} \text{ cm}^{-3}$

(nondegenerate semiconductor)

Nondegenerate semiconductor



In a nondegenerate semiconductor, the Fermi level is well below the bottom of the conduction band and well above the top of the valence band.

Compute the electron density

$$n_0 = N_C \mathcal{F}_{1/2}(\eta_F) \text{ cm}^{-3}$$
 $\eta_F = (E_F - E_C)/k_B T$ $N_C = \frac{1}{4} \left(\frac{2m^* k_B T}{\pi \hbar^2}\right)^{3/2}$

Nondegenerate semiconductors:

$$\eta_F \ll 0$$

$$E_F \ll E_C$$

$$\mathcal{F}_{1/2}(\eta_F) \to e^{\eta_F}$$

$$\eta_0 = N_C e^{\eta_F} \text{ m}^{-3}$$

Compute the hole density

$$p_0 = \int_{-\infty}^{E_V} \left[1 - f(E) \right] D_V(E) dE$$

$$f_0(E) = \frac{1}{1 + e^{(E - E_F)/k_B T}}$$

$$D_{V}(E) = g_{V} \frac{m_{p}^{*} \sqrt{2m_{p}^{*}(E_{V} - E)}}{\pi^{2} \hbar^{3}}$$

$$p_0 = N_V \mathcal{F}_{1/2}(\eta_F) \text{m}^{-3}$$
 $\eta_F = (E_V - E_F)/k_B T$ $N_V = \frac{1}{4} \left(\frac{2m_p^* k_B T}{\pi \hbar^2}\right)^{3/2}$

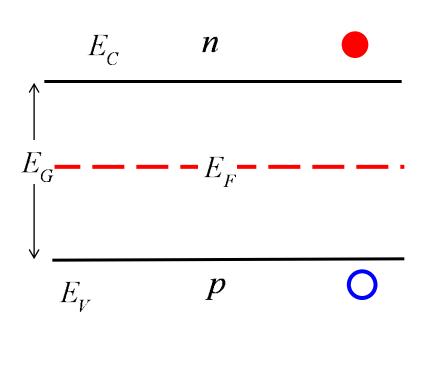
Nondegenerate semiconductor:

$$\eta_F \ll 0$$

$$E_F \gg E_U$$
 $\mathcal{F}_{1/2}(\eta_F) \rightarrow e^{\eta_F}$

$$p_0 = N_V e^{\eta_F} \text{ m}^{-3}$$

Summary



$$n_0 = N_C e^{(E_F - E_C)/k_B T}$$

$$p_0 = N_V e^{(E_V - E_F)/k_B T}$$

Nondegenerate semiconductor