

Computing matrix element of the three-body operator for parallel spins

Emmanuel Giner^{a)}

I. WHAT WE NEED TO COMPUTE

A. General matrix element

We need to compute the matrix elements of this nasty operator

$$O^3 = \frac{1}{6} \sum_{ijm} \sum_{klm} \langle ijm|klm \rangle \hat{k}^\dagger \hat{l}^\dagger \hat{n}^\dagger \hat{m} \hat{j} \hat{i}. \quad (1)$$

Therefore, it is the product of $\hat{k}^\dagger \hat{l}^\dagger \hat{n}^\dagger$ by $\hat{m} \hat{j} \hat{i}$, and each term posses a 6-fold permutation symmetry with ± 1 signature. Starting from $\hat{k}^\dagger \hat{l}^\dagger \hat{n}^\dagger$ there are two symmetric pertumations

$$\hat{k}^\dagger \hat{l}^\dagger \hat{n}^\dagger = \hat{l}^\dagger \hat{n}^\dagger \hat{k}^\dagger = \hat{n}^\dagger \hat{k}^\dagger \hat{l}^\dagger, \quad (2)$$

and three anti-symmetric permutations

$$\hat{l}^\dagger \hat{k}^\dagger \hat{n}^\dagger = \hat{n}^\dagger \hat{l}^\dagger \hat{k}^\dagger = \hat{k}^\dagger \hat{n}^\dagger \hat{l}^\dagger, \quad (3)$$

which fulfill

$$\hat{k}^\dagger \hat{l}^\dagger \hat{n}^\dagger = -\hat{l}^\dagger \hat{k}^\dagger \hat{n}^\dagger. \quad (4)$$

Therefore, one can express a 6×6 matrix that is the product all possible permutations for (k, l, n) and (i, j, m) , rearrange everything in terms of the original $\hat{k}^\dagger \hat{l}^\dagger \hat{n}^\dagger \hat{m} \hat{j} \hat{i}$ operator and obtain the following terms

$$\begin{aligned} \frac{1}{6} \mathcal{P}_{ijm}^{klm} \langle ijm|klm \rangle \hat{k}^\dagger \hat{l}^\dagger \hat{n}^\dagger \hat{m} \hat{j} \hat{i} = & \hat{k}^\dagger \hat{l}^\dagger \hat{n}^\dagger \hat{m} \hat{j} \hat{i} \left(\right. \\ & + (\langle ijm|klm \rangle + \langle ijm|lmk \rangle + \langle ijm|mkj \rangle) \\ & \left. - (\langle ijm|lkm \rangle + \langle ijm|mlk \rangle + \langle ijm|kml \rangle) \right), \end{aligned} \quad (5)$$

where \mathcal{P}_{ijm}^{klm} contains the sum over the 36 permutations of the 6 indices with the corresponding signs.

B. Diagonal term

For a diagonal term, one has that

$$\begin{aligned} i &= k \\ j &= l \\ m &= n \end{aligned} \quad (6)$$

and therefore one obtains

$$\begin{aligned} \langle I|O^3|I \rangle = & \sum_{i>j>m \in |I \rangle} + (\langle ijm|ijm \rangle + \langle ijm|jmi \rangle + \langle ijm|mij \rangle) \\ & - (\langle ijm|jim \rangle + \langle ijm|mji \rangle + \langle ijm|imj \rangle) \end{aligned} \quad (7)$$

where the summation $\sum_{i>j>m \in |I \rangle}$ means running over all orbitals occupied in $|I \rangle$ in a unique way.

^{a)}Electronic mail: emmanuel.giner@lct.jussieu.fr

C. Single excitation term

For a single excitation term, one has that

$$\begin{aligned} i &\neq k \\ j &= l \\ m &= n \end{aligned} \quad (8)$$

and then

$$\begin{aligned} \langle I|O^3 \hat{k}^\dagger \hat{i}|I \rangle = & (-1)^{\mathcal{P}_{ij}} \sum_{j>m \in |I \rangle} + (\langle ijm|kjm \rangle + \langle ijm|jmk \rangle + \langle ijm|mkj \rangle) \\ & - (\langle ijm|jkm \rangle + \langle ijm|mjk \rangle + \langle ijm|kml \rangle). \end{aligned} \quad (9)$$

D. Double excitation term

For a double excitation term, one has that

$$\begin{aligned} i &\neq k \\ j &\neq l \\ m &= n \end{aligned} \quad (10)$$

and then

$$\begin{aligned} \langle I|O^3 \hat{k}^\dagger \hat{l}^\dagger \hat{j} \hat{i}|I \rangle = & (-1)^{\mathcal{P}_{ij}} \sum_{m \in |I \rangle} + (\langle ijm|klm \rangle + \langle ijm|lmk \rangle + \langle ijm|mkj \rangle) \\ & - (\langle ijm|lkm \rangle + \langle ijm|mlk \rangle + \langle ijm|kml \rangle). \end{aligned} \quad (11)$$