

Deferred-Correction Approaches

- Size of computational molecule affects both storage requirements and effort needed to solve the algebraic system at each time-step
 - Usually, we wish to keep only the nearest neighbors of the center node P in the LHS of equations (leads to tri-diagonal matrix or something close to it) ⇒ easier to solve linear/nonlinear system
 - But, approximations that produce such molecules are often not accurate enough
- Way around this issue?
 - Leave only the terms containing the nearest neighbors in the LHS and bring all other more-remote terms to the RHS
 - This requires that these terms be evaluated with previous or old values, which may lead to divergence of the iterative scheme
- Better approach?



Better Approach

- Compute the terms that are approximated with a high-order approximation explicitly and put them in the RHS
- Take a simpler approximation to these terms (that give a small computational molecule). Insert it twice in the equation, with a + and - sign
- One of these two simpler approximations, keep it in the LHS of the equations (with unknown variables values, i.e. implicit/new). Move the other to the RHS (i.e. computing it explicitly using existing/old values)
- The RHS now contains the difference between two explicit approximations of the same term, and is likely to be small ⇒
 - Likely no convergence problems to an iteration scheme (Jacobi, GS, SOR, etc) or gradient descent (CG, etc)
- Once the iteration converges, the low order approximation terms (one explicit, the other implicit) drop out and the solution corresponds to the higher-order approximation
- ⇒ Using H & L for high & low orders:

$$\mathbf{A}^{H} \mathbf{x} = \mathbf{b} \longrightarrow \mathbf{A}^{L} \mathbf{x} = \mathbf{b} - \left[\mathbf{A}^{H} \mathbf{x} - \mathbf{A}^{L} \mathbf{x} \right]^{\text{old}}$$



- This approach can be very powerful and general
 - Used when treating higher-order approximations, non-orthogonal grids, corrections needed to avoid oscillation effects, etc
 - Since RHS can be viewed as a correction ⇒ called deferredcorrection
 - Note: both L&H terms could be implicit in time: use L&H explicit starter to get first values and then most recent old values in bracket during iterations (similar to Jacobi vs. Gauss Seidel)
 - Explicit for H (high-order) term, implicit for L (low-order) term

$$\mathbf{A}^{H} \mathbf{x} = \mathbf{b} \longrightarrow \mathbf{A}^{L} \mathbf{x}_{implicit} = \mathbf{b} - \left[\mathbf{A}^{H} \mathbf{x}_{explicit} - \mathbf{A}^{L} \mathbf{x}_{implicit} \right]^{old}$$

Implicit for both L and H terms (similar to Gauss-Seidel)

$$\mathbf{A}^{H} \mathbf{x} = \mathbf{b} \longrightarrow \mathbf{A}^{L} \mathbf{x}_{\text{implicit}} = \mathbf{b} - \left[\mathbf{A}^{H} \mathbf{x}_{\text{implicit}} - \mathbf{A}^{L} \mathbf{x}_{\text{implicit}} \right]^{\text{old}}$$



- Example 1: FD methods with High-order Pade' schemes
 - One can use the PDE itself to express implicit Pade' time derivative $\left(\frac{C\phi}{\partial t}\right)$ as a function of ϕ^{n+1} (see homework)
 - Or, use deferred-correction (within an iteration scheme of index r):

• In time:
$$\left(\frac{\partial \phi}{\partial t}\right)_n^{r+1} = \left(\frac{\phi_{n+1} - \phi_{n-1}}{2\Delta t}\right)^{r+1} + \left[\left(\frac{\partial \phi}{\partial t}\right)_n^{\text{Pade'}} - \frac{\phi_{n+1} - \phi_{n-1}}{2\Delta t}\right]^r$$

• In space:
$$\left(\frac{\partial \phi}{\partial x}\right)_{i}^{r+1} = \left(\frac{\phi_{i+1} - \phi_{i-1}}{2\Delta x}\right)^{r+1} + \left[\left(\frac{\partial \phi}{\partial x}\right)_{i}^{\text{Pade'}} - \frac{\phi_{i+1} - \phi_{i-1}}{2\Delta x}\right]^{r}$$

- The complete 2nd order CDS would be used on the LHS. The RHS would be the bracket term: the difference between the Pade' scheme and the "old" CDS. When the CDS becomes as accurate as Pade', this term in the bracket is zero
- Note: Forward/Backward DS could have been used instead of CDS, e.g. in $\left(\frac{\partial \phi}{\partial t}\right)_{n+1}^{r+1} = \left(\frac{\phi_{n+1} - \phi_n}{\Delta t}\right)^{r+1} + \left[\left(\frac{\partial \phi}{\partial t}\right)_{n+1}^{\text{Pade'}} - \frac{\phi_{n+1} - \phi_n}{\Delta t}\right]^{r+1}$



- Example 2 with FV methods: Higher-order Flux approximations
 - Higher-order flux approximations are computed with "old values" and a lower order approximation is used with "new values" (implicitly) in the linear system solver: $F_e = F_e^L + \left[F_e^H - F_e^L \right]^{\text{old}}$

where F_{ρ} is the flux. For ex., the low order approximation is a UDS or CDS

- Convergence and stability properties are close to those of the low order implicit term since the bracket is often small compared to this implicit term
- In addition, since bracket term is small, the iteration in the algebraic equation. solver can converge to the accuracy of higher-order scheme
- Additional numerical effort is explicit with "old values" and thus much smaller than the full implicit treatment of the higher-order terms
- A factor can be used to produce a mixture of pure low and pure high order. This can be used to remove undesired properties, e.g. oscillations of highorder schemes $F_e = \omega F_e^L + (1 - \omega) \left[F_e^H - F_e^L \right]^{\text{old}}$