Complex Flows - Hand in 2

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1

During the collision of two particles in the simulation there is a possibility that the particles will stick together instead of repelling. The reason being, that the simulation works in specified time steps and if the particles are in collision, they are assumed to have left the collision after an adjustment has been made to the velocity. Originally, the velocities were adjusted in such a way that a Δv was calculated using equation (2), where $\mathbf{v}_1^{\{0\}}$ and $\mathbf{v}_2^{\{0\}}$ are the particle's velocities before the collision. Then, Δv is added to both particle's velocities with opposite signs. The relative velocity, $\mathbf{G}^{\{0\}}$ is described in equation (1).

$$\mathbf{G}^{\{0\}} = \mathbf{v}_1^{\{0\}} - \mathbf{v}_2^{\{0\}} \tag{1}$$

$$\Delta v = \mathbf{n}(\mathbf{G}^{\{0\}} \cdot \mathbf{n}) \tag{2}$$

This works fine as long as during the next time step, the particles are no longer intersecting. If they are, the signs are reversed again, causing them to stick together and with higher speed, they can appear to vibrate. A simple, and perhaps a lazy solution, would be to make a matrix that records if two particles are in a collision and using that to only allow the velocities to update once per collision of two particles.

Using this method, a significant number of *sticking* can be averted. This problem will however reappear as the assignment progresses.

 $\mathbf{2}$

Before applying friction on sliding particles and restitution coefficient a few properties of the collision have to be defined.

Using $G^{\{0\}}$ from equation (1) the relative velocity of the contact point before the collision can be defined [1].

$$\mathbf{G}_c^{\{0\}} = \mathbf{G}^{\{0\}} + r_1 \omega_1^{\{0\}} \times \mathbf{n} + r_2 \omega_2^{\{0\}} \times \mathbf{n}$$
(3)

In equation (3), the radius of both particles, r_1 and r_2 , are multiplied with the corresponding angular velocity, $\omega_1^{\{0\}}$ and $\omega_2^{\{0\}}$. The cross product of these terms with the normal vector, \mathbf{n} , provides a tangential contribution to the relative velocity at the contact point, $\mathbf{G}_c^{\{0\}}$ [1]. The tangential component of $\mathbf{G}_c^{\{0\}}$ is defined in equation (4)

$$\mathbf{G}_{ct}^{\{0\}} = \mathbf{G}_c^{\{0\}} - \left(\mathbf{G}_c^{\{0\}} \cdot \mathbf{n}\right) \mathbf{n} \tag{4}$$

Using the length of the vector $\mathbf{G}_{ct}^{\{0\}}$, the tangential unit vector can be defined in equation (5)

$$\mathbf{t} = \frac{\mathbf{G}_{ct}^{\{0\}}}{|\mathbf{G}_{ct}^{\{0\}}|}.\tag{5}$$

Now, the coefficient of restition is defined in equation (6), in order to better simulate particle to particle and particle to wall collisions.

$$\mathbf{n} \cdot \mathbf{G}^{\{0\}} = -e \left(\mathbf{n} \cdot \mathbf{G}^{\{0\}} \right) \tag{6}$$

the normal and tangential components of the Impulsive force are given by equation (7) and equation (8).

$$J_n = -\frac{m_1 m_2}{m_1 + m_2} (1 + e) (\mathbf{n} \cdot \mathbf{G}^{\{0\}})$$
 (7)

$$J_t = fJ_n \tag{8}$$

During particle to particle collision, two scenarios have to be considered: If the particles slide through collision, or not. The velocities and angular velocities of the particles are determined depending on if they slide through collision or not. The condition can be described with equation (9) [1].

$$J_t > -\left(\frac{2}{7}\right) \frac{m_1 m_2}{m_1 + m_2} |\mathbf{G}_{ct}^{\{0\}}| \tag{9}$$

If the particles slide through the collision the velocities and angular velocities after the collision can be evaluated using equation (10), equation (11), equation (12) and equation (13) [1].

$$\mathbf{v_1} = \mathbf{v_1}^{\{0\}} - (\mathbf{n} + f\mathbf{t})(\mathbf{n} \cdot \mathbf{G}^{\{0\}})(1+e) \frac{m_2}{m_1 + m_2}$$
(10)

$$\mathbf{v_2} = \mathbf{v_2}^{\{0\}} + (\mathbf{n} + f\mathbf{t})(\mathbf{n} \cdot \mathbf{G}^{\{0\}})(1+e) \frac{m_1}{m_1 + m_2}$$
(11)

$$\omega_{1} = \omega_{1}^{\{0\}} - \frac{5}{2r_{1}} (\mathbf{n} \cdot \mathbf{G}^{\{0\}}) (\mathbf{n} \times \mathbf{t}) (1+e) \frac{m_{2}}{m_{1} + m_{2}}$$
(12)

$$\omega_2 = \omega_2^{\{0\}} - \frac{5}{2r_2} (\mathbf{n} \cdot \mathbf{G}^{\{0\}}) (\mathbf{n} \times \mathbf{t}) (1+e) \frac{m_2}{m_1 + m_2}$$
(13)

If the condition is not fulfilled, that is, if the particles do not slide through the condition the particles' properties are found using equation (14), equation (15), equation (16) and equation (17) [1].

$$\mathbf{v_1} = \mathbf{v_1}^{\{0\}} - \left((1+e)(\mathbf{n} \cdot \mathbf{G}^{\{0\}})\mathbf{n} + \frac{2}{7} |\mathbf{G}_{ct}^{\{0\}}|\mathbf{t} \right) \frac{m_2}{m_1 + m_2}$$
(14)

$$\mathbf{v_2} = \mathbf{v_2}^{\{0\}} + \left((1+e)(\mathbf{n} \cdot \mathbf{G}^{\{0\}})\mathbf{n} + \frac{2}{7} |\mathbf{G}_{ct}^{\{0\}}| \mathbf{t} \right) \frac{m_1}{m_1 + m_2}$$
(15)

$$\omega_{1} = \omega_{1}^{\{0\}} - \frac{5}{7r_{1}} |\mathbf{G}_{ct}^{\{0\}}| (\mathbf{n} \times \mathbf{t}) \frac{m_{2}}{m_{1} + m_{2}}$$
(16)

$$\omega_{2} = \omega_{2}^{\{0\}} - \frac{5}{7r_{2}} |\mathbf{G}_{ct}^{\{0\}}| (\mathbf{n} \times \mathbf{t}) \frac{m_{1}}{m_{1} + m_{2}}$$
(17)

This way, particle to particle collisions are modelled more accurately. For particle to wall collisions, the velocity and angular velocity of the particle is scaled by the coefficient of restitution.

Applying these conditions greatly affects the average velocity of the particles over time. In order to quantify these changes, the ratio of the average particle speed to the initial average particle speed is calculated, equation (18), for each time step.

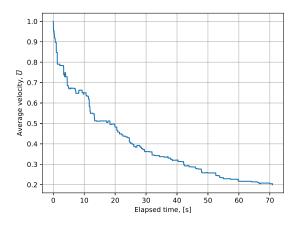
$$\overline{U} = \frac{U_{average}}{U_{initial}} \tag{18}$$

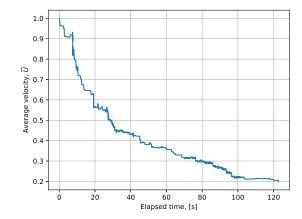
The time it takes for \overline{U} to reach 0.2 will be the variable of interest. At that point, the average velocity of the particles in the system will have lowered by 80% of the initial value. The ratio of 0.2 is chosen based on observations of the simulation and was found to be a reasonable ratio as the rate at which the average velocities decrease slows down significantly after that point. This variable of interest is defined by equation (19).

$$\tau = t_{\overline{U} = 0.2} \tag{19}$$

Since the initial velocities are randomized, there is a non predictive factor in the model, making the variable τ inconsistent to some degree. However, in general, τ is relatively consistent between models.

Comparing figure 1b and figure 1a, the improved model shows consistently a $\tau \approx 62$ s while the original model has a consistent $\tau \approx 118$ s. This shows that the implementation of the coefficients of restitution and friction significantly speed up the rate at which the kinetic energy in the system is lost. The kinetic energy of the system is determined by calculating the kinetic energy of each particle and adding them up. The kinetic energy of the system is plotted on figure 2.

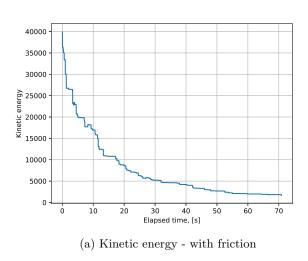


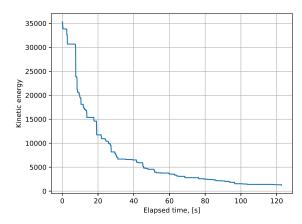


- (a) Normalized average velocity with friction
- (b) Normalized average velocity without friction

Figure 1: Normalized average velocity

As mentioned, the velocity is normalized while the kinetic energy is not which is why the y-axis on figure 2a and figure 2b does not match. Due to this, the normalized velocity is a more descriptive property of the flow to analyse than the kinetic energy of the system.





(b) Kinetic energy - without friction

Figure 2: Kinetic energy

3

In it's current state the particles are only affected by their interactions between other particles and the wall. Implementing passive forces will change the model drastically. For this assignment the velocity of each particle will be adjusted by three different accelerations in the form of their respective forces. The sum of the forces are used to adjust the velocity. Each force is represented by a matrix of the same size as the matrix that describes the velocity of the particles.

Force due to gravitational pull

The gravitational acceleration is implemented in the code by making the second column, or the y-component, of the force matrix equal to the gravitational acceleration of $-9.81 \frac{\text{m}}{\text{s}^2}$ times the mass of each particle. For each particle, i, the force due to the gravitational pull is defined in equation (20), where M is the mass of the particle. In the current model the mass is the same for each particle and it's value is 1.

$$F_g^{\{i\}} = \begin{bmatrix} 0\\ -9.81 \end{bmatrix} \cdot M \tag{20}$$

Force due to Drag

The Drag force on each particle is calculated using equation (21) [2] where u_i and v_i are the horizontal and vertical components of the particle's velocity, C_D is the coefficient of drag and A is the area normal to the flow direction, in this case $A = r^2 * \pi$ where r is the radius of the particle. The vector n_d is a directional unit vector of the velocity.

$$F_d^{\{i\}} = -\frac{1}{2} \left(\mathcal{C}_D \cdot A \left(\sqrt{u_i^2 v_i^2} \right)^2 \right) \mathbf{n_d}$$
 (21)

Force due to Coriolis acceleration

The Coriolis acceleration is calculated using equation (22) [2].

$$F_c^{\{i\}} = -2 \cdot (\omega_i \times \mathbf{n_d}) \tag{22}$$

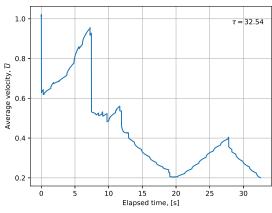
The sum of these forces, F_{sum} , are used to adjust the velocity before the particles are moved according to equation (23) where dt is the time increment by which the program moves each run.

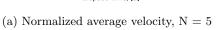
$$\mathbf{v} = \mathbf{v}^{\{0\}} + \frac{F_{sum}}{M} \cdot dt \tag{23}$$

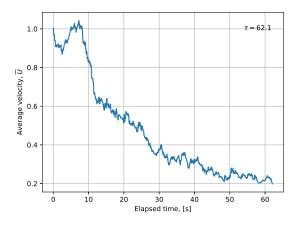
The simulation now shows the particles collect at the bottom, causing them to clip, similarly as in the first part of this assignment, however due to a different reason. Now, the particles are clipping since the particles settle at a lower velocity and are generally closer so when the velocity after collision is updated it is often not enough to expel the particles away from one another. This and the added effect of the gravitational acceleration, which accelerates the colliding particles in the same direction, while one might be at the bottom, gaining no velocity. Decreasing dt can reduce this effect.

4

Varying the number of particles provides an interesting comparison of the simulation's velocity plots. On figure 3a, the small number of particles makes particle to particle collision a rarity and with the added effect from the gravitational force, the particles bounce independently until the wall collision and other forces slow them down. The particle bounce can be seen on each plot in figure 3 where the average velocity increases as the particles accelerate towards the bottom, hit the bottom, the vertical velocity is reversed with relatively low loss, and then the particle slows down as it travels away from the bottom.







(b) Normalized average velocity, N = 35

Figure 3: Normalized average velocity with varying number of particles, N

As the number of particles increase the value of τ increases, meaning the system takes a longer time to slow lose it's energy. With increased rate of particle to particle interaction, there are more instances where the particles slide through the collision, causing them to spin. Due to the Coriolis acceleration in combination with the gravitational acceleration as the particles begin to spin more while accelerating towards the bottom the

particles begin to not only spin but travel in a circular pattern. This causes erraticity in the average velocity which can clearly be seen increasing with the number of particles on figure 3.

All previous simulations were made using the same fluid density of $\rho_{\rm air}=1.239\,{\rm kg\over m^3}$. Increasing the density of the fluid should increase the effects of the drag force. In theory this should lower the τ value. The model simulated 20 particles in a fluid with the density of air, water, corn syrup and molasses. As can be seen on figure 4, the τ value does decrease however not by a large margin. This suggests either that even still the gravitational force has a much larger effect on the particles than the drag force or that the drag force is not correctly implemented.

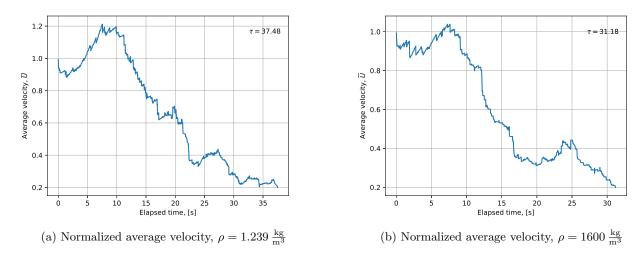


Figure 4: Normalized average velocity with varying fluid density, ρ

APPENDIX

additional plots

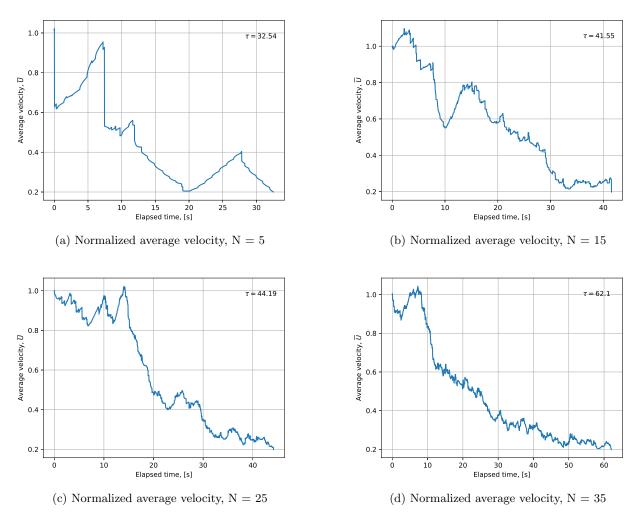
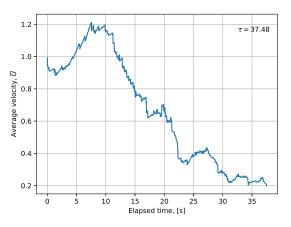
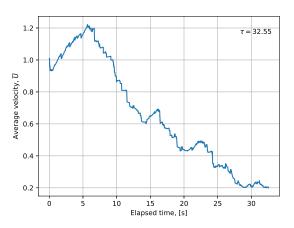


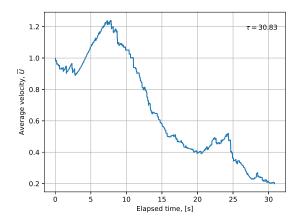
Figure 5: Normalized average velocity with varying number of particles, N



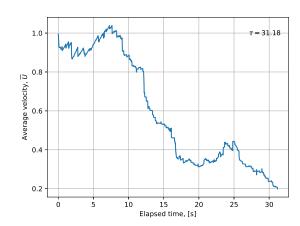
(a) Normalized average velocity, $\rho=1.239\,\frac{\mathrm{kg}}{\mathrm{m}^3}$



(c) Normalized average velocity, $\rho = 1380 \frac{\text{kg}}{\text{m}^3}$



(b) Normalized average velocity, $\rho = 997\;\frac{kg}{m^3}$



(d) Normalized average velocity, $\rho = 1600 \, \frac{\text{kg}}{\text{m}^3}$

Figure 6: Normalized average velocity with varying fluid density, ρ

Model Code

```
import numpy as np
2
     from numpy import sqrt, cos, sin, array, zeros
     import pyglet
     import matplotlib.pyplot as plt
     #parameters
     window_size = np.array((800, 800))
     nparticles = 35
     radius = 20
     max_velocity = 100 #Pixels per second
     max_omega = 0.5
11
     color1 = (20, 20, 250) # blue
12
     color2 = (200, 200, 250) # light blue
13
     prevents = 0
14
     #Tracked variables
16
     plotVel = []
17
     plotKE = []
18
     collision = np.zeros((nparticles,nparticles))
19
     timevalue = float(0)
20
     dt = 1/20000
21
22
     initialV = 0
     AvgVel = 0
23
     KE = 0
24
25
     #Physical properties
26
27
     A = (radius*10**(-3))**2*np.pi
28
    mass
    f = 0.5
                     # friction coefficient
29
```

```
30
     e = 0.8
                      # coefficient of restitution
    CD = 0.5
                     #coefficient of drag - F_D = CD*1/2*rho*U^2*A
31
    CL = 0.1
                     #coefficient of lift
     FD = np.zeros((nparticles,2) )
                                            #Drag force on particles
33
     Fg = np.zeros((nparticles,2) )
34
     FC = np.zeros((nparticles,2) )
35
     FR = np.zeros((nparticles,2) )
36
      rho = 1.239 #density of air
37
      #rho_W = 997 #density of water
38
      \#rho\_S = 1500
39
      \#rho\_M = 1200
40
      #t0 = time()
41
42
     # initialize global variables for particles
43
 44
      position = zeros((nparticles, 2))
     velocity = zeros((nparticles, 2))
45
46
     angle = zeros(nparticles)
 47
      omega = zeros(nparticles)
     circles = []
48
49
      # make window and batch
50
     window = pyglet.window.Window(window_size[0], window_size[1])
51
      batch = pyglet.graphics.Batch()
52
53
54
      # functions for particle simulation
55
56
      def avg_velocity(velocity):
57
         Sum = 0
         for i in range(len(velocity)):
58
              Sum += np.sqrt(velocity[i,0]**2+velocity[i,1]**2)
59
         avgVel = Sum/len(velocity)
60
61
         return avgVel
62
      def totKE(velocity):
63
         KE = 0
 64
         for i in range(nparticles):
65
66
             KE += 0.5*length(velocity[i])**2*mass
         return KE
67
68
      def length(vector):
69
         length = np.sqrt(vector[0]**2+vector[1]**2)
70
71
         return length
72
73
      def make_particles(position, velocity, omega, circles):
74
         # make particles in a grid at least one diameter from walls
         xstart = np.ones(2) * radius * 2
75
76
         xlength = window_size - 2 * xstart
         n = np.ceil(np.sqrt(nparticles))
77
78
         m = np.ceil(nparticles / n)
         xg, yg = np.meshgrid(np.arange(n)/(n-1), np.arange(m)/(m-1))
79
          # make random velocities and rotation
80
 81
         velangle = np.random.rand(nparticles) * 2 * np.pi
         velocity += np.random.rand(nparticles).reshape(-1,1) * max_velocity \
82
 83
                     * array([cos(velangle), sin(velangle)]).T
         omega += (np.random.rand(nparticles) - 0.5) * 2 * max_omega
 84
         global initialV
85
         initialV = avg_velocity(velocity)
         # prepare particles for simulation
87
         for i in range(nparticles):
 88
            position[i,0] = xg.flatten()[i] * xlength[0] + xstart[0]
89
             position[i,1] = yg.flatten()[i] * xlength[1] + xstart[1]
90
              # generate circle for particle
 91
              {\tt circles.append(pyglet.shapes.Circle(position[i,0], position[i,1],}\\
92
                                                  radius, color=color1, batch=batch))
93
              # generate spot on particle to track rotation
94
              x, y = position[i,:] + 0.5 * radius * array([cos(angle[i]), sin(angle[i])])
95
96
              circles.append(pyglet.shapes.Circle(x, y, 0.25*radius,
                                                  color=color2, batch=batch))
97
98
      def Forces(dt, velocity, omega, rho):
99
100
         for i in range(nparticles):
             nF = velocity[i]/length(velocity[i])
101
             nF3 = [nF[0], nF[1], 0]
102
103
              \#tF = \ np. \ array \ (nF[:,2], [-nF[:,1], \ nF[:,0]]) \ . \ T
              temp_omega = [0, 0, omega[i]]
104
              FC[i,0] = -2*np.cross(temp_omega,nF3)[0]
105
```

```
C3 = -2*np.cross(temp_omega,nF3)
                                                        #Coriolis acceleration
106
              FC[i.0] = C3[0]*mass
107
              FC[i,1] = C3[1]*mass
108
              FD = -(CD*0.5*length(velocity[i])**2*A)*nF
109
110
              Fg[i,1] = -9.81*mass
111
          Fsum = FD+Fg+FC
112
          velocity += Fsum/mass*dt
113
114
115
      def move_particles(dt, position, angle, circles):
         position += velocity * dt
116
          position = np.round(position,decimals=2)
117
          angle += omega * dt
118
          # update circle positions
119
120
          for i in range(nparticles):
              circles[i*2].position = position[i]
121
               \texttt{circles[i*2+1].position = position[i] + 0.5 * radius * array([cos(angle[i]), sin(angle[i])]) } 
122
123
      def wall collision(velocity):
124
125
          \# handle particle collition with the walls
          for i in range(nparticles):
126
127
              if position[i,0] < radius:
                  velocity[i,0] = e*abs(velocity[i,0])
128
                  omega[i] = omega[i]*-e
129
130
              if position[i,0] > window_size[0] - radius :
                  velocity[i,0] = -e*abs(velocity[i,0])
131
132
                  omega[i] = omega[i] *-e
133
              if position[i,1] < radius:</pre>
                  velocity[i,1] = e*abs(velocity[i,1])
134
135
                  omega[i] = omega[i] *-e
              if position[i.1] > window size[1] - radius :
136
                  velocity[i,1] = -e*abs(velocity[i,1])
137
                  omega[i] = omega[i] *-e
138
139
      def particle_collision(velocity):
140
          for i in range(nparticles-1):
141
              for j in range(i+1, nparticles):
142
                  distance = sqrt(((position[j] - position[i])**2).sum())
143
144
                  if distance < 2 * radius and collision[i,j] == 0:</pre>
                       # collision! - apply textbook eq. 5.14
145
                      collision[i,j] = 1
146
                      collision[j,i] = 1
147
                      n = (position[j] - position[i]) / (distance)
148
149
                      G0 = velocity[i] - velocity[j]
150
                      G03 = np.zeros(3)
151
                      G03[:-1] = G0
152
                      n3 = np.zeros((1,3))
153
                      n3[0,:-1] = n
154
                      omega3 = np.zeros((nparticles,3))
155
                      omega3[:,2] = omega
156
157
                      Gc3 = G03+np.cross(radius*omega3[i],n3) + np.cross(n3,radius*omega3[j])
                      Gc = Gc3[0, :2]
158
159
                      Gct = Gc - (np.dot(Gc, n))*n
                      Gc3 = np.squeeze(np.asarray(Gc3))
160
161
                      n3 = np.squeeze(np.asarray(n3))
                      Gct3 = Gc3 - (np.dot(Gc3, n3))*n3
162
                      lengthGct = length(Gct)
163
                      t = Gct/ lengthGct
164
                      t3 = Gct3 / length(Gct3)
165
                      Jn = -(mass**2/(2*mass))*(1+e)*(np.dot(n,G0))
166
167
                      if Jt > -(2/7)*mass**2/(mass*2)* lengthGct:
168
                          dveli = (n+f*t)*(np.dot(n,G0))*(1+e)*mass/(mass*2)
169
                          dvelj = (n+f*t)*(np.dot(n,G0))*(1+e)*mass/(mass*2)
170
                          {\tt domegi = (5)/(2*radius) *np.dot(n3,G03)*(np.cross(n3,t3))*f*(1+e)*mass/(mass*2)}
171
                          domegj = (5)/(2*radius) *np.dot(n3,G03)*(np.cross(n3,t3))*f*(1+e)*mass/(mass*2)
172
173
174
                          dveli = ((1+e)*(np.dot(n,G0))*n + 2/7*lengthGct * t)*mass/(2*mass)
175
176
                          dvelj = ((1+e)*(np.dot(n,G0))*n + 2/7*lengthGct * t)*mass/(2*mass)
                           domegi = 5/(7*radius)*length(Gct3) * (np.cross(n3,t3)) * mass/(mass*2)
177
                          domegj = 5/(7*radius)*length(Gct3) * (np.cross(n3,t3)) * mass/(mass*2)
178
179
                      velocity[i] -= dveli
180
181
                      velocity[j] += dvelj
```

```
182
                      domegi = domegi[2]
                                               #Since the simulation is two-dimensional, only spin around the z-axis affects the problem.
183
                      domegj = domegj[2]
184
185
186
                      omega[i] -= np.sign(omega[i]) * domegi
                      omega[j] += np.sign(omega[j]) * domegj
187
                  elif distance < 2 * radius and collision[i,j] == 1:
188
189
                      global prevents
190
191
                      prevents += 1
                      #print('t: ', time()-t0, ' prevented events: ', prevents)
192
193
                  else:
                      collision[i,j] = 0
194
                      collision[j,i] = 0
195
196
197
      def particle_collisionoriginal(velocity):
198
          # handle collision between particles using simple loops
199
          for i in range(nparticles-1):
              for j in range(i+1, nparticles):
200
201
                  distance = sqrt(((position[j] - position[i])**2).sum())
                  if distance < 2 * radius:
202
                       # collision! - apply textbook eq. 5.14
203
                      n = (position[j] - position[i]) / distance
204
                      G0 = velocity[i] - velocity[j]
205
206
                      dvel = n * np.dot(n, G0)
                      velocity[i] -= dvel
207
208
                      velocity[j] += dvel
209
      def plot(i, xi, yi, xlabel, ylabel, tau):
210
          tau = np.round(tau, decimals=2)
211
          plt.figure(i)
212
213
          plt.clf()
          plt.text(0.87*max(xi),0.95*max(yi), r'$\tau = {}$'.format(tau))
214
215
          plt.plot(xi,yi)
          plt.grid()
216
          plt.ylabel(ylabel)
217
218
          plt.xlabel(xlabel)
          plt.show(block = False)
219
220
221
     # modify pyglet draw command to draw our particles
222
223
      @window.event
      def on_draw():
224
225
          window.clear()
226
          batch.draw()
227
228
      # This is want we do in each time step
      def update(dt):
229
230
          global timevalue
          global initialV
231
          global AvgVel
232
233
          timevalue += dt
          Forces(dt, velocity, omega, rho)
234
235
          {\tt move\_particles(dt, position, angle, circles)}
          wall_collision(velocity)
236
          particle collision(velocity)
237
238
          AvgVel = avg_velocity(velocity)
239
240
          KE = totKE(velocity)
          normVel = AvgVel/initialV
241
242
          plotVel.append(normVel)
          plotKE.append(KE)
243
244
245
          if normVel < 0.2:</pre>
             pyglet.app.exit()
246
247
              return
248
249
250
      \# run the following if this is the main script
      if __name__ == "__main__":
251
252
          # Update the game 120 times per second
253
          pyglet.clock.schedule_interval(update, dt)
254
255
          # Create particles with random position and velocity
256
257
          make_particles(position, velocity, omega, circles)
```

```
258
            # Tell pyglet to do its thing
259
260
            pyglet.app.run()
            x = np.linspace(0, timevalue, len(plotVel))
261
            plot(1,x,plotVel, 'Elapsed time, [s]', r'Average velocity, $\overline{U}$', timevalue) plot(2,x,plotKE, 'Elapsed time, [s]', 'Kinetic energy', timevalue)
262
263
            print(r'\tau',timevalue)
264
            del window
265
            del batch
266
```

Bibliography

- $[1] \quad \hbox{Clayton T. Crow et al. } \textit{Multiphase Flows with Droplets and Particles}. \ 2012.$
- [2] Pijush K. Kundu, Ira M. Cohen, and David R. Downling. Fluid Mechanics. 2016.