

Preliminaries

Def (PDE) Eq. involving function and partial deriv.
 Not: $u_{xx} = \frac{\partial u}{\partial x_k} \rightarrow$ typically u
 well-posed:
 1. has a solution
 2. solution is unique
 3. small change eq \rightarrow s.c. sol \rightarrow stability

Theorem (Schwarz) function u cont. diff. at x
 $u_{xy}(x) = u_{yx}(x) \rightarrow$ order doesn't matter
 can be used to show that cont. sol doesn't exist

Def (strong/weak solution)
 strong: all deriv. in PDE are cont.
 \hookrightarrow otherwise weak

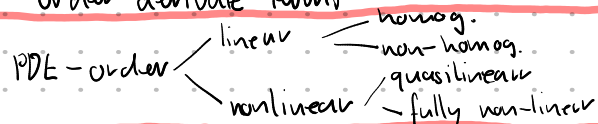
Def (order) order of PDE = highest order of partial deriv.

Def (linear) PDE is of form
 $a^{(0)}u + \sum_{i=1}^n a^{(1)}_i u_{x_i} + \sum_{i,j=1}^n a^{(2)}_{ij} u_{x_i x_j} + \dots$
 $=: \mathcal{L}(u) = f(x)$

Def (lin) homogeneous homog. $f(x) = 0$
 \hookrightarrow linear in "

Theorem: PDE $\mathcal{L}[u] = f(x)$ u_1, u_2 solutions \Rightarrow
 (Superposition) $\alpha u_1 + \beta u_2$ sol of $\mathcal{L}[u] = 0$ if $\mathcal{L}[u] = 0$
 $\alpha u_1 + \beta u_2 + u_p$ sol of $\mathcal{L}[u] = f(x)$

Def: (quasilinear) PDE is linear in its highest order derivative terms
 \hookrightarrow particular sol.



Def (Gradient & Laplacian) $u(x, y, z)$

$\nabla u := (u_{x_1}, u_{x_2}, u_{x_3})$ $\Delta u := u_{xx} + u_{yy} + u_{zz}$
 $(u: \mathbb{R}^n \rightarrow \mathbb{R})$ $(u: \mathbb{R}^n \rightarrow \mathbb{R})$
 (Hessian) $(H_f)_{ij} = \frac{\partial^2 f}{\partial x_i \partial x_j}$ (Divergence) $\text{div } v = \sum_{i=1}^n (v_i)_{x_i}$
 $(v: \mathbb{R}^n \rightarrow \mathbb{R}^n)$ $(v: \mathbb{R}^n \rightarrow \mathbb{R}^n)$

\Rightarrow Gradient/Diverge/Laplacian are linear

Method of Characteristics

Solves first order (quasilinear) PDEs
 \Rightarrow Form: $a(x, y, u) u_x + b(x, y, u) u_y = c(x, y, u)$
 $=: c(x, y, u) + c_1(x, y, u)$

1. Characteristic curve (PDE becomes ODE here)

$s \mapsto (x_0(s), y_0(s)) \Leftrightarrow u(x_0(s), y_0(s)) = \tilde{u}_0(s)$
 $\Leftrightarrow \Gamma(s) = (x_0(s), y_0(s), \tilde{u}_0(s))$

2. ODE system

$\begin{cases} \frac{dx_0}{ds} = a(x(s), y(s), \tilde{u}(s)) \rightarrow \text{dependent on } 1, 2 \\ \frac{dy_0}{ds} = b(x(s), y(s), \tilde{u}(s)) \\ \frac{d\tilde{u}}{ds} = c(x(s), y(s), \tilde{u}(s)) + c_1(x(s), y(s), \tilde{u}(s)) \end{cases}$
 $=: c(x, y, u)$

Initial conditions $x(0, s) = x_0(s)$ $y(0, s) = y_0(s)$
 $\tilde{u}(0, s) = \tilde{u}_0(s)$

3. $x(t, s), y(t, s) \Rightarrow s(x, y)$ $t(x, y)$
 4. Plug $s(x, y)$ & $t(x, y)$ into $\tilde{u}(s, t)$

Ex: Cauchy problem $\begin{cases} u_x + u_y = 1 \\ u(x, 0) = 2x^3 \end{cases}$

$\Rightarrow \Gamma(s) = (s, 0, 2s^3)$
 2. $x_t = 1$ $y_t = 1$ $\tilde{u}_t = 1$
 $x(0, s) = s$ $y(0, s) = 0$ $\tilde{u}(0, s) = 2s^3$
 $\Rightarrow x(t, s) = s+t$ $y(t, s) = t$ $\tilde{u}(t, s) = 2s^3 + t$
 3. $t = y$ $s = x - y$
 4. $u(x, y) = 2(x-y)^3 + y$

Not. (Transversality condition)

At $(0, s)$ $\det \begin{pmatrix} a(x_0(s), y_0(s), u_0(s)) & b(x_0(s), y_0(s), u_0(s)) \\ \frac{1}{ds} x_0(s) & \frac{1}{ds} y_0(s) \end{pmatrix} \neq 0$

\Leftrightarrow solution exists in neighborhood of $(x_0(s), y_0(s))$

Conservation laws & Shock waves

Def (Scalar conservation law)

$u: \mathbb{R} \times [0, \infty) \rightarrow \mathbb{R}$ $u_y + f(u)_x = 0$ f: flux
 \hookrightarrow spatial temp. $\Leftrightarrow u_y + c(u)u_x = 0$ $c(u) = f'(u)$

Ex (Transport equation)

$u_y + cu_x = 0$ $c \in \mathbb{R} \Rightarrow u(x, y) = g(x - cy)$
 $u(x, 0) = g(x)$

Prop. Implicit solution: $u(x, y) = u_0(x - c(u(x, y))y)$

- Solutions are constant along charact.
- (\tilde{u} depends only on s)
- Characteristics are straight lines
- admits strong local solution

Theorem: Scalar con. law with $c, u_0 \in C^1(\mathbb{R})$

Critical time:
 $y_c := \inf \left\{ -\frac{1}{c(u_0(s))s} : s \in \mathbb{R}, c(u_0(s))s < 0 \right\}$
 $(y_c = \infty \text{ if } c(u_0(s))s \geq 0 \forall s)$

If $y_c > 0$, then unique sol. for PDE in $[0, y_c]$ and u satisfies the imp. eq.

\hookrightarrow the first time two char. cross

Def Integral formulation no continuity needed

$\int_a^b u(x, y_2) dx - \int_a^b u(x, y_1) dx = - \int_{y_1}^{y_2} \int_a^b [f(u(b, y)) - f(u(a, y))] dy$

for all $a < b, y_1 < y_2$

Def (Weak solution) \hookrightarrow cont. diff. in each D.

$u(x, y)$ weak solution on $D = \hat{U} D$, if u

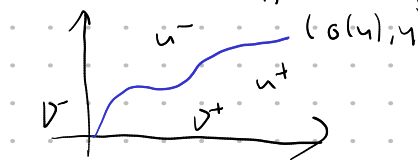
- satisfies original PDE on each D
- " integral form on D must satisfy R-H cond.

\Rightarrow boundaries between D are called shocks

(Rankine-Hugoniot condition)

$G'(u) = \frac{f^+ - f^-}{u^+ - u^-}$ for shock wave $(G(y))$

$u^\pm = \lim_{x \rightarrow G(y)}^\pm u(x, y)$ $f^\pm = f(u^\pm)$ \hookrightarrow p.e. from intersecting shock waves



Entropy condition for weak solutions:

Shock wave $x=y(y)$ satisfies $c(u^+) < y < c(u^-)$

All characteristics enter shock wave but do not emerge from it

1D Wave equation (hyperbolic / $\delta c > 0$)

Form: $u_{tt} - c^2 u_{xx} = 0$ $\alpha \mathbb{R}, x, t \in \mathbb{R} \times (0, \infty)$

General solution: $u(x,t) = F(x+ct) + G(x-ct)$

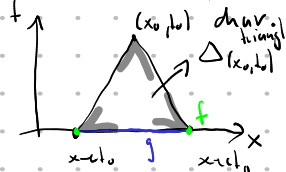
F: forward wave G: backward for $F, G \in C^2(\mathbb{R})$

Cauchy problem for homog. wave eq

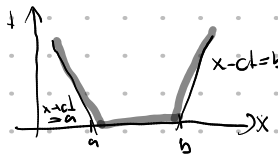
$$\begin{cases} u_{tt} - c^2 u_{xx} = 0 & (x,t) \in \mathbb{R} \times (0, \infty) \\ u(x,0) = f(x) \\ u_t(x,0) = g(x) \end{cases}$$

D'Alembert formula: $u(x,y) = \frac{f(x+ct) + f(x-ct)}{2} + \frac{1}{2c} \int_{x-ct}^{x+ct} g(y) dy$

Domain of dependence of (x_0, t_0)



Region of influence of $[a,b]$



Cauchy problem for inhomog. wave eq

$$\begin{cases} u_{tt} - c^2 u_{xx} = f(x,t) \\ u(x,0) = f(x) \\ u_t(x,0) = g(x) \end{cases}$$

Sol. (D'Alembert) unique

$$u(x,t) = \frac{f(x+ct) + f(x-ct)}{2} + \frac{1}{2c} \int_{x-ct}^{x+ct} g(y) dy + \frac{1}{2c} \int_0^t \int_{x-c(t-\tau)}^{x+c(t-\tau)} f(\xi, \tau) d\xi d\tau$$

Option F simple (depending only on x or t)

1. find part. sol $v(x,t)$ solving $v_{tt} - c^2 v_{xx} = F(x,t)$

2. Use D'Alembert (homog.) to find $w(x,t)$ with

$$\begin{cases} w_{tt} - c^2 w_{xx} = 0 \\ w(x,0) = f(x) - v(x,0) \\ w_t(x,0) = g(x) - v_t(x,0) \end{cases}$$

3. $u = v + w$

Prop: 1. Singularities of f, g get passed to u & propagated along char.
2. f, g even/odd/periodic $\Rightarrow u$ e.o.p.

Shaw with Fourier series

By uniqueness of sol, show that e.g. $-u(-x, t)$ also solves system

Separation of Variables

Homogeneous:

$$u(x,t) = X(x) \cdot T(t) \quad (x,t) \in (0,L) \times (0,\infty)$$

1. Equation: Heat: $u_t - k u_{xx} = 0$ $k \in \mathbb{R}$ parabolic
Wave: $u_{tt} - c^2 u_{xx} = 0$ $c \in \mathbb{R}$ hyperbolic

Boundary conditions

- Dirichlet: $u(0,t) = u(L,t) = 0$

- Neumann: $u_x(0,t) = u_x(L,t) = 0$

- Mixed/Robin: $\alpha_0 u(0,t) + \beta_0 u_x(0,t) = \gamma_0$
 $\alpha_L u(L,t) + \beta_L u_x(L,t) = \gamma_L$

Initial conditions: $u(x,0) = f(x)$
 $u_t(x,0) = g(x)$ Wave

2. Formulate ODEs

Heat: $T'(t) X(x) - k X''(x) T(t) = 0$

$$\Rightarrow \frac{T'(t)}{k T(t)} = \frac{X''(x)}{X(x)} = -\lambda \Rightarrow \begin{cases} X''(x) = -\lambda X(x) \\ T'(t) = -k \lambda T(t) \end{cases}$$

Wave: $X(x) T''(t) = c^2 X''(x) T(t) = 0$

$$\Rightarrow \frac{T''(t)}{c^2 T(t)} = \frac{X''(x)}{X(x)} = -\lambda \Rightarrow \begin{cases} X''(x) = -\lambda X(x) \\ T''(t) = -c^2 \lambda T(t) \end{cases}$$

3. Solve $X''(x) = -\lambda X(x)$

$$X(x) = \begin{cases} \alpha \cosh(\sqrt{\lambda} x) + \beta \sinh(\sqrt{\lambda} x) & \lambda > 0 \\ \alpha \cos(\sqrt{\lambda} x) + \beta \sin(\sqrt{\lambda} x) & \lambda < 0 \end{cases}$$

For both: With BC (other option: $T(t) = 0$ if $\lambda = 0$)

$\lambda < 0$: trivial sol.

$\lambda = 0$: Dirichlet: $X_n(x) = 0$ Neumann: $X_n(x) = \alpha_n$

$\lambda > 0$: Dirichlet: $X_n(x) = \beta_n \sin(\sqrt{\lambda_n} x)$ $\lambda_n = (\frac{n\pi}{L})^2$

Neumann: $X_n(x) = \beta_n \cos(\sqrt{\lambda_n} x)$

4. Solve $T(t)$

Heat: General sol: $T_n(t) = \beta_n e^{-k \lambda_n t}$

Wave: $T_n(t) = \gamma_n \cos(ct \sqrt{\lambda_n}) + \delta_n \sin(ct \sqrt{\lambda_n})$

For boundaries at a, b :

$$a_n \sinh(\sqrt{\lambda_n} x) + b_n \cosh(\sqrt{\lambda_n} x)$$

$$\Leftrightarrow c_n \sinh(\sqrt{\lambda_n} (x-a)) + d_n \sinh(\sqrt{\lambda_n} (x-b))$$

Second order linear PDEs

Form: $\Delta[u] = a u_{xx} + 2b u_{xy} + c u_{yy} + d u_x + e u_y + f u = g$

leading term/principal part

Def $\delta(L)(x_0, y_0) = b^2(x_0, y_0) - a(x_0, y_0) c(x_0, y_0)$

hyperbolic	$\delta(L)(x_0, y_0) > 0$	ex. $u_{tt} - u_{xx} = 0$ (wave)
parabolic	" $= 0$	$u_t - u_{xx} = 0$ (heat)
elliptic	" < 0	$u_{xx} + u_{yy} = 0$ (Laplace)

5. Putting together:

• Heat - Dirichlet: $u(x,t) = \sum_{n=1}^{\infty} C_n \sin(\sqrt{\lambda_n} x) e^{-k \lambda_n t}$

$$C_n = \frac{2}{L} \int_0^L f(x) \sin(\sqrt{\lambda_n} x) dx$$

Heat - Neumann: $u(x,t) = \left(\frac{A_0}{2} + \sum_{n=1}^{\infty} C_n \cos(\sqrt{\lambda_n} x) \right) e^{-k \lambda_n t}$

$$C_n = \frac{2}{L} \int_0^L f(x) \cos(\sqrt{\lambda_n} x) dx \quad \text{on average}$$

• Wave - Dirichlet:

$$u(x,t) = \sum_{n=1}^{\infty} \sin(\sqrt{\lambda_n} x) (A_n \cos(\sqrt{\lambda_n} t) + B_n \sin(\sqrt{\lambda_n} t))$$

$$A_n = \frac{2}{L} \int_0^L f(x) \sin(\sqrt{\lambda_n} x) dx \quad B_n = \frac{2}{c \sqrt{\lambda_n}} \int_0^L g(x) \sin(\sqrt{\lambda_n} x) dx$$

• Wave - Neumann:

$$u(x,t) = \frac{A_0 + B_0 t}{2} + \sum_{n=1}^{\infty} \cos(\sqrt{\lambda_n} x) (A_n \cos(\sqrt{\lambda_n} t) + B_n \sin(\sqrt{\lambda_n} t))$$

$$A_n = \frac{2}{L} \int_0^L f(x) \cos(\sqrt{\lambda_n} x) dx \quad B_n = \frac{2}{c \sqrt{\lambda_n}} \int_0^L g(x) \cos(\sqrt{\lambda_n} x) dx$$

Inhomogeneous

1. Find $X(x)$ like for homog
2. Write $u(x,t) = \sum_{n=1}^{\infty} T_n(t) X_n(x)$
 \hookrightarrow determine T_n in 1.
3. Plug into inhom. PDE
4. If necessary, expand $h(x,t)$ as Fourier series
5. Compare coefficient in 3 & 4 to get ODEs for $T_n(t)$
6. Use init. cond. of PDE to get init. coeffs for ODEs:
 $f(x) = u(x,0) = \sum_{n=1}^{\infty} T_n(0) X_n(x)$
7. Solve for $T_n(t)$, plug into $u(x,t)$

Fourier

• Odd: $f(x) = \sum_{n=1}^{\infty} b_n \sin\left(\frac{n\pi}{L} x\right)$ 2L-periodic on $[-L, L]$
 $\Rightarrow b_n = \frac{2}{L} \int_0^L f(x) \sin\left(\frac{n\pi}{L} x\right) dx$

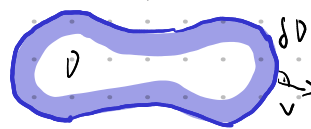
Even: $f(x) = \frac{b_0}{2} + \sum_{n=1}^{\infty} b_n \cos\left(\frac{n\pi}{L} x\right)$ 2L-periodic on $[-L, L]$
 $\Rightarrow b_n = \frac{2}{L} \int_0^L f(x) \cos\left(\frac{n\pi}{L} x\right) dx$
 $n=0 \Rightarrow 1$

Elliptic equations $\delta(x)(x_0, y_0) < 0$

Ex. Laplace's eq: $u_{xx} + u_{yy} = 0$

Poisson's eq: $u_{xx} + u_{yy} = f(x)$

• $D \subset \mathbb{R}^2$ open set, ∂D its boundary, ν outward normal to D , $d_\nu u(x,y) = \vec{\nu}(x,y) \cdot \nabla u(x,y)$



Def. (Dirichlet problem for Poisson's eq.)

$$\begin{cases} \Delta u(x,y) = p(x,y) & (x,y) \in D \\ u(x,y) = g(x,y) & (x,y) \in \partial D \end{cases}$$

$$(Neumann problem for Poisson's eq.)$$

$$\begin{cases} \Delta u(x,y) = p(x,y) & (x,y) \in D \\ d_\nu u = g & (x,y) \in \partial D \end{cases}$$

$$(Problem of third kind)$$

$$\begin{cases} \Delta u(x,y) = p(x,y) & (x,y) \in D \\ \alpha u + \beta d_\nu u = g & (x,y) \in \partial D \end{cases}$$

$$(Problem of third kind)$$

$$\begin{cases} \Delta u(x,y) = p(x,y) & (x,y) \in D \\ \alpha u + \beta d_\nu u = g & (x,y) \in \partial D \end{cases}$$

Thm. • Solution for Poisson/Laplace with Neumann cond exists if:

$$\int_{\partial D} g(x(s), y(s)) ds = \int_D p(x,y) dx dy$$

• Solution of Laplace's eq must have $\int_{\partial D} d_\nu u ds = 0$

A open subset of D , ν : outward unit of A

Def. (harmonic) $u(x,y)$ is harmonic if it solves Laplace's eq.

Ex. $e^x \sin(y)$, $\sinh(x) \cos(y)$, $\ln(x^2 + y^2)$ on $\mathbb{R}^2 \setminus \{0\}$
 c , $ax + by + c$, xy , $x^2 - y^2$

Maximum Principles

Theorem (Weak maximum) D bounded

$u \in C^2(D) \cap C(\bar{D})$ harmonic $\hookrightarrow \int_{\partial D} d_\nu u \leq 0$

• $\max_D u = \max_{\partial D} u$ Maximum is achieved on the boundary

• $\min_D u = \min_{\partial D} u$

Theorem (Mean value) Disk $B_R(x_0, y_0) \subset D$

u harmonic on D \hookrightarrow with radius R

$$u(x_0, y_0) = \frac{1}{2\pi R} \int_{\partial B_R(x_0, y_0)} u(x(s), y(s)) ds$$

Value of u at x_0, y_0 av. to value on circle around

$$= \frac{1}{2\pi} \int_0^{2\pi} u(x_0 + R \cos \theta, y_0 + R \sin \theta) d\theta$$

Theorem: (Strong maximum principle)

D connected, u harmonic. If u attains its maximum (or minimum) at an interior point, then u is constant

Theorem: (Uniqueness of Poisson equation) D bounded.

$$\begin{cases} \Delta u = f & \text{in } D \\ u = g & \text{on } \partial D \end{cases} \text{ has at most one solution.}$$

Theorem: D bounded

u_1, u_2 solve $\Delta u_1 = 0$, $\Delta u_2 = 0$ with $u_1 = g_1$, $u_2 = g_2$ on ∂D resp. $\Rightarrow \max_D |u_1 - u_2| = \max_{\partial D} |g_1 - g_2|$

Theorem (Maximum Principle for heat eq)

Domain $Q_T = [0, T] \times D$, D bounded

Parabolic boundary: $\partial_p Q_T = \{ \{0\} \times D \} \cup \{ [0, T] \times \partial D \}$

\hookrightarrow Boundary of except top corner $\{T\} \times D$

\Rightarrow If u solves $u_t = k \Delta u$ in Q_T for $k > 0$. Then

$$\max_{Q_T} u = \max_{\partial_p Q_T} u \quad \text{and} \quad \min_{Q_T} u = \min_{\partial_p Q_T} u$$

