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## Applied Numerical Methods - Lab 3

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### 0.1 Stationary heat conduction in 1-D

In a one dimensional pipe we are interested in the temperature evolution along the  $z$ -axis. We will study the behaviour of the numerical solution based on finite differences.

#### 0.1.1 Distretization to a linear system of equations

The  $z$ -axis is discretization in  $N + 1$  points spreading from 0 to  $L$ . the differential equation for the temperature is the following

$$-\kappa \frac{d^2 T}{dz^2} + v\rho C \frac{dT}{dz} = Q(z)$$

with

$$Q(z) = \begin{cases} 0 & \text{if } 0 \leq z < a \\ Q_0 \sin\left(\frac{z-a}{b-a}\pi\right) & \text{if } a \leq z \leq b \\ 0 & \text{if } b \leq z \leq L. \end{cases}$$

We assumed that  $\kappa$  as a constant value through the pipe.

The boundary conditions are

$$T(0) = T_0$$

and

$$-\kappa \frac{dT}{dz}(L) = k_v(T(L) - T_{out}).$$

For the discretization, let  $T_0, T_1, \dots, T_N$  be the unknown temperature we will solve for. The step is  $h = L/N$ . Define  $z_i = i * h$  and  $T_i \approx T(z_i)$ . It is clear that  $z_0 = 0$  and  $x_N = L$ . We also define the ghost point  $x_{N+1} = L + h$  that will be used implicitly to write the system. The real unknowns we will include in the system are  $T_1, \dots, T_N$ .  $T_0$  is known from the start.

The figure 1 illustrates the discretization and the usage of ghost points depending on the boundary conditions.

With finite difference approximation we rewrite the problem as

$$-\kappa \frac{T_{i+1} - 2T_i + T_{i-1}}{h^2} + v\rho C \frac{T_{i+1} - T_{i-1}}{2h} = Q(z_i),$$

for all the points,  $i = 1, \dots, N$ . This is equivalent to

$$\underbrace{\left(-\frac{v\rho Ch}{2} - \kappa\right)}_{\alpha} T_{i-1} + \underbrace{(2\kappa)}_{\beta} T_i + \underbrace{\left(\frac{v\rho Ch}{2} - \kappa\right)}_{\gamma} T_{i+1} = h^2 Q(z_i),$$

As for the boundary condition we have the straightforward

$$T_0 = 400,$$

Figure 1: Discretization and ghost points.

as well as

$$-\kappa \frac{T_{N+1} - T_{N-1}}{h} = k_v(T_N - T_{out}).$$

This allows to express  $T_{N+1}$  with other variables and remove it from the system of equations.

$$T_{N+1} = T_{N-1} - \frac{2hk_v}{k}T_N + \frac{2hk_v}{k}T_{out}.$$

This is a linear system of  $N$  equations. As we would like to avoid dividing by the small value  $h^2$ , we rewrite the system as follows. This system is tridiagonal and the Matlab *band solver* will be very (very) efficient if the matrix is defined as sparse.

### 0.1.2 Convergence of solution without convection

We now set  $v = 0$  and solve the equation for increasing values of  $N$ . Results on figure 2.

We see that increasing  $N$  has a great effect on the quality of the solution. Furthermore the solution corresponds visually well to the impulse  $Q(z)$ .

### 0.1.3 Increasing speed for $N = 40$

Let us set  $N = 40$ . We will now have  $v$  vary and take the values 0.1, 0.5, 1, 10. Results on figure 3.

We clearly notice that oscillations occur when  $v = 10$ .

### 0.1.4 Increasing $N$ for $v = 10$

As we want to see better how the oscillation appears, we increase the precision with  $N$  for  $v$  fixed at 10. Results on figure 4.

As the level of the discretization increases with  $N$ , we see more and more oscillation.

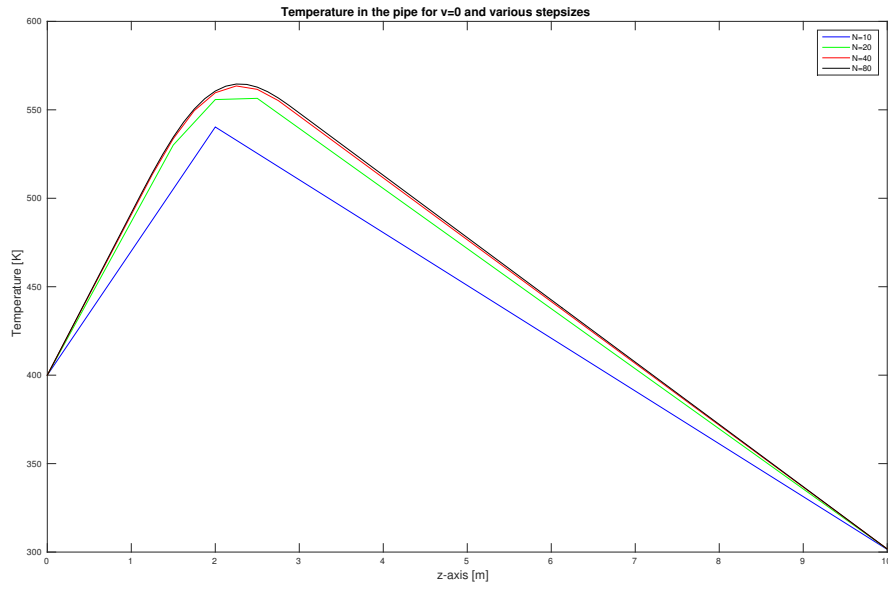


Figure 2: Convergence of solution without convection.

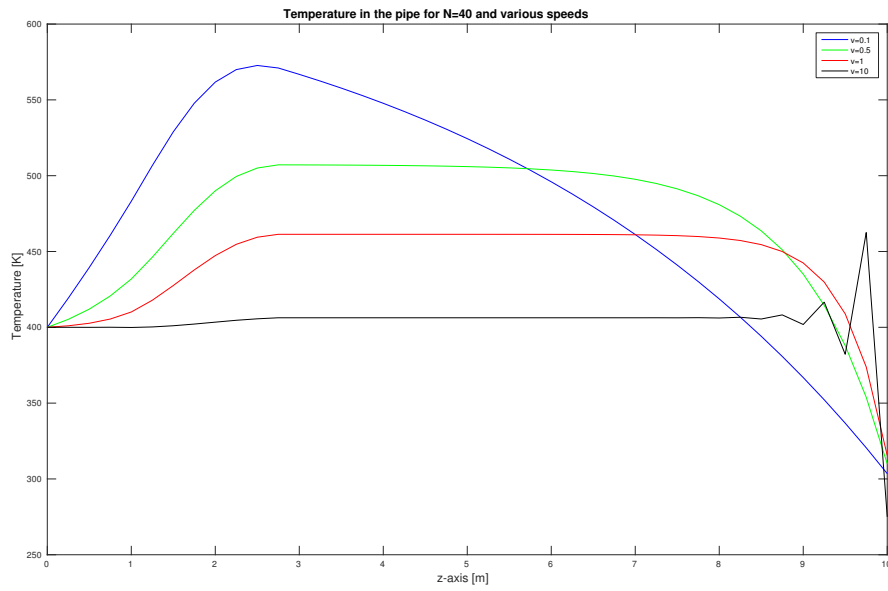


Figure 3: Solution with increasing speed  $v$ .

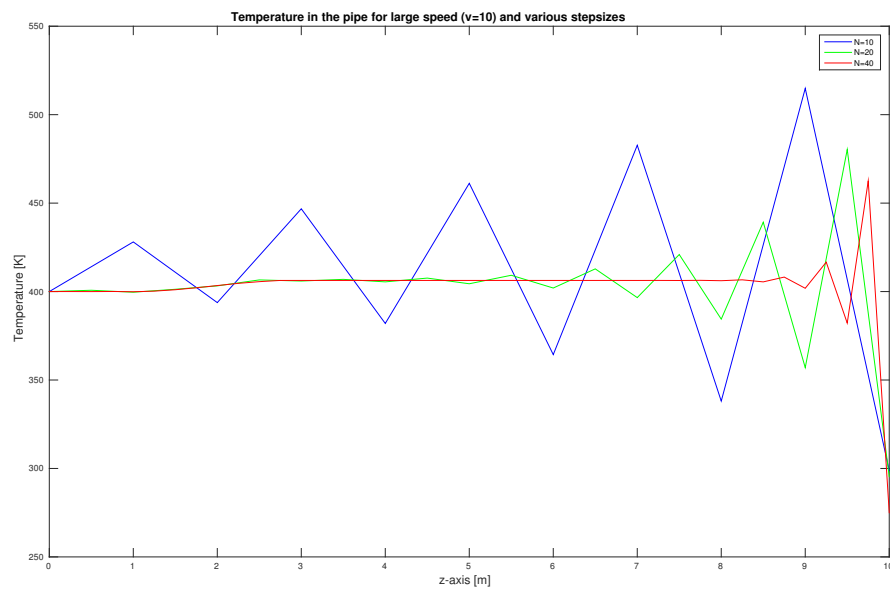


Figure 4: Capture the oscillation with increasing  $N$ .