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QFT - Lecture 15

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Contents

| | | |
|----------|--|----------|
| 1 | Correlation functions | 4 |
| 2 | Wick's Theorem | 4 |
| 2.0.1 | Summary: | 5 |
| 3 | Exponentiation and cancellation | 6 |

1 Correlation functions

$$\langle \Omega | T \phi(x) \phi(y) | \Omega \rangle = \langle 0 | T \phi(x) \phi(y) e^{-i \int dt H_I(t)} | 0 \rangle \quad (1)$$

with $\phi(x)$, $\phi(y)$ being in the Heisenberg picture on the left and in the Interaction picture on the right.

$$H_I(t) = \int \frac{\lambda}{4!} \phi^4 d^3x$$

with ϕ^4 in the interaction picture.

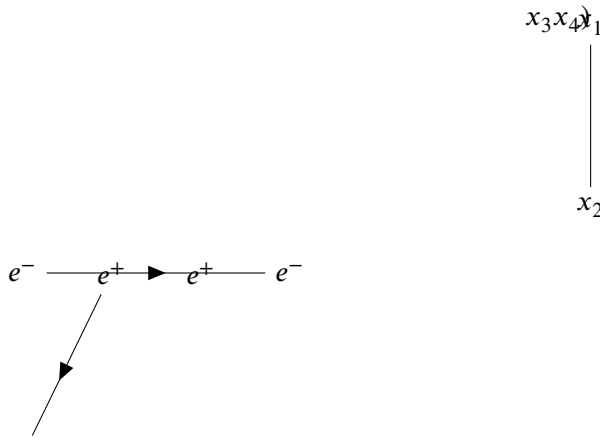
2 Wick's Theorem

$$T \phi_1 \cdots \phi_n = N(\phi_1 \cdots \phi_n + \text{all possible contractions})$$

For example:

$$\langle 0 | T \phi_1 \phi_2 \phi_3 \phi_4 | 0 \rangle = \overbrace{\phi_1 \phi_2 \phi_3 \phi_4} + \overbrace{\phi_1 \phi_2 \phi_3 \phi_4} + \overbrace{\phi_1 \phi_2 \phi_3 \phi_4} \quad (2)$$

Which can be represented as the Feynman Diagrams



Eq1 :

lowest order: Num: $\langle 0 | T \phi(x) \phi(y) | 0 \rangle = D_F(x - y) = \text{Fig2}$

first order: Num: $\langle 0 | T \phi(x) \phi(y) | 0 \rangle (-i\lambda/4!) \int d^4z \phi(z)^4 \langle 0 |$

We now have 6 fields and must find all contractions. We can contract $\phi(x)$ and $\phi(y)$ with $\phi(z)$, which leads to $4 \cdot 3 = 12$ contractions. We can also contract $\phi(x)$ with $\phi(y)$ This leads to 3 contractions.

Which leads to

$$\frac{-i\lambda}{4!} \int d^4z 3 D_F(x - y) D_F(z - z) D_F(z - z) + 12 D_F(x - z) D_F(y - z) D_F(z - z) \quad (4)$$

With the Feynman Propagator

$$D_F(x - y) = \frac{d^4p}{(2\pi)^4} \frac{i}{p^2 + m^2 + i\epsilon} e^{-ip(x-y)} \quad (5)$$

We can see that this integral becomes infinite for $x = y$, and integrating over all space again afterwards will definitely lead to infinity. We will ignore this for now, the divergences will disappear later.

For the 3 contractions we have the following Feynman Diagrams: Fig3

For the 12 contractions we have the following Feynman Diagrams: Fig4

Now considering a more complicated situation, with the third order:

$$\langle 0 | T \phi(x) \phi(y) \frac{1}{3!} \left(\frac{-i\lambda}{4!} \right)^3 \int d^4 z \phi^4 \int d^4 w \phi^4 \int d^4 u \phi(u) | 0 \rangle \quad (6)$$

We will consider one particular contraction, that being With the Feynmann Diagram Fig5

How many contractions give the same expression? We get $3!$ from interchanging verteces, $4 \cdot 3$ from the placement of contractions in the z vertex, we then get $4 \cdot 3 \cdot 2$ placements of the w vertex, and $4 \cdot 3$ from the u vertex. We have overcounted by a factor of 2, as we can interchange the contraction between the w and u fields, so we need to divide by 2.

we ignore $\frac{1}{3!}$ and we ignore $\frac{1}{4!} \cdot 4 \cdot 3$, $\frac{1}{4!}$. We then divide by the so called symmetry factor s , in our case $s = 2 \cdot 2 \cdot 2$. The first 2 comes from the line starting and ending at the z vertex, the 2nd from the line starting and ending at the u vertex, and the last because we can interchange the wu lines.

Another possibility of getting a symmetry factor is if two vertices are equivalent. Possible symmetry factors are as following

- line starting and ending at the same vertex
- equivalent lines
- equivalent vertices

These Symmetry factors are the reason why we should divide by the symmetry factor instead of leaving $\frac{1}{3!}$ and $\frac{1}{4!}$ in the third order term.

2.0.1 Summary:

$$\langle \Omega | T \phi(x) \phi(y) | \Omega \rangle = \frac{\text{Num}}{\text{Den}} \quad (7)$$

$$\text{Num: } \sum \text{all possible diagrams with two external points} \quad (8)$$

Feynman rules: Figs7

Fig8

Momentum Space Feynman Rules

Momentum Space Feynman Rules, correlation functions to Fourier space.

$$\langle \Omega | T \phi(x) \phi(y) | \Omega \rangle \mapsto \int d^4x e^{ipx} \int d^4y e^{ipy} \langle \Omega | T \phi(x) \phi(y) | \Omega \rangle \quad (9)$$

Omit exponent factors of external points, and the associated $\int \frac{d^4p}{(2\pi)^4}$ Integrals.

$$f(x) \mapsto \hat{f}(p), \quad \hat{f}(p) \mapsto f(x), \quad f(x) = \int \frac{d^4p}{(2\pi)^4} \hat{f}(p) e^{-ipx} \quad (10)$$

That way we can get $\hat{f}(p)$ by itself.

3 Exponentiation and cancellation

Consider Fig9

Then we have a product of two delta functions, which is nonsense

Let $V_i \in \{\text{The set of all different diagrams disconnected from external points}\}$.

$$A = (\text{Value of connected piece}) \cdot \prod_i \frac{1}{n_i!} V_i^{n_i}, \quad n_i = \text{number of } V_i.$$

$$\text{Num} = \sum_{\text{all connected}} (\text{value connected}) \cdot \underbrace{\prod_i \sum_{n_i} \frac{1}{n_i!} V_i^{n_i}}_{\exp(\sum_i V_i)}$$

$$\text{Den} = \exp(\sum_i V_i)$$

and therefore what we would like to have is

$$\frac{\text{Num}}{\text{Den}} = \sum \text{all connected}$$