QFT - Lecture 25

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1 Renormalization of electric charge

Consider corrections to photon propagator

Vacuum polatization diagram: Fig1

$$= (-ie)^{2}(-1) \int \frac{d^{4}k}{(2\pi)^{4}} \operatorname{tr} \left[\gamma^{\mu} \frac{i(\cancel{k} + m)}{k^{2} - m^{2} + i\epsilon} \gamma^{\nu} \frac{i(\cancel{k} + \cancel{q} + m)}{(k+q)^{2} - m^{2} + i\epsilon} \right]$$
(1)
$$= i \Pi_{2}^{\mu\nu}(q)$$
(2)

where we have (-1) and the trace because of the fermion loop.

More generally: Fig 2

 $\Pi^{\mu\nu}(q)$ is a tensor. It is comprised of the tensors that are included in the calculation. Then

$$\Pi^{\mu\nu}(q) = Cq^{\mu}q^{\nu} + Dg^{\mu\nu}, \text{ by Lorentz invariance}$$
 (3)

We use Ward's Identity:

$$q_{\mu}\Pi^{\mu\nu} = 0 \tag{4}$$

Which we apply to eq3

$$\Rightarrow Cq^2q^{\nu} + Dq\nu = 0 \tag{5}$$

$$\Rightarrow (Cq^2 + D)q^{\nu} = 0 \tag{6}$$

$$\Rightarrow Cq^2 + D = 0 \tag{7}$$

$$\Rightarrow D = -Cq^2 \tag{8}$$

Which we substitute into the previous equation

$$\Pi^{\mu\nu}(q) = C(q^{\mu}q^{\nu} - q^2g^{\mu\nu} \tag{9}$$

and we redefine $-C = \Pi(q^2)$.

$$= (q^2 g^{\mu\nu} - q^{\mu} q^{\nu}) \Pi(q^2) \tag{10}$$

2 Sum of all correction to the poton propagator

denoted by Fig 3

$$= Fig 4 \tag{11}$$

Consider a low q^2 process $(q^2 \ll m^2)$:

Fig5

What if q^2 is not small?

$$e_0 \to \frac{e_0}{\sqrt{1 - \Pi(q^2)}} = e \frac{\sqrt{1 - \Pi(0)}}{\sqrt{1 - \Pi(q^2)}}$$
 (12)

The apparent charge is dependent on q^2 .

That means there is an Effective potential from an electron that is

$$V(r) = -\frac{\alpha}{r} \left(1 + \frac{\alpha}{4\sqrt{\pi}} \frac{e^{-2mr}}{(mr)^{3/2}} + \dots \right)$$
 (13)

Plotting this, with $V(r) = \frac{-\alpha_{eff}(r)}{r}$: Fig6
When the electrons are closer together, they have a stronger attraction. Fig7