

Study guide: Numerical solution of the Navier-Stokes equations

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The physical and mathematical problem

<http://www.youtube.com/embed/P8VcZzgdfSc>

<http://www.youtube.com/embed/sI2uCHH3qIM>

Lots of physical applications involve fluid flow

- Weather (flow in the atmosphere)
- Ocean currents
- Flight
- Drag on cars
- Blood circulation
- Breathing

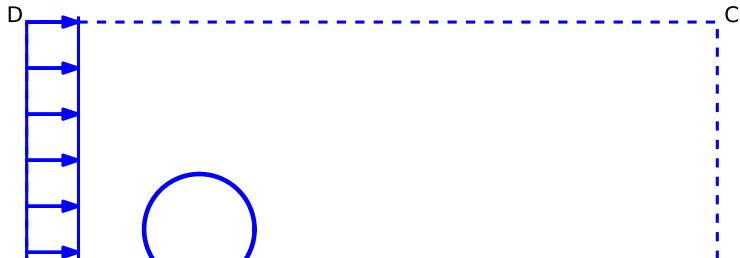
The physical assumptions behind the Navier-Stokes equations

Assumptions:

- Incompressible flow (velocity $< 1/3$ of the speed of sound)
- Laminar flow
- Simple fluids (constant viscosity ν)

Primary unknowns:

- velocity $\mathbf{u}(\mathbf{x}, t)$
- pressure $p(\mathbf{x}, t)$



The Navier-Stokes equations

Momentum balance (Newton's 2nd law):

$$\mathbf{u}_t + (\mathbf{u} \cdot \nabla) \mathbf{u} = -\frac{1}{\rho} \nabla p + \nu \nabla^2 \mathbf{u} + \mathbf{f}$$

Mass balance (eq. of continuity):

$$\nabla \cdot \mathbf{u} = 0$$

Boundary conditions

- Dirichlet conditions: components of \mathbf{u} are known
- Neumann conditions:
 - Stress condition: components of the stress vector $\boldsymbol{\sigma} \cdot \mathbf{n}$ are prescribed
 - Outflow or symmetry condition: $\partial \mathbf{u} / \partial n = 0$ (or components of this vector are zero)
- Pressure known at *a single point*

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The classical splitting method

Idea: split the N-S equations into simpler problems (*operator splitting*).

A simple, naive approach

The equation for \mathbf{u} looks like a diffusion equation...why not a Forward Euler scheme?

$$\mathbf{u}_t + (\mathbf{u} \cdot \nabla) \mathbf{u} = -\frac{1}{\rho} \nabla p + \nu \nabla^2 \mathbf{u} + \mathbf{f}$$

$$\frac{\mathbf{u}^{n+1} - \mathbf{u}^n}{\Delta t} + (\mathbf{u}^n \cdot \nabla) \mathbf{u}^n = -\frac{1}{\rho} \nabla p^n + \nu \nabla^2 \mathbf{u}^n + \mathbf{f}^n$$

$$\mathbf{u}^{n+1} = \mathbf{u}^n - \Delta t (\mathbf{u}^n \cdot \nabla) \mathbf{u}^n - \frac{\Delta t}{\rho} \nabla p^n + \Delta t \nu \nabla^2 \mathbf{u}^n + \Delta t \mathbf{f}^n$$

Two fundamental problems:

- 1 $\nabla \cdot \mathbf{u}^{n+1} \neq 0$ (that equation is not used!)
- 2 no computation of p^{n+1}

A working scheme

Idea: Forward Euler in time, but evaluate ∇p at t_{n+1} and enforce $\nabla \cdot \mathbf{u}^{n+1} = 0$.

$$\mathbf{u}^{n+1} = \mathbf{u}^n - \Delta t (\mathbf{u}^n \cdot \nabla) \mathbf{u}^n - \frac{\Delta t}{\varrho} \nabla p^{n+1} + \Delta t \nu \nabla^2 \mathbf{u}^n + \Delta t \mathbf{f}^n,$$

$$\nabla \cdot \mathbf{u}^{n+1} = 0$$

Note: *implicit system* for \mathbf{u}^{n+1} and p^{n+1}

We solve the implicit system by a splitting technique

- Use old $\beta \nabla p^n$ for ∇p^{n+1} and advance to intermediate velocity \mathbf{u}^*
- Correct the \mathbf{u}^* velocity by $\nabla \cdot \mathbf{u}^{n+1} = 0$

Intermediate velocity (Forward Euler):

$$\mathbf{u}^* = \mathbf{u}^n - \Delta t (\mathbf{u}^n \cdot \nabla) \mathbf{u}^n - \beta \frac{\Delta t}{\rho} \nabla p^n + \Delta t \nu \nabla^2 \mathbf{u}^n + \Delta t \mathbf{f}^n$$

Seek correction $\delta \mathbf{u}$ such that

$$\mathbf{u}^{n+1} = \mathbf{u}^* + \delta \mathbf{u}$$

fulfills

$$\nabla \cdot \mathbf{u}^{n+1} = 0$$

A Poisson equation must be solved to ensure $\nabla \cdot \mathbf{u} = 0$

Subtract \mathbf{u}^* equation from original \mathbf{u}^{n+1} equation to find $\delta \mathbf{u}$:

$$\delta \mathbf{u} = \mathbf{u}^{n+1} - \mathbf{u}^* = -\frac{\Delta t}{\rho} \nabla \Phi$$

where

$$\Phi = p^{n+1} - \beta p^n$$

The oldest methods had $\beta = 0$, but $\beta \neq 0$ gives in general better speed and accuracy.

$\nabla \cdot \mathbf{u}^{n+1} = 0$ implies

$$\nabla \cdot \delta \mathbf{u} = -\nabla \cdot \mathbf{u}^*$$

which gives

$$\nabla^2 \Phi = \frac{\rho}{\Delta t} \nabla \cdot \mathbf{u}^*$$

- 1 Compute the intermediate velocity \mathbf{u}^*
- 2 Solve the Poisson equation for Φ
- 3 Update the velocity: $\mathbf{u}^{n+1} = \mathbf{u}^* - \frac{\Delta t}{\rho} \nabla \Phi$
- 4 Update the pressure: $p^{n+1} = \Phi + \beta p^n$

Basically, we have $u = f$ approximation problems (1, 3, 4) and a Poisson equation to solve.

Boundary conditions

Problem: p condition at one point only in the original N-S equations. Now we need boundary conditions for Φ along the whole boundary (Poisson equation).

- Use conditions for \mathbf{u} also for \mathbf{u}^*
- Known pressure: known Φ
- Known pressure gradient: known $\partial\Phi/\partial n$
- Otherwise $\partial\Phi/\partial n = 0$

Spatial discretization by the finite element method

- $\mathbf{u}^*, \mathbf{u}^{n+1} \in V^{(u)}$ (modulo nonzero Dirichlet cond.)
- $p^{n+1} \in V^{(\Phi)}$ (modulo nonzero Dirichlet cond.)
- Test function $\mathbf{v}^{(u)} \in V^{(u)}$ for vector equations (velocity)
- Test function $v^{(\Phi)} \in V^{(\Phi)}$ for scalar equations (pressure)
- Take inner product of vector equation and $\mathbf{v}^{(u)}$
- Integrate $\nabla^2 \mathbf{u} \cdot \mathbf{v}^{(u)}$ by parts
- Integrate $\nabla p \cdot \mathbf{v}^{(u)}$ by parts (optional)
- Notation: \mathbf{u} is \mathbf{u}^{n+1} , \mathbf{u}_1 is \mathbf{u}^n , p is p^{n+1} , p_1 is p^n (as in code)

$$\begin{aligned} \int_{\Omega} (\mathbf{u}^* \cdot \mathbf{v}^{(u)} + \Delta t ((\mathbf{u}_1 \cdot \nabla) \nabla \mathbf{u}_1) \cdot \mathbf{v}^{(u)} - \frac{\Delta t}{\rho} p \nabla \cdot \mathbf{v}^{(u)} + \\ \Delta t \nu \nabla \mathbf{u}_1 \cdot \nabla \mathbf{v}^{(u)} - \Delta t f_1) \, dx + \int_{\partial \Omega_{N,u}} \left(\nu \frac{\partial \mathbf{u}}{\partial n} - p \mathbf{n} \right) \cdot \mathbf{v}^{(u)} \, ds \end{aligned} \quad (1)$$

$$\forall \mathbf{v}^{(u)} \in V^{(u)}.$$

Natural boundary condition:

Increasing the implicitness

Stability (due to Forward Euler-style scheme):

$$\Delta t \leq \frac{h^2}{2\nu + Uh}. \quad (5)$$

h : minimum element size, U : typical velocity.

Better stability by a Backward Euler scheme:

$$\mathbf{u}^{n+1} = \mathbf{u}^n - \Delta t (\mathbf{u}^{n+1} \cdot \nabla) \mathbf{u}^{n+1} - \frac{\Delta t}{\rho} \nabla p^{n+1} + \Delta t \nu \nabla^2 \mathbf{u}^{n+1} + \Delta t \mathbf{f}^n \quad (6)$$

$$\nabla \cdot \mathbf{u}^{n+1} = 0. \quad (7)$$

Intermediate velocity ($\nabla p^{n+1} \rightarrow \beta p^n$):

$$\mathbf{u}^* = \mathbf{u}^n - \Delta t (\mathbf{u}^* \cdot \nabla) \mathbf{u}^* - \beta \frac{\Delta t}{\rho} p^{n+1} + \Delta t \nu \nabla^2 \mathbf{u}^* + \Delta t \mathbf{f}^{n+1}$$

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Methods based on slight compressibility

$\nabla \cdot \mathbf{u} = 0$ is problematic. Allow slight compressibility in the fluid:

$$p_t + c^2 \nabla \cdot \mathbf{u} = 0.$$

c : speed of sound.

Now we have evolution equations for \mathbf{u} and p :

$$\mathbf{u}_t = -(\mathbf{u} \cdot \nabla) \mathbf{u} - \frac{1}{\rho} \nabla p + \nu \nabla^2 \mathbf{u} + \mathbf{f}, \quad (9)$$

$$p_t = -c^2 \nabla \cdot \mathbf{u}. \quad (10)$$

Forward Euler:

$$\mathbf{u}^{n+1} = \mathbf{u}^n - \Delta t (\mathbf{u}^n \cdot \nabla) \mathbf{u}^n - \frac{\Delta t}{\rho} \nabla p^n + \Delta t \nu \nabla^2 \mathbf{u}^n + \Delta t \mathbf{f}^n, \quad (11)$$

$$p^{n+1} = p^n - \Delta t c^2 \nabla \cdot \mathbf{u}^n. \quad (12)$$

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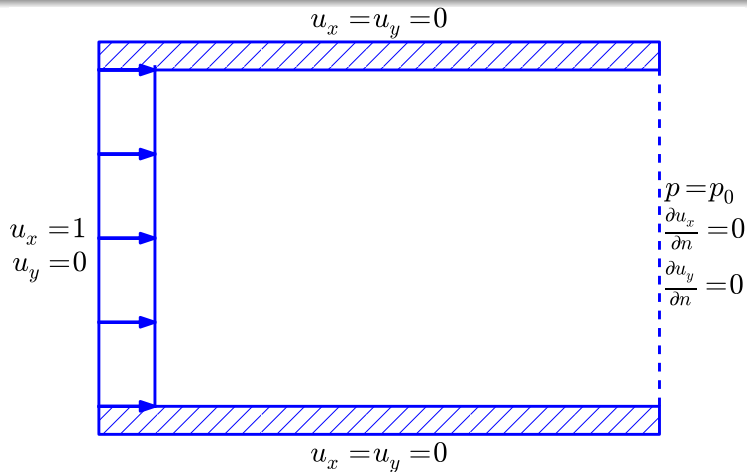


Figure : Flow in a channel.

$$u_y = 0, \frac{\partial u_x}{\partial n} = 0$$