Approximate relativistic methods

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Approximate relativistic methods

PROS:

- Approximate methods to include relativistic effects that avoid to use 2- and 4componend formalisms.
- Many variants (ZORA, IORA, DKH, NESC, X2C).
- They are very efficient.
- Spin-free one-electron versions only modify the AO basis one-electron integrals.



CONS:

- Need to use appropriate basis set (e.g. cc-pVDZ → cc-pVDZ-DK).
- Spin-orbit effects are missing.

Dirac equation

One-particle Dirac equation

$$\hat{D} \left(\begin{array}{c} \phi^L \\ \phi^S \end{array} \right) = E \left(\begin{array}{c} \phi^L \\ \phi^S \end{array} \right)$$

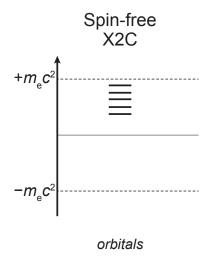
$$\hat{D} = \begin{pmatrix} \hat{V} & c (\vec{\sigma}. \vec{p}) \\ c (\vec{\sigma}. \vec{p}) & \hat{V} - 2mc^2 \end{pmatrix}$$

Dirac equation $+m_{e}c^{2}$ $-m_{e}c^{2}$ $-m_{e}c^{2}$

Exact two-component method (X2C)

$$\hat{U}^{\dagger}\hat{D}\hat{U} = \hat{U}^{\dagger} \begin{bmatrix} h_{\mathrm{LL}} & h_{\mathrm{LS}} \\ h_{\mathrm{SL}} & h_{\mathrm{SS}} \end{bmatrix} \hat{U} = \begin{bmatrix} h_{++}^{\mathrm{FW}} & 0 \\ 0 & h_{--}^{\mathrm{FW}} \end{bmatrix}$$

$$\hat{U}^{\dagger} \begin{pmatrix} \phi_{+}^{L} \\ \phi_{+}^{S} \end{pmatrix} = \begin{pmatrix} \phi^{FW} \\ 0 \end{pmatrix}$$



Computational Aspects

Solve the modified Dirac equation in a kinetically-balanced basis set

$$\phi_i^S = \sum_{\mu} C_{\mu i}^S \; \frac{\vec{\sigma} \cdot \vec{p}}{2mc} \chi_{\mu}$$

$$\phi_i^S = \sum_{\mu} C_{\mu i}^S \frac{\vec{\sigma} \cdot \vec{p}}{2mc} \chi_{\mu} \qquad \left(\begin{array}{c} V & T \\ T & \frac{W}{4m^2c^2} - T \end{array} \right) \left(\begin{array}{c} C^L \\ C^S \end{array} \right) = \left(\begin{array}{c} S & 0 \\ 0 & \frac{T}{2mc^2} \end{array} \right) \left(\begin{array}{c} C^L \\ C^S \end{array} \right) \epsilon$$

Need the following integrals

$$S_{\mu\nu} = \langle \chi_{\mu} | \chi_{\nu} \rangle$$

$$T_{\mu\nu} = \langle \chi_{\mu} | \frac{-\nabla^{2}}{2} | \chi_{\nu} \rangle$$

$$V_{\mu\nu} = \langle \chi_{\mu} | \sum_{A} \frac{Z_{A}}{|\vec{r} - \vec{R_{A}}|} | \chi_{\nu} \rangle$$

$$W_{\mu\nu} = \langle \chi_{\mu} | (\vec{\sigma} \cdot \vec{p}) V (\vec{\sigma} \cdot \vec{p}) | \chi_{\nu} \rangle$$

Spin-free version

$$W_{\mu\nu} = W_{\mu\nu}^{\rm SF} + V_{\mu\nu}^{\rm SF}$$

$$W_{\mu\nu}^{\rm SF} = \langle \chi_{\mu} | \vec{p} \cdot (V\vec{p}) | \chi_{\nu} \rangle$$

Implementation

The class RelPotentialInt computes the W integrals (no integral derivatives yet)

```
W_{\mu\nu}^{\rm SF} = \langle \chi_{\mu} | \vec{p} \cdot (V\vec{p}) | \chi_{\nu} \rangle = \langle \vec{p}\chi_{\mu} | V | \vec{p}\chi_{\nu} \rangle 
 /*! \ingroup MINTS  
   * \class RelPotentialInt  
   * \brief Computes relativistic potential integrals.  
   * Use an IntegralFactory to create this object. */ class RelPotentialInt : public OneBodyAOInt  
   {       /// Computes integrals between two shell objects.  
      void \emph{compute\_pair}(\text{const GaussianShell\&}, \text{ const GaussianShell\&});  
      /// Computes integrals between two shell objects.
```

Implementation

X2C integrals are computed by the MintsHelper class

```
void MintsHelper::integrals()
{
...

// Compute and dump one-electron SO integrals.
if (options_.get_str("RELATIVISTIC") == "NO") {
    // Overlap
    so_overlap()->save(psio_, PSIF_OEI);
    // Kinetic
    so_kinetic()->save(psio_, PSIF_OEI);
    // Potential
    so_potential()->save(psio_, PSIF_OEI);
}else if (options_.get_str("RELATIVISTIC") == "X2C") {
    outfile->Printf( " Using relativistic (X2C) overlap, kinetic,
and potential integrals.\n");
}
```

(there is another call, perhaps redundant, in one electron integrals())

Calling X2C

Calling X2C

```
set {
  basis cc-pVDZ-DK
  relativistic x2c
}
energy('x2c')
energy('ccsd(t)')
```