



Chapter 5

Linear Design Models

Equations of Motion

$$\dot{p}_n = (\cos \theta \cos \psi)u + (\sin \phi \sin \theta \cos \psi - \cos \phi \sin \psi)v + (\cos \phi \sin \theta \cos \psi + \sin \phi \sin \psi)w$$

$$\dot{p}_e = (\cos \theta \sin \psi)u + (\sin \phi \sin \theta \sin \psi + \cos \phi \cos \psi)v + (\cos \phi \sin \theta \sin \psi - \sin \phi \cos \psi)w$$

$$\dot{h} = u \sin \theta - v \sin \phi \cos \theta - w \cos \phi \cos \theta$$

$$\dot{u} = rv - qw - g \sin \theta + \frac{\rho V_a^2 S}{2m} \left[C_X(\alpha) + C_{X_q}(\alpha) \frac{cq}{2V_a} + C_{X_{\delta_e}}(\alpha) \delta_e \right] + \frac{\rho S_{\text{prop}} C_{\text{prop}}}{2m} \left[(k_{\text{motor}} \delta_t)^2 - V_a^2 \right]$$

$$\dot{v} = pw - ru + g \cos \theta \sin \phi + \frac{\rho V_a^2 S}{2m} \left[C_{Y_0} + C_{Y_\beta} \beta + C_{Y_p} \frac{bp}{2V_a} + C_{Y_r} \frac{br}{2V_a} + C_{Y_{\delta_a}} \delta_a + C_{Y_{\delta_r}} \delta_r \right]$$

$$\dot{w} = qu - pv + g \cos \theta \cos \phi + \frac{\rho V_a^2 S}{2m} \left[C_Z(\alpha) + C_{Z_q}(\alpha) \frac{cq}{2V_a} + C_{Z_{\delta_e}}(\alpha) \delta_e \right]$$

$$\dot{\phi} = p + q \sin \phi \tan \theta + r \cos \phi \tan \theta$$

$$\dot{\theta} = q \cos \phi - r \sin \phi$$

$$\dot{\psi} = q \sin \phi \sec \theta + r \cos \phi \sec \theta$$

$$\dot{p} = \Gamma_1 pq - \Gamma_2 qr + \frac{1}{2} \rho V_a^2 S b \left[C_{p_0} + C_{p_\beta} \beta + C_{p_p} \frac{bp}{2V_a} + C_{p_r} \frac{br}{2V_a} + C_{p_{\delta_a}} \delta_a + C_{p_{\delta_r}} \delta_r \right]$$

$$\dot{q} = \Gamma_5 pr - \Gamma_6 (p^2 - r^2) + \frac{\rho V_a^2 Sc}{2J_y} \left[C_{m_0} + C_{m_\alpha} \alpha + C_{m_q} \frac{cq}{2V_a} + C_{m_{\delta_e}} \delta_e \right]$$

$$\dot{r} = \Gamma_7 pq - \Gamma_1 qr + \frac{1}{2} \rho V_a^2 S b \left[C_{r_0} + C_{r_\beta} \beta + C_{r_p} \frac{bp}{2V_a} + C_{r_r} \frac{br}{2V_a} + C_{r_{\delta_a}} \delta_a + C_{r_{\delta_r}} \delta_r \right]$$

Lift and Drag Models

Nonlinear model with stall effects

$$C_L(\alpha) = (1 - \sigma(\alpha)) [C_{L_0} + C_{L_\alpha} \alpha] + \sigma(\alpha) [2 \operatorname{sign}(\alpha) \sin^2 \alpha \cos \alpha]$$

$$\sigma(\alpha) = \frac{1 + e^{-M(\alpha - \alpha_0)} + e^{M(\alpha + \alpha_0)}}{(1 + e^{-M(\alpha - \alpha_0)}) (1 + e^{M(\alpha + \alpha_0)})}$$

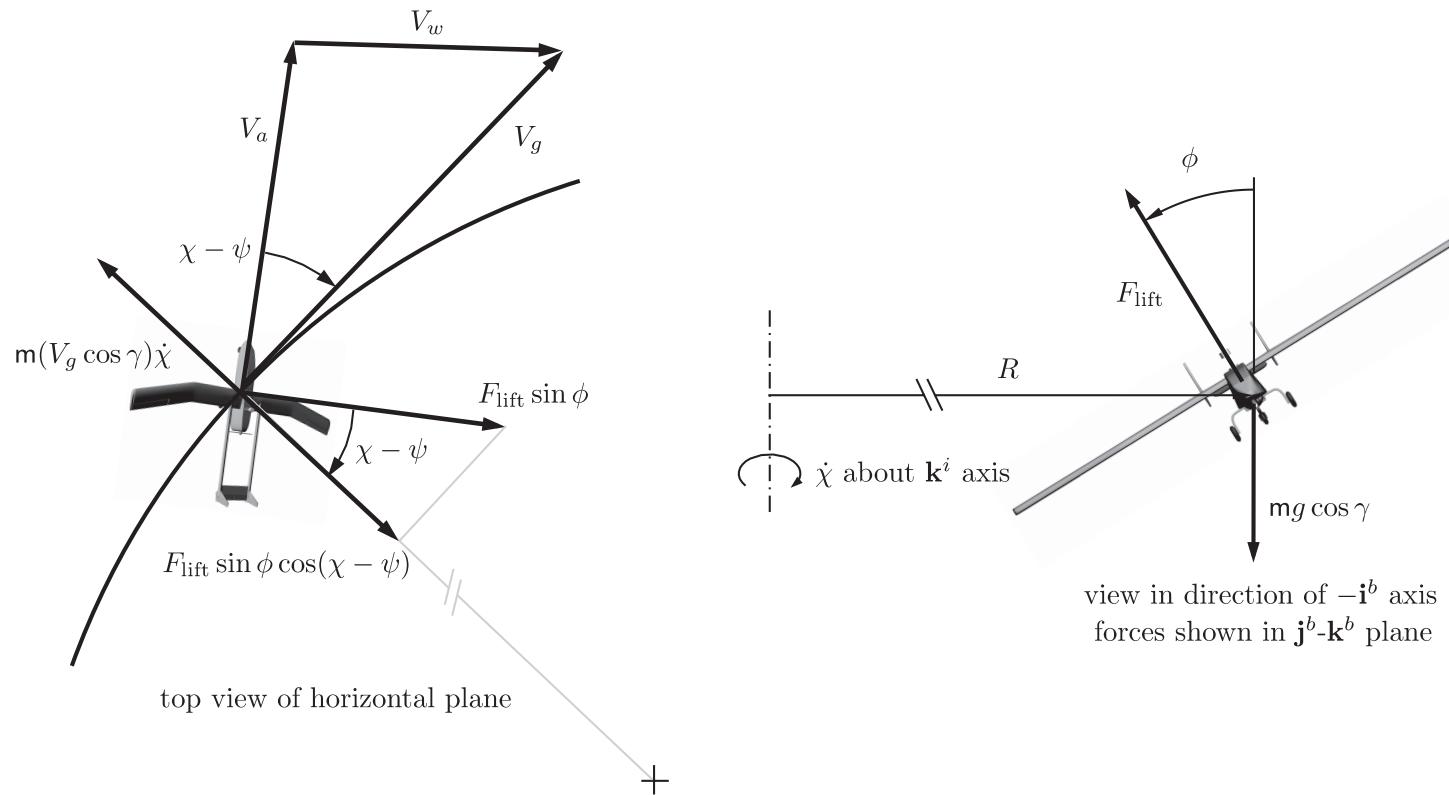
$$C_D(\alpha) = C_{D_p} + \frac{(C_{L_0} + C_{L_\alpha} \alpha)^2}{\pi e A R}$$

Linear Model

$$C_L(\alpha) = C_{L_0} + C_{L_\alpha} \alpha$$

$$C_D(\alpha) = C_{D_0} + C_{D_\alpha} \alpha$$

Coordinated Turn



$$\begin{aligned}
 F_{lift} \sin \phi \cos(\chi - \psi) &= m \frac{v^2}{R} \\
 &= mv\omega \\
 &= m(V_g \cos \gamma) \dot{\chi}
 \end{aligned}$$

$$F_{lift} \cos \phi = mg \cos \gamma$$

Coordinated Turn

Dividing the two expressions gives

$$\dot{\chi} = \frac{g}{V_g} \tan \phi \cos(\chi - \psi)$$

which is the coordinated turn condition in wind.

In the absence of wind we have $V_a = V_g$ and $\psi = \chi$ which gives

$$\dot{\psi} = \frac{g}{V_a} \tan \phi$$

which is the expression commonly seen in the literature.

The turning radius is given by

$$R = \frac{V_g \cos \gamma}{\dot{\chi}} = \frac{V_g^2 \cos \gamma}{g \tan \phi \cos(\chi - \psi)}.$$

For level flight in the absence of wind we have the standard formula

$$R = \frac{V_a^2}{g \tan \phi}.$$

Trim Conditions

Given the nonlinear systems

$$\dot{x} = f(x, u).$$

The equilibria (x^*, u^*) are defined by

$$f(x^*, u^*) = 0$$

An aircraft in equilibrium is in a trim condition. In general, trim condition may include states that are not constant. Therefore, trim conditions are given by

$$\dot{x}^* = f(x^*, u^*)$$

Calculating Trim

Objective is to compute trim states and inputs when aircraft simultaneously satisfies three conditions:

- Traveling at constant speed V_a^* ,
- Climbing at constant flight path angle γ^* ,
- In constant orbit of radius R^* .

V_a^* , γ^* , and R^* , are inputs to the trim calculations.

States: $x \triangleq (p_n, p_e, p_d, u, v, w, \phi, \theta, \psi, p, q, r)^\top$

Inputs: $u \triangleq (\delta_e, \delta_t, \delta_a, \delta_r)^\top$

Calculating Trim

For constant-climb orbit:

→ speed of aircraft not changing → $\dot{u}^* = \dot{v}^* = \dot{w}^* = 0$

→ roll and pitch angles constant → $\dot{\phi}^* = \dot{\theta}^* = \dot{p}^* = \dot{q}^* = 0$

Turn rate constant and given by

$$\dot{\psi}^* = \frac{V_a^*}{R^*} \cos \gamma^* \quad \rightarrow \quad \dot{r}^* = 0$$

Climb rate constant, and given by

$$\dot{h}^* = V_a^* \sin \gamma^*$$

Given parameters V_a^* , γ^* , and R^* , can specify \dot{x}^* as

$$\dot{x}^* = (\dot{p}_n^* \ \dot{p}_e^* \ \dot{h}^* \ \dot{u}^* \ \dot{v}^* \ \dot{w}^* \ \dot{\phi}^* \ \dot{\theta}^* \ \dot{\psi}^* \ \dot{p}^* \ \dot{q}^* \ \dot{r}^*)^\top$$

...continued...

Calculating Trim

Given parameters V_a^* , γ^* , and R^* , can specify \dot{x}^* as

$$\dot{x}^* = \begin{pmatrix} \dot{p}_n^* \\ \dot{p}_e^* \\ \dot{h}^* \\ \dot{u}^* \\ \dot{v}^* \\ \dot{w}^* \\ \dot{\phi}^* \\ \dot{\theta}^* \\ \dot{\psi}^* \\ \dot{p}^* \\ \dot{q}^* \\ \dot{r}^* \end{pmatrix} = \begin{pmatrix} [\text{don't care}] \\ [\text{don't care}] \\ V_a^* \sin \gamma^* \\ 0 \\ 0 \\ 0 \\ 0 \\ 0 \\ \frac{V_a^*}{R^*} \cos \gamma^* \\ 0 \\ 0 \\ 0 \end{pmatrix} = f(x^*, u^*)$$

Problem of finding x^* and u^* such that $\dot{x}^* = f(x^*, u^*)$, reduces to solving nonlinear algebraic systems of equations.

Can use Matlab trim command – see Appendix F

Lateral Transfer Functions - Roll

Objective: Find simple transfer function relationship between ϕ and δ_a

Start with: $\dot{\phi} = p + q \sin \phi \tan \theta + r \cos \phi \tan \theta$

Assuming that θ is small and defining the disturbance as $d_{\phi_1} \triangleq q \sin \phi \tan \theta + r \cos \phi \tan \theta$ gives

$$\dot{\phi} = p + d_{\phi_1}$$

Differentiating and substituting for $\dot{\phi}$ gives

$$\begin{aligned}\ddot{\phi} &= \dot{p} + \dot{d}_{\phi_1} \\ &= \Gamma_1 pq - \Gamma_2 qr + \frac{1}{2} \rho V_a^2 S b \left[C_{p_0} + C_{p_\beta} \beta + C_{p_p} \frac{bp}{2V_a} + C_{p_r} \frac{br}{2V_a} + C_{p_{\delta_a}} \delta_a + C_{p_{\delta_r}} \delta_r \right] + \dot{d}_{\phi_1} \\ &= \Gamma_1 pq - \Gamma_2 qr + \frac{1}{2} \rho V_a^2 S b \left[C_{p_0} + C_{p_\beta} \beta + C_{p_p} \frac{b}{2V_a} (\dot{\phi} - d_{\phi_1}) + C_{p_r} \frac{br}{2V_a} + C_{p_{\delta_a}} \delta_a + C_{p_{\delta_r}} \delta_r \right] + \dot{d}_{\phi_1} \\ &= \left(\frac{1}{2} \rho V_a^2 S b C_{p_p} \frac{b}{2V_a} \right) \dot{\phi} + \left(\frac{1}{2} \rho V_a^2 S b C_{p_{\delta_a}} \right) \delta_a \\ &\quad + \left\{ \Gamma_1 pq - \Gamma_2 qr + \frac{1}{2} \rho V_a^2 S b \left[C_{p_0} + C_{p_\beta} \beta - C_{p_p} \frac{b}{2V_a} (d_{\phi_1}) + C_{p_r} \frac{br}{2V_a} + C_{p_{\delta_r}} \delta_r \right] + \dot{d}_{\phi_1} \right\} \\ &= -a_{\phi_1} \dot{\phi} + a_{\phi_2} \delta_a + d_{\phi_2}\end{aligned}$$

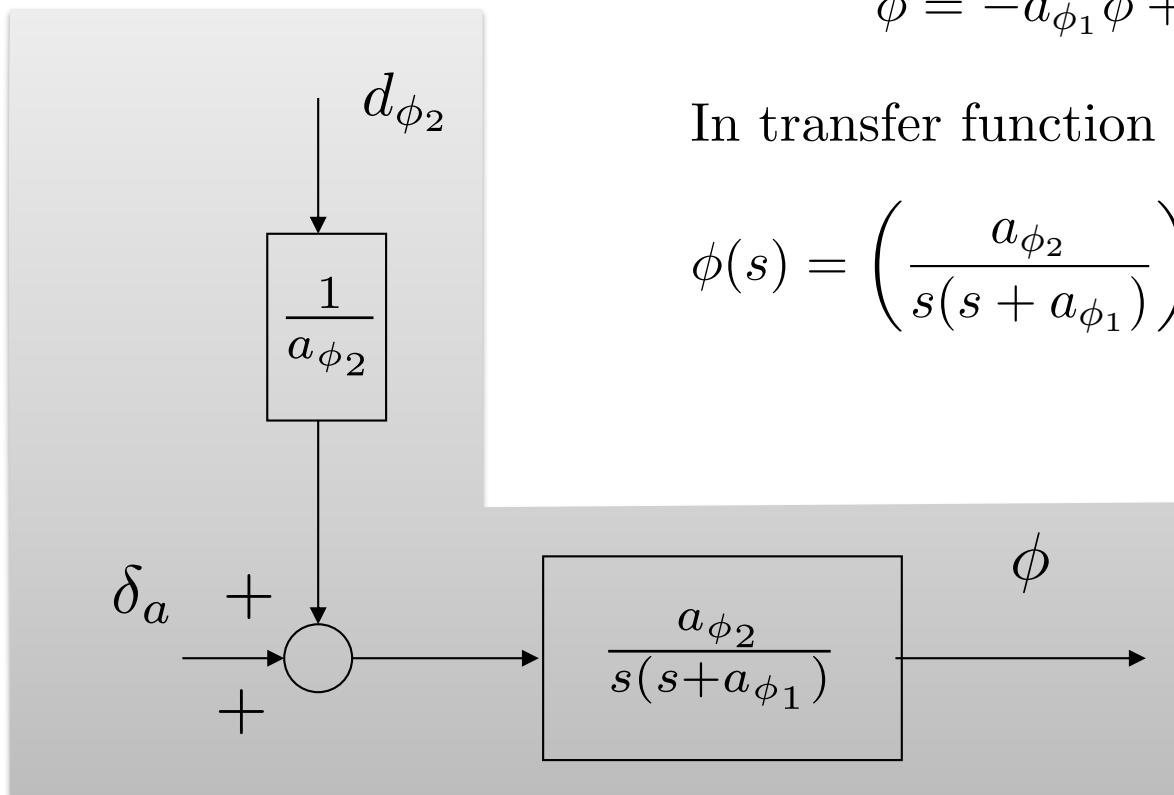
Roll Dynamics

Therefore, we have the second order differential equation

$$\ddot{\phi} = -a_{\phi_1}\dot{\phi} + a_{\phi_2}\delta_a + d_{\phi_2}.$$

In transfer function form we have

$$\phi(s) = \left(\frac{a_{\phi_2}}{s(s+a_{\phi_1})} \right) \left(\delta_a(s) + \frac{1}{a_{\phi_2}}d_{\phi_2}(s) \right).$$



Lateral Transfer Functions - Course

To derive a transfer function from roll ϕ to course χ , start with the coordinated turn condition

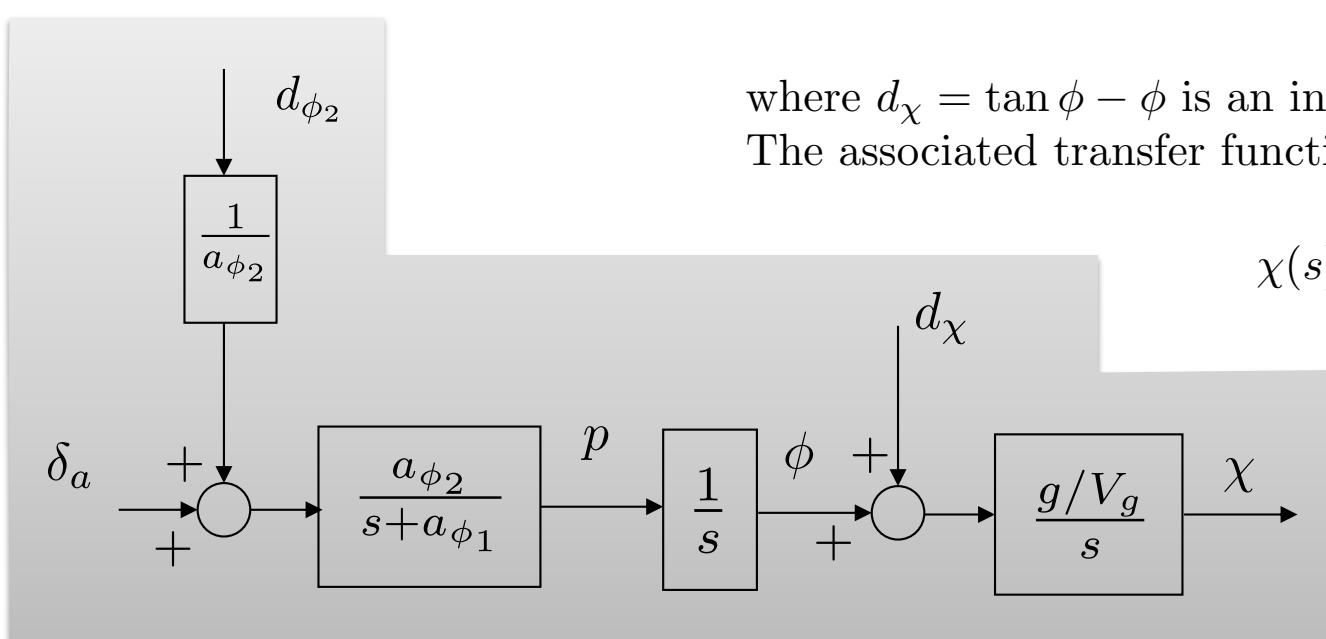
$$\dot{\chi} = \frac{g}{V_a} \tan \phi$$

add and subtract $\frac{g}{V_a}\phi$ to get

$$\begin{aligned}\dot{\chi} &= \frac{g}{V_a}\phi + \frac{g}{V_a}(\tan \phi - \phi) \\ &= \frac{g}{V_a}\phi + \frac{g}{V_a}d_\chi\end{aligned}$$

where $d_\chi = \tan \phi - \phi$ is an input disturbance.
The associated transfer function is

$$\chi(s) = \frac{g/V_a}{s} (\phi(s) + d_\chi(s))$$



Lateral Transfer Functions - Sideslip

To derive the transfer function from rudder to sideslip, note that in zero wind

$$v = V_a \sin \beta.$$

Differentiating and assuming constant airspeed gives

$$\dot{v} = (V_a \cos \beta) \dot{\beta}.$$

Substituting for \dot{v} from the dynamic gives

$$(V_a \cos \beta) \dot{\beta} = pw - ru + g \cos \theta \sin \phi \\ + \frac{\rho V_a^2 S}{2m} \left[C_{Y_0} + C_{Y_\beta} \beta + C_{Y_p} \frac{bp}{2V_a} + C_{Y_r} \frac{br}{2V_a} + C_{Y_{\delta_a}} \delta_a + C_{Y_{\delta_r}} \delta_r \right].$$

Solving for $\dot{\beta}$ and grouping terms gives:

Sideslip Dynamics

$$\dot{\beta} = -a_{\beta_1}\beta + a_{\beta_2}\delta_r + d_{\beta}$$

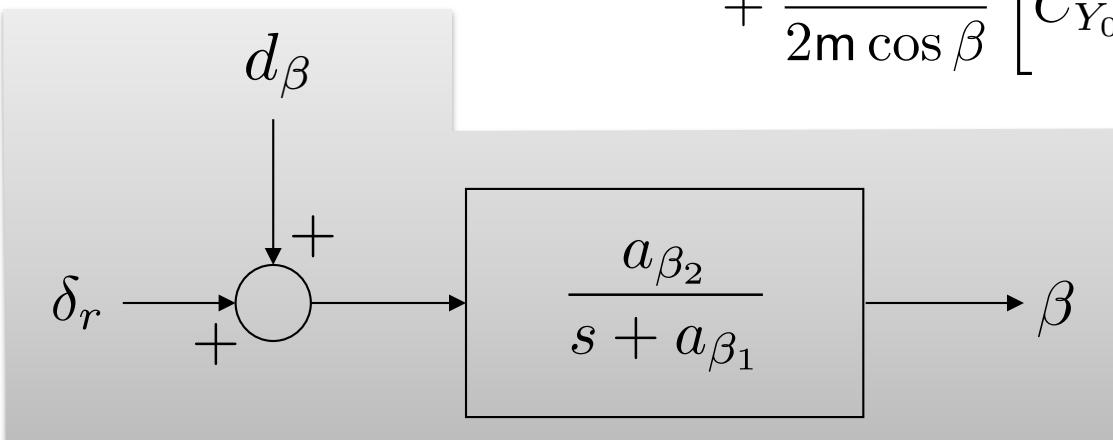
where

$$a_{\beta_1} = -\frac{\rho V_a S}{2m \cos \beta} C_{Y_{\beta}}$$

$$a_{\beta_2} = \frac{\rho V_a S}{2m \cos \beta} C_{Y_{\delta_r}}$$

$$d_{\beta} = \frac{1}{V_a \cos \beta} (pw - ru + g \cos \theta \sin \phi)$$

$$+ \frac{\rho V_a S}{2m \cos \beta} \left[C_{Y_0} + C_{Y_p} \frac{bp}{2V_a} + C_{Y_r} \frac{br}{2V_a} + C_{Y_{\delta_a}} \delta_a \right]$$



Longitudinal Transfer Functions - Pitch

To derive the transfer function from elevator to pitch angle:

$$\begin{aligned}\dot{\theta} &= q \cos \phi - r \sin \phi \\ &= q + q(\cos \phi - 1) - r \sin \phi \\ &\stackrel{\triangle}{=} q + d_{\theta_1}\end{aligned}$$

Differentiating and substituting \dot{q} from the dynamics gives

$$\begin{aligned}\ddot{\theta} &= \Gamma_6(r^2 - p^2) + \Gamma_5 pr + \frac{\rho V_a^2 c S}{2J_y} \left[C_{m_0} + C_{m_\alpha} \alpha + C_{m_q} \frac{cq}{2V_a} + C_{m_{\delta_e}} \delta_e \right] + \dot{d}_{\theta_1} \\ &= \Gamma_6(r^2 - p^2) + \Gamma_5 pr + \frac{\rho V_a^2 c S}{2J_y} \left[C_{m_0} + C_{m_\alpha}(\theta - \gamma) + C_{m_q} \frac{c}{2V_a} (\dot{\theta} - d_{\theta_1}) + C_{m_{\delta_e}} \delta_e \right] + \dot{d}_{\theta_1} \\ &= \left(\frac{\rho V_a^2 c S}{2J_y} C_{m_q} \frac{c}{2V_a} \right) \dot{\theta} + \left(\frac{\rho V_a^2 c S}{2J_y} C_{m_\alpha} \right) \theta + \left(\frac{\rho V_a^2 c S}{2J_y} C_{m_{\delta_e}} \right) \delta_e + \left\{ \Gamma_6(r^2 - p^2) \right. \\ &\quad \left. + \Gamma_5 pr + \frac{\rho V_a^2 c S}{J_y} \left[C_{m_0} - C_{m_\alpha} \gamma - C_{m_q} \frac{c}{2V_a} d_{\theta_1} \right] + \dot{d}_{\theta_1} \right\} \\ &= -a_{\theta_1} \dot{\theta} - a_{\theta_2} \theta + a_{\theta_3} \delta_e + d_{\theta_2}\end{aligned}$$

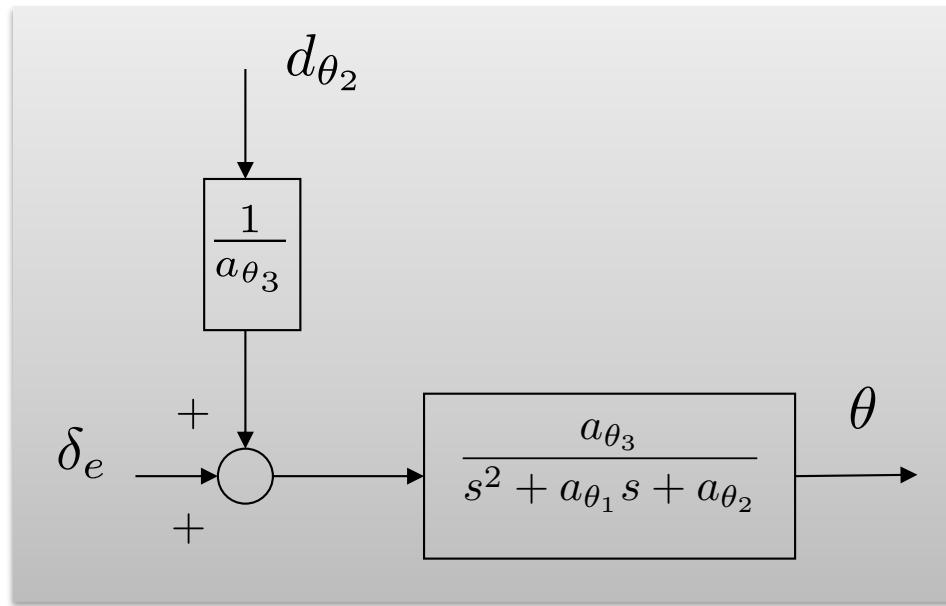
Longitudinal Transfer Functions - Pitch

The ODE model is

$$\ddot{\theta} = -a_{\theta_1}\dot{\theta} - a_{\theta_2}\theta + a_{\theta_3}\delta_e + d_{\theta_2}$$

with associated transfer function

$$\theta(s) = \left(\frac{a_{\theta_3}}{s^2 + a_{\theta_1}s + a_{\theta_2}} \right) \left(\delta_e(s) + \frac{1}{a_{\theta_3}}d_{\theta_2}(s) \right)$$



Longitudinal TF - Altitude from Pitch

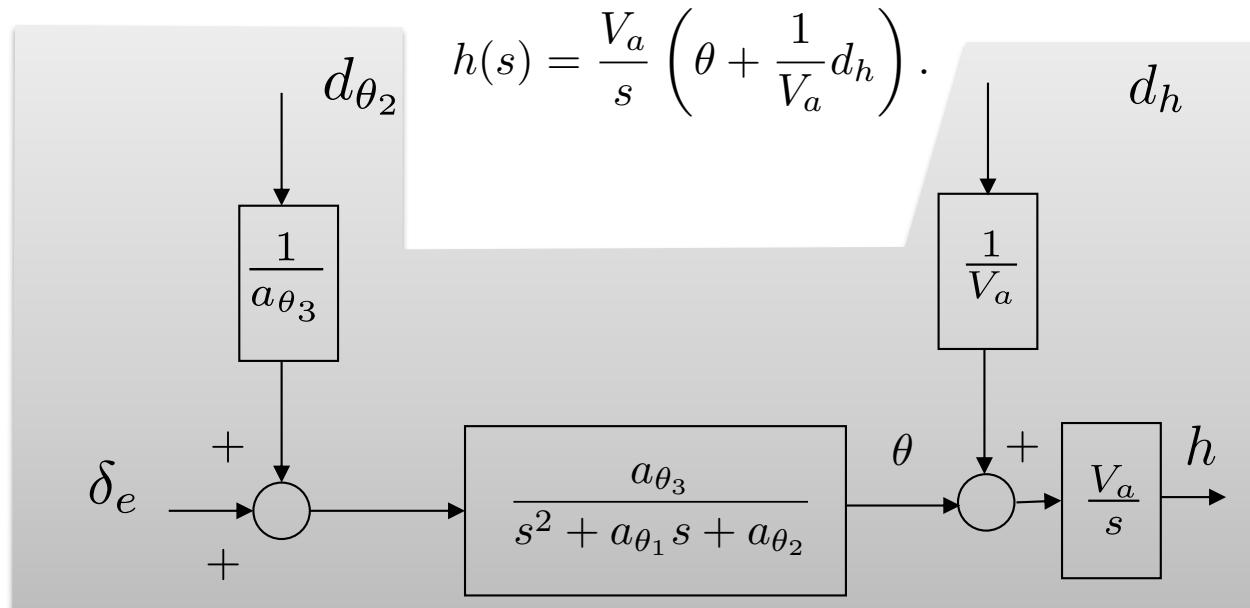
To derive a transfer function from pitch to altitude (assuming constant air-speed):

$$\begin{aligned}\dot{h} &= u \sin \theta - v \sin \phi \cos \theta - w \cos \phi \cos \theta \\ &= V_a \theta + (u \sin \theta - V_a \theta) - v \sin \phi \cos \theta - w \cos \phi \cos \theta \\ &= V_a \theta + d_h\end{aligned}$$

where

$$d_h \triangleq (u \sin \theta - V_a \theta) - v \sin \phi \cos \theta - w \cos \phi \cos \theta.$$

The associated transfer function is



Longitudinal Transfer Functions - Airspeed

Both pitch angle and throttle strongly affect airspeed. We wish to derive the relationship.

If there is no wind, then

$$V_a = \sqrt{u^2 + v^2 + w^2}.$$

Differentiating gives

$$\dot{V}_a = \dot{u} \cos \alpha \cos \beta + \dot{v} \sin \alpha \cos \beta + \dot{w} \sin \alpha \cos \beta = \dot{u} \cos \alpha + \dot{w} \sin \alpha + d_{V_1}$$

where

$$d_{V_1} = -\dot{u}(1 - \cos \beta) \cos \alpha - \dot{w}(1 - \cos \beta) \sin \alpha + \dot{v} \sin \beta.$$

Substituting from the dynamics gives

$$\begin{aligned}\dot{V}_a &= rV_a \cos \alpha \sin \beta - pV_a \sin \alpha \sin \beta - g \cos \alpha \sin \theta + g \sin \alpha \cos \theta \cos \phi \\ &\quad + \frac{\rho V_a^2 S}{2m} \left[-C_D(\alpha) - C_{D_\alpha} \alpha - C_{D_q} \frac{cq}{2V_a} - C_{D_{\delta_e}} \delta_e \right] + \frac{1}{m} T_p(\delta_t, V_a) + d_{V_1} \\ &= (rV_a \cos \alpha - pV_a \sin \alpha) \sin \beta - g \sin(\theta - \alpha) - g \sin \alpha \cos \theta (1 - \cos \phi) \\ &\quad + \frac{\rho V_a^2 S}{2m} \left[-C_{D_0} - C_{D_\alpha} \alpha - C_{D_q} \frac{cq}{2V_a} - C_{D_{\delta_e}} \delta_e \right] + \frac{1}{m} T_p(\delta_t, V_a) + d_{V_1} \\ &= -g \sin \gamma + \frac{\rho V_a^2 S}{2m} \left[-C_{D_0} - C_{D_\alpha} \alpha - C_{D_q} \frac{cq}{2V_a} - C_{D_{\delta_e}} \delta_e \right] + \frac{1}{m} T_p(\delta_t, V_a) + d_{V_2}\end{aligned}$$

Longitudinal Transfer Functions - Airspeed

Evaluate at trim: α^* , V_a^* , θ^* , δ_e^* , and let

$$\bar{V}_a = V_a - V_a^*, \quad \bar{\theta} = \theta - \theta^*, \quad \bar{\delta}_e = \delta_e - \delta_e^*$$

to get

$$\begin{aligned}\dot{\bar{V}}_a &= -g \cos(\theta^* - \alpha^*) \bar{\theta} + \left\{ \frac{\rho V_a^* S}{m} [-C_{D_0} - C_{D_\alpha} \alpha^* - C_{D_{\delta_e}} \delta_e^*] + \frac{1}{m} \frac{\partial T_p}{\partial V_a} (\delta_t^*, V_a^*) \right\} \bar{V}_a \\ &\quad + \frac{1}{m} \frac{\partial T_p}{\partial \delta_t} (\delta_t^*, V_a^*) \bar{\delta}_t + d_V \\ &= -a_{V_1} \bar{V}_a + a_{V_2} \bar{\delta}_t - a_{V_3} \bar{\theta} + d_V\end{aligned}$$

where

$$\begin{aligned}a_{V_1} &= \frac{\rho V_a^* S}{m} [C_{D_0} + C_{D_\alpha} \alpha^* + C_{D_{\delta_e}} \delta_e^*] - \frac{1}{m} \frac{\partial T_p}{\partial V_a} (\delta_t^*, V_a^*) \\ a_{V_2} &= \frac{1}{m} \frac{\partial T_p}{\partial \delta_t} (\delta_t^*, V_a^*), \\ a_{V_3} &= g \cos(\theta^* - \alpha^*),\end{aligned}$$

Longitudinal Transfer Functions - Airspeed

The associated transfer function is

$$\bar{V}_a(s) = \frac{1}{s + a_{V_1}} (a_{V_2} \bar{\delta}_t(s) - a_{V_3} \bar{\theta}(s) + d_V(s))$$

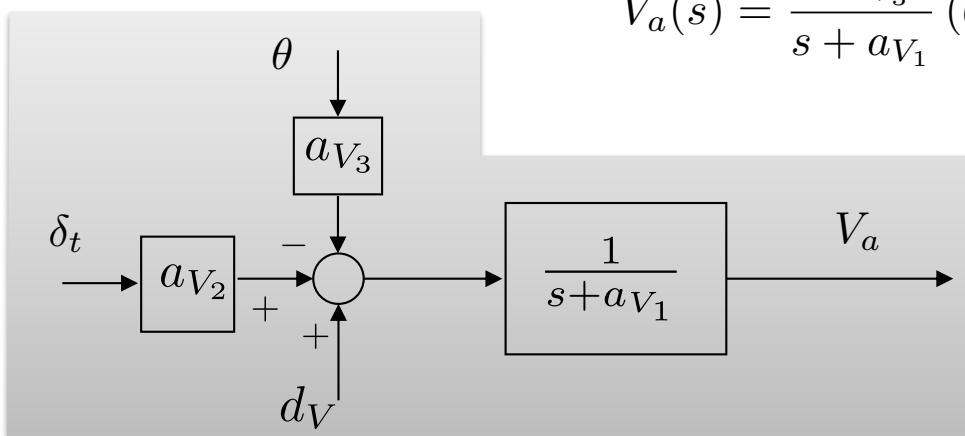
For UAVs we don't know trim values. Therefore the untrimmed variables are used and the trim values are absorbed into the disturbances.

For throttle we have

$$V_a(s) = \frac{a_{V_2}}{s + a_{V_1}} (\delta_t(s) + d_{V_1}(s))$$

and for pitch angle we have

$$V_a(s) = \frac{-a_{V_3}}{s + a_{V_1}} (\theta(s) + d_{V_2}(s))$$



Linear State-space Models

We will derive the state space equations about trim.

Recall that the nonlinear equations are given by $\dot{x} = f(x, u)$. The trim condition is $\dot{x}^* = f(x^*, u^*)$. Let $\bar{x} = x - x^*$ be the deviation from trim. Then

$$\begin{aligned}\dot{\bar{x}} &= \dot{x} - \dot{x}^* \\ &= f(x, u) - f(x^*, u^*) \\ &= f(x + x^* - x^*, u + u^* - u^*) - f(x^*, u^*) \\ &= f(x^* + \bar{x}, u^* + \bar{u}) - f(x^*, u^*)\end{aligned}$$

Using a Taylor series expansion around trim gives

$$\begin{aligned}\dot{\bar{x}} &= f(x^*, u^*) + \frac{\partial f(x^*, u^*)}{\partial x} \bar{x} + \frac{\partial f(x^*, u^*)}{\partial u} \bar{u} + H.O.T - f(x^*, u^*) \\ &\approx \frac{\partial f(x^*, u^*)}{\partial x} \bar{x} + \frac{\partial f(x^*, u^*)}{\partial u} \bar{u} \\ &\stackrel{\triangle}{=} A\bar{x} + B\bar{u}\end{aligned}$$

Lateral State-space Equations

The lateral states and inputs are defined as

$$\dot{x}_{\text{lat}} \triangleq (v, p, r, \phi, \psi)^\top, \quad u_{\text{lat}} \triangleq (\delta_a, \delta_r)^\top.$$

The nonlinear equations of motion are

$$\begin{aligned} \dot{v} &= pw - ru + g \cos \theta \sin \phi + \frac{\rho \sqrt{u^2 + v^2 + w^2} S}{2m} \frac{b}{2} [C_{Y_p} p + C_{Y_r} r] \\ &\quad + \frac{\rho(u^2 + v^2 + w^2)S}{2m} [C_{Y_0} + C_{Y_\beta} \tan^{-1} \left(\frac{v}{\sqrt{u^2 + w^2}} \right) + C_{Y_{\delta_a}} \delta_a + C_{Y_{\delta_r}} \delta_r] \\ \dot{p} &= \Gamma_1 pq - \Gamma_2 qr + \frac{\rho \sqrt{u^2 + v^2 + w^2} S}{2} \frac{b^2}{2} [C_{p_p} p + C_{p_r} r] \\ &\quad + \frac{1}{2} \rho(u^2 + v^2 + w^2) S b [C_{p_0} + C_{p_\beta} \tan^{-1} \left(\frac{v}{\sqrt{u^2 + w^2}} \right) + C_{p_{\delta_a}} \delta_a + C_{p_{\delta_r}} \delta_r] \\ \dot{r} &= \Gamma_7 pq - \Gamma_1 qr + \frac{\rho \sqrt{u^2 + v^2 + w^2} S}{2} \frac{b^2}{2} [C_{r_p} p + C_{r_r} r] \\ &\quad + \frac{1}{2} \rho(u^2 + v^2 + w^2) S b [C_{r_0} + C_{r_\beta} \tan^{-1} \left(\frac{v}{\sqrt{u^2 + w^2}} \right) + C_{r_{\delta_a}} \delta_a + C_{r_{\delta_r}} \delta_r] \\ \dot{\phi} &= p + q \sin \phi \tan \theta + r \cos \phi \tan \theta \\ \dot{\psi} &= q \sin \phi \sec \theta + r \cos \phi \sec \theta \end{aligned}$$

where we have used $\beta = \tan^{-1} \left(\frac{v}{\sqrt{u^2 + w^2}} \right)$ and $V_a = \sqrt{u^2 + v^2 + w^2}$.

Jacobian Matrices

$$A_{\text{lat}} = \frac{\partial f_{\text{lat}}}{\partial x_{\text{lat}}} = \begin{pmatrix} \frac{\partial \dot{v}}{\partial v} & \frac{\partial \dot{v}}{\partial p} & \frac{\partial \dot{v}}{\partial r} & \frac{\partial \dot{v}}{\partial \phi} & \frac{\partial \dot{v}}{\partial \psi} \\ \frac{\partial \dot{p}}{\partial v} & \frac{\partial \dot{p}}{\partial p} & \frac{\partial \dot{p}}{\partial r} & \frac{\partial \dot{p}}{\partial \phi} & \frac{\partial \dot{p}}{\partial \psi} \\ \frac{\partial \dot{r}}{\partial v} & \frac{\partial \dot{r}}{\partial p} & \frac{\partial \dot{r}}{\partial r} & \frac{\partial \dot{r}}{\partial \phi} & \frac{\partial \dot{r}}{\partial \psi} \\ \frac{\partial \dot{\phi}}{\partial v} & \frac{\partial \dot{\phi}}{\partial p} & \frac{\partial \dot{\phi}}{\partial r} & \frac{\partial \dot{\phi}}{\partial \phi} & \frac{\partial \dot{\phi}}{\partial \psi} \\ \frac{\partial \dot{\psi}}{\partial v} & \frac{\partial \dot{\psi}}{\partial p} & \frac{\partial \dot{\psi}}{\partial r} & \frac{\partial \dot{\psi}}{\partial \phi} & \frac{\partial \dot{\psi}}{\partial \psi} \end{pmatrix}$$
$$B_{\text{lat}} = \frac{\partial f_{\text{lat}}}{\partial u_{\text{lat}}} = \begin{pmatrix} \frac{\partial \dot{v}}{\partial \delta_a} & \frac{\partial \dot{v}}{\partial \delta_r} \\ \frac{\partial \dot{p}}{\partial \delta_a} & \frac{\partial \dot{p}}{\partial \delta_r} \\ \frac{\partial \dot{r}}{\partial \delta_a} & \frac{\partial \dot{r}}{\partial \delta_r} \\ \frac{\partial \dot{\phi}}{\partial \delta_a} & \frac{\partial \dot{\phi}}{\partial \delta_r} \\ \frac{\partial \dot{\psi}}{\partial \delta_a} & \frac{\partial \dot{\psi}}{\partial \delta_r} \end{pmatrix}$$

Take partial derivatives and evaluate at trim to get

Lateral State-space Equations

$$\begin{pmatrix} \dot{\bar{v}} \\ \dot{\bar{p}} \\ \dot{\bar{r}} \\ \dot{\bar{\phi}} \\ \dot{\bar{\psi}} \end{pmatrix} = \begin{pmatrix} Y_v & Y_p & Y_r & g \cos \theta^* \cos \phi^* & 0 \\ L_v & L_p & L_r & 0 & 0 \\ N_v & N_p & N_r & 0 & 0 \\ 0 & 1 & \cos \phi^* \tan \theta^* & q^* \cos \phi^* \tan \theta^* - r^* \sin \phi^* \tan \theta^* & 0 \\ 0 & 0 & \cos \phi^* \sec \theta^* & p^* \cos \phi^* \sec \theta^* - r^* \sin \phi^* \sec \theta^* & 0 \end{pmatrix} \begin{pmatrix} \bar{v} \\ \bar{p} \\ \bar{r} \\ \bar{\phi} \\ \bar{\psi} \end{pmatrix} + \begin{pmatrix} Y_{\delta_a} & Y_{\delta_r} \\ L_{\delta_a} & L_{\delta_r} \\ N_{\delta_a} & N_{\delta_r} \\ 0 & 0 \\ 0 & 0 \end{pmatrix} \begin{pmatrix} \bar{\delta}_a \\ \bar{\delta}_r \end{pmatrix}$$

Lateral	Formula
Y_v	$\frac{\rho S b v^*}{4m V_a^*} [C_{Y_p} p^* + C_{Y_r} r^*] + \frac{\rho S v^*}{m} [C_{Y_0} + C_{Y_\beta} \beta^* + C_{Y_{\delta_a}} \delta_a^* + C_{Y_{\delta_r}} \delta_r^*] + \frac{\rho S C_{Y_\beta}}{2m} \sqrt{u^{*2} + w^{*2}}$
Y_p	$w^* + \frac{\rho V_a^* S b}{4m} C_{Y_p}$
Y_r	$-u^* + \frac{\rho V_a^* S b}{4m} C_{Y_r}$
Y_{δ_a}	$\frac{\rho V_a^{*2} S}{2m} C_{Y_{\delta_a}}$
Y_{δ_r}	$\frac{\rho V_a^{*2} S}{2m} C_{Y_{\delta_r}}$
L_v	$\frac{\rho S b^2 v^*}{4V_a^*} [C_{p_p} p^* + C_{p_r} r^*] + \rho S b v^* [C_{p_0} + C_{p_\beta} \beta^* + C_{p_{\delta_a}} \delta_a^* + C_{p_{\delta_r}} \delta_r^*] + \frac{\rho S b C_{p_\beta}}{2} \sqrt{u^{*2} + w^{*2}}$
L_p	$\Gamma_1 q^* + \frac{\rho V_a^* S b^2}{4} C_{p_p}$
L_r	$-\Gamma_2 q^* + \frac{\rho V_a^* S b^2}{4} C_{p_r}$
L_{δ_a}	$\frac{\rho V_a^{*2} S b}{2} C_{p_{\delta_a}}$
L_{δ_r}	$\frac{\rho V_a^{*2} S b}{2} C_{p_{\delta_r}}$
N_v	$\frac{\rho S b^2 v^*}{4V_a^*} [C_{r_p} p^* + C_{r_r} r^*] + \rho S b v^* [C_{r_0} + C_{r_\beta} \beta^* + C_{r_{\delta_a}} \delta_a^* + C_{r_{\delta_r}} \delta_r^*] + \frac{\rho S b C_{r_\beta}}{2} \sqrt{u^{*2} + w^{*2}}$
N_p	$\Gamma_7 q^* + \frac{\rho V_a^* S b^2}{4} C_{r_p}$
N_r	$-\Gamma_1 q^* + \frac{\rho V_a^* S b^2}{4} C_{r_r}$
N_{δ_a}	$\frac{\rho V_a^{*2} S b}{2} C_{r_{\delta_a}}$
N_{δ_r}	$\frac{\rho V_a^{*2} S b}{2} C_{r_{\delta_r}}$

Alternative Form - Lateral

The lateral dynamics are often expressed using β instead of v . Recall that

$$v = V_a \sin \beta.$$

Therefore

$$\bar{v} = V_a^* \cos \beta^* \bar{\beta}.$$

Differentiating gives

$$\dot{\bar{\beta}} = \frac{1}{V_a^* \cos \beta^*} \dot{\bar{v}}.$$

The associated state space equations are

$$\begin{pmatrix} \dot{\bar{\beta}} \\ \dot{\bar{p}} \\ \dot{\bar{r}} \\ \dot{\bar{\phi}} \\ \dot{\bar{\psi}} \end{pmatrix} = \begin{pmatrix} Y_v & \frac{Y_p}{V_a^* \cos \beta^*} & \frac{Y_r}{V_a^* \cos \beta^*} & \frac{g \cos \theta^* \cos \phi^*}{V_a^* \cos \beta^*} & 0 \\ L_v V_a^* \cos \beta^* & L_p & L_r & 0 & 0 \\ N_v V_a^* \cos \beta^* & N_p & N_r & 0 & 0 \\ 0 & 1 & \cos \phi^* \tan \theta^* & q^* \cos \phi^* \tan \theta^* & 0 \\ 0 & 0 & \cos \phi^* \sec \theta^* & -r^* \sin \phi^* \tan \theta^* & 0 \end{pmatrix} \begin{pmatrix} \bar{\beta} \\ \bar{p} \\ \bar{r} \\ \bar{\phi} \\ \bar{\psi} \end{pmatrix} + \begin{pmatrix} \frac{Y_{\delta_a}}{V_a^* \cos \beta^*} & \frac{Y_{\delta_r}}{V_a^* \cos \beta^*} \\ L_{\delta_a} & L_{\delta_r} \\ N_{\delta_a} & N_{\delta_r} \\ 0 & 0 \\ 0 & 0 \end{pmatrix} \begin{pmatrix} \bar{\delta}_a \\ \bar{\delta}_r \end{pmatrix}$$

Longitudinal State-space Equations

The longitudinal states and inputs are defined as

$$\dot{x}_{\text{lon}} \triangleq (u, w, q, \theta, h)^\top, \quad u_{\text{lon}} \triangleq (\delta_e, \delta_t)^\top.$$

The nonlinear equations of motion are

$$\begin{aligned} \dot{u} &= -qw - g \sin \theta + \frac{\rho(u^2 + w^2)S}{2m} \left[C_{X_0} + C_{X_\alpha} \tan^{-1} \left(\frac{w}{u} \right) + C_{X_{\delta_e}} \delta_e \right] \\ &\quad + \frac{\rho\sqrt{u^2 + w^2}S}{4m} C_{X_q} cq + \frac{\rho S_{\text{prop}}}{2m} C_{\text{prop}} \left[(k\delta_t)^2 - (u^2 + w^2) \right] \\ \dot{w} &= qu + g \cos \theta + \frac{\rho(u^2 + w^2)S}{2m} \left[C_{Z_0} + C_{Z_\alpha} \tan^{-1} \left(\frac{w}{u} \right) + C_{Z_{\delta_e}} \delta_e \right] \\ &\quad + \frac{\rho\sqrt{u^2 + w^2}S}{4m} C_{Z_q} cq \\ \dot{q} &= \frac{1}{2J_y} \rho(u^2 + w^2)cS \left[C_{m_0} + C_{m_\alpha} \tan^{-1} \left(\frac{w}{u} \right) + C_{m_{\delta_e}} \delta_e \right] + \frac{1}{4J_y} \rho\sqrt{u^2 + w^2} S C_{m_q} c^2 q \\ \dot{\theta} &= q \\ \dot{h} &= u \sin \theta - w \cos \theta \end{aligned}$$

where we have used $\alpha = \tan^{-1} \left(\frac{w}{u} \right)$ and $V_a = \sqrt{u^2 + w^2}$.

Jacobian Matrices

$$A_{\text{lon}} = \frac{\partial f_{\text{lon}}}{\partial x_{\text{lon}}} = \begin{pmatrix} \frac{\partial \dot{u}}{\partial u} & \frac{\partial \dot{u}}{\partial w} & \frac{\partial \dot{u}}{\partial q} & \frac{\partial \dot{u}}{\partial \theta} & \frac{\partial \dot{u}}{\partial h} \\ \frac{\partial \dot{w}}{\partial u} & \frac{\partial \dot{w}}{\partial w} & \frac{\partial \dot{w}}{\partial q} & \frac{\partial \dot{w}}{\partial \theta} & \frac{\partial \dot{w}}{\partial h} \\ \frac{\partial \dot{q}}{\partial u} & \frac{\partial \dot{q}}{\partial w} & \frac{\partial \dot{q}}{\partial q} & \frac{\partial \dot{q}}{\partial \theta} & \frac{\partial \dot{q}}{\partial h} \\ \frac{\partial \dot{\theta}}{\partial u} & \frac{\partial \dot{\theta}}{\partial w} & \frac{\partial \dot{\theta}}{\partial q} & \frac{\partial \dot{\theta}}{\partial \theta} & \frac{\partial \dot{\theta}}{\partial h} \\ \frac{\partial \dot{h}}{\partial u} & \frac{\partial \dot{h}}{\partial w} & \frac{\partial \dot{h}}{\partial q} & \frac{\partial \dot{h}}{\partial \theta} & \frac{\partial \dot{h}}{\partial h} \end{pmatrix}$$

$$B_{\text{lon}} = \frac{\partial f_{\text{lon}}}{\partial u_{\text{lon}}} = \begin{pmatrix} \frac{\partial \dot{u}}{\partial \delta_e} & \frac{\partial \dot{u}}{\partial \delta_t} \\ \frac{\partial \dot{w}}{\partial \delta_e} & \frac{\partial \dot{w}}{\partial \delta_t} \\ \frac{\partial \dot{q}}{\partial \delta_e} & \frac{\partial \dot{q}}{\partial \delta_t} \\ \frac{\partial \dot{\theta}}{\partial \delta_e} & \frac{\partial \dot{\theta}}{\partial \delta_t} \\ \frac{\partial \dot{h}}{\partial \delta_e} & \frac{\partial \dot{h}}{\partial \delta_t} \end{pmatrix}$$

Take partial derivatives and evaluate at trim to get

Longitudinal State-space Equations

$$\begin{pmatrix} \dot{\bar{u}} \\ \dot{\bar{w}} \\ \dot{\bar{q}} \\ \dot{\bar{\theta}} \\ \dot{\bar{h}} \end{pmatrix} = \begin{pmatrix} X_u & X_w & X_q & -g \cos \theta^* & 0 \\ Z_u & Z_w & Z_q & -g \sin \theta^* & 0 \\ M_u & M_w & M_q & 0 & 0 \\ 0 & 0 & 1 & 0 & 0 \\ \sin \theta^* & -\cos \theta^* & 0 & u^* \cos \theta^* + w^* \sin \theta^* & 0 \end{pmatrix} \begin{pmatrix} \bar{u} \\ \bar{w} \\ \bar{q} \\ \bar{\theta} \\ \bar{h} \end{pmatrix} + \begin{pmatrix} X_{\delta_e} & X_{\delta_t} \\ Z_{\delta_e} & 0 \\ M_{\delta_e} & 0 \\ 0 & 0 \\ 0 & 0 \end{pmatrix} \begin{pmatrix} \bar{\delta}_e \\ \bar{\delta}_t \end{pmatrix}$$

Longitudinal	Formula
X_u	$\frac{u^* \rho S}{m} [C_{X_0} + C_{X_\alpha} \alpha^* + C_{X_{\delta_e}} \delta_e^*] - \frac{\rho S w^* C_{X_\alpha}}{2m} + \frac{\rho S c C_{X_q} u^* q^*}{4m V_a^*} - \frac{\rho S_{prop} C_{prop} u^*}{m}$
X_w	$-q^* + \frac{w^* \rho S}{m} [C_{X_0} + C_{X_\alpha} \alpha^* + C_{X_{\delta_e}} \delta_e^*] + \frac{\rho S c C_{X_q} w^* q^*}{4m V_a^*} + \frac{\rho S C_{X_\alpha} u^*}{2m} - \frac{\rho S_{prop} C_{prop} w^*}{m}$
X_q	$-w^* + \frac{\rho V_a^* S C_{X_q} c}{4m}$
X_{δ_e}	$\frac{\rho V_a^{*2} S C_{X_{\delta_e}}}{2m}$
X_{δ_t}	$\frac{\rho S_{prop} C_{prop} k^2 \delta_t^*}{m}$
Z_u	$q^* + \frac{u^* \rho S}{m} [C_{Z_0} + C_{Z_\alpha} \alpha^* + C_{Z_{\delta_e}} \delta_e^*] - \frac{\rho S C_{Z_\alpha} w^*}{2m} + \frac{u^* \rho S C_{Z_q} c q^*}{4m V_a^*}$
Z_w	$\frac{w^* \rho S}{m} [C_{Z_0} + C_{Z_\alpha} \alpha^* + C_{Z_{\delta_e}} \delta_e^*] + \frac{\rho S C_{Z_\alpha} u^*}{2m} + \frac{\rho w^* S c C_{Z_q} q^*}{4m V_a^*}$
Z_q	$u^* + \frac{\rho V_a^* S C_{Z_q} c}{4m}$
Z_{δ_e}	$\frac{\rho V_a^{*2} S C_{Z_{\delta_e}}}{2m}$
M_u	$\frac{u^* \rho S c}{J_y} [C_{m_0} + C_{m_\alpha} \alpha^* + C_{m_{\delta_e}} \delta_e^*] - \frac{\rho S c C_{m_\alpha} w^*}{2J_y} + \frac{\rho S c^2 C_{m_q} q^* u^*}{4J_y V_a^*}$
M_w	$\frac{w^* \rho S c}{J_y} [C_{m_0} + C_{m_\alpha} \alpha^* + C_{m_{\delta_e}} \delta_e^*] + \frac{\rho S c C_{m_\alpha} u^*}{2J_y} + \frac{\rho S c^2 C_{m_q} q^* w^*}{4J_y V_a^*}$
M_q	$\frac{\rho V_a^* S c^2 C_{m_q}}{4J_y}$
M_{δ_e}	$\frac{\rho V_a^{*2} S c C_{m_{\delta_e}}}{2J_y}$

Alternative Form – Longitudinal

The longitudinal dynamics are often expressed using α instead of w . Recall that

$$w = V_a \sin \alpha \cos \beta.$$

At trim, when $\beta = 0$ we have

$$\bar{w} = V_a^* \cos \alpha^* \bar{\alpha}.$$

Differentiating gives

$$\dot{\bar{\alpha}} = \frac{1}{V_a^* \cos \alpha^*} \dot{\bar{w}}.$$

The associated state space equations are

$$\begin{pmatrix} \dot{\bar{u}} \\ \dot{\bar{\alpha}} \\ \dot{\bar{q}} \\ \dot{\bar{\theta}} \\ \dot{\bar{h}} \end{pmatrix} = \begin{pmatrix} X_u & X_w V_a^* \cos \alpha^* & X_q & -g \cos \theta^* & 0 \\ \frac{Z_u}{V_a^* \cos \alpha^*} & Z_w & \frac{Z_q}{V_a^* \cos \alpha^*} & \frac{-g \sin \theta^*}{V_a^* \cos \alpha^*} & 0 \\ M_u & M_w V_a^* \cos \alpha^* & M_q & 0 & 0 \\ 0 & 0 & 1 & 0 & 0 \\ \sin \theta^* & -V_a^* \cos \theta^* \cos \alpha^* & 0 & u^* \cos \theta^* + w^* \sin \theta^* & 0 \end{pmatrix} \begin{pmatrix} \bar{u} \\ \bar{\alpha} \\ \bar{q} \\ \bar{\theta} \\ \bar{h} \end{pmatrix} + \begin{pmatrix} X_{\delta_e} & X_{\delta_t} \\ \frac{Z_{\delta_e}}{V_a^* \cos \alpha^*} & 0 \\ M_{\delta_e} & 0 \\ 0 & 0 \\ 0 & 0 \end{pmatrix} \begin{pmatrix} \bar{\delta}_e \\ \bar{\delta}_t \end{pmatrix}$$

Reduced Order Modes

Longitudinal Modes

- Short Period Mode
 - Short period mode is the fast mode seen in pitch rate q and pitch angle theta.
- Phugoid Mode
 - Phugoid mode is slow mode seen in theta and u .
- Induced by impulse on elevator

Lateral Modes

- Roll Mode
 - First order fast mode between aileron and pitch rate.
- Spiral-divergence Mode
 - First order (really) slow mode between aileron and yaw/course
- Dutch-roll Mode
 - Second order mode: coupling between roll, yaw, and side slip. Like a duck wagging its tail.
- Induced by doublet on aileron or rudder

Simulation Project

- Find trim states and inputs. For the aerosonde model use $V_a = 17$ m/s, and $R = 150$ m.
- Find the transfer functions derived in this chapter.
- Find the state space models derived in this chapter.
- Compute the eigenvalues of A_{lon} and A_{lat} and associate them with the phugoid mode, the short period mode, the roll mode, the dutch roll mode, and the spiral mode. Show these modes in the simulator.