

# 量子计算与机器学习作业

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**3.1** The *fidelity*  $F$  of two quantum states  $|\psi_1\rangle$  and  $|\psi_2\rangle$  is defined by  $F \equiv |\langle\psi_1|\psi_2\rangle|^2$ . It is a measure of the distance between the two quantum states: We have  $0 \leq F \leq 1$ , with  $F = 1$  when  $|\psi_1\rangle$  coincides with  $|\psi_2\rangle$  and  $F = 0$  when  $|\psi_1\rangle$  and  $|\psi_2\rangle$  are orthogonal. Show that  $F = \cos^2 \frac{\alpha}{2}$ , with  $\alpha$  the angle between the Bloch vectors corresponding to the quantum states  $|\psi_1\rangle$  and  $|\psi_2\rangle$ .

假设两个量子态  $|\psi_1\rangle$  和  $|\psi_2\rangle$  在 Bloch 球面上的位置分别为  $(\theta_1, \varphi_1)$  和  $(\theta_2, \varphi_2)$ , 即

$$|\psi_1\rangle = \cos \frac{\theta_1}{2} |0\rangle + e^{i\varphi_1} \sin \frac{\theta_1}{2} |1\rangle = \begin{pmatrix} \cos \frac{\theta_1}{2} \\ e^{i\varphi_1} \sin \frac{\theta_1}{2} \end{pmatrix},$$
$$|\psi_2\rangle = \cos \frac{\theta_2}{2} |0\rangle + e^{i\varphi_2} \sin \frac{\theta_2}{2} |1\rangle = \begin{pmatrix} \cos \frac{\theta_2}{2} \\ e^{i\varphi_2} \sin \frac{\theta_2}{2} \end{pmatrix}.$$

则这两个量子态的内积为

$$\begin{aligned} \langle\psi_1|\psi_2\rangle &= \left( \cos \frac{\theta_1}{2}, e^{-i\varphi_1} \sin \frac{\theta_1}{2} \right) \begin{pmatrix} \cos \frac{\theta_2}{2} \\ e^{i\varphi_2} \sin \frac{\theta_2}{2} \end{pmatrix} \\ &= \cos \frac{\theta_1}{2} \cos \frac{\theta_2}{2} + \sin \frac{\theta_1}{2} \sin \frac{\theta_2}{2} e^{i(\varphi_2 - \varphi_1)}. \end{aligned}$$

因此两个量子态的 fidelity 为

$$\begin{aligned} F = |\langle\psi_1|\psi_2\rangle|^2 &= \left( \cos \frac{\theta_1}{2} \cos \frac{\theta_2}{2} + \sin \frac{\theta_1}{2} \sin \frac{\theta_2}{2} \cos(\varphi_2 - \varphi_1) \right)^2 + \left( \sin \frac{\theta_1}{2} \sin \frac{\theta_2}{2} \sin(\varphi_2 - \varphi_1) \right)^2 \\ &= \cos^2 \frac{\theta_1}{2} \cos^2 \frac{\theta_2}{2} + \sin^2 \frac{\theta_1}{2} \sin^2 \frac{\theta_2}{2} + 2 \cos \frac{\theta_1}{2} \cos \frac{\theta_2}{2} \sin \frac{\theta_1}{2} \sin \frac{\theta_2}{2} \cos(\varphi_2 - \varphi_1) \\ &= \frac{1}{2} (1 + \cos \theta_1 \cos \theta_2 + \sin \theta_1 \sin \theta_2 \cos(\varphi_2 - \varphi_1)). \end{aligned}$$

而两个量子态在 Bloch 球面上的夹角  $\alpha$  即为向量  $(\sin \theta_1 \cos \varphi_1, \sin \theta_1 \sin \varphi_1, \cos \theta_1)^\top$  和  $(\sin \theta_2 \cos \varphi_2, \sin \theta_2 \sin \varphi_2, \cos \theta_2)^\top$  之间的夹角, 有

$$\begin{aligned} \cos \alpha &= (\sin \theta_1 \cos \varphi_1, \sin \theta_1 \sin \varphi_1, \cos \theta_1) \begin{pmatrix} \sin \theta_2 \cos \varphi_2 \\ \sin \theta_2 \sin \varphi_2 \\ \cos \theta_2 \end{pmatrix} \\ &= \sin \theta_1 \sin \theta_2 (\cos \varphi_1 \cos \varphi_2 + \sin \varphi_1 \sin \varphi_2) + \cos \theta_1 \cos \theta_2 \\ &= \cos \theta_1 \cos \theta_2 + \sin \theta_1 \sin \theta_2 \cos(\varphi_2 - \varphi_1) \\ &= 2F - 1. \end{aligned}$$

因此,  $F = \frac{1}{2}(1 + \cos \alpha) = \cos^2 \frac{\alpha}{2}$ .

**3.2** Show that the unitary operator moving the state parametrized on the Bloch sphere by the angles  $(\theta_1, \varphi_1)$  into the state  $(\theta_2, \varphi_2)$  is given by

$$R_z\left(\frac{\pi}{2} + \varphi_2\right) H R_z(\theta_2 - \theta_1) H R_z\left(-\frac{\pi}{2} - \varphi_1\right).$$

The *phase-shift gate* is defined as

$$R_z(\delta) = \begin{bmatrix} 1 & 0 \\ 0 & e^{i\delta} \end{bmatrix}.$$

角度  $(\theta_1, \varphi_1)$  对应的量子态为

$$|\psi_1\rangle = \cos \frac{\theta_1}{2} |0\rangle + e^{i\varphi_1} \sin \frac{\theta_1}{2} |1\rangle = \begin{pmatrix} \cos \frac{\theta_1}{2} \\ e^{i\varphi_1} \sin \frac{\theta_1}{2} \end{pmatrix}.$$

经过给出的各个酉操作后，量子态依次变为：

$$\begin{aligned} |\psi_1\rangle &= \begin{pmatrix} \cos \frac{\theta_1}{2} \\ e^{i\varphi_1} \sin \frac{\theta_1}{2} \end{pmatrix} \xrightarrow{R_z(-\frac{\pi}{2}-\varphi_1)} \begin{pmatrix} \cos \frac{\theta_1}{2} \\ -i \sin \frac{\theta_1}{2} \end{pmatrix} \xrightarrow{H} \frac{1}{\sqrt{2}} \begin{pmatrix} e^{-i\frac{\theta_1}{2}} \\ e^{i\frac{\theta_1}{2}} \end{pmatrix} \xrightarrow{R_z(\theta_2-\theta_1)} \frac{1}{\sqrt{2}} e^{-i\frac{\theta_1}{2}} \begin{pmatrix} 1 \\ e^{i\theta_2} \end{pmatrix} \\ &\xrightarrow{H} \frac{1}{2} e^{-i\frac{\theta_1}{2}} \begin{pmatrix} 1 + e^{i\theta_2} \\ 1 - e^{i\theta_2} \end{pmatrix} \xrightarrow{R_z(\frac{\pi}{2}+\varphi_2)} \frac{1}{2} e^{-i\frac{\theta_1}{2}} \begin{pmatrix} 1 + e^{i\theta_2} \\ e^{\frac{\pi}{2}+\varphi_2}(1 - e^{i\theta_2}) \end{pmatrix} \end{aligned}$$

而

$$\begin{aligned} \frac{1}{2} e^{-i\frac{\theta_1}{2}} \begin{pmatrix} 1 + e^{i\theta_2} \\ e^{\frac{\pi}{2}+\varphi_2}(1 - e^{i\theta_2}) \end{pmatrix} &= \frac{1}{2} e^{i(\theta_2-\theta_1)} \begin{pmatrix} e^{-i\frac{\theta_2}{2}} + e^{i\frac{\theta_2}{2}} \\ ie^{\varphi_2} (e^{-i\frac{\theta_2}{2}} - e^{i\frac{\theta_2}{2}}) \end{pmatrix} \\ &= e^{i\frac{\theta_2-\theta_1}{2}} \begin{pmatrix} \cos \frac{\theta_2}{2} \\ e^{i\varphi_2} \sin \frac{\theta_2}{2} \end{pmatrix} \\ &= e^{i\frac{\theta_2-\theta_1}{2}} |\psi_2\rangle \\ &= |\psi_2\rangle. \end{aligned}$$

其中  $|\psi_2\rangle$  是角度  $(\theta_2, \varphi_2)$  对应的量子态。综上所述， $R_z\left(\frac{\pi}{2} + \varphi_2\right) H R_z(\theta_2 - \theta_1) H R_z\left(-\frac{\pi}{2} - \varphi_1\right)$  是将角度  $(\theta_1, \varphi_1)$  对应的量子态变为角度  $(\theta_2, \varphi_2)$  的酉操作。

**4.1** 证明贝尔态  $|\Phi^+\rangle = \frac{1}{\sqrt{2}}(|00\rangle + |11\rangle)$  可以等效表达为  $|\Phi^+\rangle = \frac{1}{\sqrt{2}}(|aa\rangle + |bb\rangle)$ ，其中  $|a\rangle$  和  $|b\rangle$  是任意一组正交归一基。

因为  $|a\rangle$  和  $|b\rangle$  是任意一组正交归一基，所以可以将  $|0\rangle$  和  $|1\rangle$  表示为

$$\begin{aligned} |0\rangle &= \alpha|a\rangle + \beta|b\rangle, \\ |1\rangle &= -\beta|a\rangle + \alpha|b\rangle. \end{aligned}$$

其中  $\alpha^2 + \beta^2 = 1$ 。因此， $|\Phi^+\rangle$  可以表示为

$$\begin{aligned}
|\Phi^+\rangle &= \frac{1}{\sqrt{2}}(|00\rangle + |11\rangle) \\
&= \frac{1}{\sqrt{2}}((\alpha|a\rangle + \beta|b\rangle) \otimes (\alpha|a\rangle + \beta|b\rangle) + (-\beta|a\rangle + \alpha|b\rangle) \otimes (-\beta|a\rangle + \alpha|b\rangle)) \\
&= \frac{1}{\sqrt{2}}((\alpha^2 + \beta^2)|aa\rangle + (\alpha^2 + \beta^2)|bb\rangle) \\
&= \frac{1}{\sqrt{2}}(|aa\rangle + |bb\rangle).
\end{aligned}$$

命题得证。

**5.1** Let  $|x\rangle$  be a basis state of  $n$  qubits. Prove that

$$H^{\otimes n}|x\rangle = \frac{\sum_z (-1)^{x \cdot z} |z\rangle}{\sqrt{2^n}}$$

where  $x \cdot z$  is the bitwise inner product of  $x$  and  $z$ , modulo 2, and the sum is over all  $z \in \{0, 1\}^n$ .

如下所示，有

$$\begin{aligned}
H^{\otimes n}|x\rangle &= H|x_1\rangle \otimes H|x_2\rangle \otimes \cdots \otimes H|x_n\rangle \\
&= \frac{1}{\sqrt{2}}(|0\rangle + (-1)^{x_1}|1\rangle) \otimes \frac{1}{\sqrt{2}}(|0\rangle + (-1)^{x_2}|1\rangle) \otimes \cdots \otimes \frac{1}{\sqrt{2}}(|0\rangle + (-1)^{x_n}|1\rangle) \otimes \\
&= \frac{1}{\sqrt{2^n}} \sum_{z_1 \in \{0,1\}} (-1)^{x_1 z_1} |z_1\rangle \otimes (|0\rangle + (-1)^{x_2}|1\rangle) \otimes \cdots \otimes (|0\rangle + (-1)^{x_n}|1\rangle) \\
&= \frac{1}{\sqrt{2^n}} \sum_{z_1 \in \{0,1\}} \sum_{z_2 \in \{0,1\}} (-1)^{x_1 z_1} |z_1\rangle \otimes (-1)^{x_2 z_2} |z_2\rangle \otimes (|0\rangle + (-1)^{x_3}|1\rangle) \otimes \cdots \otimes (|0\rangle + (-1)^{x_n}|1\rangle) \otimes \\
&= \cdots \\
&= \frac{1}{\sqrt{2^n}} \sum_{z_1 \in \{0,1\}} \sum_{z_2 \in \{0,1\}} \cdots \sum_{z_n \in \{0,1\}} (-1)^{x_1 z_1} |z_1\rangle \otimes (-1)^{x_2 z_2} |z_2\rangle \otimes \cdots \otimes (-1)^{x_n z_n} |z_n\rangle \\
&= \frac{1}{\sqrt{2^n}} \sum_{z_1, z_2, \dots, z_n \in \{0,1\}} (-1)^{x_1 z_1 + x_2 z_2 + \cdots + x_n z_n} |z_1 z_2 \cdots z_n\rangle \\
&= \frac{1}{\sqrt{2^n}} \sum_{z \in \{0,1\}^n} (-1)^{x \cdot z} |z\rangle.
\end{aligned}$$

即  $H^{\otimes n}|x\rangle = \frac{\sum_z (-1)^{x \cdot z} |z\rangle}{\sqrt{2^n}}$ ，命题得证。