

Homework 1

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1 Problem 4.2

Use separation of variable in *cartesian* coordinates to solve the infinite *cubical* well (or particle in a box):

$$V(x, y, z) = \begin{cases} 0, & \forall x, y, z \in [0, a] \\ \infty, & \forall x, y, z \notin [0, a] \end{cases}$$

1. Find the stationary states, and the corresponding energies.
2. Call the distinct energies E_1, E_2, \dots in order of increasing energy. Find E_1, E_2, E_3, E_4, E_5 and E_6 . Determine their degeneracies (that is, the number of different states that share the same energy).
3. What is the degeneracy of E_{14} , and why is this case interesting?

Solution 1: Stationary states

To find the stationary states of the infinite cubical well, we are going to solve the time independent Schrödinger equation,

$$-\frac{\hbar^2}{2m} \nabla^2 \psi = E\psi, \quad \forall x, y, z \in [0, a],$$

with the following boundary conditions $\psi(0, 0, 0) = \psi(a, a, a) = 0$. To solve the equation we are going to use the method of separation of variables, that is, that we assume that the solution of the differential equation has the following form $\psi(x, y, z) = X(x)Y(y)Z(z)$. Substituting this solution to the differential equation, we can perform

some algebraic manipulation,

$$\begin{aligned}
-\frac{\hbar^2}{2m}\nabla^2\psi &= E\psi \\
\left(\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} + \frac{\partial^2}{\partial z^2}\right)X(x)Y(y)Z(z) &= -\frac{2m}{\hbar^2}EX(x)Y(y)Z(z) \\
Y(y)Z(z)\frac{\partial^2}{\partial x^2}X(x) + X(x)Z(z)\frac{\partial^2}{\partial y^2}Y(y) + X(x)Y(y)\frac{\partial^2}{\partial z^2}Z(z) &= -\frac{2m}{\hbar^2}EX(x)Y(y)Z(z) \\
\frac{1}{X(x)}\frac{\partial^2}{\partial x^2}X(x) + \frac{1}{Y(y)}\frac{\partial^2}{\partial y^2}Y(y) + \frac{1}{Z(z)}\frac{\partial^2}{\partial z^2}Z(z) &= -\frac{2m}{\hbar^2}E.
\end{aligned}$$

Now we can re-write this partial differential equation into three differential equations assuming that $E = \frac{\hbar^2}{2m}(k_x^2 + k_y^2 + k_z^2)$,

$$\begin{aligned}
\frac{d^2X(x)}{dx^2} &= -k_x^2X(x) \rightarrow X(x) = A_x \sin[k_x x] + B_x \cos[k_x x], \\
\frac{d^2Y(y)}{dy^2} &= -k_y^2Y(y) \rightarrow Y(y) = A_y \sin[k_y y] + B_y \cos[k_y y], \\
\frac{d^2Z(z)}{dz^2} &= -k_z^2Z(z) \rightarrow Z(z) = A_z \sin[k_z z] + B_z \cos[k_z z].
\end{aligned}$$

In order to find the expression for the coefficients A_n , B_n and k_n , we start by applying the boundary conditions. Since sin and cos are periodic functions, they satisfy $f(0) = f(a)$, however only the sin function satisfy the condition $f(0) = f(a) = 0$, hence, we set $B_x = B_y = B_z = 0$ leading to,

$$X(x) = A_x \sin[k_x x], \quad Y(y) = A_y \sin[k_y y], \quad Z(z) = A_z \sin[k_z z].$$

Now we recall the fact that x, y and z have units of distance and that the argument of the sin function must be dimensionless, combining this restriction with the property of periodicity we can define the constants k_n as, $k_x = n_x\pi/a$, $k_y = n_y\pi/a$, $k_z = n_z\pi/a$, where $(n_x, n_y, n_z) \in \mathbb{Z}^+$. With this information we can re-write the solution as,

$$\psi(x, y, z) = A_x A_y A_z \sin\left[\frac{n_x\pi}{a}x\right] \sin\left[\frac{n_y\pi}{a}y\right] \sin\left[\frac{n_z\pi}{a}z\right],$$

with

$$E = \frac{\pi^2 \hbar^2}{2ma^2}(n_x^2 + n_y^2 + n_z^2), \quad (n_x, n_y, n_z) \in \mathbb{Z}^+.$$

Finally, in order to get the expression for A_x, A_y and A_z we apply the normalization restriction to each spatial dimension,

$$\int_0^a A_l^2 \sin^2\left[\frac{n_l\pi}{a}s\right] ds = A_l^2 \frac{a}{4} \left(2 - \frac{1}{\pi n} \sin[2\pi n]\right) = 1,$$

since $n \in \mathbb{Z}^+$ we get that $A_l = \sqrt{2/a}$, therefore,

$$\psi(x, y, z) = \sqrt{\frac{8}{a^3}} \sin\left[\frac{n_x\pi}{a}x\right] \sin\left[\frac{n_y\pi}{a}y\right] \sin\left[\frac{n_z\pi}{a}z\right], \quad (n_x, n_y, n_z) \in \mathbb{Z}^+$$

Solution 2: Energy analysis

2 Problem 4.3

Use

$$P_l^m(x) \equiv (1-x^2)^{|m|/2} \left(\frac{d}{dx} \right)^{|m|} P_l(x)$$

$$P_l(x) \equiv \frac{1}{2^l l!} \left(\frac{d}{dx} \right)^l (x^2-1)^l$$

$$Y_l^m(\theta, \phi) = \epsilon \sqrt{\frac{(2l+1)}{4\pi} \frac{(l-|m|)!}{(l+|m|)!}} e^{im\phi} P_l^m(\cos[\theta])$$

to construct Y_0^0 and Y_2^l . Check that they are normalized and orthogonal.

3 Problem 4.13

- Find $\langle r \rangle$ and $\langle r^2 \rangle$ for an electron in the ground state of hydrogen. Express your answers in terms of the Bohr radius.
- Find $\langle x \rangle$ and $\langle x^2 \rangle$ for an electron in the ground state of hydrogen. *Hint:* this requires no new integration-note that $r^2 = x^2 + y^2 + z^2$, and exploit the symmetry of the ground state.
- Find $\langle x^2 \rangle$ in the state $n=2, l=1, m=1$. *Warning:* This state is not symmetrical in x, y, z . Use $x = r \sin \theta \cos \phi$.

4 Problem 4.14

What is the *most probable* value of r , in the ground state of hydrogen? (The answer is not zero!) *Hint:* First you must figure out the probability that the electron would be found between r and $r + dr$.

5 Problem 4.23

In problem 4.3 you showed that

$$Y_2^l(\theta, \phi) = -\sqrt{\frac{15}{8\pi}} \sin \theta \cos \theta e^{i\phi}.$$

Apply the raising operator to find $Y_2^2(\theta, \phi)$. Use equation $A_l^m = \hbar\sqrt{l(l+1) - m(m \pm 1)} = \hbar\sqrt{(l \mp m)(l \pm m + 1)}$ to get the normalization.

[heading=bibintoc,title=References]