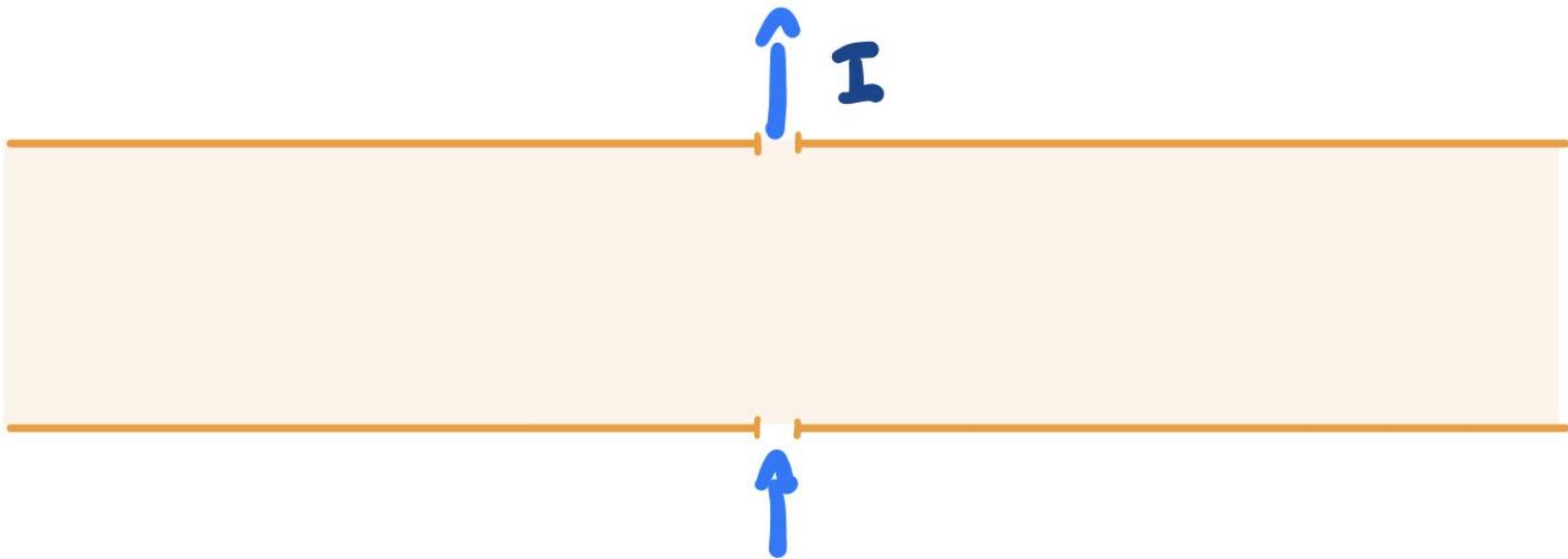


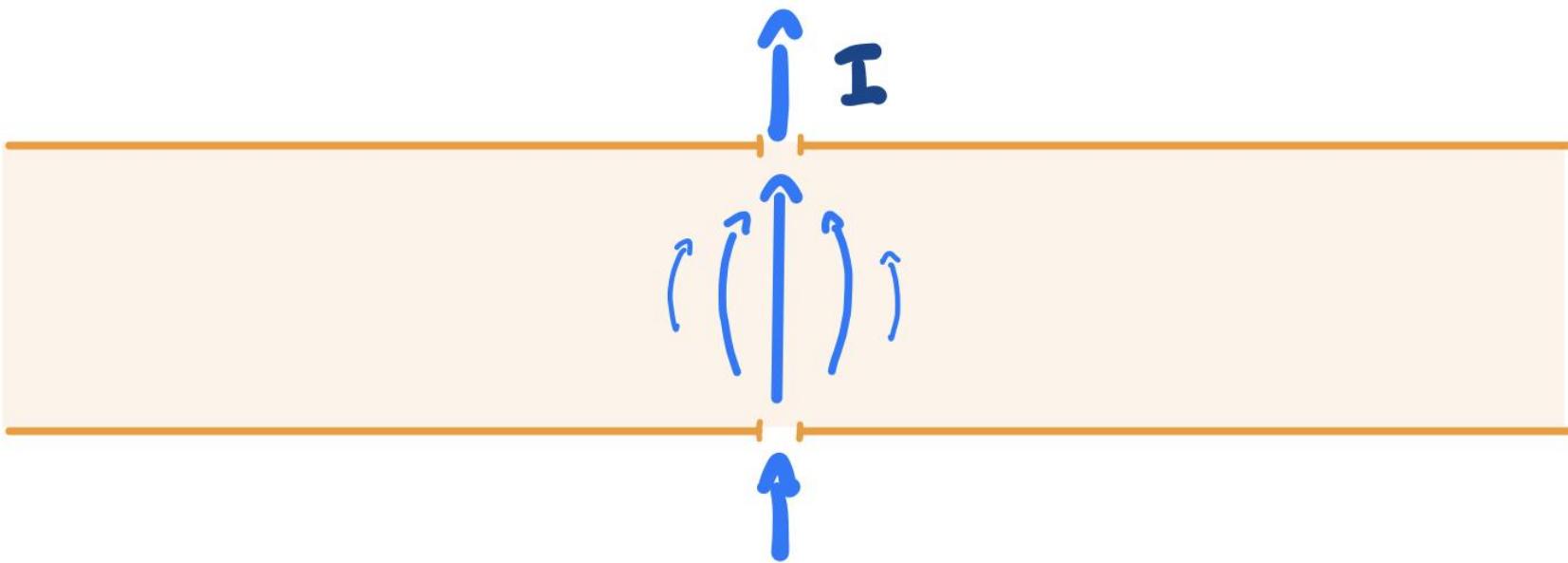
Trasporto non-locale di valle nel grafene a doppio strato

Relatore:
Marco Polini

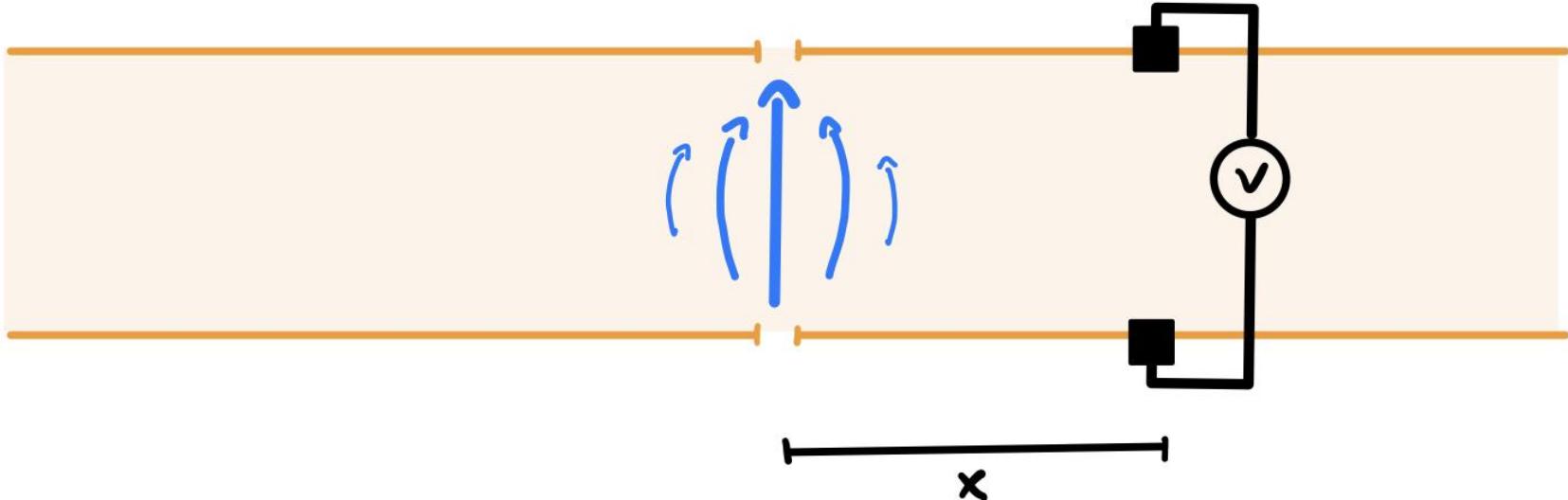
Candidato:
Francesco Sacco

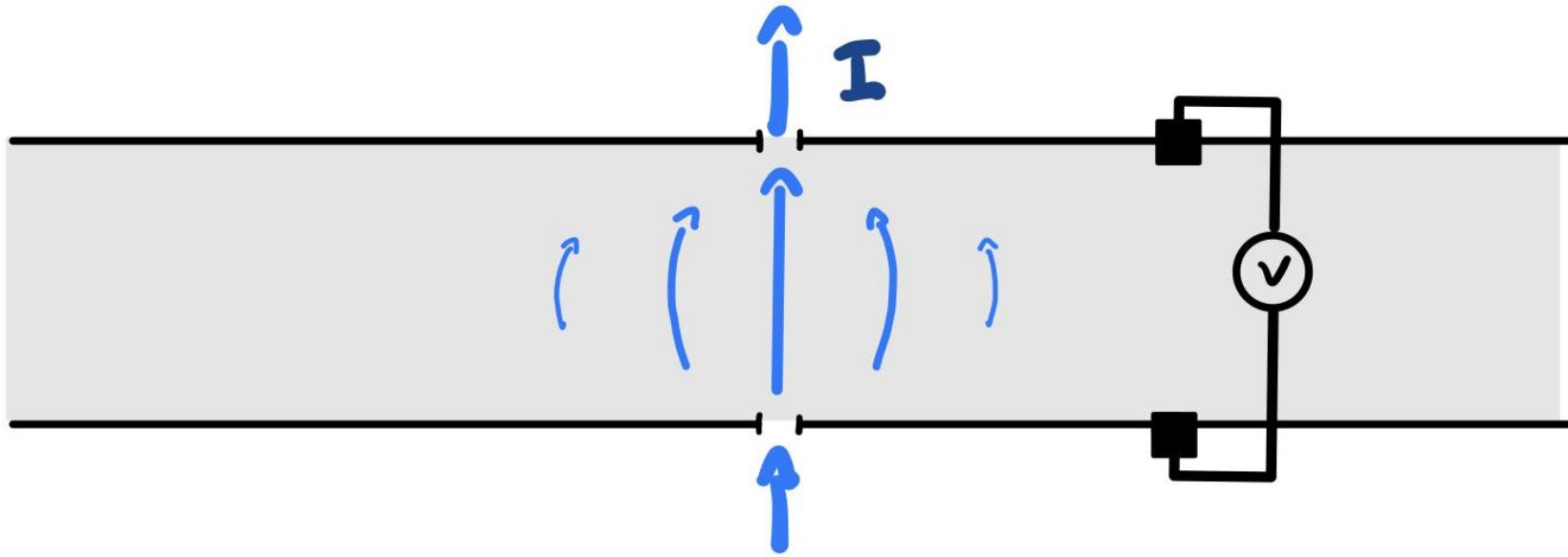




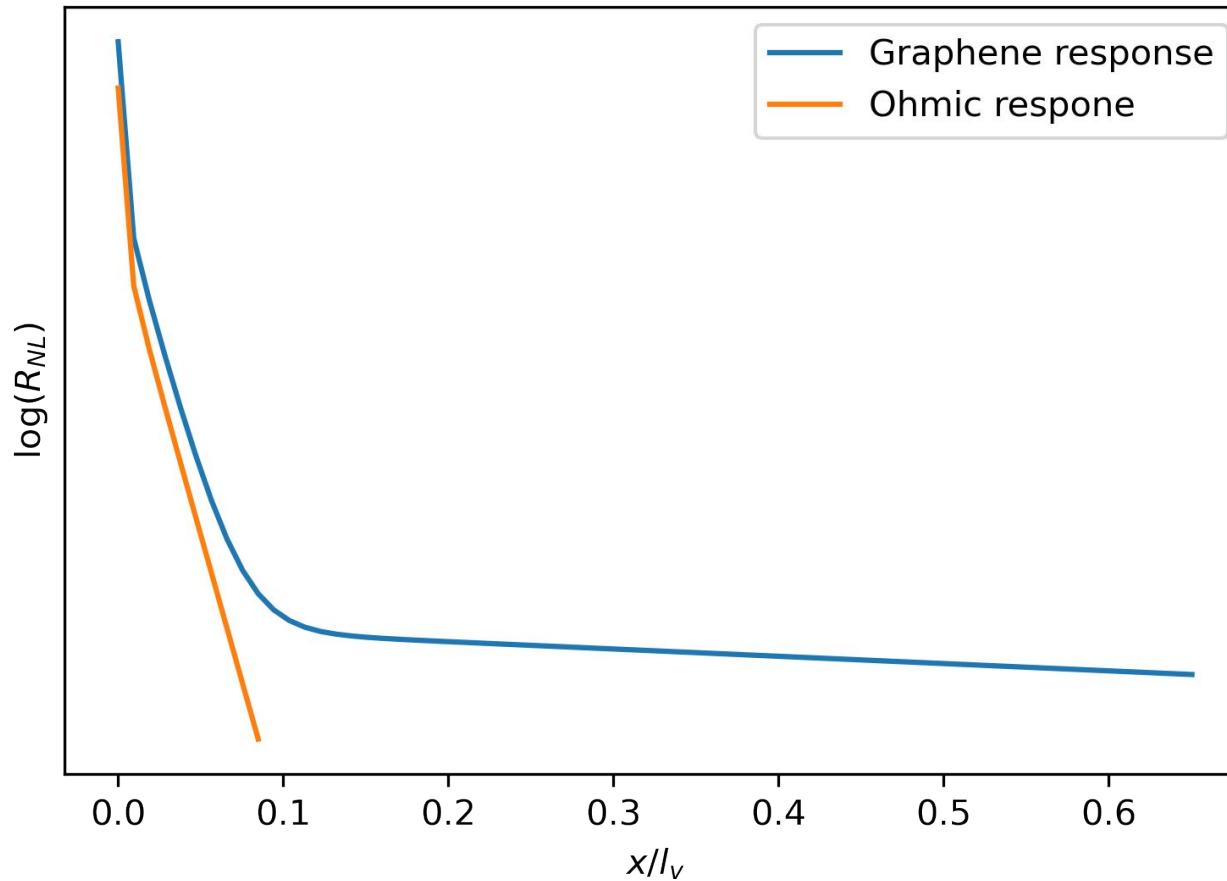


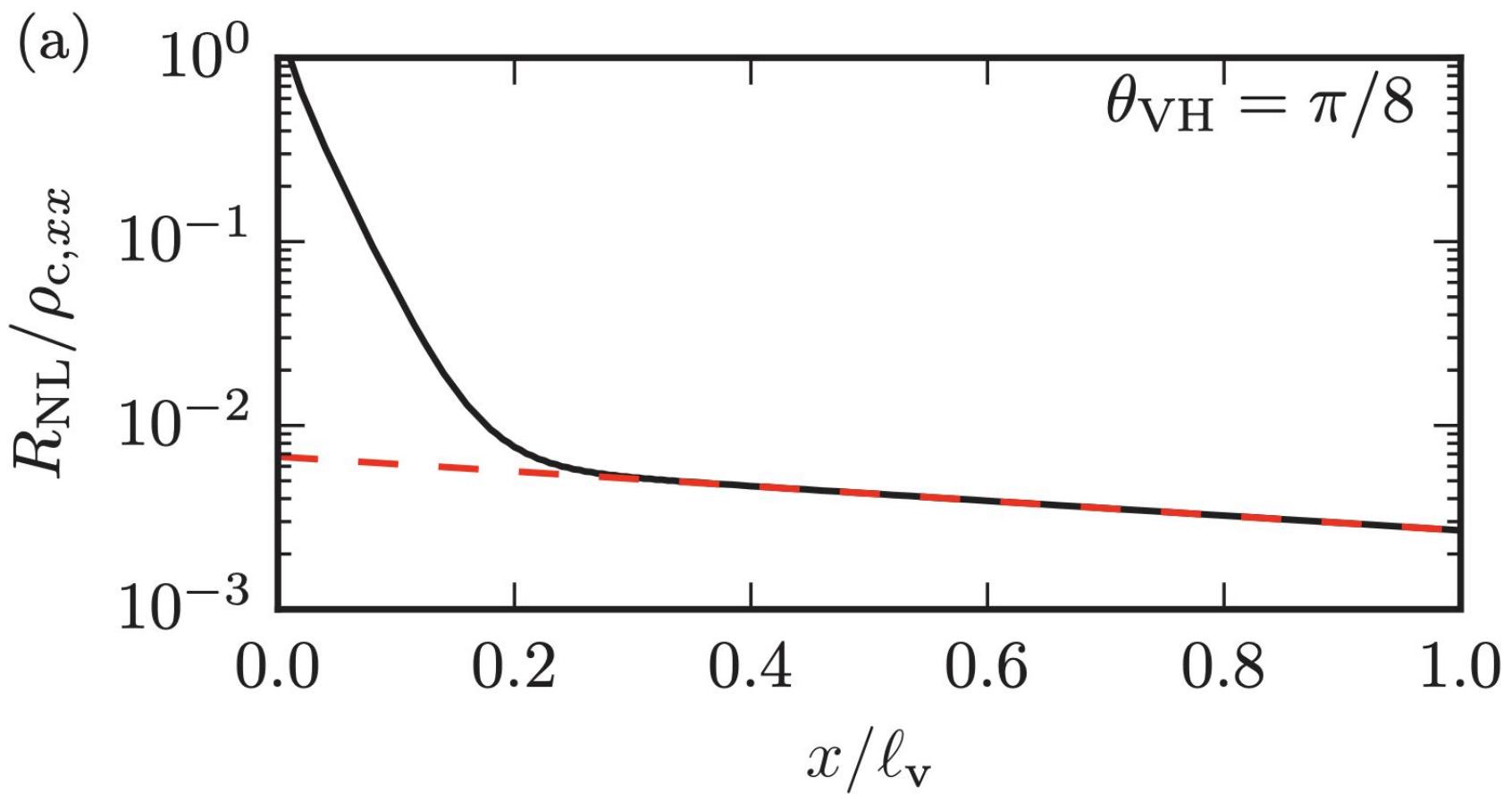
$$\frac{V(x)}{I} \equiv R_{\text{NL}}(x) = \frac{2\rho}{\pi} \ln \left| \coth \left(\frac{\pi x}{2W} \right) \right| \approx \frac{4\rho}{\pi} e^{-\pi|x|/W}$$



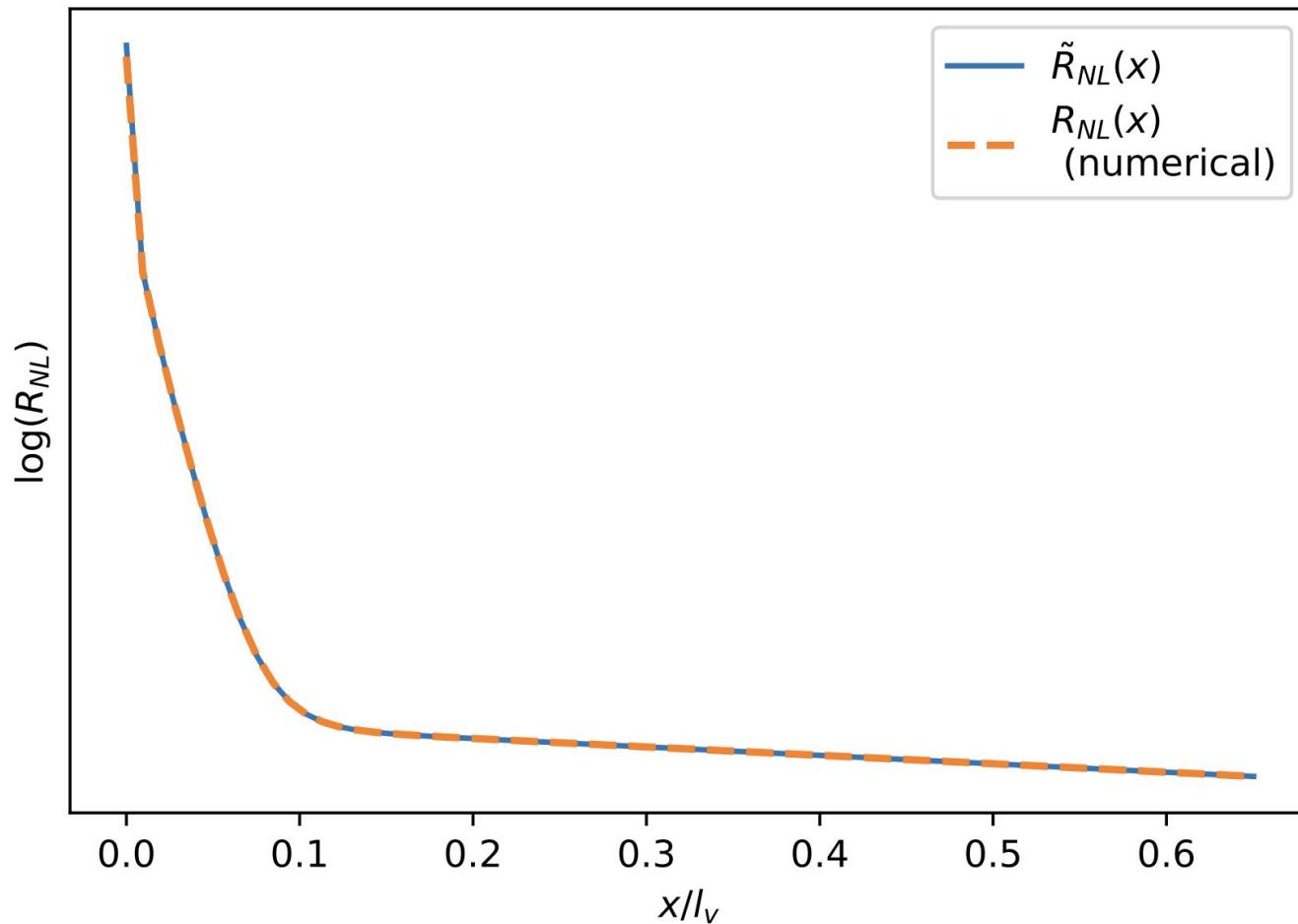


$$\frac{R_{xx}}{2L_v} e^{-|x|/L_v} \sin^2(\theta_{\text{VH}}) - \frac{2R_{xx}}{\pi W} \ln \left| \tanh \left(\frac{\pi x}{2W} \right) \right| \cos^2(\theta_{\text{VH}})$$

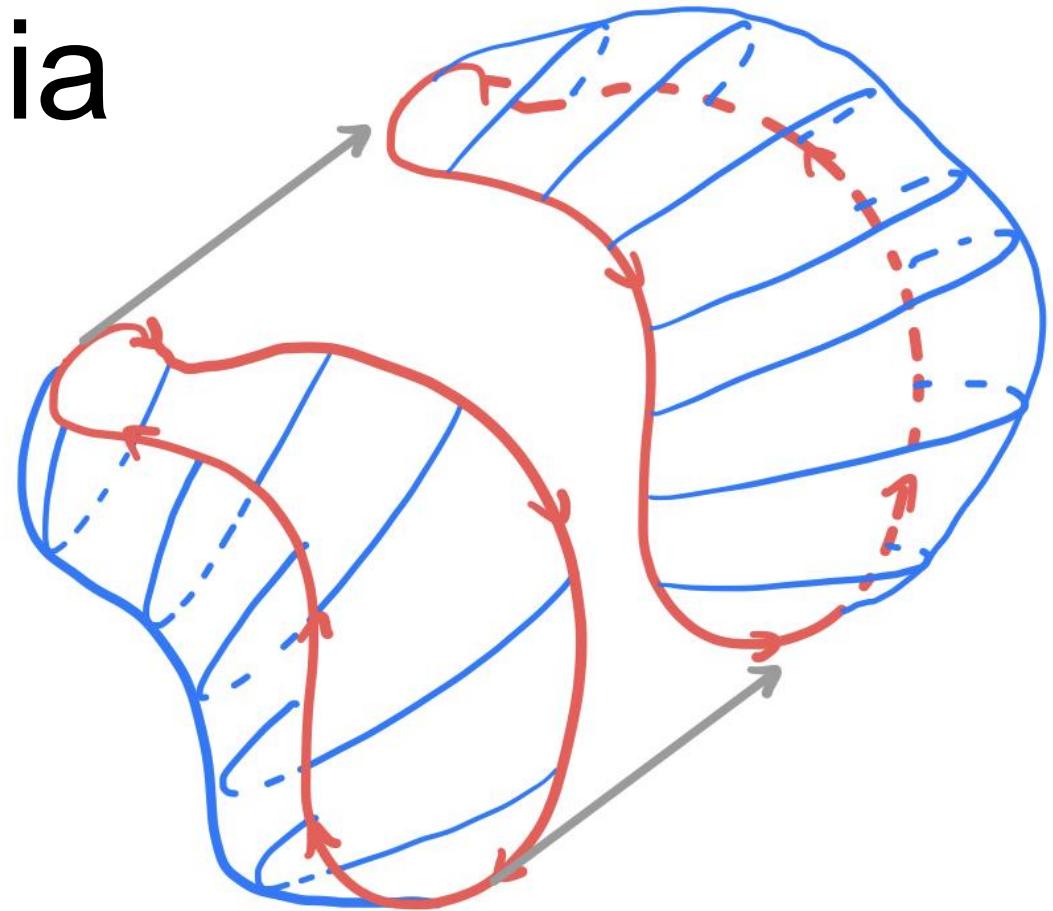




M. Beconcini, F. Taddei, and M. Polini, “Nonlocal topological valley transport at large valley hall angles,” *Physical Review B*, vol. **94**, no. 12, p. 121408, 2016.



Berryologia

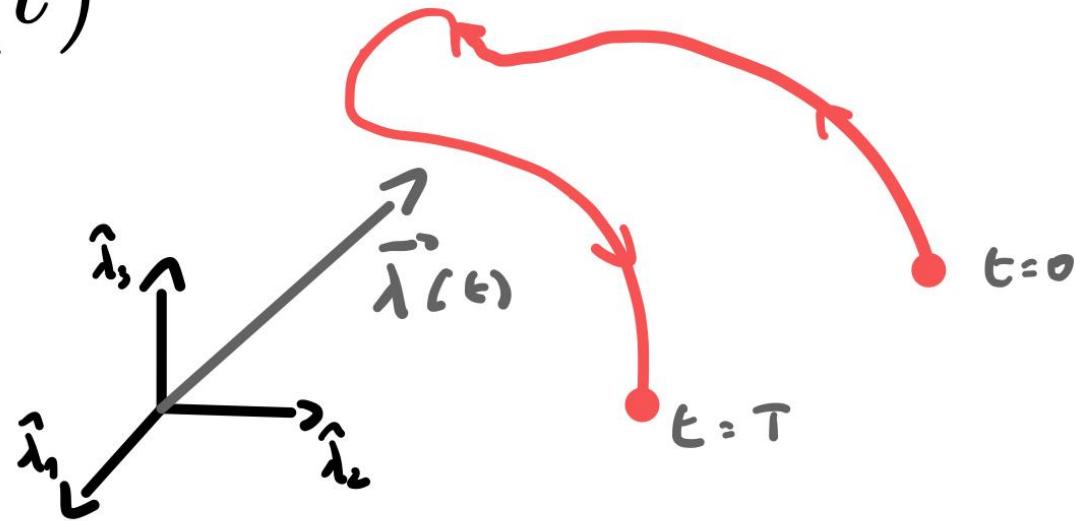


$$H(\boldsymbol{\lambda}) |n, \boldsymbol{\lambda}\rangle = E_n(\boldsymbol{\lambda}) |n, \boldsymbol{\lambda}\rangle$$

$$\boldsymbol{\lambda} = (\lambda_1, \lambda_2, \dots)$$

$$|n, \boldsymbol{\lambda}\rangle \rightarrow \underbrace{e^{i\gamma_n(\boldsymbol{\lambda})}}_{\text{Berry phase}} |n, \boldsymbol{\lambda}\rangle$$

$$\lambda = \lambda(t)$$



$$\lambda = \lambda(t) \quad |n, \lambda\rangle \rightarrow e^{i\gamma_n(\lambda(t))} |n, \lambda\rangle$$

$$|\psi_n(t)\rangle = \underbrace{e^{i\gamma_n(\lambda(t))}}_{\text{Berry phase}} \cdot \underbrace{e^{-\frac{i}{\hbar} \int_0^t E_n(\lambda(t')) dt'}}_{\text{dynamical phase}} |n, \lambda(t)\rangle,$$

$$\dot{\gamma}_n(t) = i \langle n, t | \partial_t | n, t \rangle$$

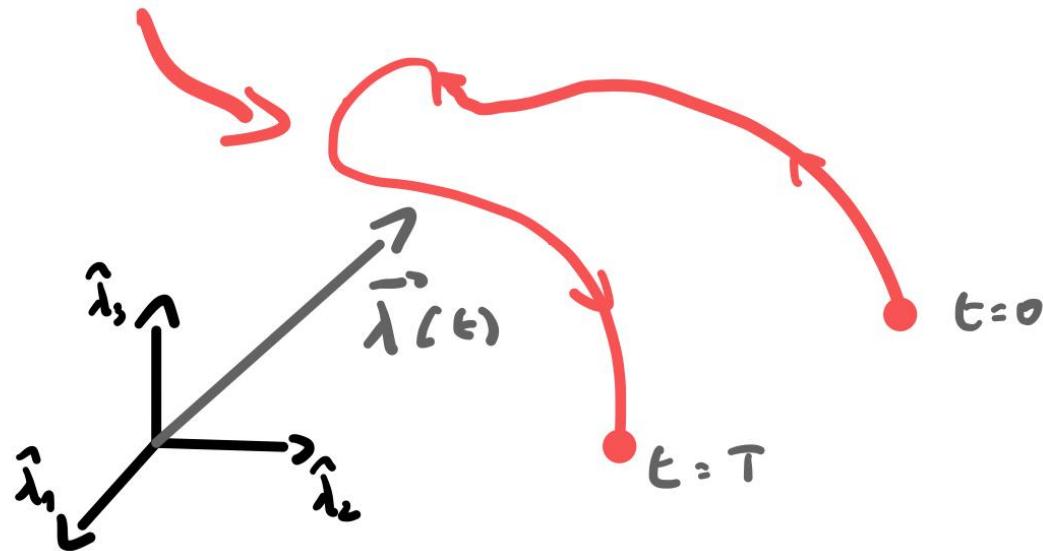
$$\dot{\gamma}_n(t) = \dot{\boldsymbol{\lambda}} \cdot \boldsymbol{A}_n(\boldsymbol{\lambda})$$

$$\boldsymbol{A}_n(\boldsymbol{\lambda}) = i \langle n, t | \partial_{\boldsymbol{\lambda}} | n, t \rangle$$

$$\dot{\gamma}_n(t) = \dot{\lambda} \cdot A_n(\lambda)$$

$$\gamma_n = \int_{\mathcal{P}} \mathbf{A}_n(\lambda) \cdot d\lambda$$

$$\vec{x}_n = \int_{\lambda} \vec{A}_n(\vec{\lambda}) \cdot d\vec{\lambda}$$

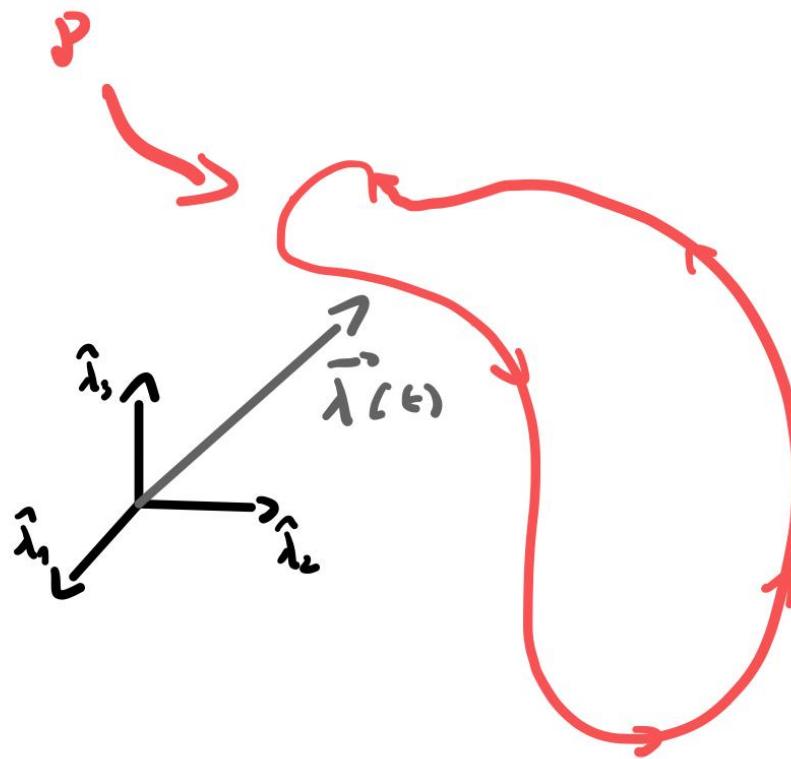


$$|n,\lambda\rangle\rightarrow e^{if_n(\lambda)}\,|n,\lambda\rangle$$

$$\mathbf{A}_n \rightarrow \mathbf{A}_n - \partial_{\pmb{\lambda}} f_n$$

$$\gamma_n = \int_{\mathcal{P}} A_n(\pmb{\lambda}) \cdot d\pmb{\lambda}$$

$$\delta_n = \oint \vec{A}_n(\vec{\lambda}) \cdot d\vec{\lambda}$$



$$\gamma_n = \oint_{\mathcal{P}} \mathbf{A}_n(\lambda) \cdot d\lambda + 2n\pi$$

$$\Omega_{\mu\nu}^n = \partial_\mu A_\nu^n(\lambda) - \partial_\nu A_\mu^n(\lambda)$$

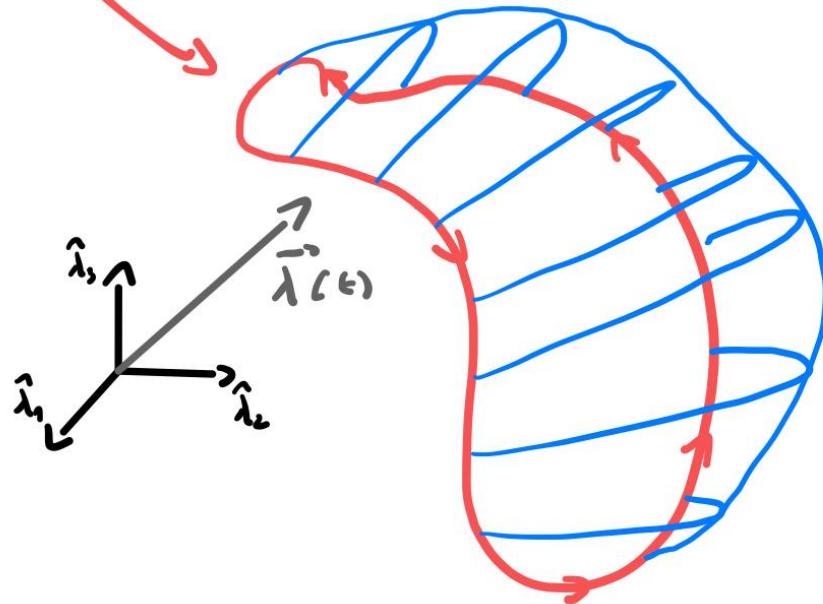
Berry Curvature

$$\Omega_{\mu\nu}^n = \partial_\mu A_\nu^n(\lambda) - \partial_\nu A_\mu^n(\lambda)$$

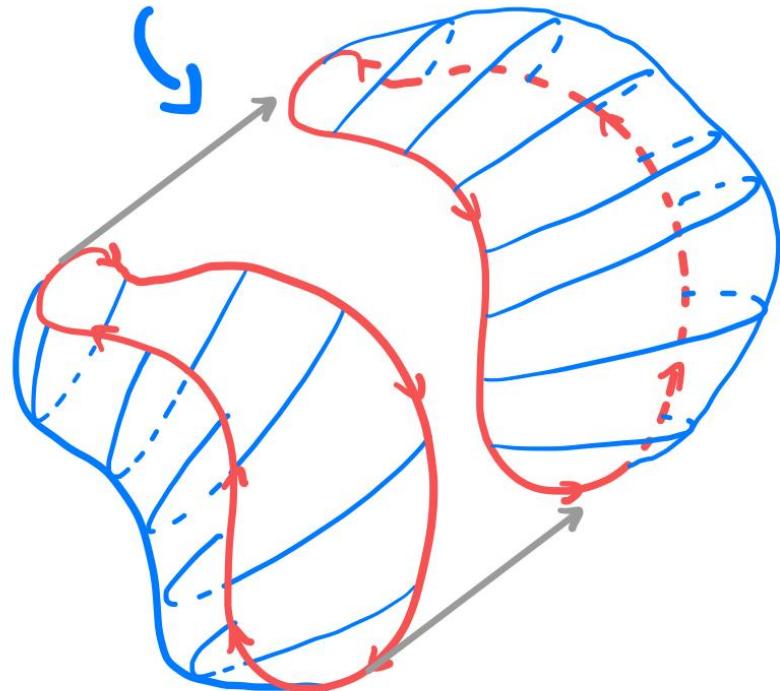
$$\Omega_{\mu\nu}^n = i \sum_{n' \neq n} \frac{\langle n | \partial_\mu H | n' \rangle \langle n' | \partial_\nu H | n \rangle}{(E_{n'} - E_n)^2} - (\mu \leftrightarrow \nu)$$

$$\mathcal{J}_n = \oint \vec{\lambda}_n(\vec{\lambda}) \cdot d\vec{\lambda} = \frac{1}{2} \int_S \Omega_{uv}^n d\lambda^u d\lambda^v$$

Curva di integrazione } Superficie di integrazione



$$\oint_s \Omega_{\alpha\nu}^n d\lambda^\alpha \wedge d\lambda^\nu$$



$$\oint_{\mathcal{S}} \Omega^n_{\mu\nu} d\lambda^\mu \wedge d\lambda^\nu = 2\pi C_n \quad \quad C_n \in \mathbb{Z}$$

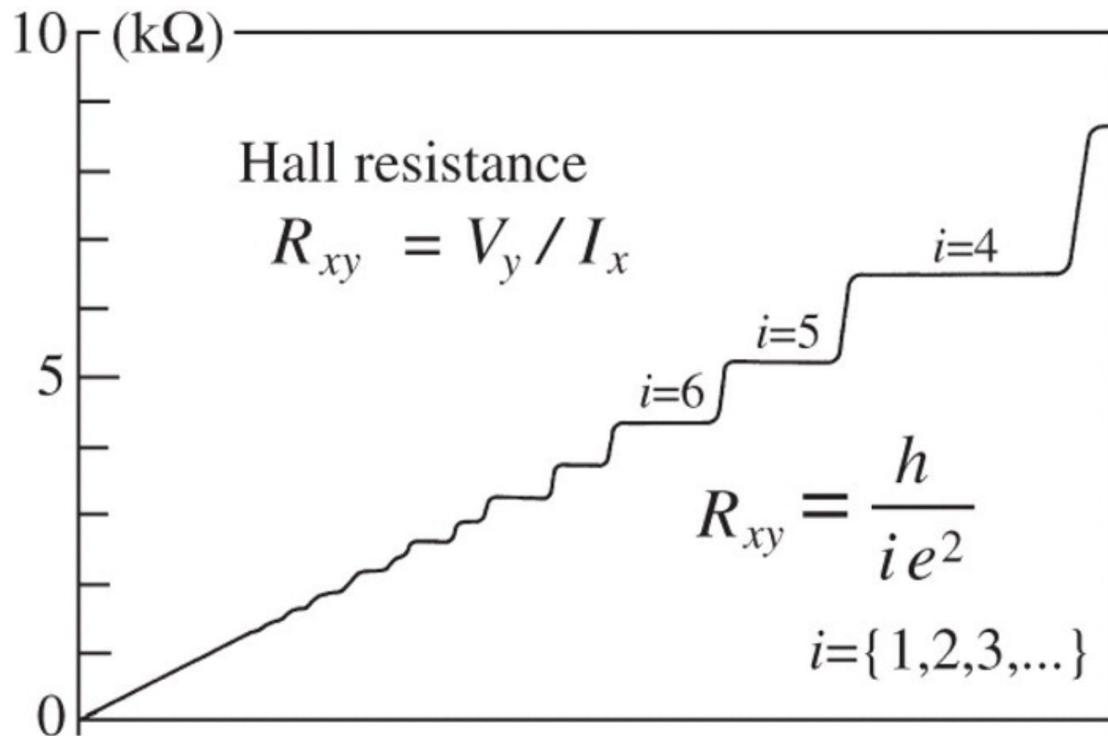
$$\sigma_{xy} = i \frac{e^2}{\hbar} \sum_{m, E_n < E_F} \int_{\mathcal{B}} \frac{\langle n, \mathbf{k} | \partial_y H | m, \mathbf{k} \rangle \langle m, \mathbf{k} | \partial_x H | n, \mathbf{k} \rangle - (n \leftrightarrow m)}{E_m(\mathbf{k}) - E_n(\mathbf{k})} \frac{d^2 \mathbf{k}}{(2\pi)^2}$$

$$\sigma_{xy} = \frac{e^2}{\hbar} \sum_{E_n < E_F} \int_{\mathcal{B}} \Omega_{k_x k_y}^n \frac{d^2 \mathbf{k}}{(2\pi)^2}$$

D. J. Thouless, M. Kohmoto, M. P. Nightingale, and M. den Nijs,
 “Quantized hall conductance in a two-dimensional periodic potential,”
Physical review letters, vol. **49**, no. 6, p. 405, 1982.

$$\sigma_{xy} = \frac{e^2}{2\pi\hbar} \sum_n C_n$$

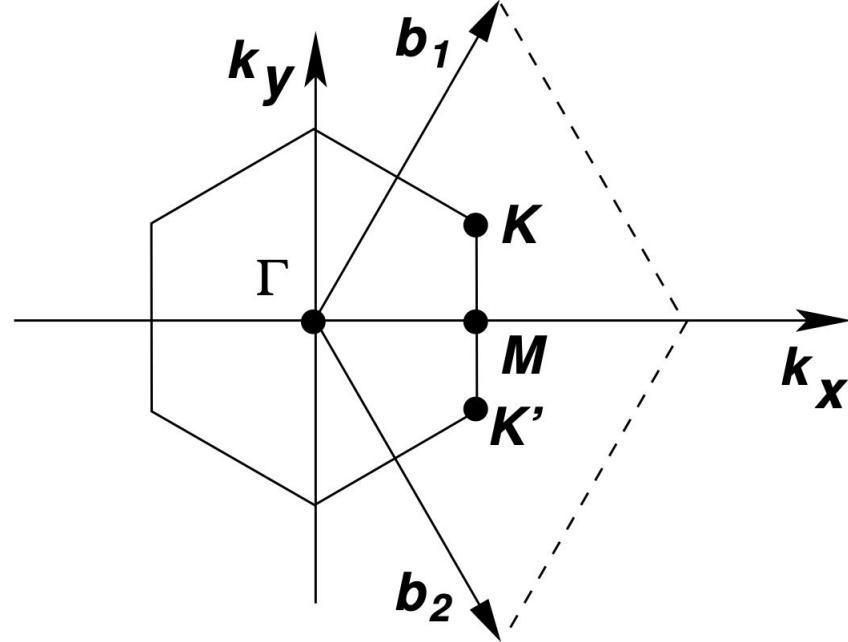
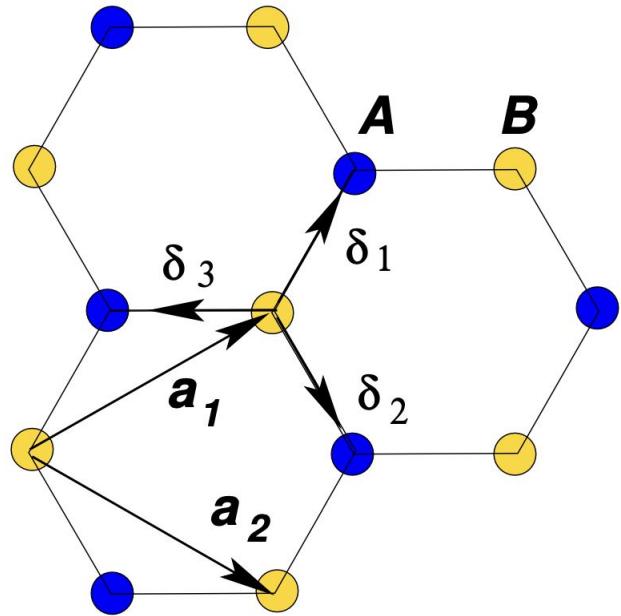
Quantized conductivity

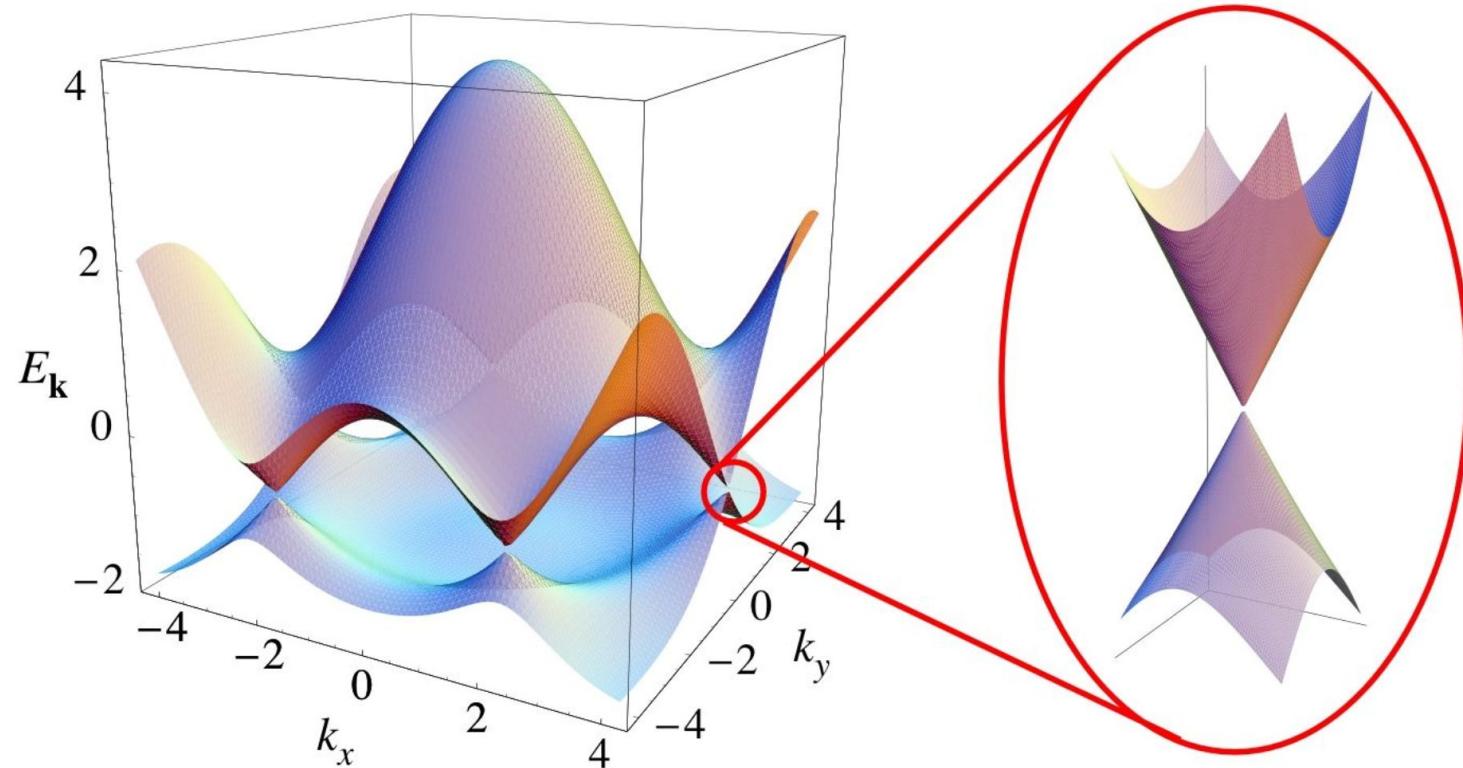


$$t \rightarrow -t \quad \Omega(\mathbf{k}) \rightarrow -\Omega(-\mathbf{k})$$

$$r \rightarrow -r \quad \Omega(\mathbf{k}) \rightarrow \Omega(-\mathbf{k})$$

Grafene





A. C. Neto, F. Guinea, N. M. Peres, K. S. Novoselov, and A. K. Geim,
“The electronic properties of graphene,” *Reviews of modern physics*,
vol. **81**, no. 1, p. 109, 2009.

$$H_{\boldsymbol{K}_0}(\boldsymbol{k}) = -H_{\boldsymbol{K}_1}^*(\boldsymbol{k}) = \begin{bmatrix} 0 & v_F\hbar(k_x - ik_y) \\ v_F\hbar(k_x + ik_y) & 0 \end{bmatrix}$$

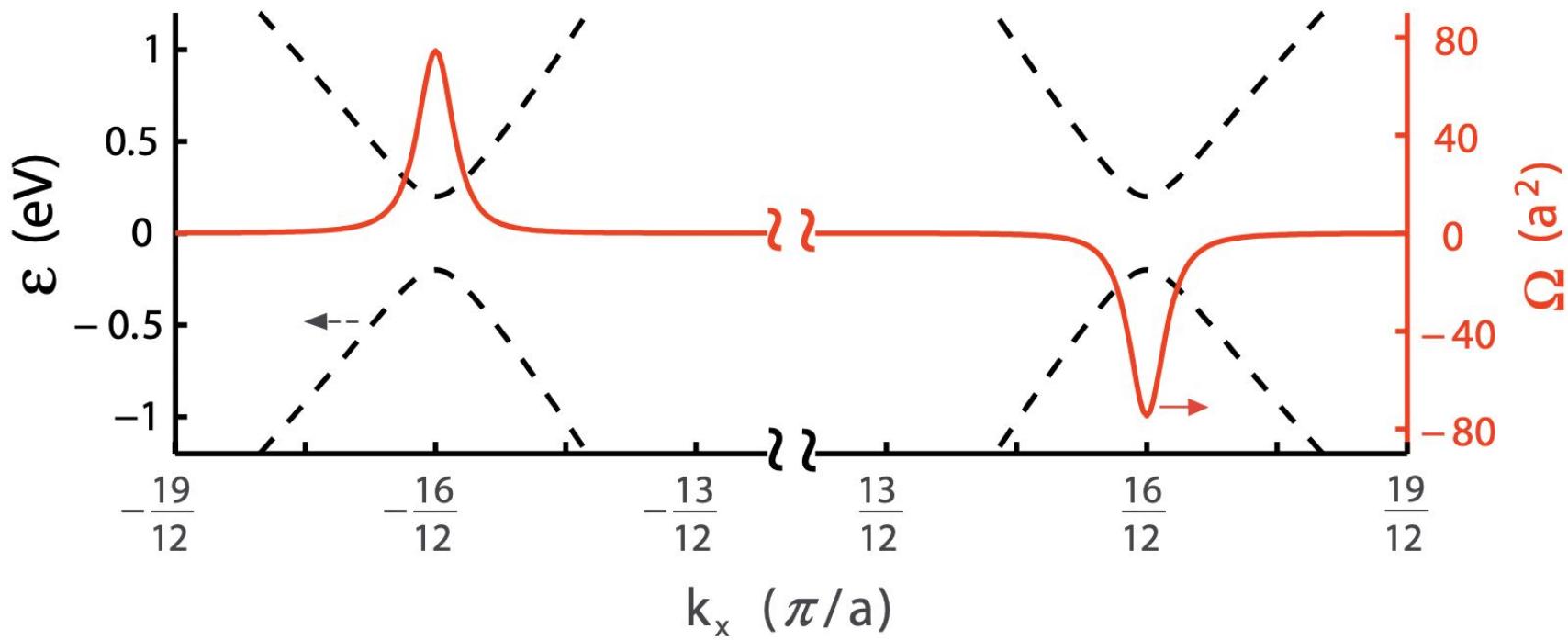
$$H(\boldsymbol{k})=\hbar v_{\mathrm{F}}(\sigma_x\tau_zk_x+\sigma_yk_y)$$

$$H(\mathbf{k}) = \begin{bmatrix} \Delta & v_F \hbar (\tau_z k_x - i k_y) \\ v_F \hbar (\tau_z k_x + i k_y) & -\Delta \end{bmatrix}$$

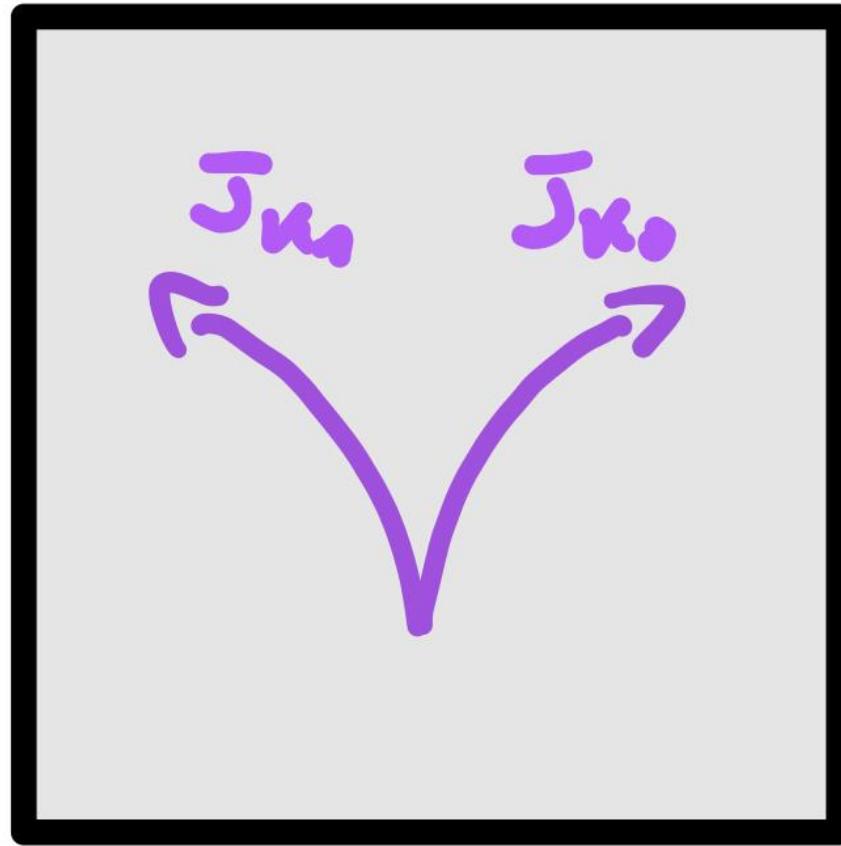
$$H(\mathbf{k}) = \hbar v_F (\sigma_x \tau_z k_x + \sigma_y k_y) + \Delta \sigma_z$$

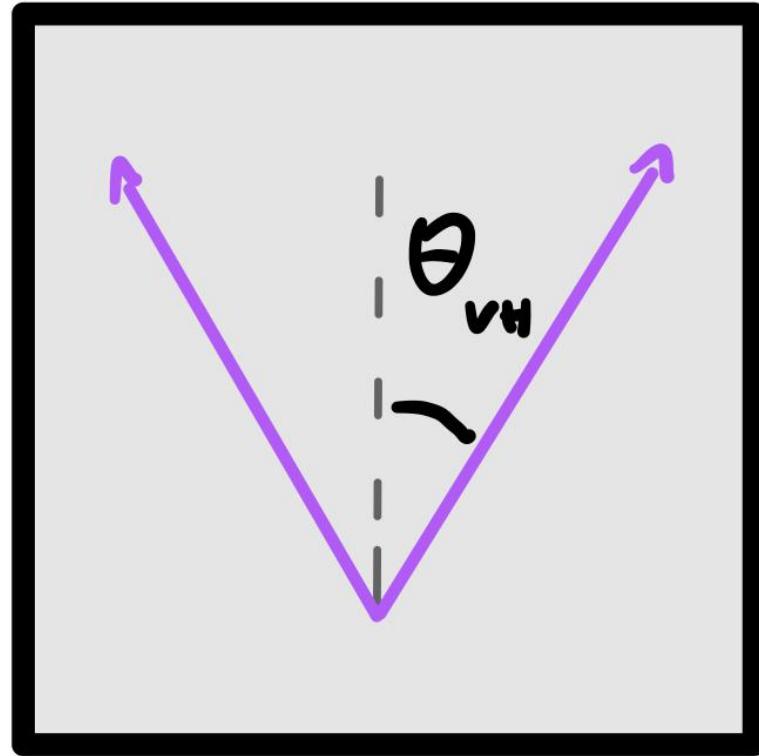
$E(k)$ k

- $\Delta = 0$
- $\Delta \neq 0$



E ↑





$$\nabla \psi_{K_\alpha}(\mathbf{r}) = \nabla V(\mathbf{r}) - \frac{1}{e}\frac{\partial}{\partial n_{K_\alpha}}\mu_\alpha(n_{K_\alpha}(\mathbf{r}))\nabla n_{K_\alpha}$$

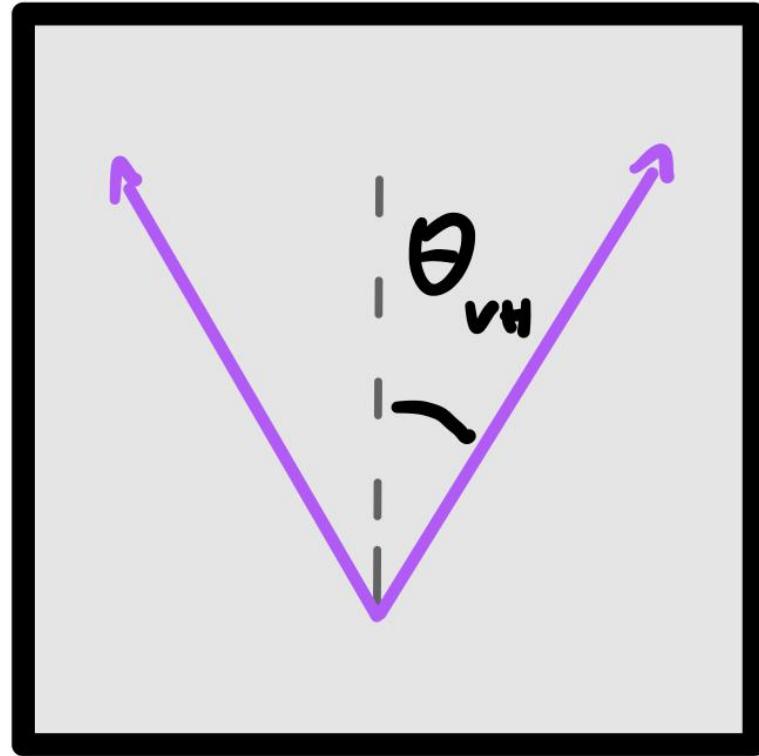
$$-e\mathbf{J}_{K_\alpha}(\mathbf{r})=\sigma_{K_\alpha}(\mathbf{r})\nabla\psi_{K_\alpha}(\mathbf{r})$$

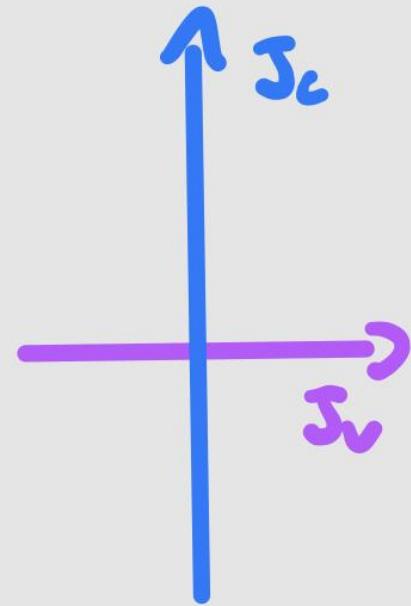
$$\sigma_{\mathbf{K}_0,ij} = \sigma_{xx} \begin{bmatrix} 1 & \tan(\theta_{\text{VH}}) \\ -\tan(\theta_{\text{VH}}) & 1 \end{bmatrix}$$

$$\sigma_{\mathbf{K}_1,ij} = \sigma_{xx} \begin{bmatrix} 1 & -\tan(\theta_{\text{VH}}) \\ \tan(\theta_{\text{VH}}) & 1 \end{bmatrix}$$

$$-eJ_{K_\alpha,i}(\mathbf{r})=\sigma_{K_\alpha,ij}E_j(\mathbf{r})-eD_{K_\alpha,ij}\partial_jn_{K_\alpha}(\mathbf{r})$$

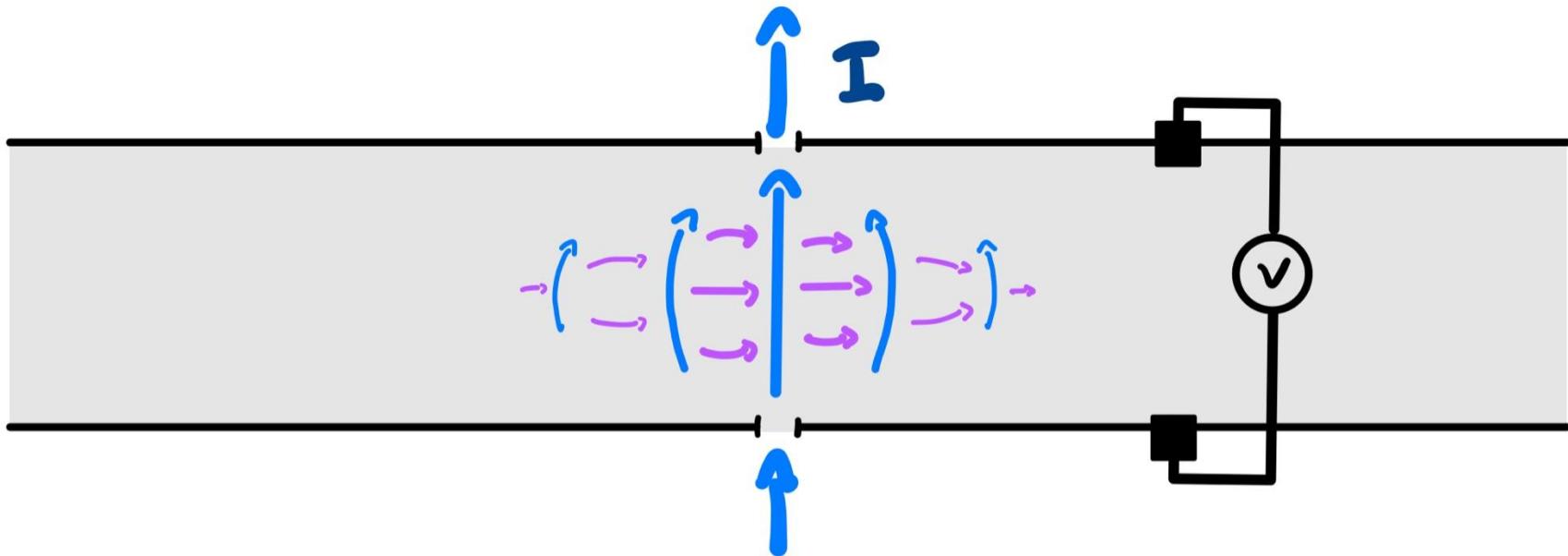
$$e^2 D_{K_\alpha,ij} = \sigma_{K_\alpha,ij} \frac{\partial \mu_\alpha}{\partial n_{K_\alpha}}[n_{K_\alpha}(\mathbf{r})]$$





$$J_{c,i}(\mathbf{r}) = \sum_j \sigma_{c,xx} \delta_{ij} E_i(\mathbf{r}) + e D_{cv,xy} \epsilon_{ij} \partial_j n_v(\mathbf{r})$$

$$J_{v,i}(\mathbf{r}) = \sum_j \sigma_{c,xy} \epsilon_{ij} E_j(\mathbf{r}) + e D_{v,xx} \delta_{ij} \partial_j n_v(\mathbf{r})$$

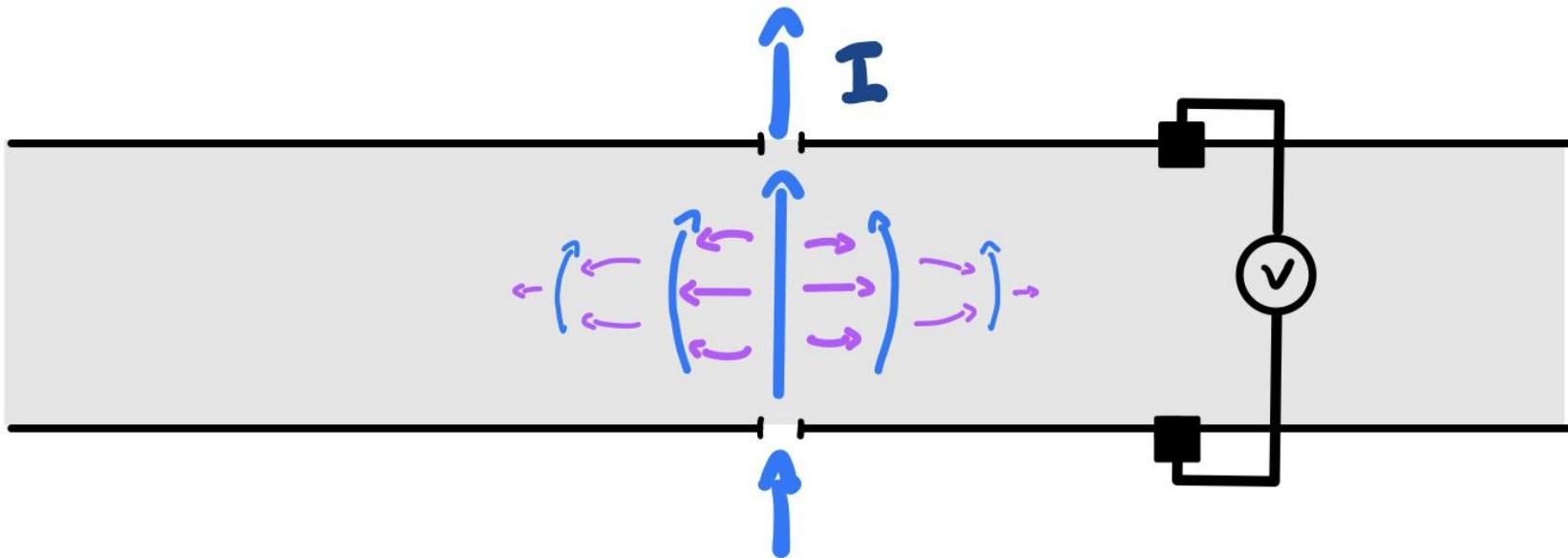


$$\nabla^2 J_c = 0$$

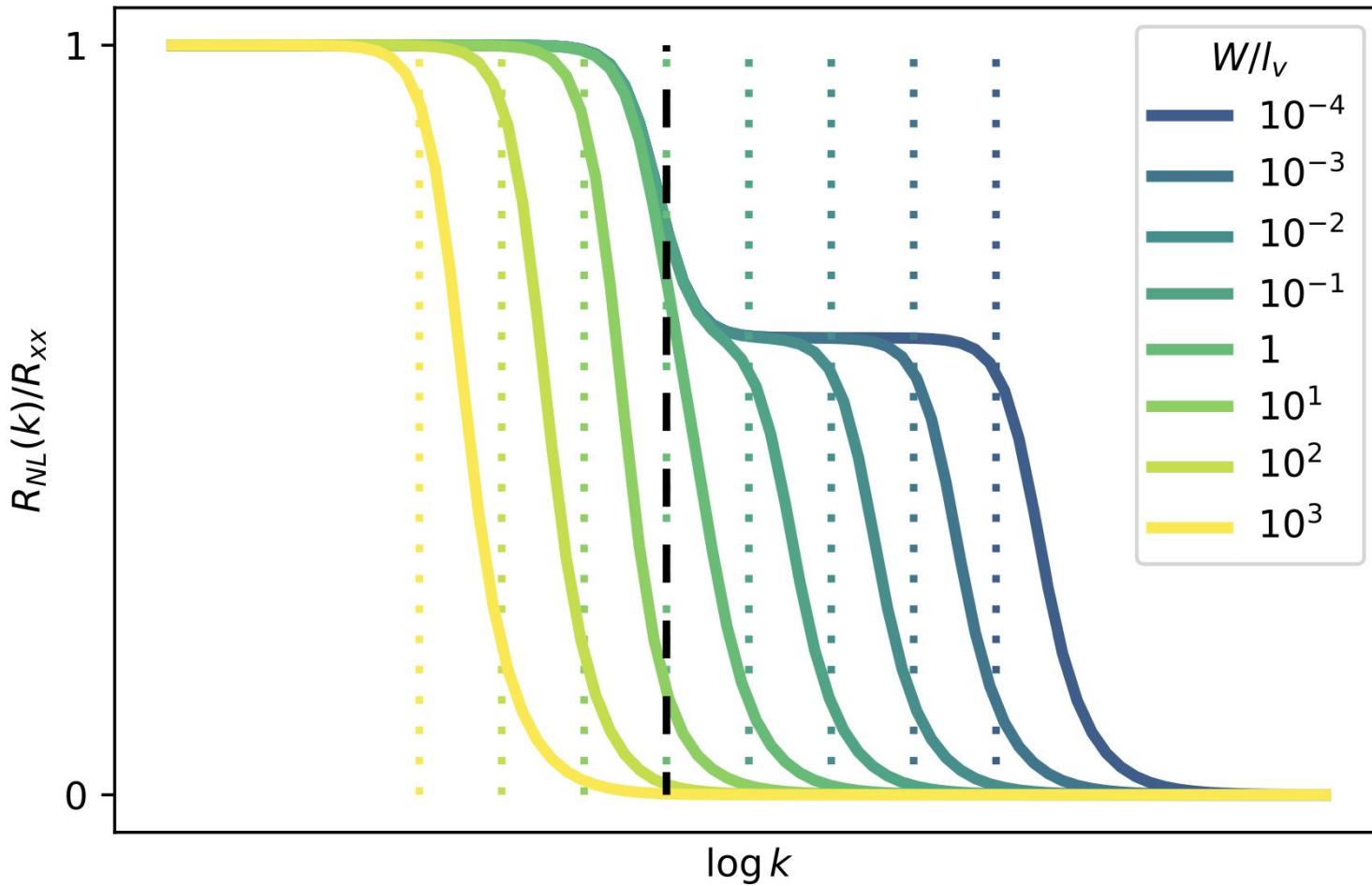
$$\nabla^2 J_v = \frac{\partial n_v}{\partial t}$$

$$\nabla^2 V(x,y) = 0$$

$$\left(\nabla^2 - l_{\rm v}^{-2}\right)n_v(r)=0$$

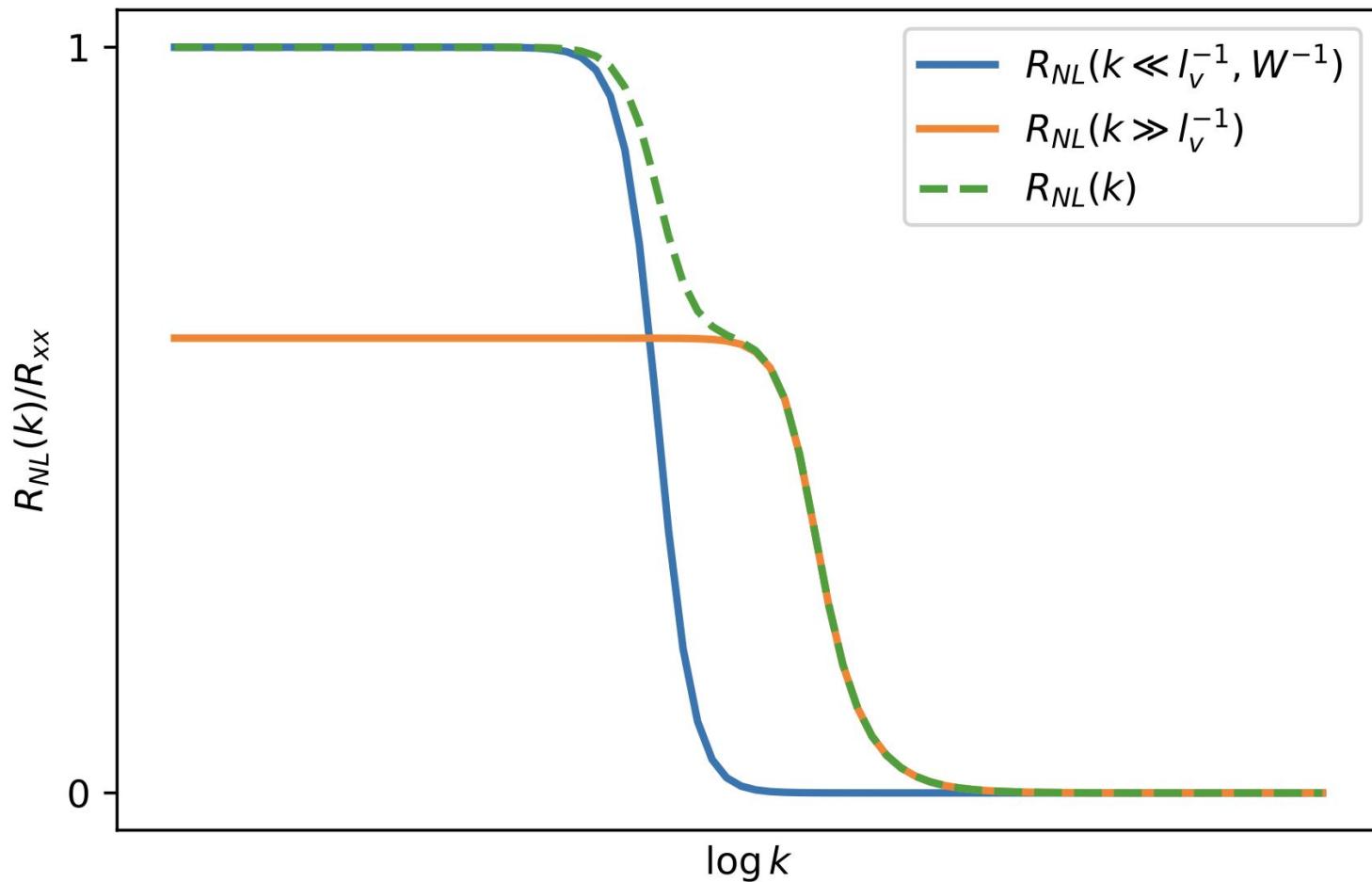


$$R_{\text{NL}}(k) = \frac{2\omega(k)}{k\sigma_c} \left\{ \frac{\omega(k)}{\tanh(kW/2)} + \frac{k\tan^2(\theta_{\text{VH}})}{\tanh[\omega(k)W/2]} \right\}^{-1}$$



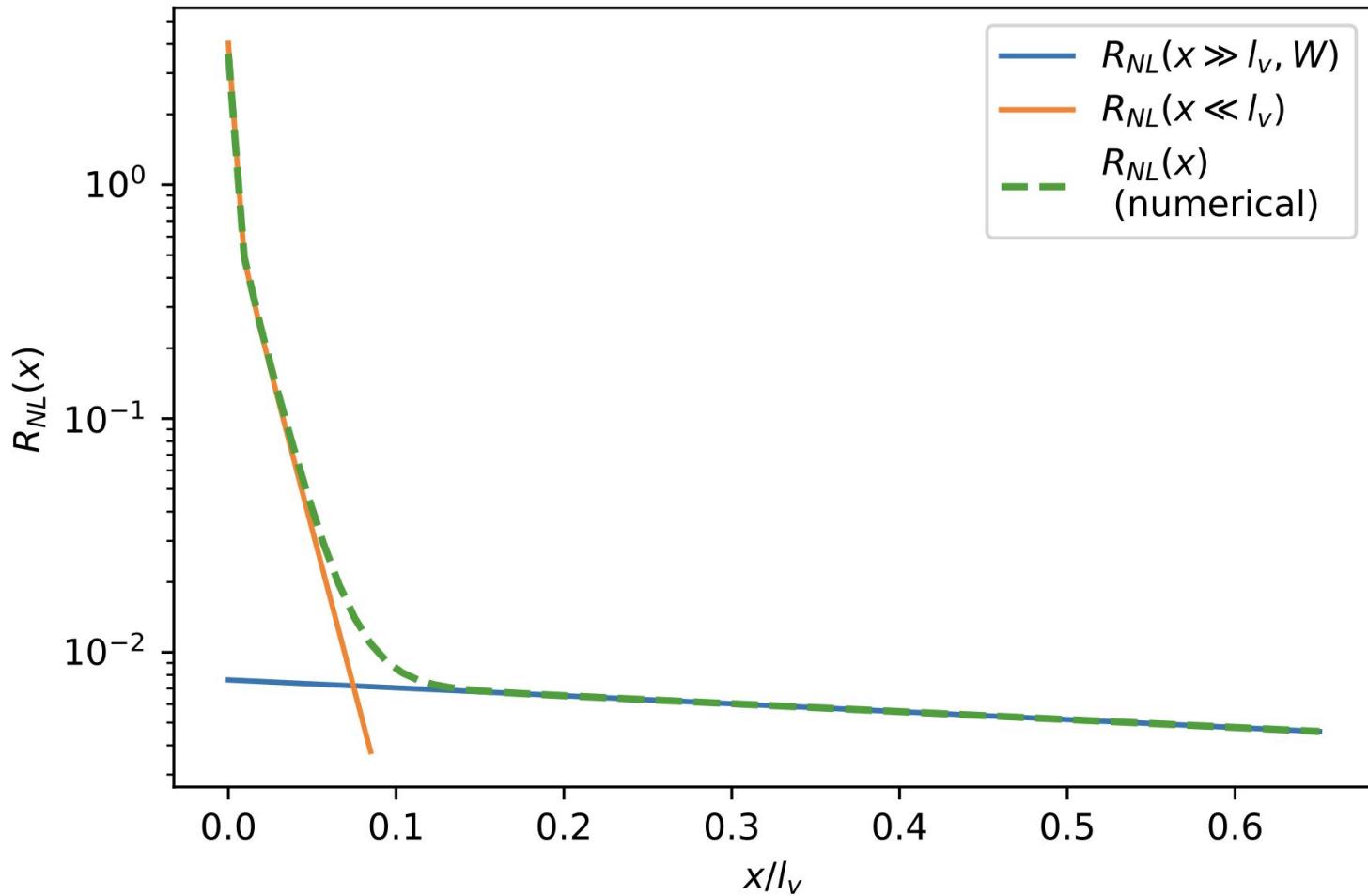
$$R_{\rm NL}(k \gg l_{\rm v}^{-1}) \approx \cos^2(\theta_{\rm VH}) R_{\rm NL}^{(0)}(k)$$

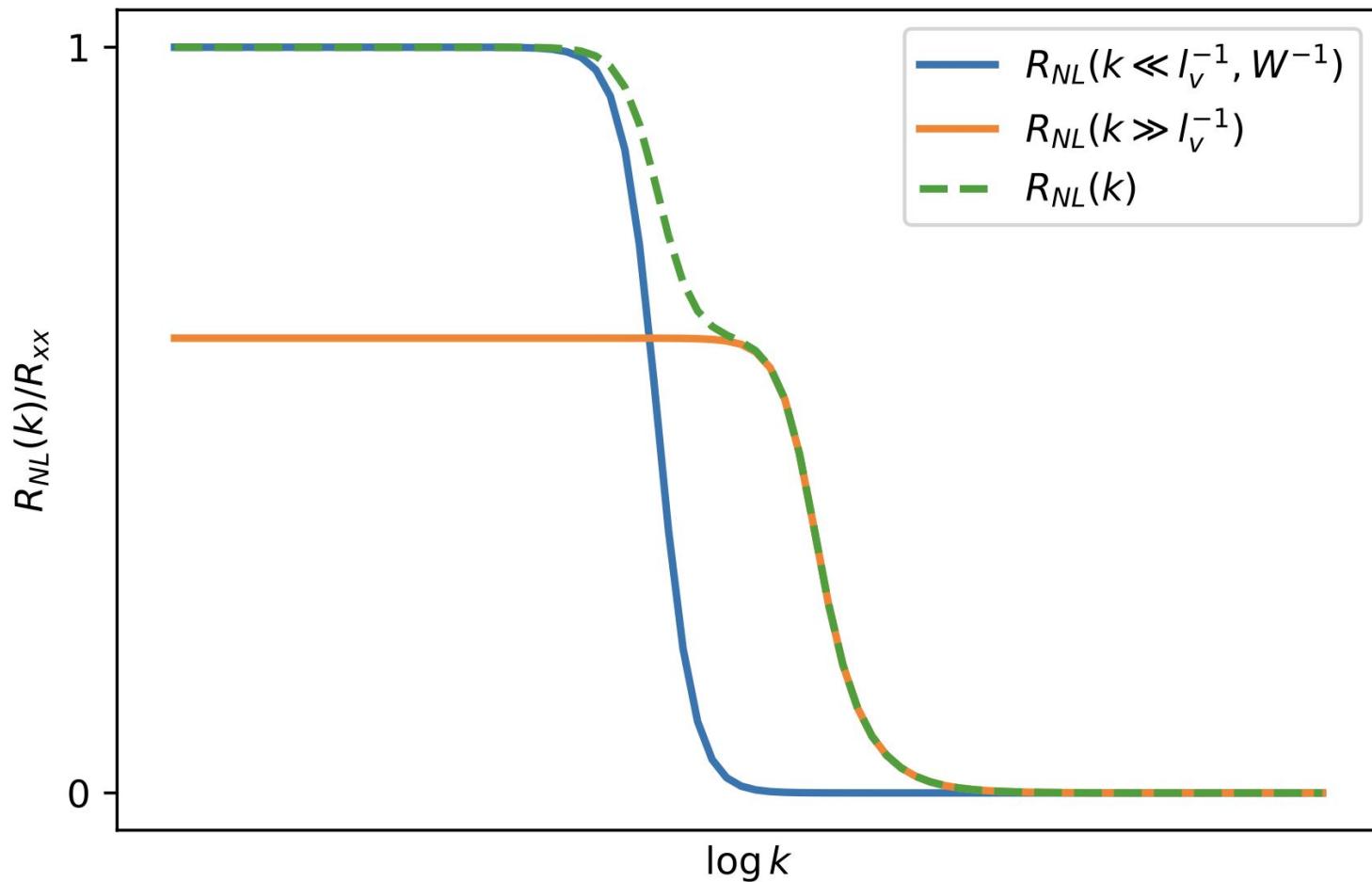
$$R_{\rm NL}(k \ll l_{\rm v}^{-1}, W^{-1}) \approx R_{\rm NL}^{\rm T}(k) = \frac{R_{xx}}{1 + L_{\rm v}^2 k^2}$$



$$\cos^2(\theta_{\rm VH}) R_{\rm NL}^0(x) = - \cos^2(\theta_{\rm VH}) \frac{2\rho}{\pi} \ln \left| \tanh \left(\frac{\pi x}{2W} \right) \right|$$

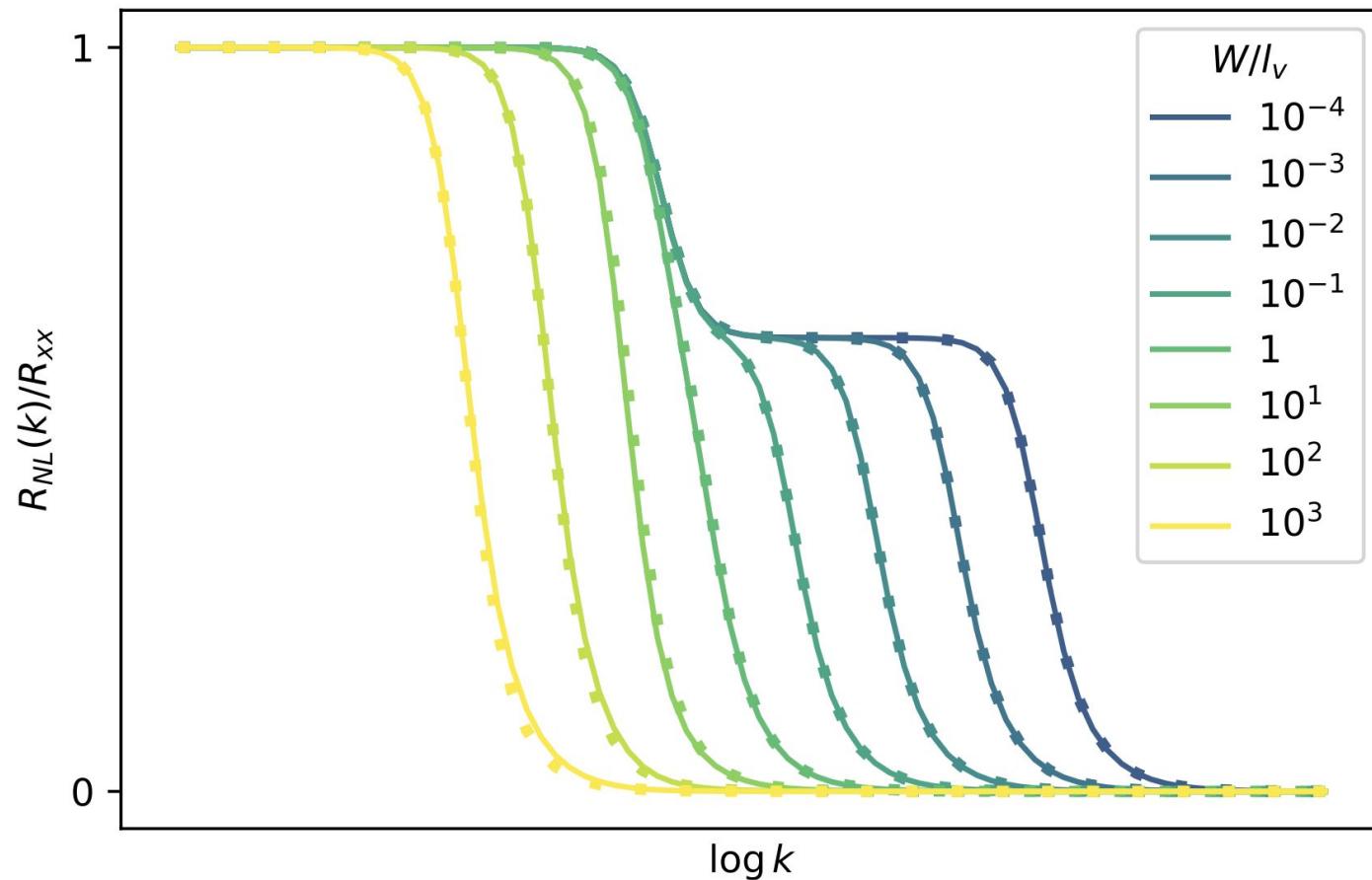
$$R_{\rm NL}^T(x)=R_{xx}\int_{-\infty}^{+\infty}\frac{e^{-ikx}}{1+L_{\rm v}^2k^2}\frac{dk}{2\pi}=\frac{R_{xx}}{2L_{\rm v}}e^{-\frac{|x|}{L_{\rm v}}}$$





$$R_{\rm NL}(k) \approx \alpha R_{\rm NL}(k \ll l_{\rm v}^{-1}, W^{-1}) + \beta R_{\rm NL}(k \gg l_{\rm v}^{-1})$$

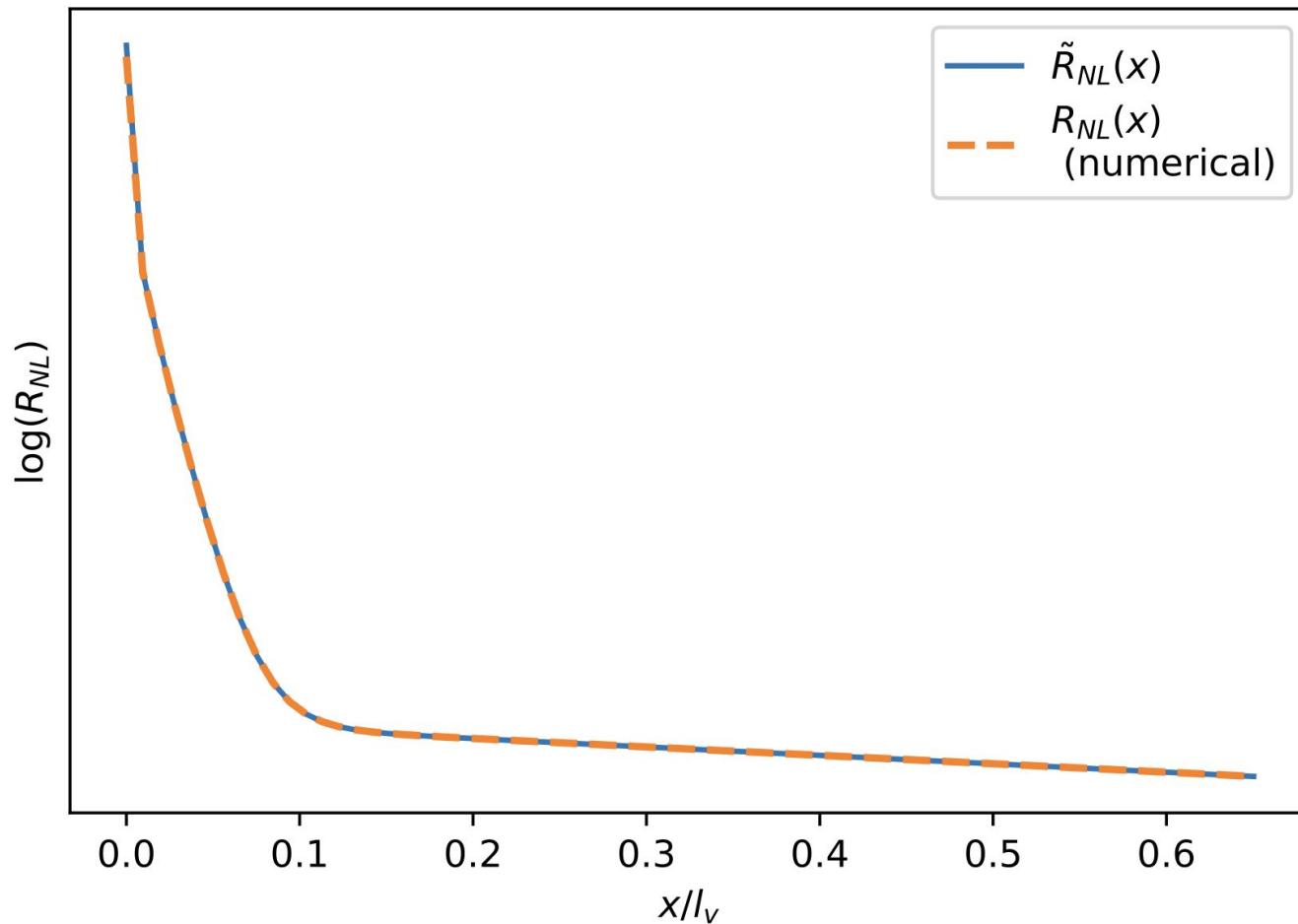
$$\tilde{R}_{\rm NL}(k)\equiv \sin^2(\theta_{\rm VH})R_{\rm NL}^T(k)+\cos^2(\theta_{\rm VH})R_{\rm NL}^{(0)}(k)$$

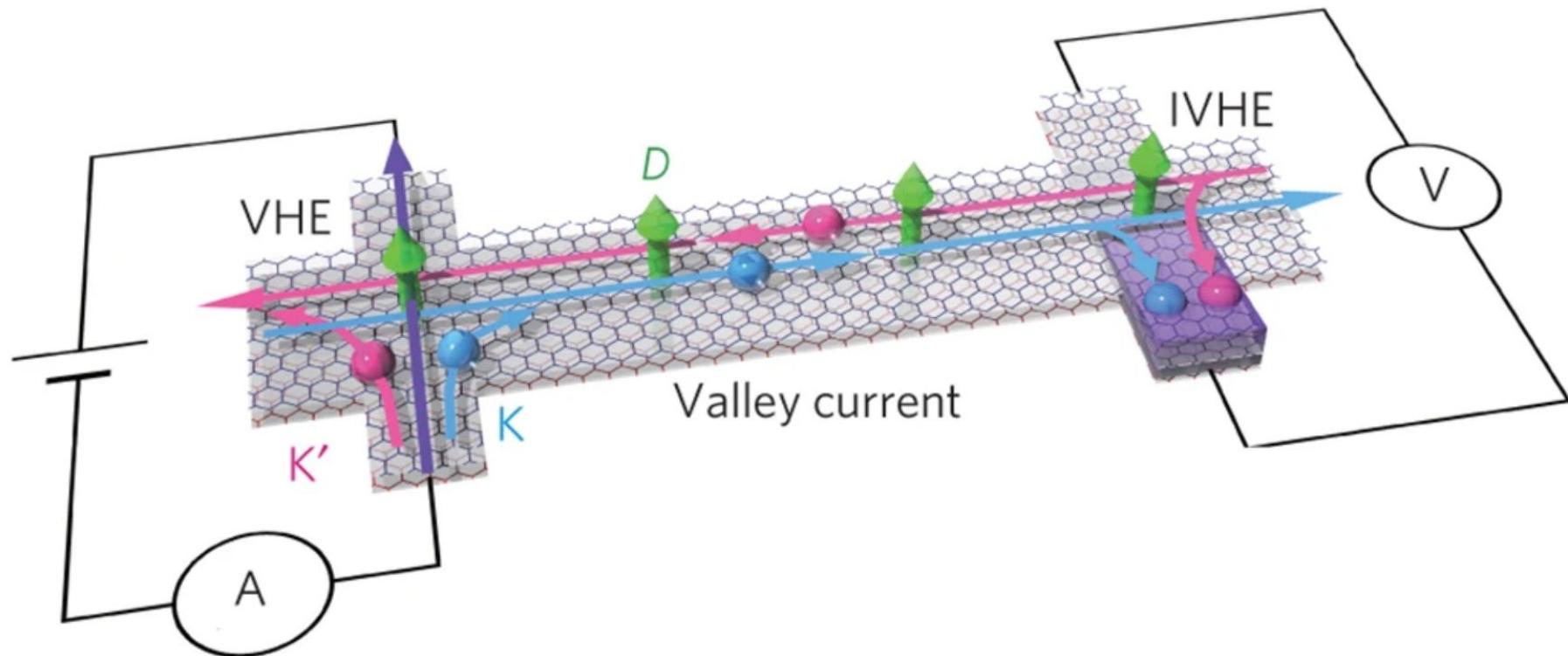


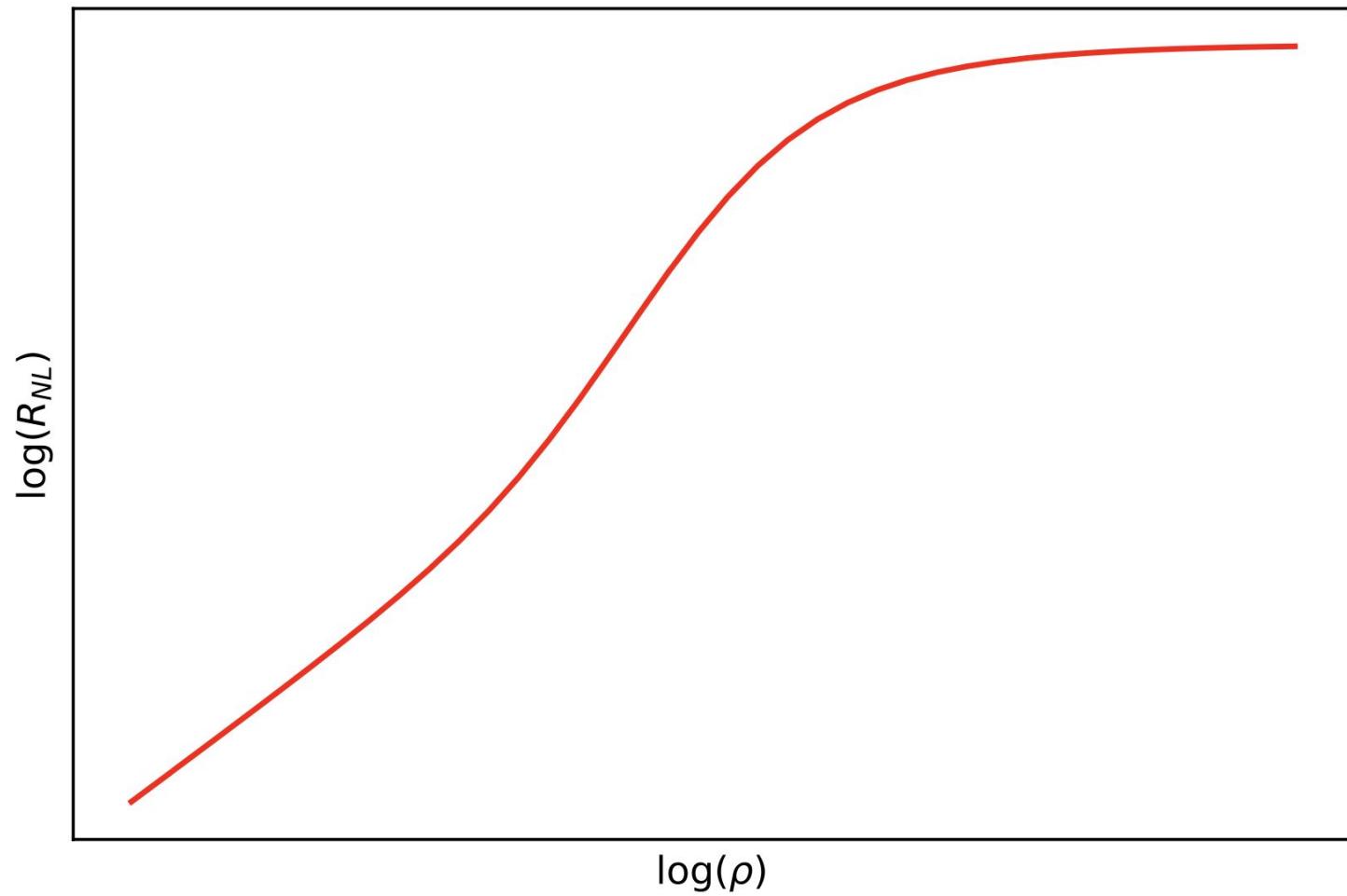
$$\tilde{R}_{\text{NL}}(k) \equiv \sin^2(\theta_{\text{VH}}) R_{\text{NL}}^T(k) + \cos^2(\theta_{\text{VH}}) R_{\text{NL}}^{(0)}(k)$$

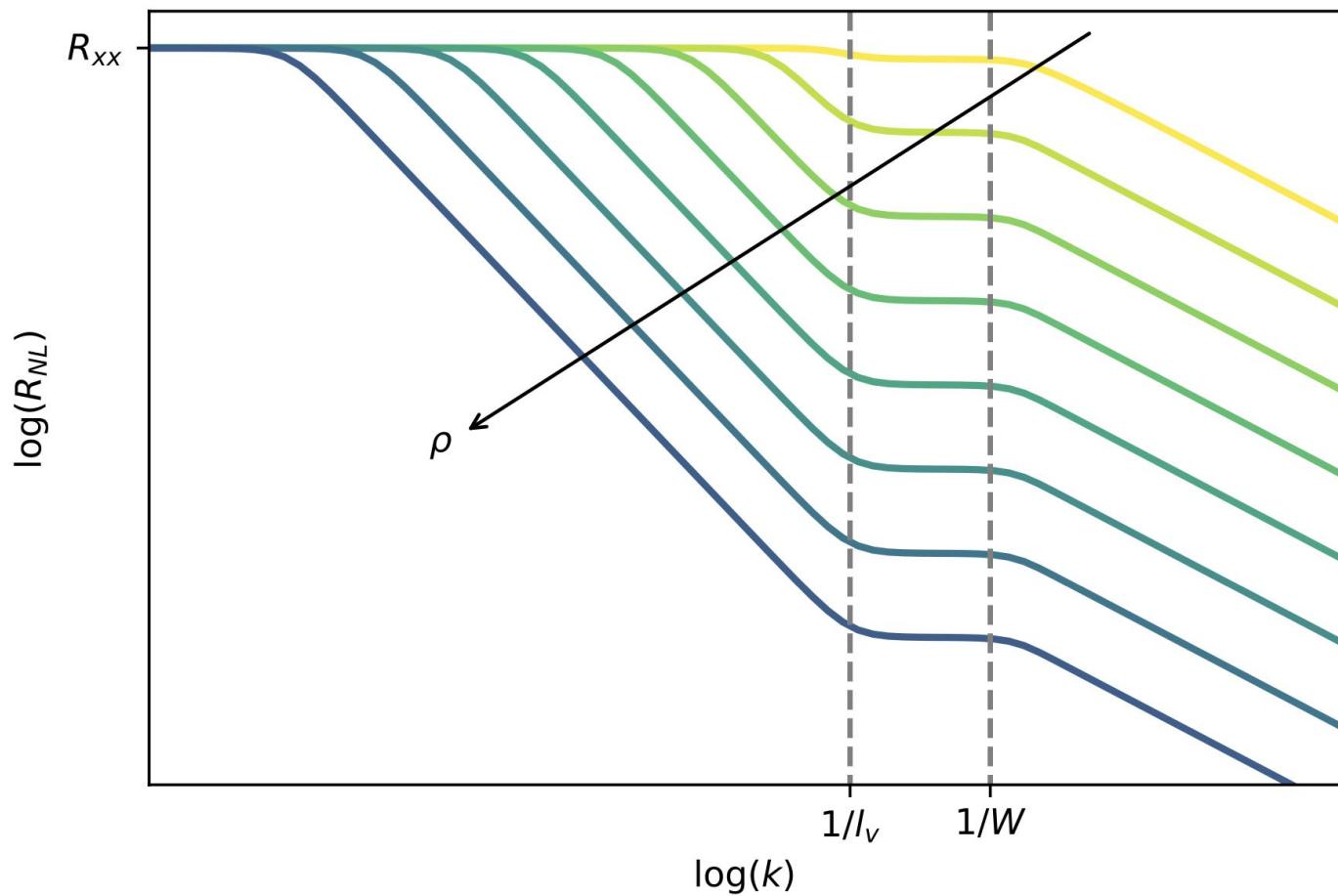
$$\tilde{R}_{\text{NL}}(x) = \sin^2(\theta_{\text{VH}}) R_{\text{NL}}^T(x) + \cos^2(\theta_{\text{VH}}) R_{\text{NL}}^{(0)}(x)$$

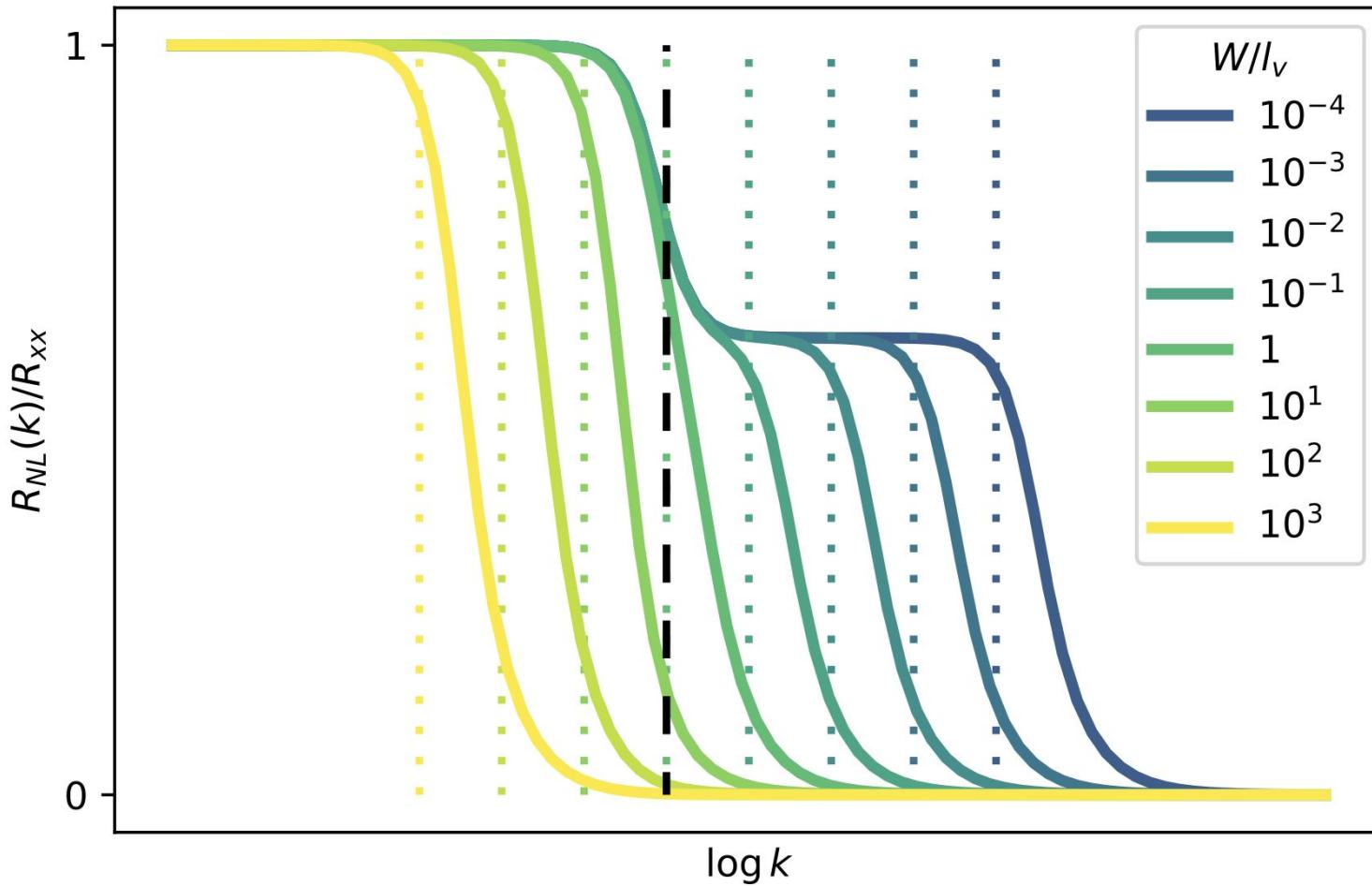
$$\tilde{R}_{\text{NL}}(x) = \frac{R_{xx}}{2L_{\text{v}}} e^{-|x|/L_{\text{v}}} \sin^2(\theta_{\text{VH}}) - \frac{2R_{xx}}{\pi W} \ln \left| \tanh \left(\frac{\pi x}{2W} \right) \right| \cos^2(\theta_{\text{VH}})$$

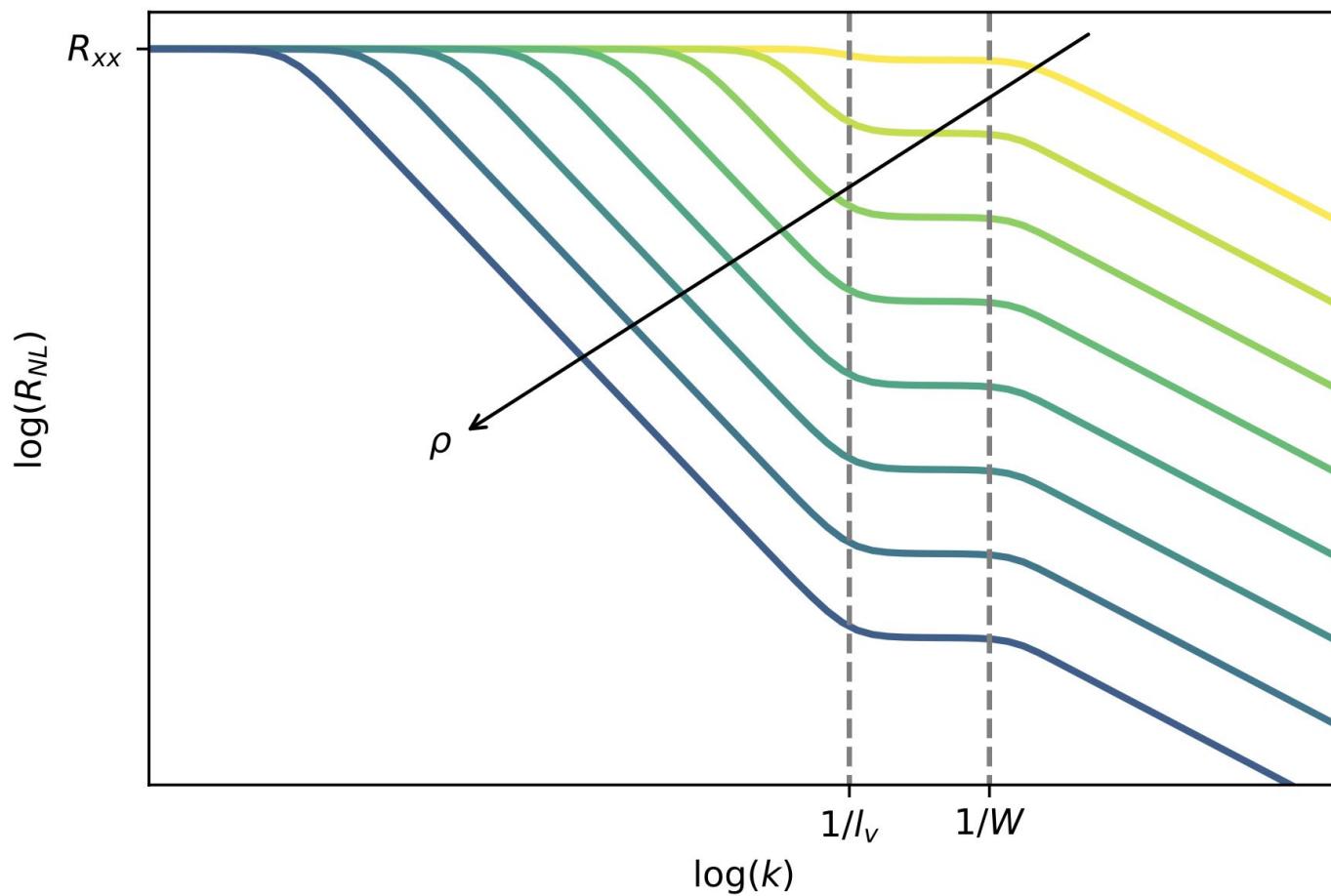






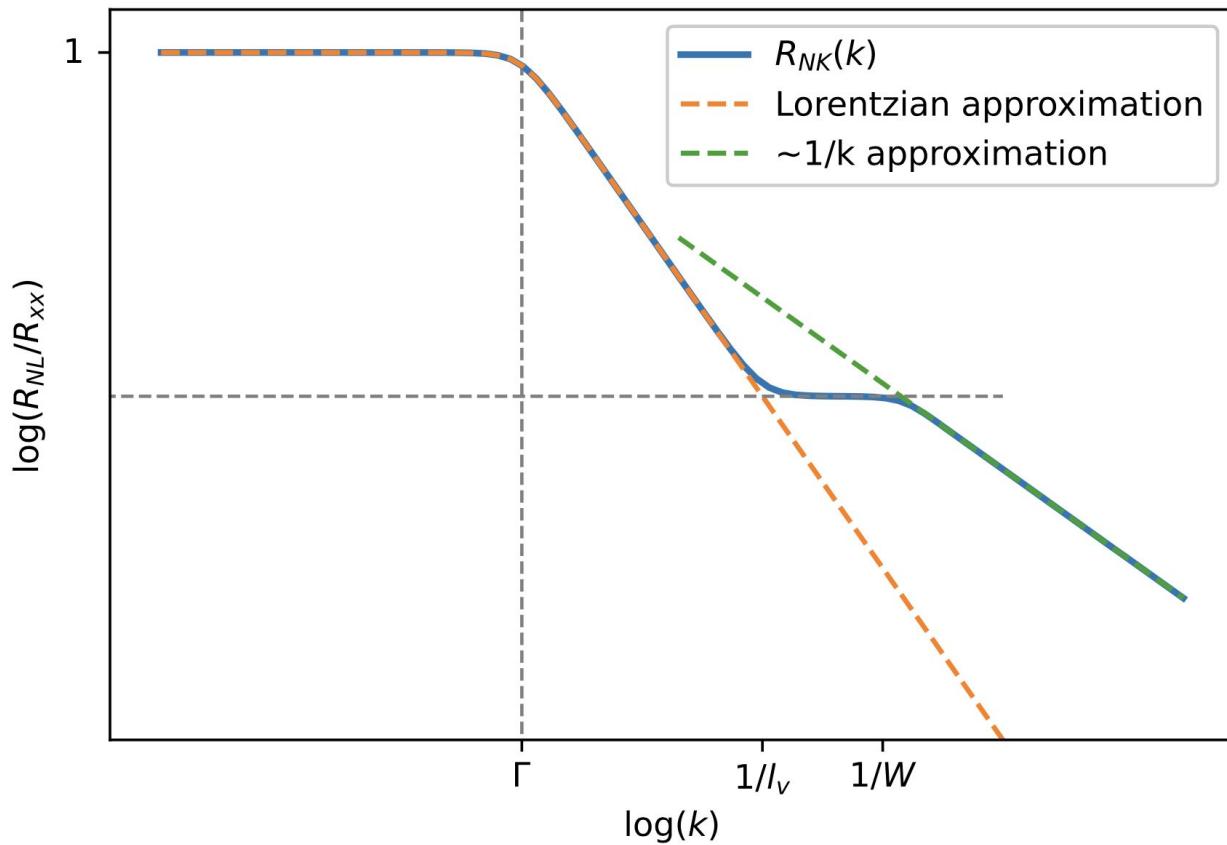






$$R_{\rm NL}(k)=\frac{2\omega(k)}{k\sigma_c}\bigg\{\frac{\omega(k)}{\tanh(kW/2)}+\frac{k\tan^2(\theta_{\rm VH})}{\tanh[\omega(k)W/2]}\bigg\}^{-1}$$

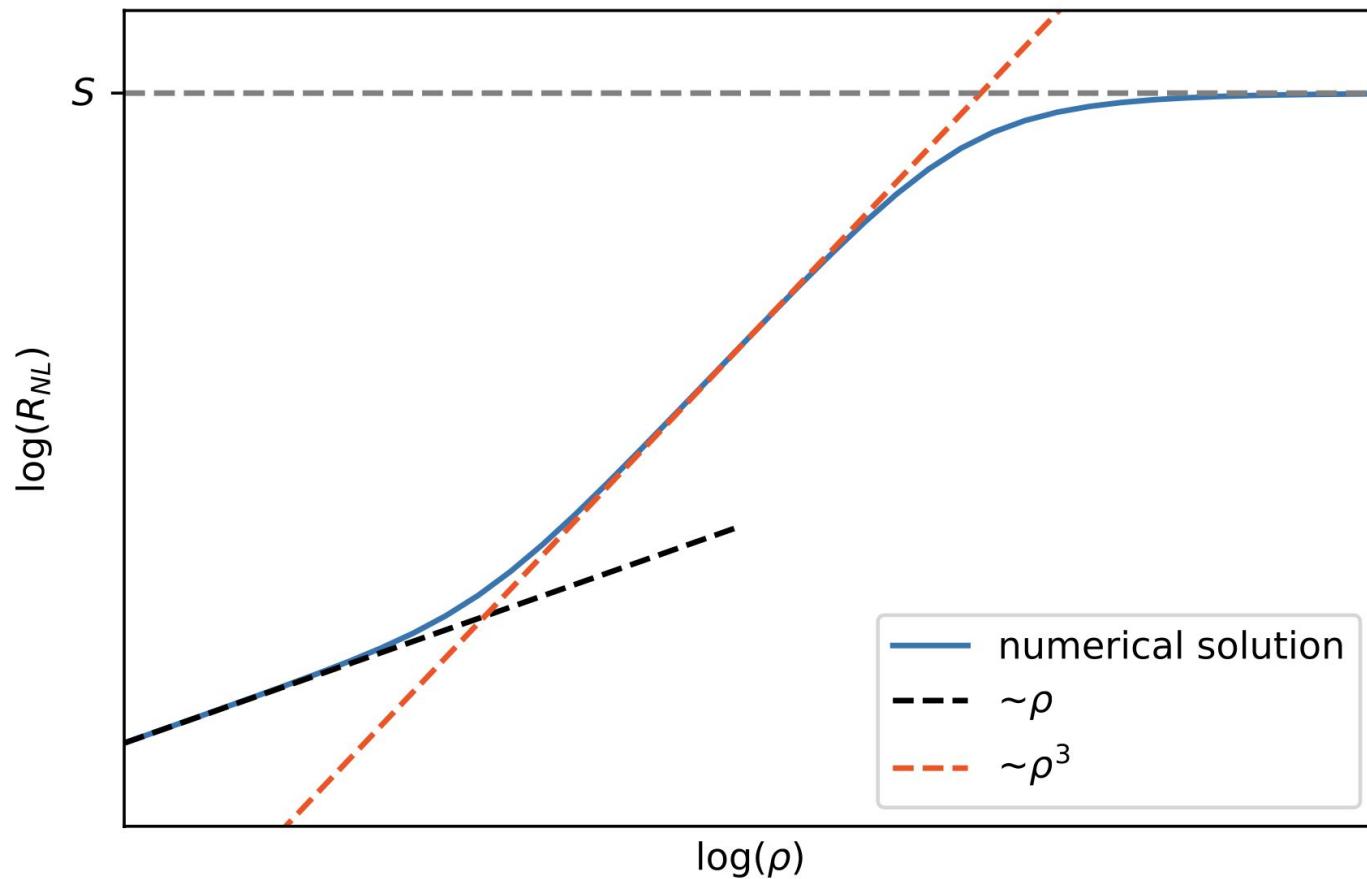
$$\lim_{\rho \rightarrow 0} R_{\rm NL}(x) = \frac{2\rho}{\pi} \ln \left| \coth \left(\frac{\pi x}{2W} \right) \right| + \rho^3 F(x) + o(\rho^5)$$

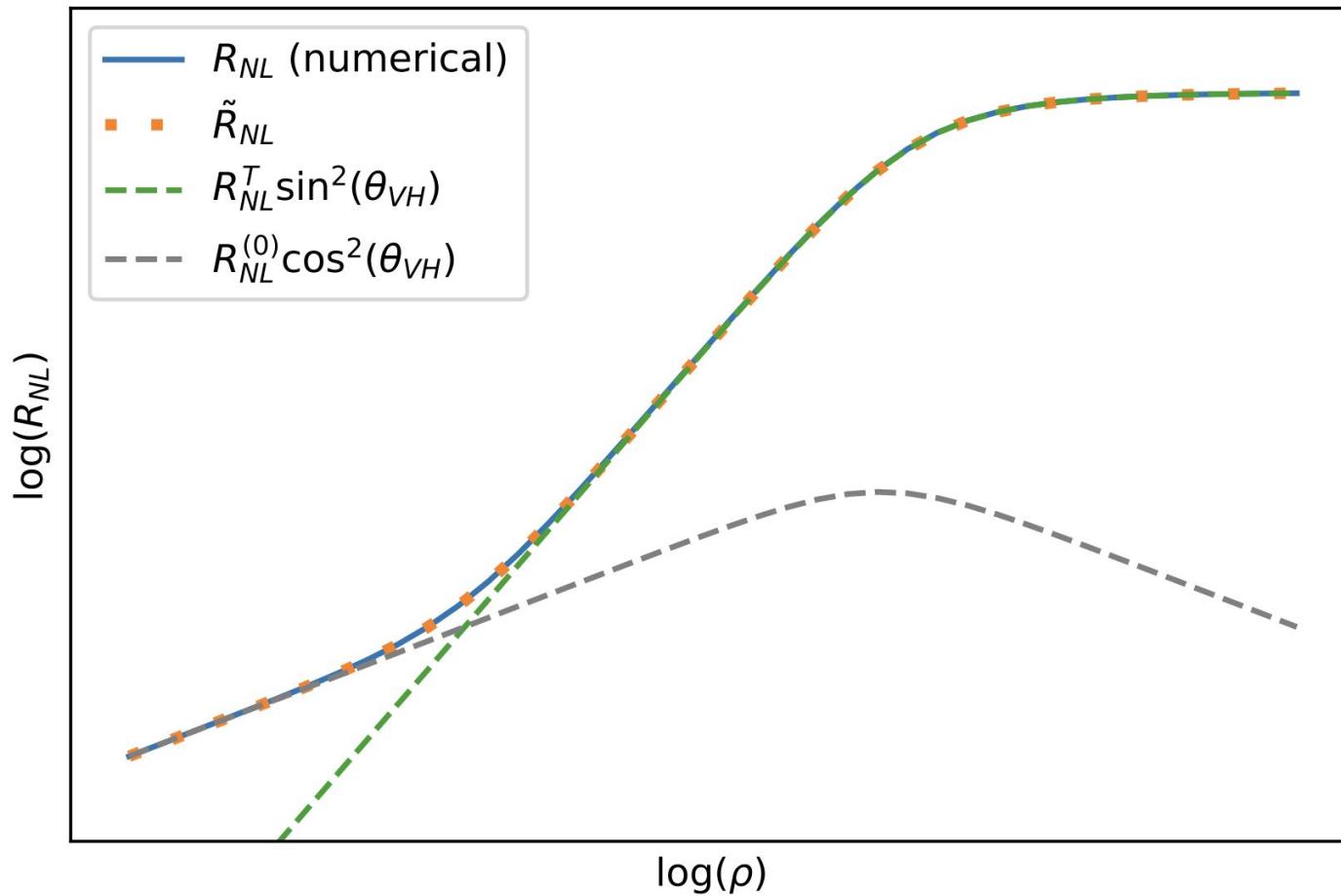


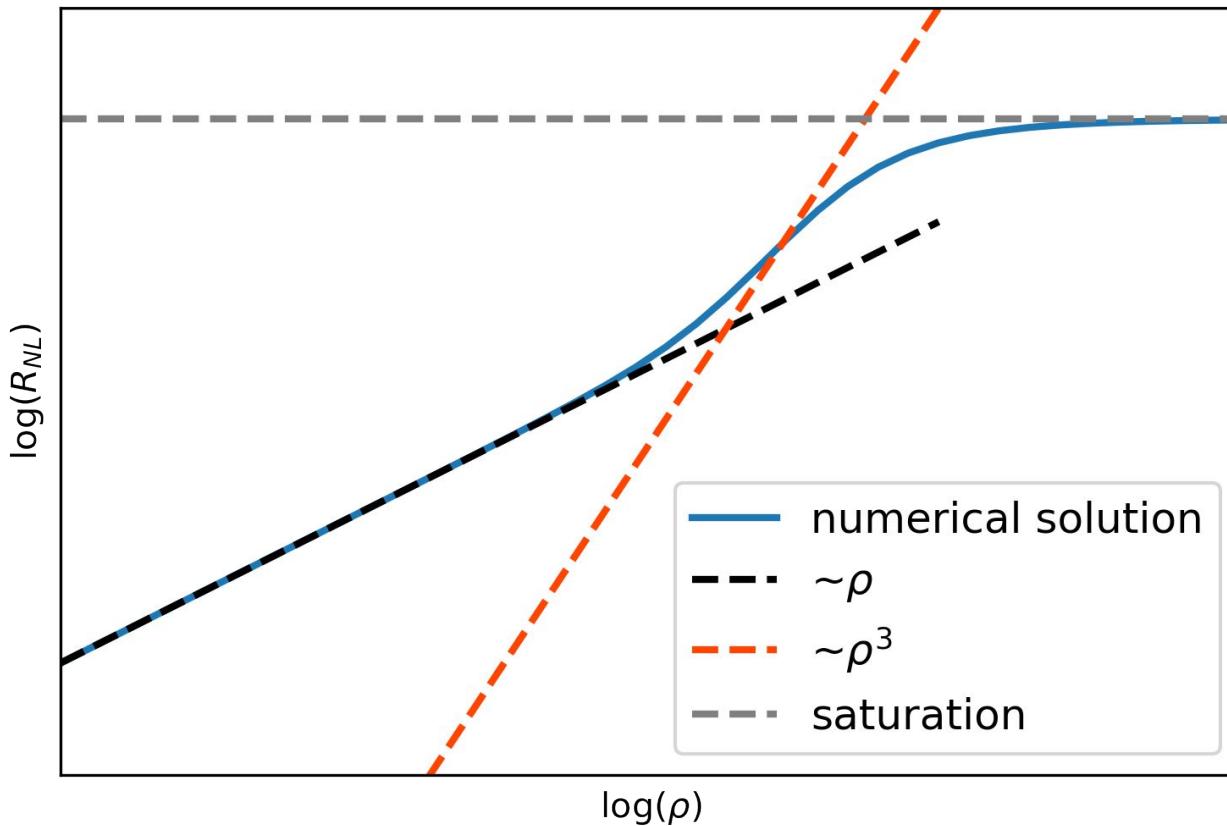
$$\lim_{\rho \rightarrow \infty} R_{NL}(k) = \frac{\rho W}{1 + (k\rho/\Gamma)^2} + \frac{1}{\rho} G(k)$$

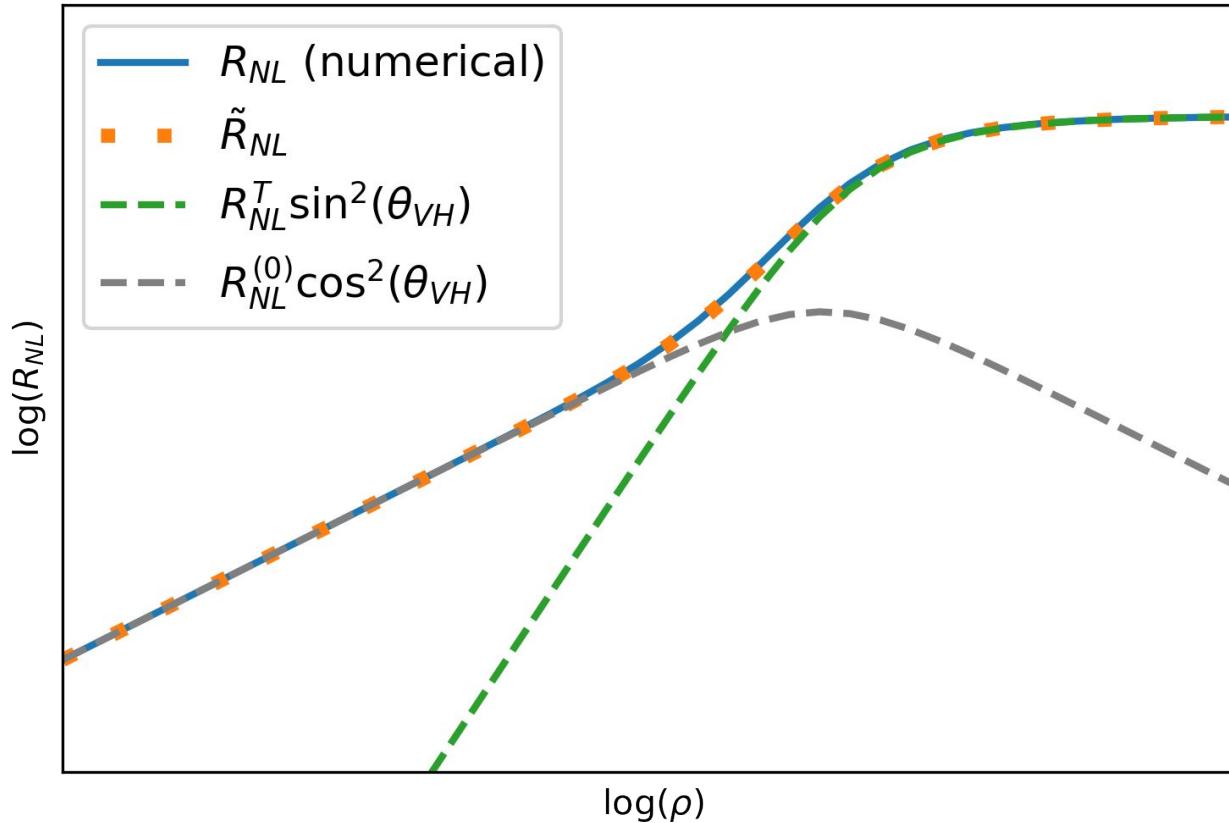
$$\lim_{\rho \rightarrow \infty} R_{\text{NL}}(k) = \frac{\rho W}{1 + (k\rho/\Gamma)^2} + \frac{1}{\rho} G(k)$$

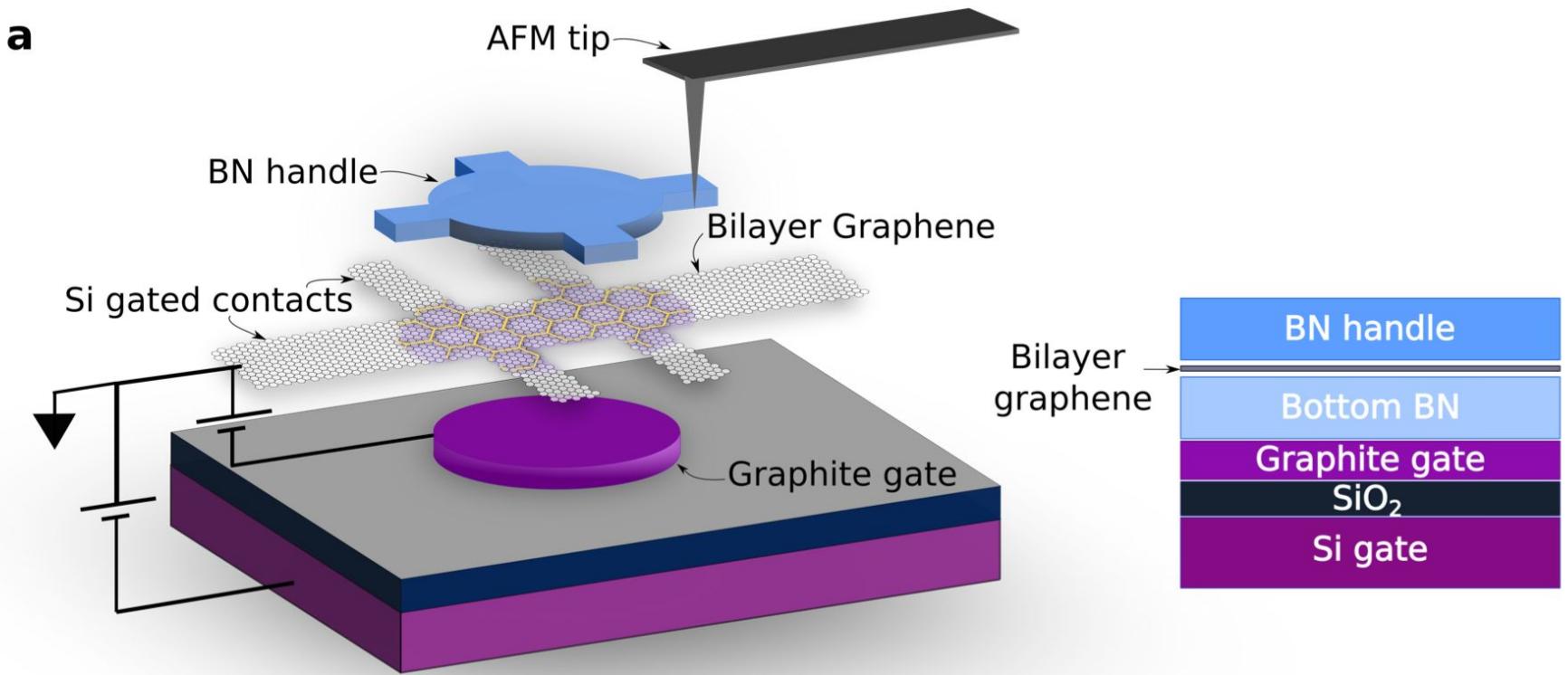
$$\lim_{\rho \rightarrow \infty} R_{\text{NL}}(x) = \frac{1}{\sigma_v} \sqrt{\frac{2W}{l_v} \tanh \left(\frac{W}{2l_v} \right)}$$

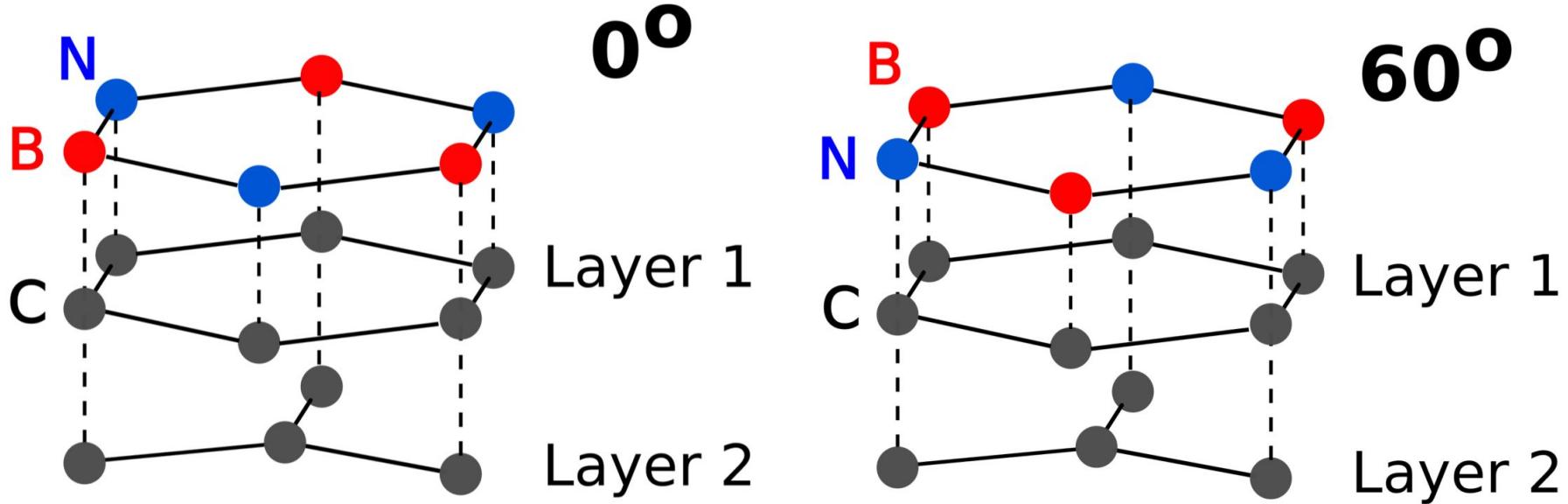


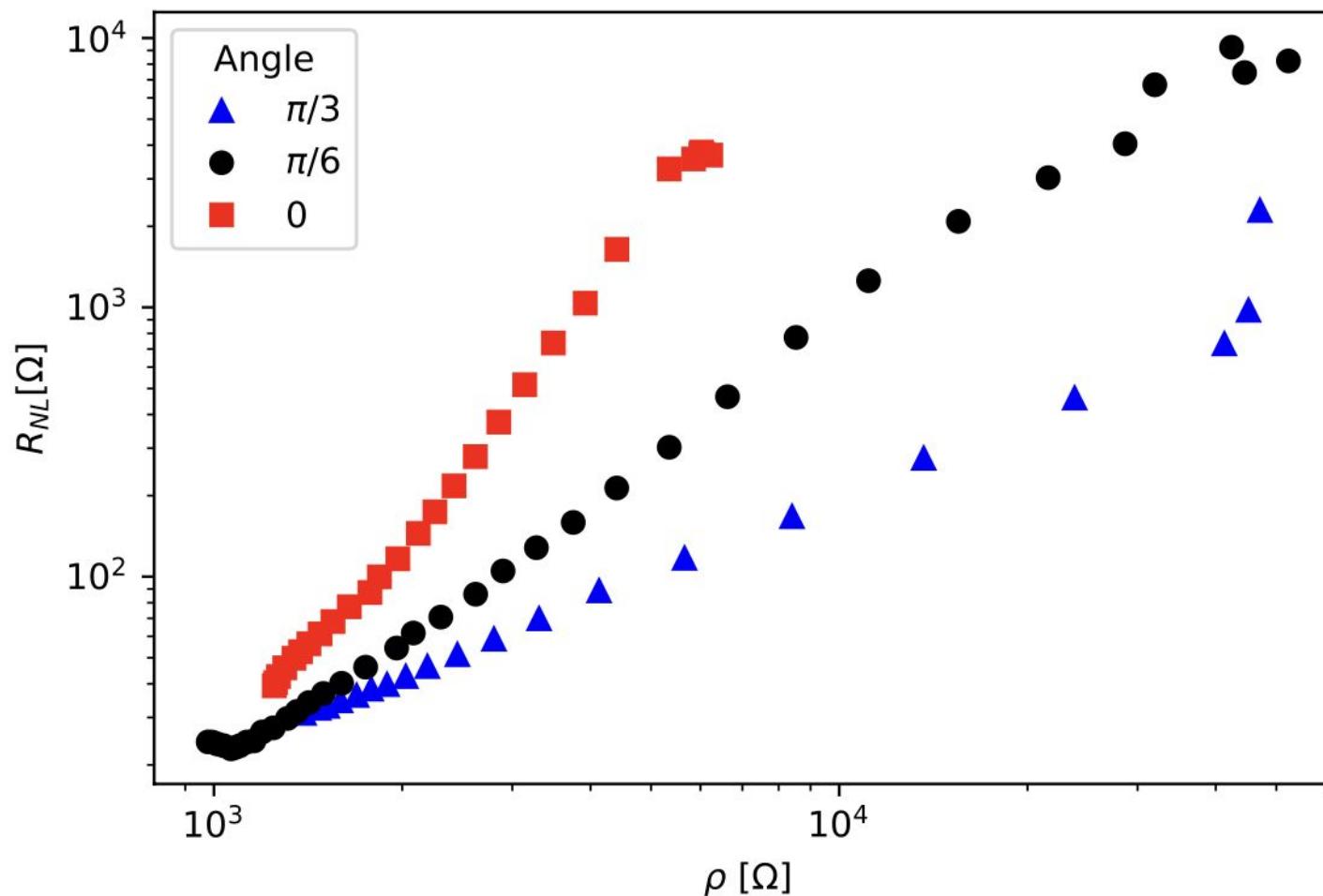


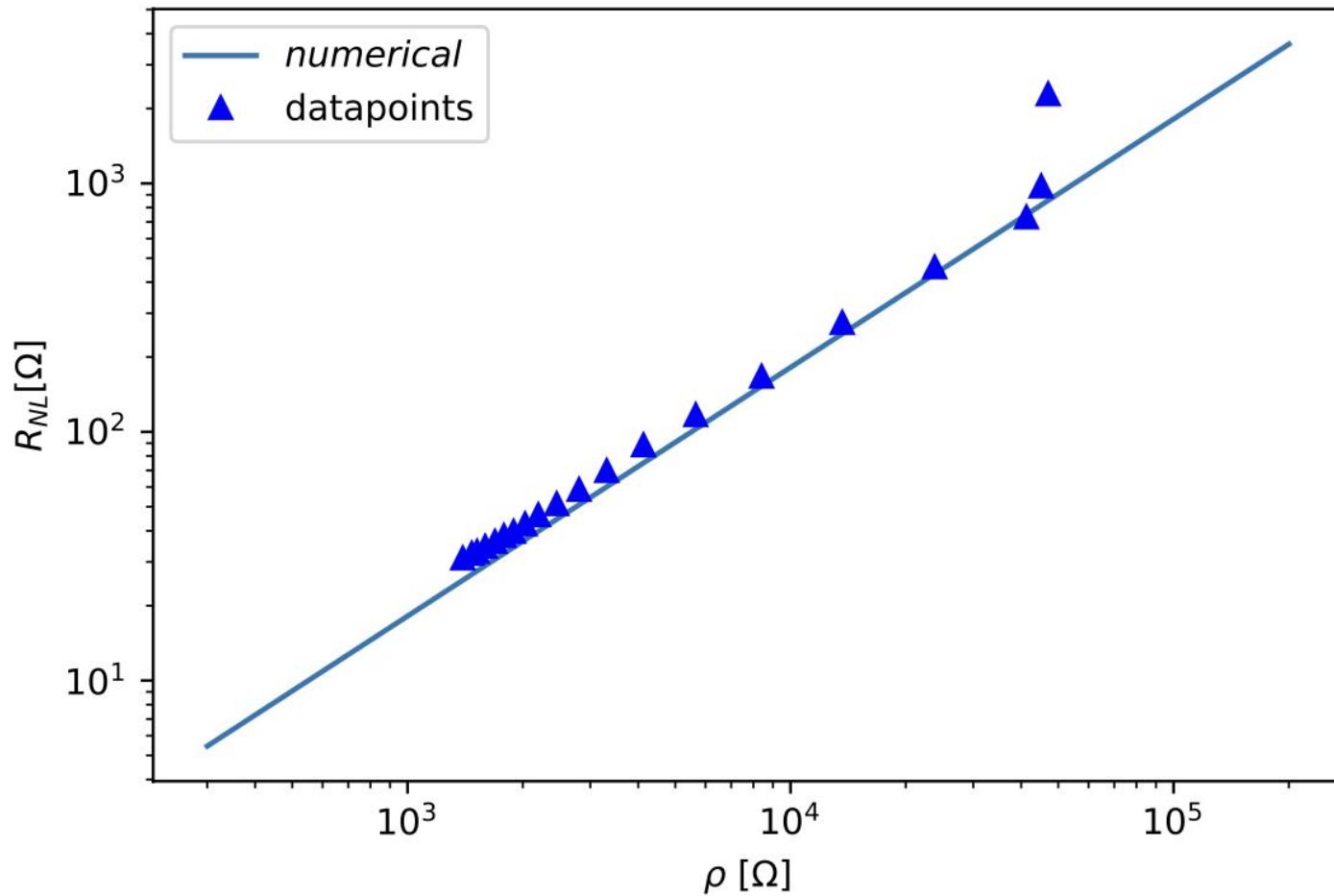


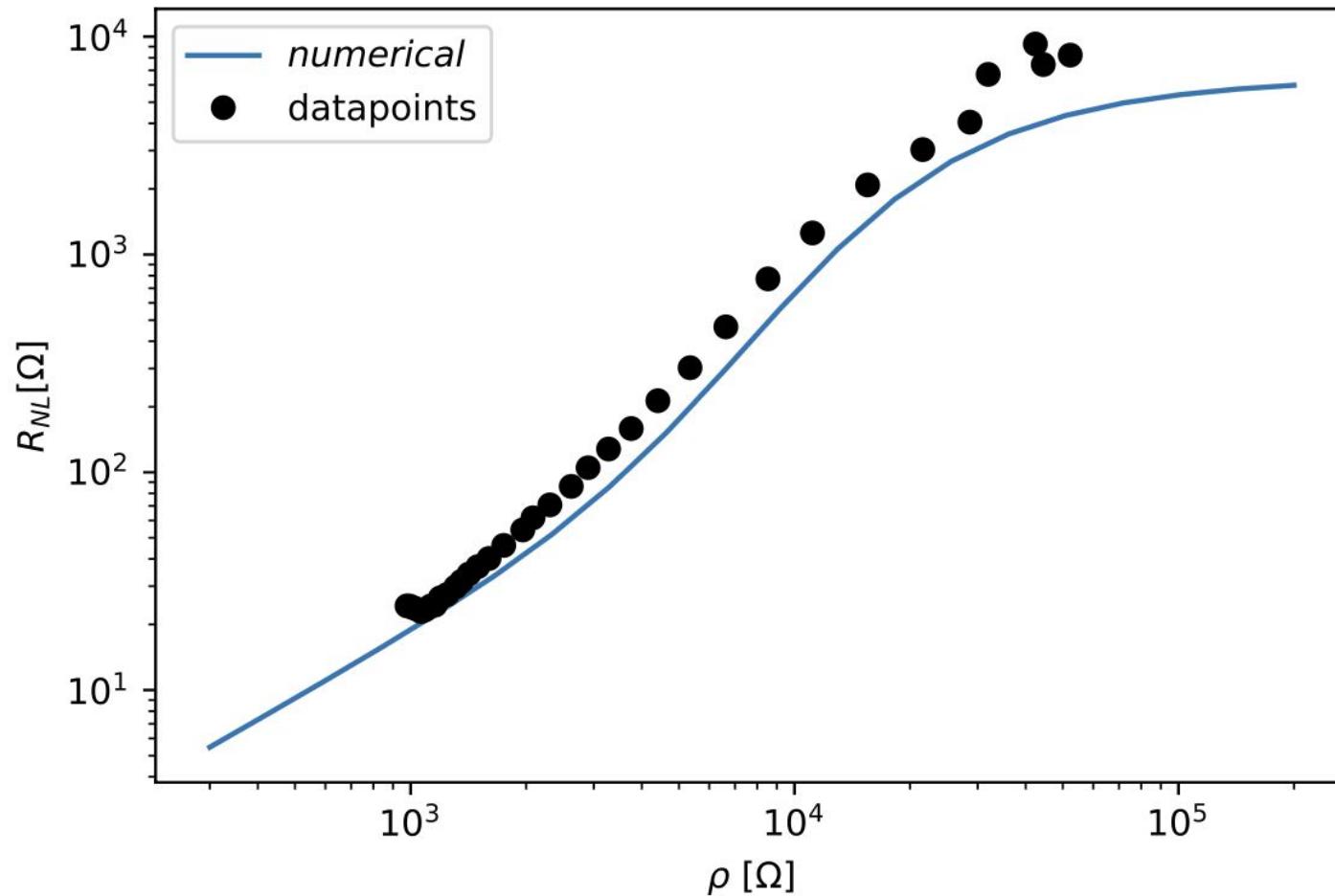


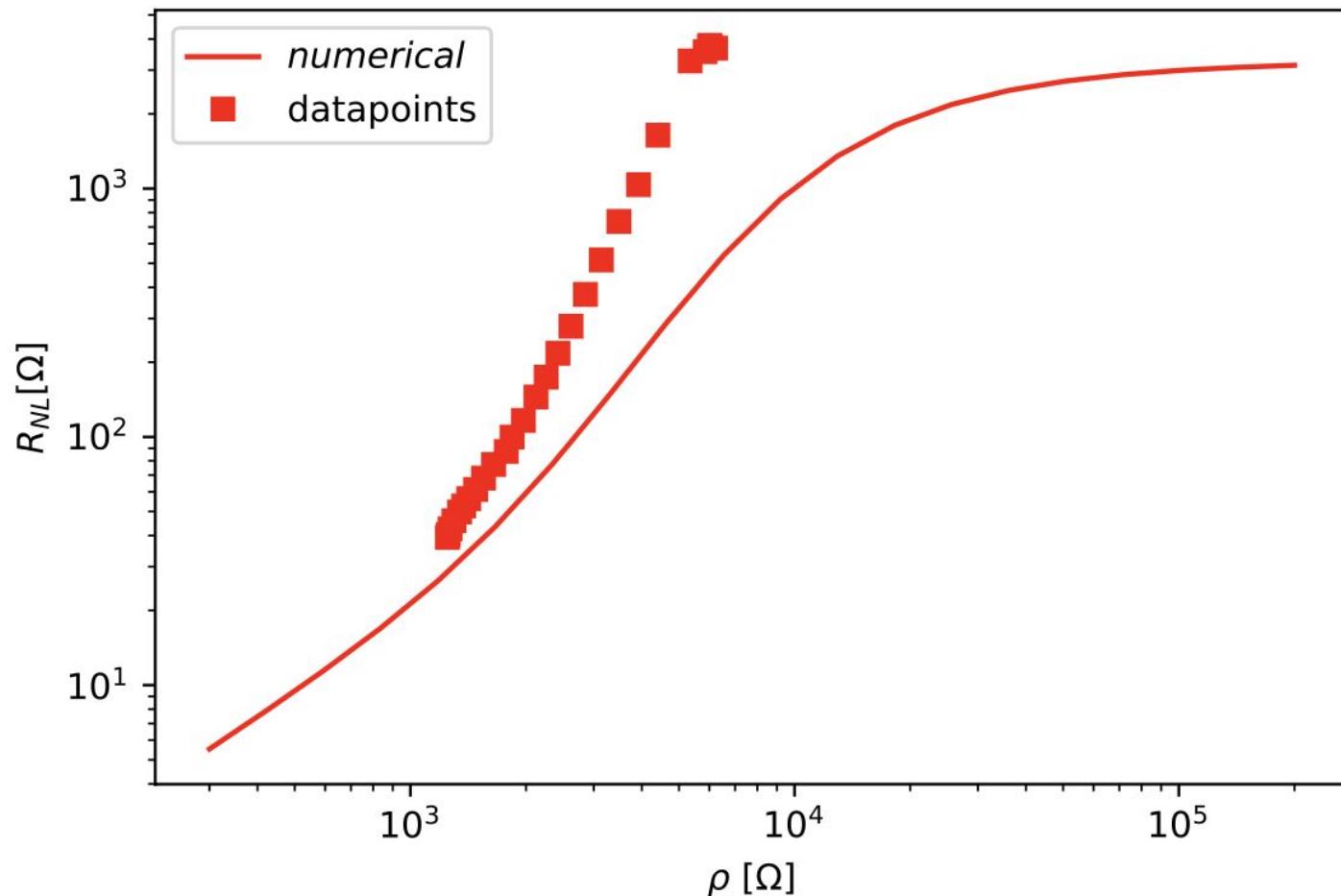












Domande?