## The symmetrized Bregman divergence as dual geodesic energy functionals\*

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The Bregman divergence [2] between two vector parameters  $\theta_1$  and  $\theta_2$  induced by strictly convex function  $F(\theta)$  is  $B_F(\theta_1:\theta_2) := F(\theta_1) - F(\theta_2) - (\theta_1 - \theta_2)^{\top} \nabla F(\theta_2)$ . Let  $\eta = \nabla F(\theta)$  and  $\theta = \nabla F^*(\eta)$  denote the dual parameterizations obtained by the Legendre-Fenchel convex conjugate  $F^*(\eta)$  of  $F(\theta)$ . The Jeffreys-type symmetrization<sup>1</sup> (i.e., arithmetic symmetrization of the sided Bregman divergences [5]) of the Bregman divergence is defined by

$$S_F(\theta_1; \theta_2) := B_F(\theta_1 : \theta_2) + B_F(\theta_2 : \theta_1) = (\theta_2 - \theta_1)^\top (\eta_2 - \eta_1) = S_{F^*}(\eta_1; \eta_2),$$

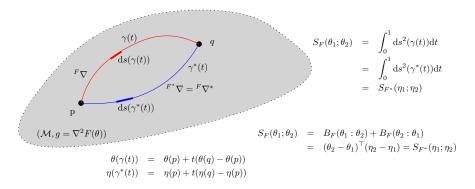


Figure 1: The symmetrized Bregman divergence  $S_F(\theta_p; \theta_q)$  can be interpreted as the energy of the Hessian metric along the primal or dual geodesics linking p to q.

Proposition 1 (Theorem 3.2 of [1], illustration in Figure 1) The Jeffreys-Bregman divergence  $S_F(\theta_1; \theta_2)$  is interpreted as the energy induced by the Hessian metric  $\nabla^2 F(\theta)$  on the dual geodesics:

$$S_F(\theta_1; \theta_2) = \int_0^1 ds^2(\gamma(t))dt = \int_0^1 ds^2(\gamma^*(t))dt.$$

<sup>\*</sup>Extracted from ongoing working notes [4].

<sup>&</sup>lt;sup>1</sup>To distinguish it with the Jensen-Shannon type symmetrization [3].

**Proof:** The proof is based on the first-order and second-order directional derivatives. The first-order directional derivative  $\nabla_u F(\theta)$  with respect to vector u is defined by

$$\nabla_u F(\theta) = \lim_{t \to 0} \frac{F(\theta + tv) - F(\theta)}{t} = v^{\top} \nabla F(\theta).$$

The second-order directional derivatives  $\nabla^2_{u,v}F(\theta)$  is

$$\nabla_{u,v}^{2} F(\theta) = \nabla_{u} \nabla_{v} F(\theta),$$

$$= \lim_{t \to 0} \frac{v^{\top} \nabla F(\theta + tu) - v^{\top} \nabla F(\theta)}{t},$$

$$= u^{\top} \nabla^{2} F(\theta) v.$$

Now consider the squared length element  $ds^2(\gamma(t))$  on the primal geodesic  $\gamma(t)$  expressed using the primal coordinate system  $\theta$ :  $ds^2(\gamma(t)) = d\theta(t)^{\top} \nabla^2 F(\theta(t)) d\theta(t)$  with  $\theta(t) = \theta_1 + t(\theta_2 - \theta_1)$  and  $d\theta(t) = \theta_2 - \theta_1$ . Let us express the  $ds^2(\gamma(t))$  using the second-order directional derivative:

$$ds^{2}(\gamma(t)) = \nabla_{\theta_{2}-\theta_{1}}^{2} F(\theta(t)).$$

Thus we have  $\int_0^1 ds^2(\gamma(t))dt = [\nabla_{\theta_2-\theta_1}F(\theta(t))]_0^1$ , where the first-order directional derivative is  $\nabla_{\theta_2-\theta_1}F(\theta(t)) = (\theta_2-\theta_1)^\top \nabla F(\theta(t))$ . Therefore we get  $\int_0^1 ds^2(\gamma(t))dt = (\theta_2-\theta_1)^\top (\nabla F(\theta_2) - \nabla F(\theta_1)) = S_F(\theta_1;\theta_2)$ .

Similarly, we express the squared length element  $ds^2(\gamma^*(t))$  using the dual coordinate system  $\eta$  as the second-order directional derivative of  $F^*(\eta(t))$  with  $\eta(t) = \eta_1 + t(\eta_2 - \eta_1)$ :

$$ds^{2}(\gamma^{*}(t)) = \nabla_{\eta_{2}-\eta_{1}}^{2} F^{*}(\eta(t)).$$

Therefore, we have  $\int_0^1 \mathrm{d} s^2(\gamma^*(t)) \mathrm{d} t = [\nabla_{\eta_2 - \eta_1} F^*(\eta(t))]_0^1 = S_{F^*}(\eta_1; \eta_2)$ . Since  $S_{F^*}(\eta_1; \eta_2) = S_F(\theta_1; \theta_2)$ , we conclude that

$$S_F(\theta_1; \theta_2) = \int_0^1 \mathrm{d}s^2(\gamma(t)) \mathrm{d}t = \int_0^1 \mathrm{d}s^2(\gamma^*(t)) \mathrm{d}t$$

In 1D, both pregeodesics  $\gamma(t)$  and  $\gamma^*(t)$  coincide. We have  $ds^2(t) = (\theta_2 - \theta_1)^2 f''(\theta(t)) = (\eta_2 - \eta_1) f^{*''}(\eta(t))$  so that we check that  $S_F(\theta_1; \theta_2) = \int_0^1 ds^2(\gamma(t)) dt = (\theta_2 - \theta_1) [f'(\theta(t))]_0^1 = (\eta_2 - \eta_1) [f^{*'}(\eta(t))]_0^1 = (\eta_2 - \eta_1) (\theta_2 - \theta_2)$ .

**Remark 1** In Riemannian geometry, a curve  $\gamma$  minimizes the energy  $E(\gamma) = \int_0^1 |\dot{\gamma}(t)|^2 dt$  if it minimizes the length  $L(\gamma) = \int_0^1 |\dot{\gamma}(t)| dt$  and  $||\dot{\gamma}(t)||$  is constant. Using Cauchy-Schwartz inequality, we can show that  $L(\gamma) \leq E(\gamma)$ .

## References

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