

# Accelerated Dynamics in HMC Simulations of Lattice Field Theory

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# Introduction

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- ▶ Why are we doing it?
  - ▶ Can be used for calculations in lattice field theory.

## Theory - The Path Integral

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### Transition Amplitude

$$\langle x_b, t_b | x_a, t_a \rangle = \int_{x_a}^{x_b} \mathcal{D}x \exp(iS[x(t)]/\hbar) \quad (1)$$

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$$\int_{x_a}^{x_b} \mathcal{D}x = \lim_{N \rightarrow \infty} A_N \prod_{n=1}^{N-1} \int_{-\infty}^{\infty} dx_n \quad (2)$$

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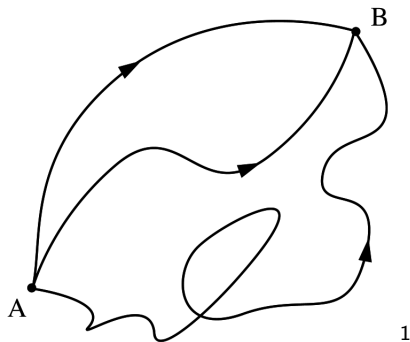
$$\langle x_b, t_b | x_a, t_a \rangle = \int_{x_a}^{x_b} \mathcal{D}x \exp(iS[x(t)]/\hbar) \quad (1)$$

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### Minkowski Action

$$S_M = \int_{t_a}^{t_b} dt \left[ \frac{1}{2} m \left( \frac{dx}{dt} \right)^2 - V(x) \right] \quad (3)$$

## Theory - The Path Integral

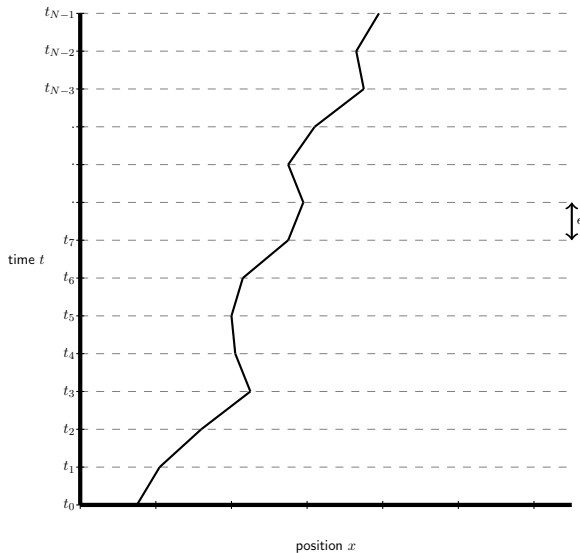


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<sup>1</sup>The Free Encyclopedia Wikipedia. *Path Integral Formulation*. 2017.  
URL: [https://en.wikipedia.org/wiki/Path\\_integral\\_formulation#/media/File:Three\\_paths\\_from\\_A\\_to\\_B.png](https://en.wikipedia.org/wiki/Path_integral_formulation#/media/File:Three_paths_from_A_to_B.png).

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$$S_M = \sum_{j=0}^{N-1} \epsilon \left[ \frac{1}{2} m \left( \frac{x_{j+1} - x_j}{\epsilon} \right)^2 - V(x_j) \right] \quad (1)$$

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### Discrete Path Integral

$$\langle x_b, t_b | x_a, t_a \rangle \sim \int_{-\infty}^{\infty} \prod_{j=1}^{N-1} dx_j \exp \left( \frac{i}{\hbar} S_M \{x_j\} \right) \quad (2)$$

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$$S_M = i \sum_{j=0}^{N-1} a \left[ \frac{1}{2} m \left( \frac{x_{j+1} - x_j}{a} \right)^2 + V(x_j) \right] := iS_E \quad (5)$$

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### Discrete Euclidean Path Integral

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### Partition Function

$$Z \sim \int_{-\infty}^{+\infty} \prod_{j=1}^{N-1} dx_i \exp \left( -\beta H (\{x_i\}) \right) \quad (4)$$



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### Monte Carlo Estimate

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### Boltzmann Distribution

$$p(\mathbf{x}_{\nu}) \mathcal{D}\mathbf{x} = \frac{\exp(-S(\mathbf{x}_{\nu})) \mathcal{D}\mathbf{x}}{\int \mathcal{D}\mathbf{x} \exp(-S(\mathbf{x}))} \quad (6)$$

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### HMC Hamiltonian

$$H_{hmc} := \sum_{i=0}^{N-1} \frac{p_i^2}{2m} + S(\{x_i\}) \quad (8)$$

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3. Accept configuration  $\{q_i^*\}$  with probability  $\min[1, \exp(-H_{HMC}(\{q_i^*\}, \{p_i^*\}) + H_{HMC}(\{q_i\}, \{p_i\}))]$  (Metropolis update).

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4. Return to step 1.

## Results - Harmonic Oscillator Expectation Values

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Value	Measured	Discrete Theory <sup>1</sup>	Continuum Theory
$\langle x \rangle$	0.00015(20)	0	0

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<sup>1</sup>M Creutz and B Freedman. "A statistical approach to quantum mechanics". In: *Annals of Physics* 132.2 (1981), pp. 427–462. DOI: [10.1016/0003-4916\(81\)90074-9](https://doi.org/10.1016/0003-4916(81)90074-9).

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$E_0$	0.44723(14)	0.4472135955	0.5

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$E_1$	0.9679(90)	<i>FILL</i>	1

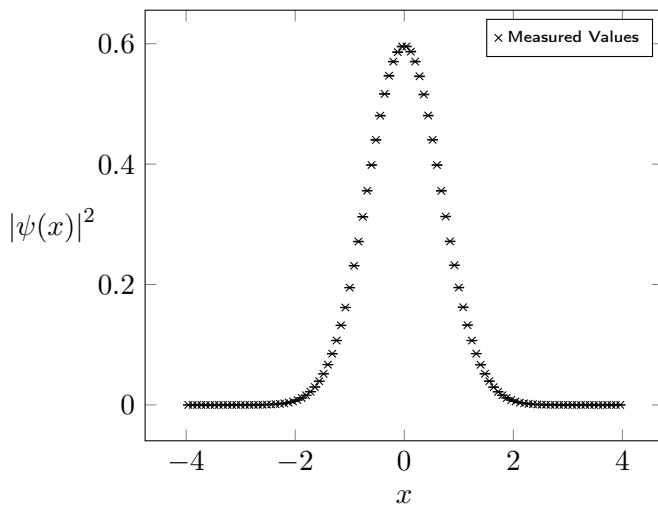
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### Discrete Wave Function<sup>1</sup>

$$\psi_{disc.}(x) = \left(\frac{\omega}{\pi}\right)^{\frac{1}{4}} \exp\left(-\frac{1}{2}\omega x^2\right) \quad (7)$$

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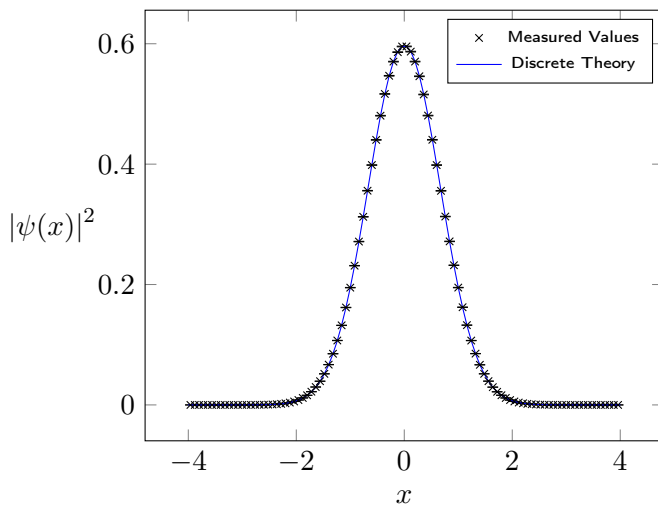
$$\psi_{disc.}(x) = \left(\frac{\omega}{\pi}\right)^{\frac{1}{4}} \exp\left(-\frac{1}{2}\omega x^2\right) \quad (7)$$

$$\omega^2 = \mu^2 \left(1 + \frac{a^2 \mu^2}{4}\right) \quad (8)$$

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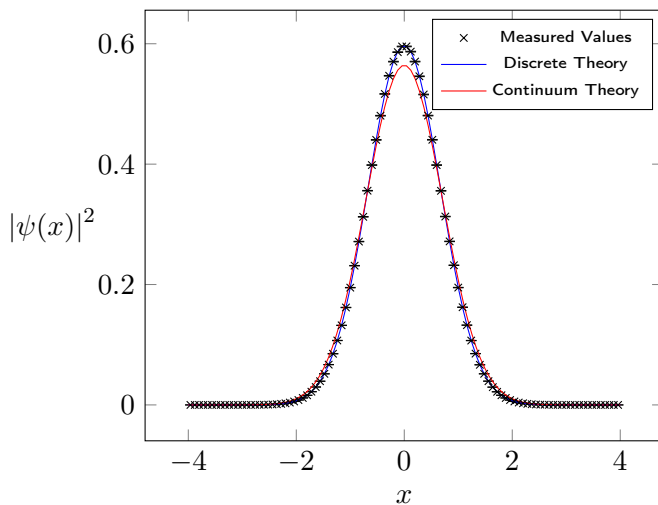


## Results - Harmonic Oscillator Wave Function

### Continuous Wave Function

$$\psi_{cont.}(x) = \left(\frac{\mu}{\pi}\right)^{\frac{1}{4}} \exp\left(-\frac{1}{2}\mu x^2\right) \quad (7)$$

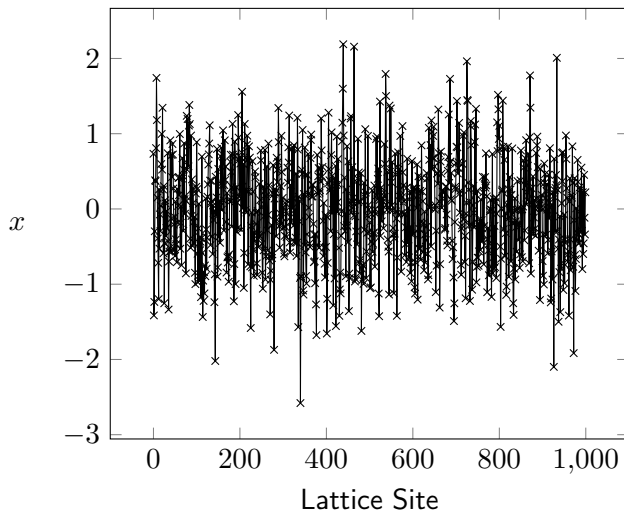
## Results - Harmonic Oscillator Wave Function



## Results - Harmonic Oscillator Typical Trajectory



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## Results - Anharmonic Oscillator Expectation Values

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Value	Measured	Reference Values <sup>1</sup>
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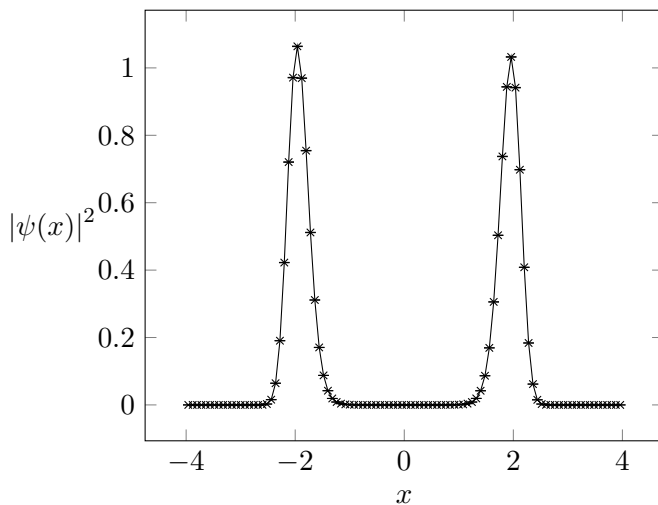
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## Results - Anharmonic Oscillator Wave Function

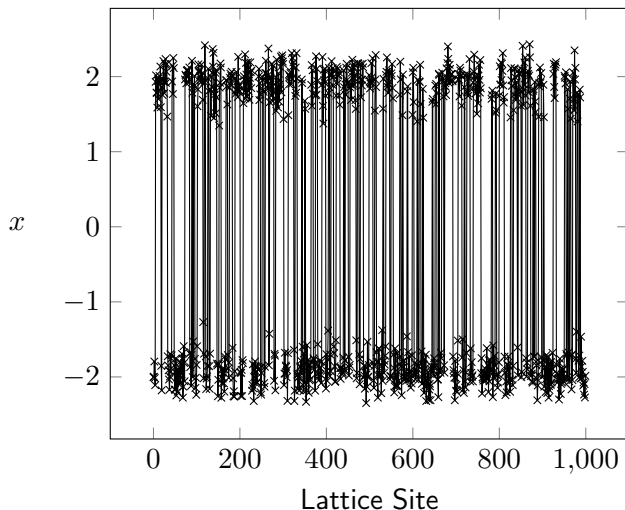
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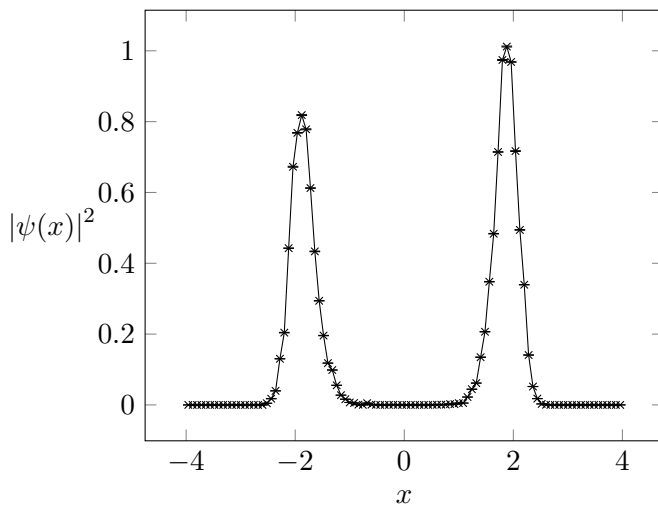
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## Results - Isolated Modes Wave Function

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  - ▶ Introduce “tempering” into the dynamics to sample from isolated modes.
- ▶ Applications of tempering?
  - ▶ Potentially applicable to lattice field theory where computation time is far more costly.