

Spurious Quasi-Resonances for Stabilized BIE-Volume Formulations for the Helmholtz Transmission Problem

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Abstract

A particular *regularised variational formulation* of the Helmholtz transmission problem is studied on a two-dimensional disk for varying frequencies. In particular, the operator norm of its associated inverse operator is investigated. In scenarios where the inner refractive index is bigger than the outer one, the operator norm associated with the considered formulation resonates and grows as a function of the wave number. The solution operator did not expose this resonance and growth behavior. This behavior, studied in previous research for other operators, is called *spurious quasi-resonances*. The occurring resonances, however, were much weaker in the formulation considered here than in others. The origin of these resonances is explained.



1 Introduction

LET $\Omega^- \subset \mathbb{R}^d, d > 0$ be a bounded Lipschitz domain and define $\Omega^+ := \mathbb{R}^d \setminus \overline{\Omega^-}$, $\Gamma := \partial\Omega^-$. For any function f on \mathbb{R}^d define $f^\pm := f|_{\Omega^\pm}$.

Let $H_{\text{loc}}^1(\Omega^\pm, \Delta) := \{v : \chi v \in H^1(\Omega^\pm), \Delta(\chi v) \in L^2(\Omega^\pm) \text{ for } \forall \chi \in C_{\text{comp}}^\infty(\mathbb{R}^d)\}$, as defined in [1]. Consider the Neumann and Dirichlet trace operators

$$\begin{aligned}\gamma_D^\pm : H_{\text{loc}}^1(\Omega^\pm) &\rightarrow H^{\frac{1}{2}}(\Gamma), (\gamma_D^\pm f)(x) := f(x) \\ \gamma_N^\pm : H_{\text{loc}}^1(\Omega^\pm, \Delta) &\rightarrow H^{-\frac{1}{2}}(\Gamma), (\gamma_N^\pm f)(x) := \text{grad } f(x) \cdot \mathbf{n}(x)\end{aligned}$$

where $H_{\text{loc}}^1(\Omega^\pm)$, $H_{\text{loc}}(\Delta, \Omega^\pm)$, $H^{\frac{1}{2}}(\Gamma)$, and $H^{-\frac{1}{2}}(\Gamma)$ are defined in chapter 3 of [2]. Now, the Cauchy trace is $\gamma_C^\pm : H_{\text{loc}}^1(\Omega^\pm, \Delta) \rightarrow H^{1/2}(\Gamma) \times H^{-1/2}(\Gamma)$ with values given by $\gamma_C^\pm := (\gamma_D^\pm, \gamma_N^\pm)$. This concludes all required definitions to formulate the Helmholtz transmission problem.

Definition 1.1 (Helmholtz transmission problem). *For $\tilde{\kappa}, c_i, c_o > 0$ and $f = (f^1, f^2) \in H^{1/2}(\Gamma) \times H^{-1/2}(\Gamma)$ find $U \in H_{\text{loc}}^1(\mathbb{R}^d \setminus \Gamma)$ such that*

$$\begin{aligned}(\Delta + \tilde{\kappa}^2 c_i) U^- &= 0 & \text{in } \Omega^- \\ (\Delta + \tilde{\kappa}^2 c_o) U^+ &= 0 & \text{in } \Omega^+ \\ \gamma_C^+ U^+ - \gamma_C^- U^- &= f & \text{on } \Gamma.\end{aligned}\tag{1}$$

Additionally, u must satisfy the Sommerfeld radiation condition $\lim_{r \rightarrow \infty} r^{\frac{d-1}{2}} \left(\frac{\partial U}{\partial r} - i\sqrt{c_o} \tilde{\kappa} U \right) = 0$ where r refers to the radial spherical coordinate.

This problem is well-posed and the solution is unique as shown in Lemma 2.2 of [3]. Hiptmair et al. considered the *single-trace formulations* (STF) of this problem [1]. This is a reformulation of this problem in terms of *boundary integral equations* (BIEs). When investigating the case $c_i < c_o$ they found that the involved *boundary integral operators* (BIOs) as a function of $\tilde{\kappa}$ exposed a nonphysical resonance behavior. More specifically, they found the operator norm of the inverse STF BIOs to have resonances (called *spurious quasi-resonances*) that the norm of the solution operator did not. Formally, Hiptmair et al. defined the solution operator as follows [1].

Definition 1.2. Given positive real numbers k, c_i , and c_o , $S : H^{\frac{1}{2}}(\Gamma) \times H^{-\frac{1}{2}}(\Gamma) \rightarrow H^{\frac{1}{2}}(\Gamma) \times H^{-\frac{1}{2}}(\Gamma)$ is the solution operator if $S(c_i, c_o) f := \gamma_{\bar{C}} u$ where u solves eq 1.

Hiptmair et al. removed these *spurious quasi-resonances* using an augmented formulation of the BIEs. P. Meury's doctoral thesis contained a *regularized variational formulation* of the Helmholtz transmission problem [4]. The goal of this report is to investigate the occurrence of *spurious quasi-resonance* in this regularized variational formulation of the Helmholtz transmission problem. For easier notation, let $\kappa = \tilde{\kappa} \sqrt{c_o}$ and $\tilde{c} := \frac{c_i}{c_o}$ on the following pages. Moreover, for the rest of the report we consider the simple representational example $d = 2$ and $\Omega^- = B_1(0)$.

Before reviewing this regularized variational formulation, we need to review a few concepts and introduce some notations.

2 Definitions

We review some basic definitions which are necessary in the coming sections. The following definition was introduced in [4].

Definition 2.1 (Interior Dirichlet-to-Neumann map). *The interior Dirichlet-to-Neumann map $\text{DtN}_{\kappa}^- : H^{\frac{1}{2}}(\Gamma) \times H^{-\frac{1}{2}}(\Gamma)$ is the operator that returns $\gamma_N U$ if U solves the Dirichlet problem defined by a vector in $H^{\frac{1}{2}}(\Gamma)$.*

Moreover, we will need the following BIODs defined in [5] [4].

Definition 2.2. Let $\{\gamma_i V\}_{\Gamma} := \frac{1}{2} (\gamma_i^+ V + \gamma_i^- V)$ for $i = D, N$. Moreover, introduce the single and double layer potential $\Psi_{\text{SL}}^{\kappa}(\vartheta)(\mathbf{x}) := \int_{\Gamma} G_{\kappa}(|\mathbf{x} - \mathbf{y}|) \vartheta(\mathbf{y}) dS(\mathbf{y})$ and $\Psi_{\text{DL}}^{\kappa}(v)(\mathbf{x}) := \int_{\Gamma} \frac{\partial G_{\kappa}(|\mathbf{x} - \mathbf{y}|)}{\partial \mathbf{n}(\mathbf{y})} v(\mathbf{y}) dS(\mathbf{y})$ with $G_{\kappa}(z) := \frac{1}{4\pi} \frac{\exp(i\kappa z)}{z}$. Then for $|s| < \frac{1}{2}$ we define four BIODs:

$$\begin{aligned} V_{\kappa} : H^{s-\frac{1}{2}}(\Gamma) &\rightarrow H^{s+\frac{1}{2}}(\Gamma), & V_{\kappa} &:= \{\gamma_D \Psi_{\text{SL}}^{\kappa}\}_{\Gamma} \\ K_{\kappa} : H^{s+\frac{1}{2}}(\Gamma) &\rightarrow H^{s+\frac{1}{2}}(\Gamma), & K_{\kappa} &:= \{\gamma_D \Psi_{\text{DL}}^{\kappa}\}_{\Gamma} \\ K'_{\kappa} : H^{s-\frac{1}{2}}(\Gamma) &\rightarrow H^{s-\frac{1}{2}}(\Gamma), & K'_{\kappa} &:= \{\gamma_N \Psi_{\text{SL}}^{\kappa}\}_{\Gamma} \\ W_{\kappa} : H^{s+\frac{1}{2}}(\Gamma) &\rightarrow H^{s-\frac{1}{2}}(\Gamma), & W_{\kappa} &:= -\{\gamma_N \Psi_{\text{DL}}^{\kappa}\}_{\Gamma}. \end{aligned}$$

3 Regularized Variational Formulation

P. Meury provides a *regularized variational formulation* of the Helmholtz transmission problem as follows [4].¹

Definition 3.1 (Regularized variational formulation). *Find $U \in H^1(\Omega)$, $\theta \in H^{-1/2}(\Gamma)$ and $p \in H^1(\Gamma)$ such that for all $V \in H^1(\Omega)$, $\varphi \in H^{-1/2}(\Gamma)$ and $q \in H^1(\Gamma)$ there holds*

$$\begin{aligned} q_{\kappa}(U, V) + \left(W_{\kappa}(\gamma_D^- U), \gamma_D^- V \right)_{\Gamma} - \left(\left(\frac{1}{2} \text{Id} - K'_{\kappa} \right)(\theta), \gamma_D^- V \right)_{\Gamma} &= g_1(V) \\ \left(\left(\frac{1}{2} \text{Id} - K_{\kappa} \right)(\gamma_D^- U), \varphi \right)_{\Gamma} + (f V_{\kappa}(\theta), \varphi)_{\Gamma} + i \bar{\eta}(p, \varphi)_{\Gamma} &= g_2(\varphi) \\ - \left(W_{\kappa}(\gamma_D^- U), q \right)_{\Gamma} - \left((K'_{\kappa} + \frac{1}{2} \text{Id})(\theta), q \right)_{\Gamma} + b(p, q) &= g_3(q). \end{aligned} \tag{2}$$

1. The formulation provided here is equivalent to the case $U_i = 0, f = 0, n(x) = c_i/c_o$ in section 3 of [4].

where we have

$$\begin{aligned}
g_1(V) &:= -\left(f^2, \gamma_D^- V\right)_\Gamma - \left(W_\kappa(f^1), \gamma_D^- V\right)_\Gamma \\
g_2(\varphi) &:= \left(\left(K_\kappa - \frac{1}{2}\text{Id}\right)(f^1), \varphi\right)_\Gamma \\
g_3(q) &:= \left(W_\kappa(f^1), q\right)_\Gamma \\
q_\kappa(U, V) &:= \int_\Omega \text{grad } U \cdot \text{grad } \bar{V} - \kappa^2 n(x) U \bar{V} \, dx \\
b(p, q) &:= (\text{grad}_\Gamma p, \text{grad}_\Gamma q)_\Gamma + (p, q)_\Gamma.
\end{aligned}$$

This formulation is derived from the Helmholtz transmission problem by partially integrating the Helmholtz equation, applying Green's first formula, coupling the resulting variational problem to the BIEs using Dirichlet-to-Neumann maps, and transforming the Cauchy trace [4].

4 Operator Formulation

The form of the variational problem in eq. 2 can be rewritten in a simpler operator formulation. We start by rewriting the expressions for q_κ and b .

Lemma 4.1. *Let $U, V \in H^1(\Omega^-)$ such that U solves the Helmholtz equation. Then $q_\kappa(U, V) = (\text{DtN}_\kappa^- \gamma_D^- U, V)_\Gamma$ with the Dirichlet-to-Neumann map DtN_κ^- .*

Proof. Greens first formula $\int_U (\psi \text{grad } \varphi + \text{grad } \psi \cdot \text{grad } \varphi) dx = \oint_{\partial U} \psi \text{grad } \varphi \cdot n dS$ with normal vector n implies

$$\begin{aligned}
q_\kappa(U, V) &= \int_{\Omega^-} (\text{grad } U \text{ grad } \bar{V} - \tilde{c} \kappa^2 U \bar{V}) dx \\
&= \int_{\partial \Omega^-} \bar{V} \text{grad } U \cdot n dS - \int_{\Omega^-} (\tilde{c} \kappa^2 U + \Delta U) \bar{V} \cdot n dS \\
&= \int_{\partial \Omega^-} \bar{V} \text{grad } U \cdot n dS
\end{aligned}$$

with normal vector n where we used the Helmholtz equation in the third equation. Plugging in the definition for Γ and the Dirichlet-to-Neumann map yields the result. \square

Lemma 4.2. *Let ϕ be the angular polar coordinate in two dimensions. There exists a unique adjoint map grad'_Γ such that $(\hat{\phi} p, \text{grad}_\Gamma q)_\Gamma = (\text{grad}'_\Gamma p, q)_\Gamma$ for all $p \in H^1(\Gamma)$, $q \in H^1(\Gamma)$.*

Proof. We will start by proving boundedness of the operator $\hat{\phi} \text{grad}_\Gamma$. Note that for $q \in H^1(\Gamma)$ we have

$$\begin{aligned}
\|\hat{\phi} \text{grad}_\Gamma q\|_{L^2(\Gamma)}^2 &= (\text{grad}_\Gamma q, \text{grad}_\Gamma q)_\Gamma \\
&\leq (\text{grad}_\Gamma q, \text{grad}_\Gamma q)_\Gamma + (q, q)_\Gamma \\
&= \|q\|_{H^1(\Gamma)}^2 < \infty.
\end{aligned}$$

So in particular $\hat{\phi} \text{grad}_\Gamma$ is a map $H^1(\Gamma) \rightarrow L^2(\Gamma)$. We have

$$\|\hat{\phi} \text{grad}_\Gamma q\|_{L^2(\Gamma)} \leq \sqrt{\|\hat{\phi} \text{grad}_\Gamma q\|_{L^2(\Gamma)}^2 + \|q\|_{L^2(\Gamma)}^2} = \|u\|_{H^1(\Gamma)},$$

so $\hat{\phi} \text{grad}_\Gamma$ is bounded.

As $\hat{\phi} \text{grad}_\Gamma$ is bounded, according to the Riesz representation theorem [6], for every p there exists a unique p' such that for the functional $q \mapsto (p, \hat{\phi} \text{grad}_\Gamma q)$ we have $(p, \hat{\phi} \text{grad}_\Gamma q) = (p', q)$. Therefore $\text{grad}'_\Gamma : p \mapsto p'$ is the unique adjoint operator. \square

We define Fourier spaces \mathcal{H}^s to the spaces H^s similar to section 3 in [7]. Using these spaces will simplify calculations.

Definition 4.3. For $0 \leq s < \infty$ and $\tilde{\kappa} > 0$ the space $\mathcal{H}_{\tilde{\kappa}}^s(X)$ is defined as the subspace of all functions $\varphi \in L^2(X)$ such that

$$\sum_{n \in \mathbb{Z}} (\tilde{\kappa}^2 + n^2)^s |\varphi_n|^2 < \infty.$$

for the Fourier coefficients φ_n of φ . We define an inner product in this space by

$$(\varphi, \psi)_{\mathcal{H}^s(\mathbb{S})} := \sum_{n \in \mathbb{Z}} (\tilde{\kappa}^2 + n^2)^s \varphi_n \overline{\psi_n}.$$

We used the $\tilde{\kappa}$ -weighted norm that were also used in [1] for dimensional reasons.² The following lemma justifies the use of this space.

Lemma 4.4. Let $s \in \mathbb{R}$. Then the space $\mathcal{H}^s(\Gamma)_{\tilde{\kappa}}$ is a Hilbert space. Moreover, $\mathcal{H}_{\tilde{\kappa}}^s(\Gamma) = H^s(\Gamma)$ and the norms generated from their respective inner products are equivalent.

Proof. The case $\tilde{\kappa} = 1$ is proven in Theorem 2 of [7]. Now consider a general $\tilde{\kappa}$. We have to show that $\mathcal{H}_{\tilde{\kappa}}^s(\Gamma) = \mathcal{H}_1^s(\Gamma)$, equivalence of their norms and completeness of $\mathcal{H}_{\tilde{\kappa}}^s(\Gamma)$.

By the limit comparison test convergence of $\sum_{n \in \mathbb{Z}} (\tilde{\kappa}^2 + n^2)^s |\varphi_n|^2$ and $\sum_{n \in \mathbb{Z}} (1 + n^2)^s |\varphi_n|^2$ are equivalent, because

$$\lim_{n \rightarrow \infty} \frac{(\tilde{\kappa}^2 + n^2)^s |\varphi_n|^2}{(1 + n^2)^s |\varphi_n|^2} = 1.$$

Thus $\mathcal{H}_1^s = \mathcal{H}_{\tilde{\kappa}}^s$.

For equivalence of norms we have to find c_1, c_2 such that $c_1 \|\varphi\|_{\mathcal{H}_1^s} \leq \|\varphi\|_{\mathcal{H}_{\tilde{\kappa}}^s} \leq c_2 \|\varphi\|_{\mathcal{H}_1^s}$. This is satisfied by $c_1 = 1, c_2 = \tilde{\kappa}^s$ for $\tilde{\kappa} > 1, s > 0$, $c_1 = \tilde{\kappa}^s, c_2 = 1$ for $\tilde{\kappa} < 1, s > 0$, $c_1 = \tilde{\kappa}^s, c_2 = 1$ for $\tilde{\kappa} > 1, s < 0$, and $c_1 = 1, c_2 = \tilde{\kappa}^s$ for $\tilde{\kappa} < 1, s < 0$.

Completeness of $\mathcal{H}_{\tilde{\kappa}}^s(\Gamma)$ is implied by completeness of $\mathcal{H}_1^s(\Gamma)$: $\mathcal{H}_1^s(\Gamma)$ and $\mathcal{H}_{\tilde{\kappa}}^s(\Gamma)$ have the same Cauchy sequences because of equivalence of their norms. Since the spaces are equal, if the Cauchy sequence converges in one it does in the other. $\mathcal{H}_1^s(\Gamma)$ is complete, so $\mathcal{H}_{\tilde{\kappa}}^s(\Gamma)$ is also as their norms are equivalent and they contain the same vectors. \square

This isomorphism between the Hilbert spaces $\mathcal{H}_{\tilde{\kappa}}^s$ and H^s justifies the use of the lemmata 4.1 and 4.2 for the spaces $\mathcal{H}_{\tilde{\kappa}}^s$. Using these results, we can rewrite the formulation in eq. 2 in an operational formulation: Define the operator

$$A : \mathcal{H}_{\tilde{\kappa}}^1(\Omega) \times \mathcal{H}_{\tilde{\kappa}}^{-\frac{1}{2}}(\Gamma) \times \mathcal{H}_{\tilde{\kappa}}^1(\Gamma) \rightarrow \mathcal{H}_{\tilde{\kappa}}^{-\frac{1}{2}}(\Gamma) \times \mathcal{H}_{\tilde{\kappa}}^{\frac{1}{2}}(\Gamma) \times \mathcal{H}^{-\frac{1}{2}}(\Gamma)$$

$$A = \begin{pmatrix} (\text{DtN}_{\tilde{\kappa}}^- + W_{\tilde{\kappa}})\gamma_D^- & -(\frac{1}{2} - K'_{\tilde{\kappa}}) & 0 \\ (\frac{1}{2} - K_{\tilde{\kappa}})\gamma_D^- & V_{\tilde{\kappa}} & i\bar{\eta} \\ -W_{\tilde{\kappa}}\gamma_D^- & -(K'_{\tilde{\kappa}} + \frac{1}{2}) & 1 + \text{grad}'_{\Gamma} \hat{\phi} \text{grad}_{\Gamma} \end{pmatrix}$$

Moreover, define

$$\mathcal{B} : \mathcal{H}_{\tilde{\kappa}}^{-\frac{1}{2}} \times \mathcal{H}_{\tilde{\kappa}}^{\frac{1}{2}} \rightarrow \mathcal{H}_{\tilde{\kappa}}^{-\frac{1}{2}} \times \mathcal{H}_{\tilde{\kappa}}^{\frac{1}{2}} \times \mathcal{H}_{\tilde{\kappa}}^{-\frac{1}{2}}, \mathcal{B} = \begin{pmatrix} -1 & -W_{\kappa} \\ 0 & (K_{\kappa} - \frac{1}{2}) \\ 0 & W_{\kappa} \end{pmatrix}.$$

2. Note that n implicitly has a dimension $[\frac{1}{r}]$ here. Since we did not write down dimensions for it ($r = 1$), we do not make this dependency explicit in our calculations.

Proposition 4.5. *If $U \in \mathcal{H}^1(\Omega)$, $\theta \in \mathcal{H}^{-1/2}(\Gamma)$ and $p \in \mathcal{H}^1(\Gamma)$ solve the following problem*

$$\begin{pmatrix} (\text{DtN}_\kappa^- + W_\kappa)\gamma_D^- & -(\frac{1}{2} - K'_\kappa) & 0 \\ (\frac{1}{2} - K_\kappa)\gamma_D^- & V_\kappa & i\bar{\eta} \\ -W_\kappa\gamma_D^- & -(K'_\kappa + \frac{1}{2}) & 1 + \text{grad}'_\Gamma \hat{\phi} \text{grad}_\Gamma \end{pmatrix} \begin{pmatrix} U \\ \theta \\ p \end{pmatrix} = \mathcal{B} \begin{pmatrix} f_2 \\ f_1 \end{pmatrix} \quad (3)$$

they also solve the problem in eq. 2.

Proof. We start our proof from the equations presented in eq. 2. After plugging in the results presented in lemma 4.1 and 4.2 for q_κ and b , we have to show that for arbitrary $V \in \mathcal{H}^1(\Gamma)$, $\varphi \in \mathcal{H}^{-\frac{1}{2}}(\Gamma)_\kappa$, $p \in H_\kappa^1(\Gamma)$ eq. 3 implies

$$\begin{aligned} (\text{DtN}_\kappa^- U, V)_\Gamma + (W_\kappa(\gamma_D^- U), \gamma_D^- V)_\Gamma - \left(\left(\frac{1}{2} \text{Id} - K'_\kappa \right)(\theta), \gamma_D^- V \right)_\Gamma &= g_1(V) \\ \left(\left(\frac{1}{2} \text{Id} - K_\kappa \right)(\gamma_D^- U), \varphi \right)_\Gamma + (V_\kappa(\theta), \varphi)_\Gamma + i\bar{\eta}(p, \varphi)_\Gamma &= g_2(\varphi) \\ - (W_\kappa(\gamma_D^- U), q)_\Gamma - \left((K'_\kappa + \frac{1}{2} \text{Id})(\theta), q \right)_\Gamma + (\text{grad}'_\Gamma \hat{\phi} \text{grad}_\Gamma p, q)_\Gamma + (p, q)_\Gamma &= g_3(q). \end{aligned}$$

This follows directly from eq. 3 by taking the inner product with V for the first row, with φ for the second row, and with q for the third row of A . \square

Proposition 4.5 allows us to restrict our search of the solution from $U \in H^1(\Gamma)$ to $\gamma_D^- U \in H^{\frac{1}{2}}(\Gamma)$, as all entries in the first column of the operator matrix include a factor γ_D^- . Therefore, we will consider the problem with $U \in H^{\frac{1}{2}}(\Gamma)$ in all following sections.³ Moreover, we will also consider $p \in \mathcal{H}^{\frac{1}{2}}$ since this simplifies the problem by reducing the number of different vector spaces and makes it more consistent with general theoretical results in [4]. The following lemma shows that a solution of this augmentation remains a solution of eq. 3.

Lemma 4.6. *Let $s \leq t$. Then $\mathcal{H}_\kappa^t \subset \mathcal{H}_\kappa^s$.*

Proof. This follows directly by applying the majorant criterion to definition 4.3. \square

Considering these augmentations, we consider the problem $A \begin{pmatrix} U \\ \theta \\ p \end{pmatrix} = \mathcal{B} \begin{pmatrix} f_2 \\ f_1 \end{pmatrix}$ with operators

$$A : \mathcal{H}_\kappa^{\frac{1}{2}} \times \mathcal{H}_\kappa^{-\frac{1}{2}} \times \mathcal{H}_\kappa^{\frac{1}{2}} \rightarrow \mathcal{H}_\kappa^{-\frac{1}{2}} \times \mathcal{H}_\kappa^{\frac{1}{2}} \times \mathcal{H}_\kappa^{-\frac{1}{2}}, \quad \mathcal{B} : \mathcal{H}_\kappa^{-\frac{1}{2}} \times \mathcal{H}_\kappa^{\frac{1}{2}} \rightarrow \mathcal{H}_\kappa^{\frac{1}{2}} \times \mathcal{H}_\kappa^{-\frac{1}{2}} \times \mathcal{H}_\kappa^{\frac{1}{2}}$$

5 Spectral Discretisation

Similar to Hiptmair et al. in [1], we want to investigate the operator norm of the inverse operator A^{-1} . This norm indicates the invertibility of the operator A . We will investigate whether the stabilized BIE-volume formulation regularises any spurious quasi-resonances in the operator norm that might appear as a spectral phenomenon. For the following computations, we will restrict our solution space to a finite subspace $\mathcal{S}_N^{\frac{1}{2}} \times \mathcal{S}_N^{-\frac{1}{2}} \times \mathcal{S}_N^1$ where $N \in \mathbb{N}$ and \mathcal{S}_N^s is the restriction to $\text{span}(y_{-N}, y_{-N+1}, \dots, y_N)$ where $y_n = e^{in\phi}$.

As the following Lemma shows, we can calculate the norm of the component-wise restricted operator A_S from the smallest singular value of the matrix representation of A_S , if we choose an orthonormal basis.

3. As in the chapter 9 the solution in all of Ω^- the Helmholtz PDE allows reconstruction of the solution in all of Ω^- from the Fourier coefficients on Γ .

Lemma 5.1. *Let $F : V \rightarrow W$ be a linear operator between normed vector spaces V, W of finite dimension with norms $\|\cdot\|_V, \|\cdot\|_W$. Let $\mathcal{B}_V, \mathcal{B}_W$ be orthonormal bases of these spaces and $C := M_{\mathcal{B}_V}^{\mathcal{B}_W}(F)$ the representation matrix of F with respect to these bases. Then*

$$\|F\|_{op} := \sup_{\|v\|_V=1} \|F(v)\|_W = \|M\|_{op} := \sup_{|x|=1} |Cx|$$

where $|\cdot|$ is the euclidian norm.

Proof. Let $(v_i)_{i=1}^N = \mathcal{B}_V, (w_i)_{i=1}^M = \mathcal{B}_W$. Then by definition $Fv_i = \sum_{j=0}^M C_{ij}w_j$.

$$\begin{aligned} \|F\|_{op} &:= \sup_{\|v\|_V=1} \|F(v)\|_W = \sup_{\|v\|_V=1} \|F(v)\|_W = \sup_{\|v\|_V=1} \left\| \sum_{i=0}^N \sum_{j=0}^M C_{ij}x_iw_j \right\|_W \\ &= \sup_{|x|=1} \left| \sum_{i=0}^N \sum_{j=0}^M C_{ij}x_i \right| = \sup_{|x|=1} |Cx| \end{aligned} \quad (4)$$

where we used orthogonality $\|v\|_V = \left\| \sum_{i=0}^N x_i v_i \right\|_V = |x|$ (and similar for w_j) in the fourth equation. \square

Together with the results from linear algebra, that the biggest singular value of a matrix corresponds to its operator norm and that inverting a matrix yields to inversion of its singular values we obtain that the operator norm of A_S^{-1} is

$$\|A_S^{-1}\|_{op} = \frac{1}{\sigma_{min}(A^{num})}$$

where A^{num} is a representation of A_S with respect to an orthonormal basis and $\sigma_{min}(A_S)$ its smallest singular value.

An orthonormal basis of \mathcal{S}_N^s is given by

$$\left(d_n^s e^{in\phi} \right)_{n=-N}^N \text{ where } d_n^s := \frac{1}{(\tilde{\kappa}^2 + n^2)^{\frac{s}{2}}} \quad (5)$$

We will use the following variables as coefficients.

$$v_n = d_n^{\frac{1}{2}}, w_n = d_n^{-\frac{1}{2}}, l_n = d_n^{\frac{1}{2}}. \quad (6)$$

Moreover, define $c_n = \sqrt{\tilde{\kappa}^2 + n^2}$.

Now we can construct the Fourier Galerkin matrix. The following lemma from Theorem 2 of [7] establishes a simple form of the BIOs applied to Fourier monomials.

Lemma 5.2. *The following eigenvalue equations hold for $V_\kappa, K_\kappa, K'_\kappa$, and \mathcal{W}_κ .*

$$\begin{aligned} V_\kappa e^{in\phi} &= \lambda^{(V)} e^{in\phi}, \quad \lambda^{(V)} := \frac{i\pi}{2} J_n(\kappa) H_n^{(1)}(\kappa) \\ K_\kappa e^{in\phi} &= \lambda^{(K)} e^{in\phi}, \quad \lambda^{(K)} := \frac{i\pi\kappa}{2} J_n(\kappa) H_n^{(1)'}(\kappa) + \frac{1}{2} = \frac{i\pi\kappa}{2} J'_n(\kappa) H_n^{(1)}(\kappa) - \frac{1}{2} \\ K'_\kappa e^{in\phi} &= \lambda^{(K')} e^{in\phi}, \quad \lambda^{(K')} := \frac{i\pi\kappa}{2} J_n(\kappa) H_n^{(1)'}(\kappa) + \frac{1}{2} = \frac{i\pi\kappa}{2} J'_n(\kappa) H_n^{(1)}(\kappa) - \frac{1}{2} \\ \mathcal{W}_\kappa e^{in\phi} &= \lambda^{(W)} e^{in\phi}, \quad \lambda^{(W)} := -\frac{i\pi\kappa^2}{2} J'_n(\kappa) H_n^{(1)'}(\kappa). \end{aligned}$$

Furthermore, we have to find explicit expressions for DtN_κ^- and $\text{grad}'_\Gamma \hat{\phi} \text{grad}_\Gamma$.

Lemma 5.3. *We have $\text{DtN}_\kappa^- v_l e^{il\phi} = \alpha_l v_l e^{il\phi}$ and $\text{grad}'_\Gamma \hat{\phi} \text{grad}_\Gamma v_l e^{il\phi} = \beta_l v_l e^{il\phi}$ where*

$$\alpha_l = \sqrt{\tilde{c}} \kappa \frac{J'_l(\sqrt{\tilde{c}} \kappa)}{J_l(\sqrt{\tilde{c}} \kappa)}, \quad \beta_l = l^2 \quad (7)$$

Proof. To derive the Dirichlet-to-Neumann map we can consider the original problem in eq. 1. We make the Fourier Ansatz $V_l = V_l^r e^{il\phi}$. As U must satisfy $(\Delta + \tilde{c}\kappa^2)U = 0$ from eq. 1, we have

$$r^2 \partial_r^2 V_l^r + r \partial_r V_l^r + (r^2 \tilde{c}\kappa^2 - l^2) V_l^r = 0.$$

This is Bessel's differential equation. Since we require convergence at the origin this implies $V_l^r(r) = J_l(\sqrt{\tilde{c}\kappa}r)$. Converting this to the same scaling as $v_l e^{il\phi}$ on the boundary, we get $n \cdot \text{grad} \left(\frac{J_l(\sqrt{\tilde{c}\kappa}r) v_l e^{il\phi}}{J_l(\sqrt{\tilde{c}\kappa})} \right) = \sqrt{\tilde{c}\kappa} \frac{J_l'(\sqrt{\tilde{c}\kappa}r) v_l e^{il\phi}}{J_l(\sqrt{\tilde{c}\kappa})}$ where n is the normal vector and we used $n \cdot \text{grad} = \frac{\partial}{\partial r}$.

To derive the expression for the composite gradient, consider the inner product of the expression with $e^{im\phi}$:

$$\begin{aligned} (\text{grad}'_\Gamma \hat{\phi} \text{grad}_\Gamma v_l e^{il\phi}, e^{im\phi})_\Gamma &= (\text{grad}_\Gamma v_l e^{il\phi}, \text{grad}_\Gamma e^{im\phi})_\Gamma \\ &= v_l l m (2\pi \delta_{lm}) \end{aligned}$$

This implies $\beta_n = l^2$. □

Using these insights, we can find the Fourier Galerkin matrix. Since the Fourier modes are eigenvectors of each of the entries in the operator A we obtain diagonal blocks of the form in the representation with respect to the basis $e^{in\phi}$ for each \mathcal{S}_N^s component

$$A_n^{num,0} \begin{pmatrix} (\alpha_n + \lambda^{(W)}) & -(\frac{1}{2} - \lambda^{(K')}) & 0 \\ (\frac{1}{2} - \lambda^{(K)}) & \lambda^{(V)} & i\bar{\eta} \\ -\lambda^{(W)} & -(\lambda^{(K')} + \frac{1}{2}) & (1 + \beta_n) \end{pmatrix}.$$

Then we obtain the Galerkin matrix scaled to the selected bases of $\mathcal{S}_N^{\frac{1}{2}} \times \mathcal{S}_N^{-\frac{1}{2}} \times \mathcal{S}_N^{\frac{1}{2}}$ and $\mathcal{S}_N^{-\frac{1}{2}} \times \mathcal{S}_N^{\frac{1}{2}} \times \mathcal{S}_N^{-\frac{1}{2}}$ defined in eq. 5:

$$A_n^{num} := T_1 A_N^{num,0} T_2 = \begin{pmatrix} (\alpha_n + \lambda^{(W)})(c_n)^{-1} & -(\frac{1}{2} - \lambda^{(K')}) & 0 \\ (\frac{1}{2} - \lambda^{(K)}) & \lambda^{(V)} c_n & i\bar{\eta} \\ -\lambda^{(W)}(c_n)^{-1} & -(\lambda^{(K')} + \frac{1}{2}) & (1 + \beta_n)(c_n)^{-1} \end{pmatrix}. \quad (8)$$

with the basis scaling blocks

$$T_1 = \begin{pmatrix} (c_n)^{-\frac{1}{2}} & & \\ & (c_n)^{\frac{1}{2}} & \\ & & (c_n)^{-\frac{1}{2}} \end{pmatrix}, \quad T_2 = \begin{pmatrix} (c_n)^{-\frac{1}{2}} & & \\ & (c_n)^{\frac{1}{2}} & \\ & & (c_n)^{-\frac{1}{2}} \end{pmatrix}.$$

For the overall matrix we write $A^{num} = \text{diag}(A_{-N}^{num}, A_{-N+1}^{num}, \dots, A_N^{num})$.

6 Validation

6.1 Operator solution

To validate the correctness of our derived matrix in eq. 8 we demonstrate that it yields the correct numerical solution for a simple example. Consider the special case

$$\vec{f} = \begin{pmatrix} H_n^{(1)}(\kappa) - J_n(\sqrt{\tilde{c}\kappa}) \\ \kappa H_n'^{(1)}(\kappa) - \sqrt{\tilde{c}\kappa} J_n'(\sqrt{\tilde{c}\kappa}) \end{pmatrix} e^{in\phi}$$

where $n = -N, -N+1, \dots, N$. Then the solution to eq. 1 is

$$U = J_n(\sqrt{\tilde{c}\kappa}r) e^{in\phi} \text{ for } r < 1, U = H_n^{(1)}(\kappa r) e^{in\phi} \text{ for } r > 1.$$

This can be seen directly by plugging in. If eq. 8 is correct, it should yield the complete analytical solution as the restricted space contains the analytical solution.

As pointed out on p. 33 of [4], if U is a solution we have $\vartheta = \gamma_N^+ U$. Moreover, on page 38 of [4] it is explained that in the solution vector we must have $p = 0$. Therefore, it follows that the boundary-restricted interior analytical solution can be written as $(U, \theta, p) = (J_n(\sqrt{\tilde{c}\kappa})e^{in\phi}, \kappa H_n'^{(1)}(\kappa)e^{in\phi}, 0)$. So overall, the components solution vector should be

$$(C_j^{ana})^U = \delta_{nj} \frac{J_n(\sqrt{\tilde{c}\kappa})}{v_n}, \quad (C_j^{ana})^\theta = \delta_{nj} \frac{\kappa H_n'^{(1)}(\kappa)}{w_n}, \quad (C_j^{ana})^p = 0, \quad \forall j.$$

We can now validate whether we get the same numerical solution using our matrix A^{num} . To do this, also extend the right hand side of eq. 3, namely $\mathcal{B} \begin{pmatrix} f_2 \\ f_1 \end{pmatrix}$ into the Fourier coefficients of f^1, f^2 . We obtain blocks of the form

$$f_n^1 e^{in\phi} \begin{pmatrix} -\lambda_n^{(W)} \\ (\lambda_n^{(K)} - 0.5) \\ \lambda_n^{(W)} \end{pmatrix} + f_2^n e^{in\phi} \begin{pmatrix} -1 \\ 0 \\ 0 \end{pmatrix}. \quad (9)$$

In this particular case, the boundary conditions translate to

$$f_j^1 = \delta_{nj} (H_n^{(1)}(\kappa) - J_n(\sqrt{\tilde{c}\kappa})), \quad f_j^2 = \delta_{nj} (\kappa H_n'^{(1)}(\kappa) - \sqrt{\tilde{c}\kappa} J_n'(\sqrt{\tilde{c}\kappa})).$$

So, writing our right hand side in the $\mathcal{H}_{\tilde{\kappa}}^{-\frac{1}{2}} \times \mathcal{H}_{\tilde{\kappa}}^{\frac{1}{2}} \times \mathcal{H}_{\tilde{\kappa}}^{-\frac{1}{2}}$ basis:

$$g_n^{num} := \underbrace{f_n^1 \begin{pmatrix} -(c_n)^{-\frac{1}{2}} \lambda_n^{(W)} \\ (c_n)^{\frac{1}{2}} (\lambda_n^{(K)} - 0.5) \\ (c_n)^{-\frac{1}{2}} \lambda_n^{(W)} \end{pmatrix}}_{\vec{x}_1} + \underbrace{f_2^n \begin{pmatrix} -(c_n)^{-\frac{1}{2}} \\ 0 \\ 0 \end{pmatrix}}_{\vec{x}_2}.$$

We call the overall vector $g^{num} = (g_{-N}^{num}, g_{-N+1}^{num}, \dots, g_N^{num})$. To measure how good the solution is, we introduce the ζ -number:

Definition 6.1 (ζ -number). *For a fixed n and a fixed κ , let $\vec{C}_n^{num}(\kappa)$ be the numerical solution vector to the problem $A_n^{num} C_n^{num}(\kappa) = g^{num}$. Let $C_n^{ana}(\kappa)$ be the analytical solution vector for the same problem. The ζ -number is defined as*

$$\zeta(\kappa) = \max_{n \in [-N, \dots, N]} \frac{\|C_n^{num}(\kappa) - C_n^{ana}(\kappa)\|}{\|C_n^{ana}(\kappa)\|}$$

As seen in fig. 1, the ζ -number is negligible across all values of κ and n as deviations from 0 are in the magnitude of computer precision. This validates our derivation of A_n^{num} and g_n^{num} .

6.2 Projector Properties

We validate our matrix A with another method. As mentioned earlier, we must have $p = 0$, if (U, θ) solves the problem according to [4]. Assuming A_n^{num} is invertible⁴, we can use the remark to validate that our derived matrix A_n^{num} is correct. We can formulate the property $p = 0$ with regards to A_n^{num} . Let $V_{\vec{b}} := \text{span}(\vec{x}_1, \vec{x}_2)$. Then $p = 0$ means that for $\vec{b} \in V_{\vec{b}}$, every solution of $A_n^{num} \vec{x} = \vec{b}_n$ satisfies $x_3 = 0$. Now let $P_{V_{\vec{b}}}$ be the projector onto $V_{\vec{b}}$ and P_3 be the projector onto $(0, 0, 1)$. Then $P_3 A_n^{-1} P_{V_{\vec{b}}} = 0$. We calculate the euclidian norm of $P_3 (A_n^{num})^{-1} P_{V_{\vec{b}}}$ for a range of κ values to validate this. Again, as shown in fig. 2, this is satisfied. This is another validation of our matrix A_n^{num} .

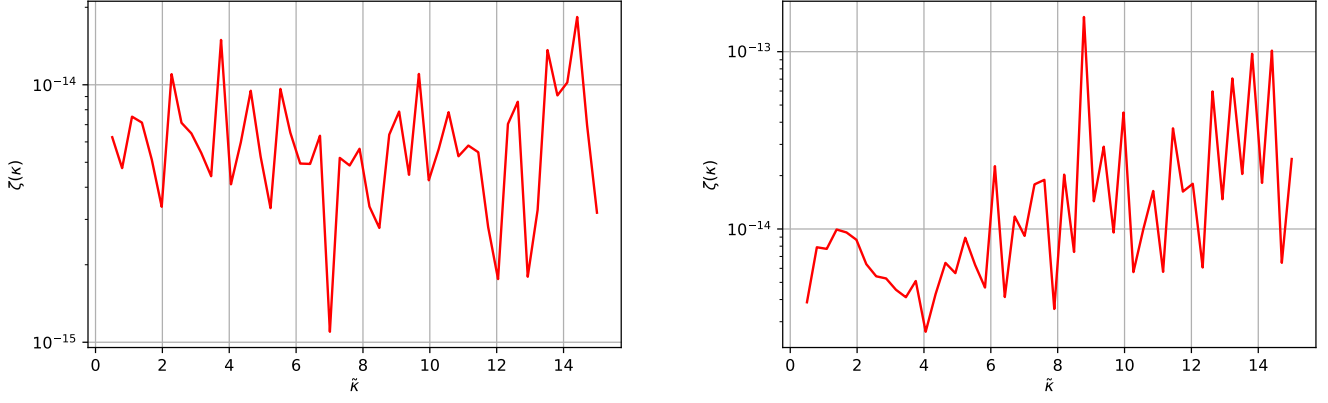


Fig. 1: Maximum relative residuum of the numerical solution ζ by wave number. On the left plot we have $c_i = 1, c_o = 3$ and on the right plot we have $c_i = 3, c_o = 1$. The number of considered Fourier modes was $N = 30$.

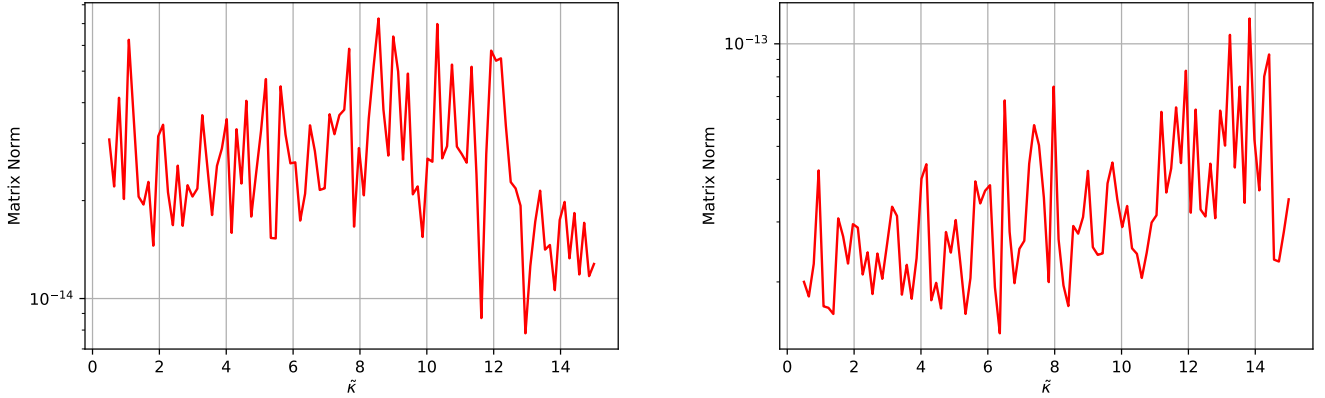


Fig. 2: Euclidian matrix norm of composed matrix $P_3(A_n^{num})^{-1}P_{V_b}$. On the left plot we have $c_i = 1, c_o = 3$ and on the right plot we have $c_i = 3, c_o = 1$. The number of considered Fourier modes was $N = 100$.

7 Numerical Results

Now, we numerically investigate the operator norm of the inverse operator A^{-1} . As established in section 5, we can compute the inverse of the smallest singular value of $\text{diag}((A_n^{num})_{n=-N}^N)$ for this. The results of the simulation are presented in fig. 3 alongside the results for the maximum euclidean norm of the solution operator.⁵ A derivation of the solution operator is provided in the [Appendix](#).

The solution operator exposes resonance frequencies for the case $c_i = 3, c_o = 1$, called quasi-resonances in [1]. This resonance behavior is expected physically: If $c_i > c_o$, total internal reflection can undergo. For certain angles of internal reflection, the solution becomes very localised around the boundary Γ , meaning that the solution operator norm peaks. As we see, peaks in the solution operators coincide with the peaks of the considered inverse boundary integral operator. However, in the case $c_i = 1, c_o = 3$ the operator A^{-1} exposes bad condition for some frequencies while the solution operator is regular for all frequencies. This shows that spurious quasi-resonances also occur for the regularised variational

4. This is fair to assume for most frequencies since the variational problem has a unique solution.

5. We picked particular representative cases. More cases are shown in section 9.

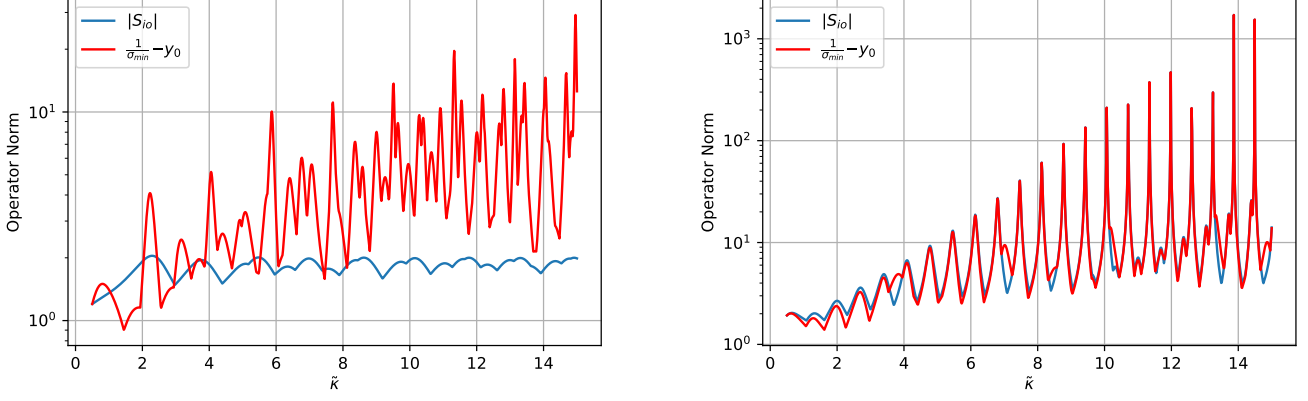


Fig. 3: Inverted minimum singular value $\|A_S^{-1}\|_{op} = \frac{1}{\sigma_{\min}(M)}$ of the matrix $M = \text{diag}((A_n^{num})_{n=-N}^N)$ by wave number κ . Only the first 30 Fourier blocks were considered, as the operator norms were not affected by higher modes. In fact, the highest selected mode was $n = 20$ for the left and $n = 21$ for the right plot. When plotting the operator norm against N for a fixed κ , we see that it approaches a constant value for big N , justifying the cutoff of $N = 30$. A more detailed explanation is given in the section [Appendix](#). Each of plots is shifted by an absolute value y_0 to allow for better comparison with the solution operator. In the left plot we have $c_i = 1, c_o = 3$, and $y_0 = 0.712$ and on the right plot we have $c_i = 3, c_o = 1$, and $y_0 = 0.450$.

formulation. However, it does affect the condition much less than other examples: the STF BIO considered by Hiptmair et al. exposed resonance peaks that were three orders of our magnitude higher than our considered case [1]. The following proposition helps explain this irregular behavior.

Proposition 7.1. Define $\mathcal{L} : \mathcal{H}_{\tilde{\kappa}}^{\frac{1}{2}} \times \mathcal{H}_{\tilde{\kappa}}^{-\frac{1}{2}} \rightarrow \mathcal{H}_{\tilde{\kappa}}^{-\frac{1}{2}} \times \mathcal{H}_{\tilde{\kappa}}^{\frac{1}{2}} \times \mathcal{H}_{\tilde{\kappa}}^{-\frac{1}{2}}$ as the map that maps the boundary conditions of eq. 1 to the right hand side of eq. 3

$$\mathcal{L} = \begin{pmatrix} -W_{\kappa} & -id \\ (K_{\kappa} - \frac{1}{2}) & 0 \\ W_{\kappa} & 0 \end{pmatrix}$$

and let $E : \mathcal{H}_{\tilde{\kappa}}^{\frac{1}{2}} \times \mathcal{H}_{\tilde{\kappa}}^{-\frac{1}{2}} \times \mathcal{H}_{\tilde{\kappa}}^{\frac{1}{2}} \rightarrow \mathcal{H}_{\tilde{\kappa}}^{\frac{1}{2}} \times \mathcal{H}_{\tilde{\kappa}}^{-\frac{1}{2}}, E = \begin{pmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \end{pmatrix}$. Then we can write the solution operator (see definition 1.2) as $S = EA^{-1}\mathcal{L}$.

Proof. Consider the boundary data $f \in \mathcal{H}_{\tilde{\kappa}}^{\frac{1}{2}} \times \mathcal{H}_{\tilde{\kappa}}^{-\frac{1}{2}}$. Then by construction we have that the right hand side of eq. 3 is $g = \mathcal{L}f$. Applying A^{-1} to the same equation and projecting onto the first two components yields $\begin{pmatrix} U \\ \theta \end{pmatrix} = EA^{-1}\mathcal{L}$. This concludes the proof. \square

This relationship between A and S implies that the lack of either E or L causes the different resonance behavior between A and S . However, it cannot be E : analytically, the third solution component $p = 0$ vanishes as mentioned and numerically validated in section 6.2. Therefore, L must be the origin of the irregularity. To show this, we lift the irregularity by introducing the operator L_{ε} for $\varepsilon > 0$:

$$L_{\varepsilon} : \mathcal{H}_{\tilde{\kappa}}^{-\frac{1}{2}} \times \mathcal{H}_{\tilde{\kappa}}^{\frac{1}{2}} \times \mathcal{H}_{\tilde{\kappa}}^{-\frac{1}{2}} \rightarrow \mathcal{H}_{\tilde{\kappa}}^{-\frac{1}{2}} \times \mathcal{H}_{\tilde{\kappa}}^{\frac{1}{2}} \times \mathcal{H}_{\tilde{\kappa}}^{-\frac{1}{2}}, L_{\varepsilon} = \begin{pmatrix} -id & -W_{\kappa} & 0 \\ 0 & (K_{\kappa} - \frac{1}{2}) & 0 \\ 0 & W_{\kappa} & \varepsilon id \end{pmatrix}. \quad (10)$$

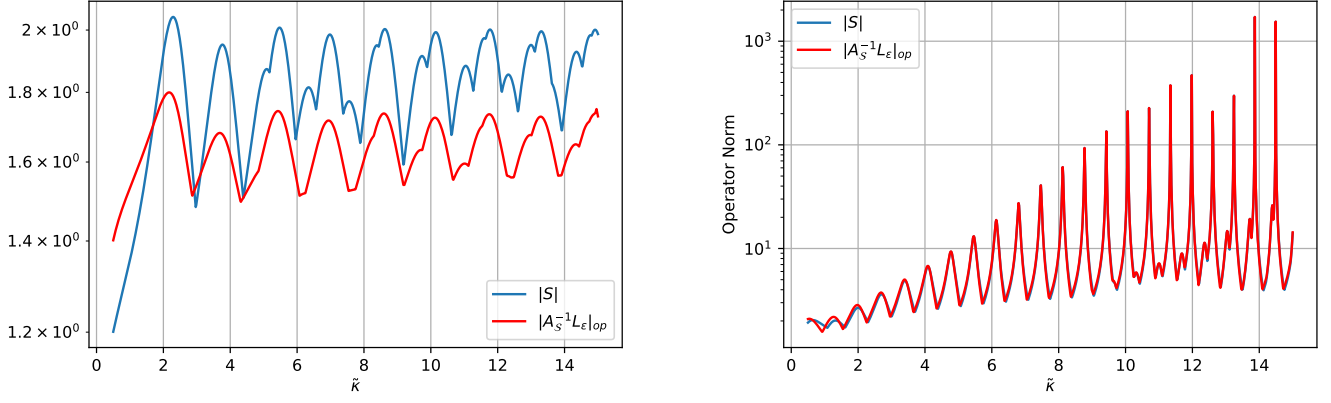


Fig. 4: Inverted minimum singular value $\|A_S^{-1}L_\varepsilon\|_{op} = \frac{1}{\sigma_{\min}(M)}$ by wave number κ . The highest selected mode out of the considered first $N = 30$ modes was $n = 15$ for the left and $n = 22$ for the right plot. In the left plot we have $c_i = 1$ and $c_o = 3$ and on the right plot we have $c_i = 3$ and $c_o = 1$.

Here we introduced the parameter ε to formally introduce a new boundary condition parameter, allowing us to invert the operator formally. As this new component is strongly suppressed by ε , it will not impact the outcome. We will use $\varepsilon = 0.01$ in the following calculations. Assume that L_ε is invertible and consider the operator $L_\varepsilon^{-1}A : \mathcal{H}_{\tilde{\kappa}}^{-\frac{1}{2}} \times \mathcal{H}_{\tilde{\kappa}}^{\frac{1}{2}} \times \mathcal{H}_{\tilde{\kappa}}^{-\frac{1}{2}} \rightarrow \mathcal{H}_{\tilde{\kappa}}^{\frac{1}{2}} \times \mathcal{H}_{\tilde{\kappa}}^{-\frac{1}{2}} \times \mathcal{H}_{\tilde{\kappa}}^{\frac{1}{2}}$. Note that while we introduced an additional boundary condition parameter in the third argument, it is strongly suppressed by ε . Invertibility is satisfied for the considered case of Γ , as L_ε becomes a blockdiagonal operator in the Fourier basis with non-vanishing determinant in each block. As fig. 4 shows, the irregular resonance behavior is indeed removed in the case $c_i = 1$, $c_o = 3$ while in the case $c_i = 3$, $c_o = 1$ the resonance behavior still agrees. This implies that the spurious quasi-resonances of the considered regularised formulation operator in eq. 3 were indeed caused by the operator applied to the boundary data on the right hand side of the problem.

However, we did not prove the invertibility of L_ε for general Ω^- . In particular, we did not provide an explicit expression for this inverse.⁶ Future investigations into this area could include other augmentations of the considered operator that remove these spurious quasi-resonances while also being explicitly expressed for general bounded Lipschitz domains $\Omega^- \subset \mathbb{R}^d$.

8 Conclusion

This report considers a regularised variational formulation of the Helmholtz transmission problem on a two-dimensional disk (eq. 1). We converted this variational formulation to an operator formulation (eq. 3) to simplify the calculations and theoretical results following. After discretizing this formulation (eq. 8), we calculated the inverse operator norm of the considered operator A (left side operator of eq. 3) numerically. We found, that for $c_i < c_o$ the operator exposed unphysical spurious quasi-resonances and growth that the solution operator did not (fig. 3). However, the occurring resonances, were smaller than for other operator formulations ([1]). By composing with another operator $\mathcal{L}_\varepsilon^{-1}$ (eq. 10), we were able to remove these spurious quasi-resonances and explain their origin by relating the domain of A^{-1} and the solution operator S . However, more general augmentation techniques are required as the considered operator \mathcal{L}_ε was not proven to be invertible generally, and no explicit inverse was derived for it.

6. While a formal inverse (treating the entries of L_ε as matrix entries) can be computed symbolically, expressions in the symbolic results may not be well-defined.

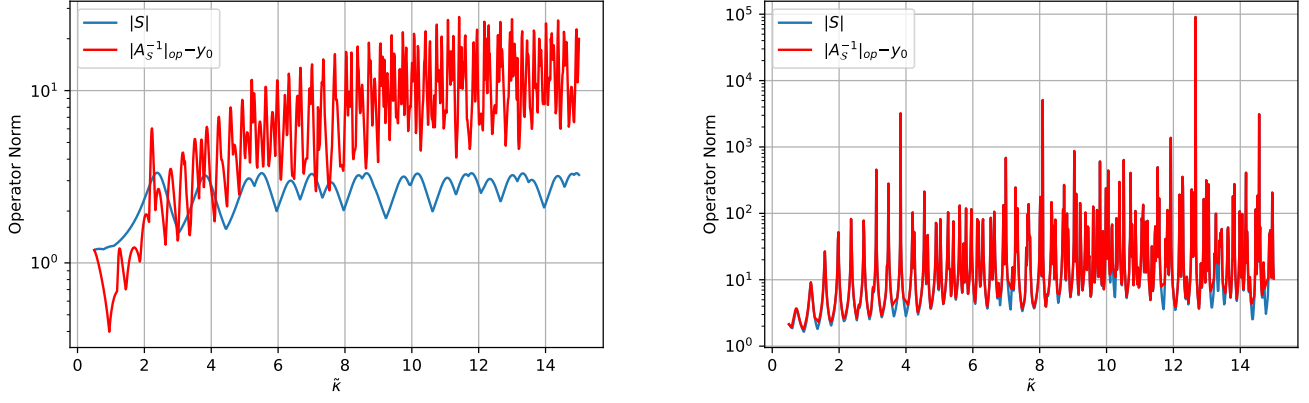


Fig. 5: Inverted minimum singular value $\|A_S^{-1}\|$ by wave number κ . In the upper left plot we have $c_i = 1$, $c_o = 10$ and $y_0 = 1.336$ and on the upper right plot we have $c_i = 10$, $c_o = 1$ and $y_0 = -0.159$.

9 Appendix

9.1 Numerical Results for other examples

We considered the example $c_i = 1$, $c_o = 3$ and vice versa in section 7. In fig. 5 we present other examples for the refractive indices, showing that the overall trends and occurrence of spurious quasi-resonances is very similar to the considered example.

Moreover, we only studied $\eta = 1.0$ in section 7. As shown in fig. 7, the operator growth is smallest around $\eta = 1.0$ out of the selected values. This is in agreement with the results of P. Meury's doctoral thesis, where he achieved optimal stability around $\eta = 1.0$ [4].

Lastly, we will justify the introduction of a cutoff N for the Fourier modes. In fig. 7 the singular value of the operator block A_n^{num} inverse are plotted as a function of the mode n . We clearly see, that for n big enough the singular value converges. There is also a simple theoretical explanation to justify the cutoff in n : consider the matrix A_n^{num} as defined in eq. 8. We have the approximate form $J_n(x) \sim \frac{1}{\Gamma(n+1)} \left(\frac{x}{2}\right)^n$ for $x \ll \sqrt{n+1}$ [8]. This implies that $\lim_{n \rightarrow \infty} \frac{\alpha_n}{n} = 1$. Moreover, we know that $\beta_n = O(n^2)$. The following limits can be validated numerically or symbolically:⁷

$$\begin{aligned} \lim_{n \rightarrow \infty} \lambda_n^{(V)} &= \frac{1}{2} \\ \lim_{n \rightarrow \infty} \lambda_n^{(K)} &= 0 \\ \lim_{n \rightarrow \infty} \frac{\lambda_n^{(W)}}{n} &= \frac{1}{2}. \end{aligned}$$

Therefore, the matrix A_n^{num} behaves the following for big n :

$$A_n^{num} \sim \begin{pmatrix} 1 + \frac{1}{2} & -\frac{1}{2} & 0 \\ \frac{1}{2} & \frac{1}{2} & i\bar{\eta} \\ -\frac{1}{2} & -\frac{1}{2} & O(n) \end{pmatrix} \text{ for } n \rightarrow \infty.$$

The smallest singular value of this matrix is given by $\inf_{|x|=1} \|A_n^{num}x\|$. Since one of the entries in the third column has order n , The extremal value of x must have the form $(x_n^1, x_n^2, 0)$. However, the matrix in the first two columns converges, implying convergence of the minimal singular value.

7. Note that the n -dependence is not written out explicitly.

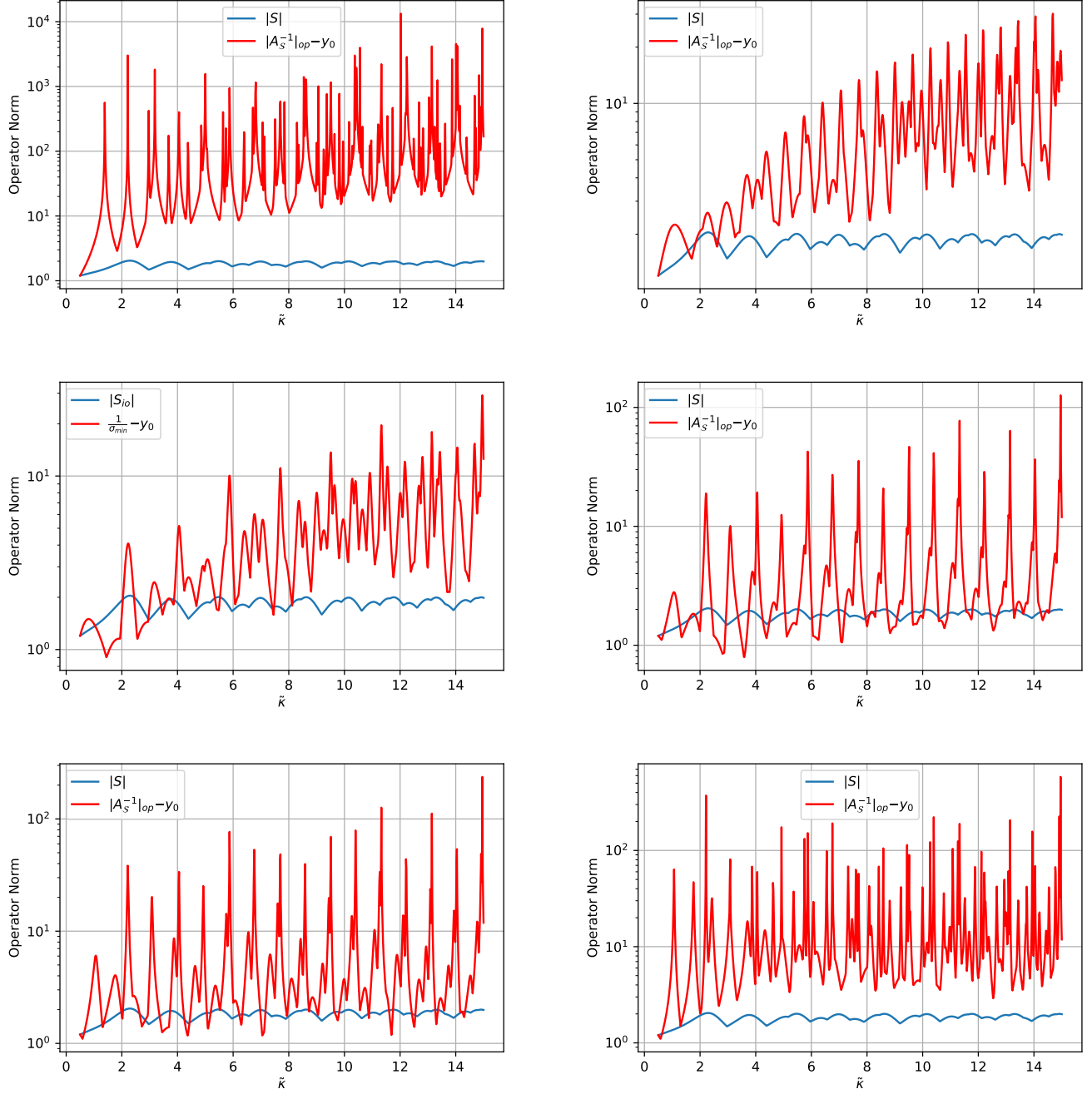


Fig. 6: Inverted minimum singular value $\|A_S^{-1}\|$ by wave number κ . In all six scenarios we have $c_i = 1.0, c_o = 3.0$ and we vary η . The values of η are as follows. Top left: $\eta = 0$, top right: $\eta = 0.5$, middle left: $\eta = 1.0$, middle right: $\eta = 5.0$, lower left: $\eta = 10.0$, and lower right: $\eta = 100.0$.

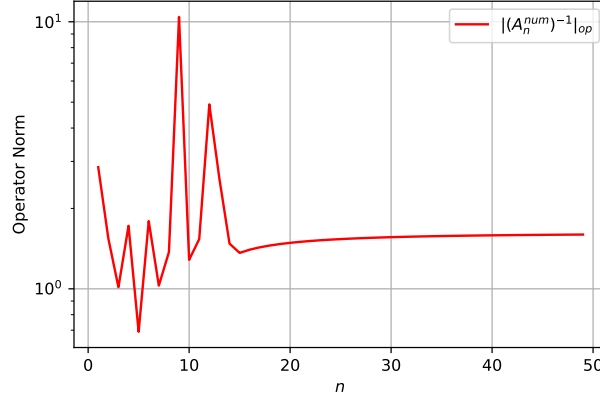


Fig. 7: Inverted singular value $\|A_n^{num}\|$ for fixed κ as a function of the Fourier mode n .

9.2 Constructing the Solution Operator

Plugging the Fourier ansatz into eq. 1 and imposing convergence at the origin and the Sommerfeld radiation condition results in

$$\begin{aligned} u^- &= \sum_{n=-\infty}^{\infty} u_n^- \frac{J_n(\sqrt{c_i}\tilde{\kappa}r)}{J_n(\sqrt{c_i}\tilde{\kappa})} e^{in\phi} \\ u^+ &= \sum_{n=-\infty}^{\infty} u_n^+ \frac{H_n(\sqrt{c_o}\tilde{\kappa}r)}{H_n(\sqrt{c_o}\tilde{\kappa})} e^{in\phi} \end{aligned}$$

where u_n^- and u_n^+ are the restrictions of u to Ω^- and Ω^+ .

Extend $f^i = \sum_{n=-\infty}^{\infty} f_n^i e^{in\phi}$ for $i = 1, 2$. The transmission condition $\gamma_C^+ u^+ = \gamma_C^- u^- + f$ implies

$$\begin{pmatrix} \frac{1}{\sqrt{c_i}\tilde{\kappa}} \frac{J'_n(\sqrt{c_i}\tilde{\kappa})}{J_n(\sqrt{c_i}\tilde{\kappa})} & -\frac{1}{\sqrt{c_o}\tilde{\kappa}} \frac{H'_n(\sqrt{c_o}\tilde{\kappa})}{H_n(\sqrt{c_o}\tilde{\kappa})} \end{pmatrix} \begin{pmatrix} u_n^- \\ u_n^+ \end{pmatrix} = \begin{pmatrix} f_n^1 \\ f_n^2 \end{pmatrix}. \quad (11)$$

By using the well-known inverse of a 2x2 matrix we obtain

$$u_n^- = \zeta(-\sqrt{c_o}\tilde{\kappa}J_n(\sqrt{c_i}\tilde{\kappa})H'_n(\sqrt{c_i}\tilde{\kappa})f_n^1 + J_n(\sqrt{c_i}\tilde{\kappa})H_n(\sqrt{c_o}\tilde{\kappa})f_n^2). \quad (12)$$

where $\zeta := \frac{1}{\sqrt{c_i}\tilde{\kappa}J'_n(\sqrt{c_i}\tilde{\kappa})H_n(\sqrt{c_o}\tilde{\kappa}) - \sqrt{c_o}\tilde{\kappa}H'_n(\sqrt{c_o}\tilde{\kappa})J_n(\sqrt{c_i}\tilde{\kappa})}$. This presents the solution operator, as we get γ_D^- from restriction to the boundary. γ_N^- can be directly obtained by restriction of the normal derivative of u_n^- to the boundary which is equivalent to multiplying the Fourier coefficients by $\sqrt{c_i}\tilde{\kappa} \frac{J'_n(\sqrt{c_i}\tilde{\kappa})}{J_n(\sqrt{c_i}\tilde{\kappa})}$. After rescaling the u_n , f_n^i to the complete orthonormal system defined in eq. 5, we obtain the solution operator matrix:

$$S_{io}^n = \zeta \begin{pmatrix} -\sqrt{c_o}\tilde{\kappa}J_n(\sqrt{c_i}\tilde{\kappa})H'_n(\sqrt{c_i}\tilde{\kappa}) & \sqrt{n^2 + \tilde{\kappa}^2}J_n(\sqrt{c_i}\tilde{\kappa})H_n(\sqrt{c_o}\tilde{\kappa}) \\ -\frac{1}{\sqrt{n^2 + \tilde{\kappa}^2}}\sqrt{c_o}\sqrt{c_i}\tilde{\kappa}^2J'_n(\sqrt{c_i}\tilde{\kappa})H'_n(\sqrt{c_i}\tilde{\kappa}) & \sqrt{c_i}\tilde{\kappa}J'_n(\sqrt{c_i}\tilde{\kappa})H_n(\sqrt{c_o}\tilde{\kappa}) \end{pmatrix} \quad (13)$$

that maps the Fourier coefficients of f in the complete orthonormal system for $\mathcal{H}^{\frac{1}{2}} \times \mathcal{H}^{-\frac{1}{2}}$ to the Fourier coefficients of $\gamma_C^- u$ in the same basis.

Symbols

This is an index of symbols repeatedly used in this report with references where the corresponding terms are defined (if applicable).

- $H^1(\Omega)$: Sobolev space, [9].
- $H_{\text{loc}}^1(\Omega^\pm)$: $H_{\text{loc}}^1(\Omega^\pm, \Delta) := \{v : \chi v \in H^1(\Omega^\pm), \Delta(\chi v) \in L^2(\Omega^\pm) \text{ for all } \chi \in C_{\text{comp}}^\infty(\mathbb{R}^d)\}$, section 1
- $H^{\frac{1}{2}}(\Gamma)$: Dirichlet trace space, [10].
- $H^{-\frac{1}{2}}(\Gamma)$: Neumann trace space, [10].
- Ω^- : Generally a bounded Lipschitz domain. $\Omega^0 = B_1(0) \subset \mathbb{R}^2$ for most of the report, section 1.
- Ω^+ : $\Omega^+ := \mathbb{R}^d \setminus \overline{\Omega^-}$, section 1.
- Γ : $\Gamma := \partial\Omega^- = \partial\Omega^+$.
- grad_Γ : surface gradient, [11].
- grad_Γ : adjoint map of $\hat{\phi} \text{grad}_\Gamma$, 4.
- φ^\pm : $\varphi^\pm := \varphi|_{\Omega^\pm}$, section 1.
- $C_{\text{comp}}^\infty(\mathbb{R}^d)$: Smooth functions with compact support
- γ_D^\pm : Dirichlet trace operators, section 1.
- γ_N^\pm : Neumann trace operator, section 1.
- γ_C^\pm : Cauchy trace, section 1.
- f : Boundary conditions of Helmholtz transmission problem, section 1.
- $\kappa, \tilde{\kappa}$: Wave number, section 1.
- c_i, c_o, \tilde{c} : Refractive indices, section 1.
- U, θ, p : Solution of Helmholtz transmission problem, section 1 & section 3.
- S : Solution operator, section 1.
- DtN_κ^- : Interior Dirichlet-to-Neumann map, section 2.
- $V_\kappa, K_\kappa, K'_\kappa, W_\kappa$: Boundary integral operators, section 2.
- $q_\kappa(\cdot, \cdot), b(\cdot, \cdot)$: Bilinear forms in regularized variational formulation, section 3.
- g_i : Right hand side of regularized variational formulation, section 3.
- $\mathcal{H}^s(\Gamma)_{\tilde{\kappa}}$: Fourier Sobolev Hilbert space, section 4.
- A : Left hand side operator in the regularized operator formulation, section 4.
- \mathcal{B} : Right hand side boundary value operator in the regularized operator formulation, section 4.
- \mathcal{S}_N^s : Restricted Sobolev space for numerical calculations, section 5.
- $\lambda^{(V)}, \lambda^{(K)}, \lambda^{(K')}, \lambda^{(W)}$: Fourier eigenvalues of boundary value operators, section 5.
- α_n, β_n : Fourier eigenvalues of operators occurring in regularised operator, section 5.
- A_n^{num} : Galerkin matrix, section 5.
- \mathcal{L}_ε : augmenting operator for A , section 7.

Code

Calculations and plots were implemented in Python. The corresponding files are available [here](#).

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