

# Simulating N particles in a box

January 2023

## 1 Introduction

The  $N$ -particle Hamiltonian in a box has the form [RKT03]:

$$H = \sum_{n=1,\dots,N} \sum_{i=1}^d \left[ \frac{(p_i^{(n)})^2}{2} + V_i(q_i^{(n)}) \right] + \sum_{\substack{n,m=1,\dots,N \\ n \neq m}} W(q^{(n)} - q^{(m)}), \quad (1)$$

where  $W$  is a repelling potential,  $C^\infty$  for  $\|q^{(n)} - q^{(m)}\| > \rho$ . We consider the limit of large energy per particle, and look for motions which are fast only in the last coordinate: we study the motion at the energy level  $H = Nh$  for a fixed  $N$  and large  $h$ , with most of the particles' energy is at the vertical motion. We scale the vertical momenta  $p_d^{(n)}$  by  $\sqrt{2h}$ , and the Hamiltonian transforms to

$$H = \sum_{n=1,\dots,N} \left[ \frac{(p_d^{(n)})^2}{2} + \delta V_d(q_d^{(n)}) \right] + \delta \sum_{n=1,\dots,N} \sum_{i=1}^{d-1} \left[ \frac{(p_i^{(n)})^2}{2} + V_i(q_i^{(n)}) \right] + \delta \sum_{\substack{n,m=1,\dots,N \\ n \neq m}} W(q^{(n)} - q^{(m)}), \quad (2)$$

where  $\delta = \frac{1}{2h}$ ; we study the behavior on the fixed energy level  $H = \frac{N}{2}$ , with

$$V_i(q_i) = \frac{1}{Q_i(q_i)^\alpha}, \quad \alpha > 0 \quad (3)$$

where the  $C^\infty$  function  $Q_i$  measures the distance to the box boundary in the  $i$ -th coordinate direction, i.e.,  $Q_i(0) = 0$ ,  $Q_i(l_i) = 0$ ,  $Q_i'(0) > 0$ ,  $Q_i'(l_i) < 0$ , and  $Q_i(q_i) > 0$  for  $q_i \in (0, l_i)$ .

## 2 Specific example to be studied

Take  $d = 2$  (so  $q_2$  is the fast motion and  $q_1$  is fixed for the periodic solutions) and small  $N$  (e.g. start with  $N = 2$ , then 3, and then maybe more.), and

$$Q_i(q_i) = q_i(l_i - q_i)(1 + P(q_i)), \quad P_i(0) = 0, P_i(q_i) \geq 0, q_i \in (0, l_i). \quad (4)$$

so

$$Q'_i(q_i) = (l_i - q_i)(1 + P_i(q_i)) - q_i(1 + P_i(q_i)) + P'_i(q_i)q_i(l_i - q_i) = \begin{cases} l_i & q_i = 0 \\ -l_i(1 + P_i(l_i)) & q_i = l_i \end{cases} \quad (5)$$

and

$$W(q^{(n)} - q^{(m)}) = \frac{a}{\|q^{(n)} - q^{(m)}\|^\beta}, \quad a, \beta > 0. \quad (6)$$

### 3 Assignment

1. Find the fixed points of (2) for  $N = 2, 3$  (start with  $P_i(q_i) = 0$  and  $a = 0$ , but also check what can you do for  $a > 0$ , and see if by preparing  $P_1(q_1)$  to get more fixed points in the  $q_1$  direction helps.
2. Write a code to find numerically the solutions of (2).
3. Look for periodic solutions in the  $q_2$  direction.
4. Write a code for calculating the averaged potential and find the "minimizing phases" for it. See if we can observe islands of stability.

### 4 Answers

#### 4.1 Finding the Fixed Points

Since there are  $N$  particles each living in  $d$ -dimensions, the system is of  $Nd$  degrees of freedom (each particle is described by  $d$  scalars). The fixed points of Hamiltonian systems are the points where the gradient of  $H$  vanishes. We find the partial derivatives of  $H$ . For  $j = 1, \dots, d$ ,  $k = 1, \dots, N$ ,

$$\frac{\partial H}{\partial p_j^{(k)}} = p_j^{(k)}$$

Recall that:

$$V_j(q_j^{(k)}) = \frac{1}{Q_j(q_j^{(k)})^\alpha}$$

so

$$V'_j(q_j^{(k)}) = \frac{-\alpha \cdot Q'_j(q_j^{(k)})}{Q_j(q_j^{(k)})^{\alpha+1}}$$

$$W(q^{(n)} - q^{(m)}) = \frac{a}{\|q^{(n)} - q^{(m)}\|^\beta} = a \left( \sum_{i=1}^d (q_i^{(n)} - q_i^{(m)})^2 \right)^{-\beta/2}$$

$$\begin{aligned}
& \frac{\partial}{\partial q_j^{(k)}} \sum_{\substack{n,m=1,\dots,N \\ n \neq m}} W(q^{(n)} - q^{(m)}) = \\
& = 2 \frac{\partial}{\partial q_j^{(k)}} \sum_{\substack{n=1 \\ n \neq k}}^N W(q^{(k)} - q^{(n)}) = \\
& = 2 \frac{\partial}{\partial q_j^{(k)}} \sum_{\substack{n=1 \\ n \neq k}}^N a \left( \sum_{i=1}^d (q_i^{(k)} - q_i^{(n)})^2 \right)^{-\beta/2} = \\
& = 2a \sum_{\substack{n=1 \\ n \neq k}}^N \left( \sum_{i=1}^d (q_i^{(k)} - q_i^{(n)})^2 \right)^{-(\beta+2)/2} \cdot \frac{-\beta}{2} \cdot 2 (q_j^{(k)} - q_j^{(n)}) \\
& = -2a\beta \sum_{\substack{n=1 \\ n \neq k}}^N \frac{q_j^{(k)} - q_j^{(n)}}{\|q^{(k)} - q^{(n)}\|^{\beta+2}}
\end{aligned}$$

(The factor of 2 comes here, because we sum up the potential inflicted on the  $k$ -th particle and inflicted by the  $k$ -th particle, which are equal).

To summarize, the fixed points are where the momenta vanish and the particles' positions  $q_j^{(k)}$  satisfy:

$$\frac{\alpha \cdot Q_j'(q_j^{(k)})}{Q_j(q_j^{(k)})^{\alpha+1}} + 2a\beta \sum_{\substack{n=1 \\ n \neq k}}^N \frac{q_j^{(k)} - q_j^{(n)}}{\|q^{(k)} - q^{(n)}\|^{\beta+2}} = 0$$

where

$$Q_j(q) = q(l_j - q)(1 + P(q))$$

and

$$Q_j'(q) = (l_j - q)(1 + P_j(q)) - q(1 + P_j(q)) + q(l_j - q)P_j'(q)$$

This system of equations is generally hard to solve. However, when taking  $a = 0$ , (i.e no repelling potential), the system is just a set of  $Nd$  separate equations, and the solution boils down to finding where the derivatives of the polynomials  $Q_j$  vanish.

**Example 1** For  $a = 0$ ,  $P_j(q) = 0$ ,

$$Q_j(q) = q(l_j - q) \implies Q_j'(q) = l_j - 2q \implies q = \frac{l_j}{2}$$

**Example 2** For  $a = 0$ ,  $P_j(q) = q$ ,

$$\begin{aligned}
Q_j(q) &= q(l_j - q)(1 + q) \implies Q_j'(q) = -3q^2 + 2(l_j - 1)q + l_j \\
&\implies q = -\frac{2(1 - l_j) \pm \sqrt{4(l_j - 1)^2 + 12l_j}}{6}
\end{aligned}$$

**Example 3** For  $a = 0$ ,  $P_j(q) = 24q^3 - 30q^2 + 10$ ,  $l_j = 1$ ,  $Q_j(q)$  has 3 critical points as shown in the Figure 1:

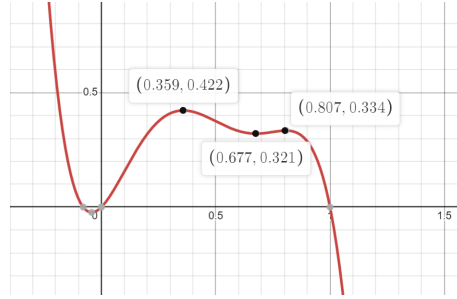


Figure 1:  $P_j(q) = 24q^3 - 30q^2 + 10$