## Numerical Modeling of Kirchhoff-Love Plates with Rotational Inertia

#### Francesco Bonaldi

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Université de Montpellier

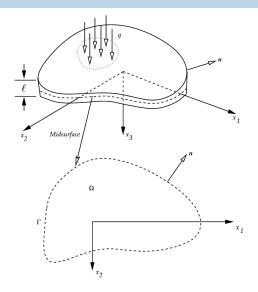
Journées FreeFEM++

Paris, December 16th, 2015

### Outline

- 1 Introduction
  - Problem Statement
  - Mathematical Setting
- 2 Numerical Treatment
  - General Method
  - Dynamics
  - Statics
  - Remarks
  - Conclusions

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■ This problem derives from a simplified model of a magneto-electro-thermo-elastic 3D plate-like body (sensor/actuator) [F. B., G. Geymonat, F. Krasucki, M. Serpilli, Math. Mech. Solids, 2015].

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- *D* is the *flexural rigidity* of the plate; it takes into account all these couplings and does not depend only on the mechanical behavior.

■ Let  $q \in L^2(0, T; L^2(\Omega))$  and

$$V=H_0^2(\Omega)=\{u\in H^2(\Omega): u=\partial_n u=0 \text{ on } \partial\Omega\}, \quad H=\frac{H_0^1(\Omega)}{(0)}=\{u\in H^1(\Omega): u=0 \text{ on } \partial\Omega\}.$$

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$$a(u,v) := \int_{\Omega} \left( D \mathbf{v} \, \Delta u \, \Delta v + D(1-\mathbf{v}) \, \nabla \nabla u : \nabla \nabla v \right) \mathrm{d}\Omega.$$

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■ Linear form on *H* to be identified with (f(t), v) for some  $f \in L^2(0, T; H)$ :

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■ The following problem is well-posed (J. L. Lions):

Find 
$$u \in C^0([0, T]; V) \cap C^1([0, T]; H)$$
 such that 
$$\forall v \in V, \quad \frac{d^2}{dt^2}(u(t), v) + a(u(t), v) = L_t(v),$$
$$u(0) = u_0 \in V, \quad \frac{du}{dt}(0) = u_1 \in H.$$

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- On setting  $\xi(t) := (\xi_1(t), \dots, \xi_N(t))$  and  $\mathcal{F}(t) := (L_t(\varphi_1), \dots, L_t(\varphi_N))$ , the semi-discrete version of the problem reads

$$\mathcal{M}\frac{\mathrm{d}^2 \boldsymbol{\xi}}{\mathrm{d}t^2}(t) + \mathcal{K}\boldsymbol{\xi}(t) = \mathcal{F}(t),$$

$$\mathcal{M}_{ij} = (\varphi_i, \varphi_j)_{1 \le i, j \le N}, \quad \mathcal{K}_{ij} = a(\varphi_i, \varphi_j)_{1 \le i, j \le N}.$$

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- Space discretization: **HCT finite elements** ( $C^1$ -class finite elements).

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- Space discretization: **HCT finite elements** ( $C^1$ -class finite elements).
- Time discretization: Newmark's (constant) average acceleration method.

# Dynamics I

 $\quad \blacksquare \quad q \equiv 0.$ 

#### Clamped circular plate

■ Initial conditions:

$$u_0(x, y) = 0,$$
  
 $u_1(x, y) = \alpha(1 - (x^2 + y^2)).$ 

■ Boundary conditions:

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#### Simply supported circular plate

Initial conditions:

$$u_0(x, y) = 0,$$
  
 $u_1(x, y) = \alpha(1 - (x^2 + y^2)).$ 

■ Boundary conditions:

$$u = 0$$
 and  $\mathbf{Mn} \cdot \mathbf{n} = 0$  on  $\partial \Omega$ ,

$$\mathbf{M} := -D((1 - \nu)\nabla\nabla u + \nu\Delta u \mathbf{I}).$$

# Dynamics II

What occurs upon refining the mesh?



## **Dynamics II**

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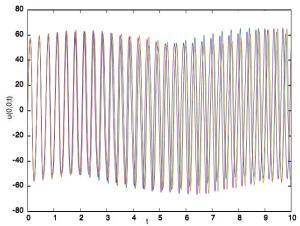
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Number N of boundary mesh nodes: N = 60, N = 90, N = 120.



## **Dynamics III**

■ What are these inconsistencies due to? **HCT or Newmark?** 

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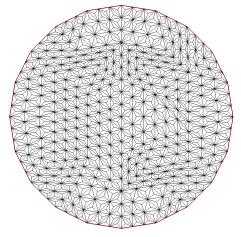
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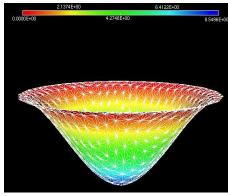
- What are these inconsistencies due to? HCT or Newmark?
- Let us consider a test case in the static regime.

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- Data set:  $D = 1831.5 \, kN \cdot m$ ,  $\ell = 0.1 \, m$ ,  $R = 10 \, m$ ,  $q = 1 \, kN/m^2$ .
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Table: Percentage error between  $u_h(0,0)$  and u(0,0).

Boundary mesh nodes	N = 75	N = 90	N = 105	N = 120	N = 135
$u_h(0,0)$	8.691 cm	8.850 cm	8.811 cm	8.849 cm	8.876 cm
Percentage error	1.89065%	3.74634%	3.28311%	3.73569%	4.0511%

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- There is something wrong with the implementation of HCT in FreeFEM++ 3.41.

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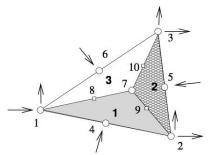
### **HCT** Element

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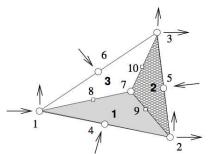
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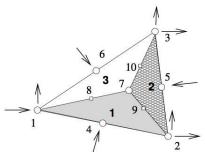
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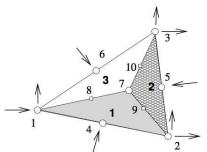
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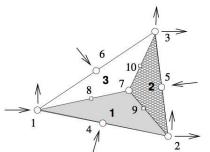
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- **Solution**: instead of integrating on the element as a whole, integrate on each sub-triangle using such a quadrature formula and then sum up.
- Corrections have been carried out in FreeFEM++ 3.42, available online since December 10th.

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Boundary mesh nodes	N = 75	N = 90	N = 105	N = 120	N = 135
$u_h(0,0)$	8.315 cm	8.385 cm	8.427 cm	8.452 cm	8.469 cm
Percentage error	2.5295%	1.70338%	1.2146%	0.924838%	0.719373%

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■ Percentage error **decreases** as mesh is refined.

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- Using FreeFEM++, we have been able to solve an evolution problem with a non-standard inertia term.
- We discovered technical issues with integration in the implementation of HCT finite elements in FreeFEM++ 3.41 which have been fixed in the latest release, available at <a href="http://www.freefem.org/ff++/">http://www.freefem.org/ff++/</a>.

### Conclusions

- Using FreeFEM++, we have been able to solve an evolution problem with a non-standard inertia term.
- We discovered technical issues with integration in the implementation of HCT finite elements in FreeFEM++ 3.41 which have been fixed in the latest release, available at http://www.freefem.org/ff++/.

#### Future work

As HCT finite elements are computationally expensive, we would like to use a different space discretization (e.g., HHO methods [Di Pietro, Ern & Lemaire, Comput. Meth. Appl. Math., 2014]).

# Acknowledgements

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### THANKS FOR YOUR ATTENTION

### Motivation

 The problem results from a dimension-reduction procedure carried out on a problem posed in the 3D domain

$$\mathcal{P}^{\varepsilon} := \Omega \times \left( -\frac{\ell^{\varepsilon}}{2}, \frac{\ell^{\varepsilon}}{2} \right), \quad \ell^{\varepsilon} = \varepsilon \ell, \quad \ell := \operatorname{diam}(\mathcal{P}^{\varepsilon}) \simeq \operatorname{diam}(\Omega).$$

■ Thus,  $\ell^{\varepsilon}$  is the real thickness of the plate-like domain under study, whereas  $\ell$  is a scaled thickness, i.e. the thickness of the domain

$$\mathcal{P} := \Omega \times \left(-\frac{\ell}{2}, \frac{\ell}{2}\right).$$

■ In other words,  $\ell$  is not *small*.



Rotational inertia is not negligible