Non linear mechanical simulation of thin films on soft substrate in the FreeFem++ environment

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Outlines

- Introduction
- **■** Asymptotic Numerical Method
- Application of the ANM to the non linear mechanical study of stiff thin films on soft substrates
- 3D implementation in the FreeFem++ environment
- Numerical results
- **■** Conclusions

Introduction

- Wrinkles of stiff thin layers on soft materials
 - Wrinkles in nature:



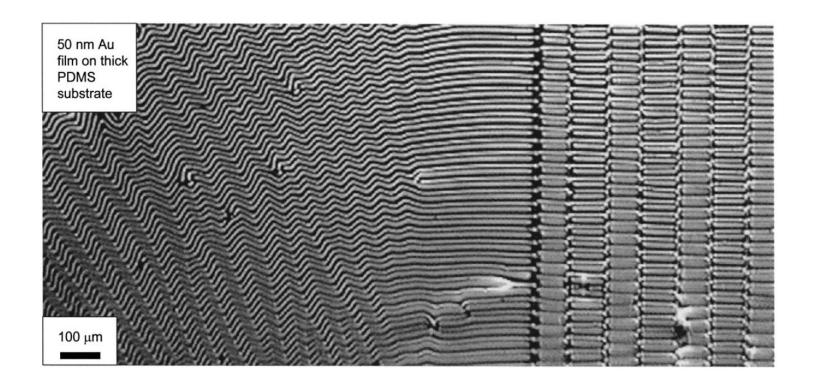




leaf Human finger Land form

Introduction

Buckling of a 50 nm gold film on a thick PDMS substrate



Introduction

- Non linear mechanical problems are generally solved using predictor-corrector algorithms such as Newton Raphson scheme (Riks, 1972, 1984), (Ramm, 1981), ...
- A family of Asymptotic-Numerical-Method based on Taylor series and Finite Element Method has been developed for non linear elastic structures
 - Damil and Potier-Ferry (1990), compute perturbed bifurcations
 - Azrar, Cochelin, Damil, Potier-Ferry (1993), compute the post-buckling behavior of elastic plate
 - Cochelin, Damil, Potier-Ferry (1995), show the extension to any non linear elastic solutions

Consider the non linear problem :

Linear operator Control parameter
$$\mathbf{R}(\mathbf{U},\lambda) = \mathbf{L}(\mathbf{U}) + \mathbf{Q}(\mathbf{U},\mathbf{U}) - \lambda \mathbf{F} = \mathbf{0}$$
 (1) Unknown vector Quadratic operator Vector

■ Principles of the ANM: from a known solution $(\mathbf{U}^j, \lambda^j)$, the new solution is expanded into truncated power serie of a perturbation parameter a:

(2)
$$\begin{cases} \mathbf{U}(a) = \mathbf{U}^{j} + \sum_{p=1}^{N} a^{p} \mathbf{U}_{p} \\ \lambda(a) = \lambda^{j} + \sum_{p=1}^{N} a^{p} \lambda_{p} \end{cases}$$

A good choice for the parameter a is the linearized arc-length parameter defined by the projection of the pair $(\mathbf{U} - \mathbf{U}^j, \lambda - \lambda^j)$ on the tangent direction $(\mathbf{U}_1, \lambda_1)$ $a = \langle \mathbf{U} - \mathbf{U}^j, \mathbf{U}_1 \rangle + (\lambda - \lambda^j) \lambda_1$ (3)

$$a = \left\langle \mathbf{U} - \mathbf{U}^{j}, \mathbf{U}_{1} \right\rangle + \left(\lambda - \lambda^{j}\right) \lambda_{1}$$
 (3)

By estimating the radius of convergence of the series, the validity range of a for each branch is:

$$a_{\text{max}} = \left(\delta \frac{\|U_1\|}{\|U_N\|}\right)^{\frac{1}{N-1}} \tag{4}$$

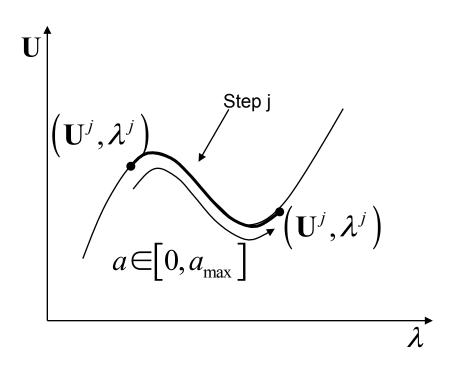
- Truncation order: $15 \le N \le 50$ (5)
- tolerance δ (affects the residual): $10^{-6} \le \delta \le 10^{-3}$ (6)

■ By substituting (2) into (1) and (3), and equating terms with the same power of a, we obtain the following sequence of linear problems:

• order 1:
$$L_t(\mathbf{U}_1) = \lambda_1 \mathbf{F}$$
, and, $\langle \mathbf{U}_1, \mathbf{U}_1 \rangle + \lambda_1 \lambda_1 = 1$ (7)

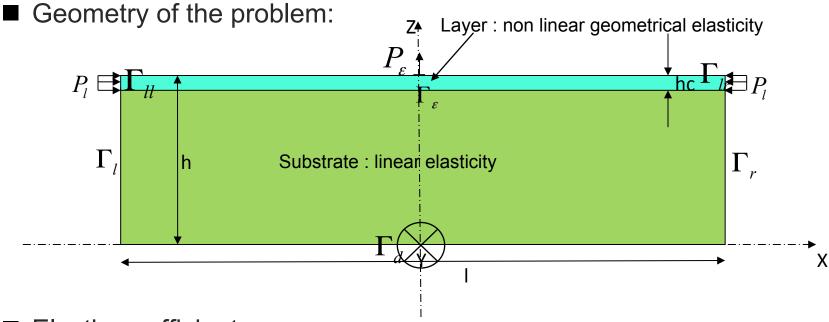
• order 2:
$$L_t(\mathbf{U}_2) = \lambda_2 \mathbf{F} - \mathbf{Q}(\mathbf{U}_1, \mathbf{U}_1)$$
, and, $\langle \mathbf{U}_2, \mathbf{U}_1 \rangle + \lambda_2 \lambda_1 = 0$ (8)

• order p:
$$L_t(\mathbf{U}_p) = \lambda_p \mathbf{F} - \sum_{r=1}^{p-1} \mathbf{Q}(\mathbf{U}_r, \mathbf{U}_{p-r})$$
, and, $\langle \mathbf{U}_p, \mathbf{U}_1 \rangle + \lambda_p \lambda_1 = 0$ (9) With, the tangent operator $\mathbf{L}_t(.)$ defined by $\mathbf{L}_t(.) = \mathbf{L}(.) + 2\mathbf{Q}(., \mathbf{U}^j)$



In case of bifurcation points:

► accumulation of small steps indicates the presence of a bifurcation



- Elastic coefficients:
 - Substrate: $E_s = 1.8 \text{ MPa}$, $v_s = 0.48$
 - Layer: $E_l = 1.3 \cdot 10^5 \text{ MPa}$ $v_l = 0.3$
- Boundary conditions: on Γ_l , Γ_{ll} , Γ_r , Γ_{rl} , $u_Y = u_Z = 0$, on Γ_d $u_Z = 0$ only model half structure, symmetry with respect to the (x=0) plane

- \blacksquare Assumptions for the mechanical model ($E_{L}\gg E_{S}$)
 - Non linear geometrical elasticity in the layer
 - Linear elasticity in the substrate
- Instabilities appear for small deformations
- Fan Xu's PhD thesis Dec 2014: "numerical study of instability patterns of film substrate patterns", based on a plate finite element formulation for the film.

■ For elastic bodies, the governing equations can be stated from the mixed Hellinger-Reissner functional.

Green-Lagrange strain tensor

$$HR(\mathbf{u}, \mathbf{S}, \lambda) = \int_{\Omega} \left(\mathbf{S} : \gamma - \frac{1}{2} \mathbf{S} : \mathbf{D}^{-1} : \mathbf{S} \right) d\Omega - \lambda P_e(\mathbf{u})$$
 Elasticity matrix (10)

Piola-Kirchhoff stress tensor of the second kind

with,
$$P_{e}(\mathbf{u}) = -\int_{\Gamma_{lr}} P_{l} u_{x} + \int_{\Gamma_{ll}} P_{l} u_{x} + \int_{\Gamma_{\varepsilon}} P_{\varepsilon} u_{z}$$
 Small perturbed force

■ Let us compute the variational formulation:

$$\int_{\Omega} \left(\mathbf{S} : \delta \gamma + \delta \mathbf{S} : \gamma - \mathbf{S} : \mathbf{D}^{-1} : \delta \mathbf{S} \right) d\Omega - \lambda P_e \left(\delta \mathbf{u} \right) = 0$$
 (12)

■ The Green Lagrange tensor is:

$$\gamma(\mathbf{u}) = \underbrace{\frac{1}{2} \left(\nabla \mathbf{u} + \nabla^T \mathbf{u} \right)}_{\gamma_l(\mathbf{u})} + \underbrace{\frac{1}{2} \left(\nabla^T \mathbf{u} \cdot \nabla \mathbf{u} \right)}_{\gamma_{nl}(\mathbf{u}, \mathbf{u})} = \gamma_l(\mathbf{u}) + \gamma_{nl}(\mathbf{u}, \mathbf{u})$$
(13)

which gives:

$$\delta \gamma = \gamma_{l} \left(\delta \mathbf{u} \right) + 2 \gamma_{nl} \left(\mathbf{u}, \delta \mathbf{u} \right) = \gamma_{l} \left(\delta \mathbf{u} \right) + \delta \gamma_{nl} \left(\mathbf{u}, \delta \mathbf{u} \right)$$

$$\underbrace{\delta \gamma_{nl} \left(\mathbf{u}, \delta \mathbf{u} \right)}_{\delta \gamma_{nl} \left(\mathbf{u}, \delta \mathbf{u} \right)} = \gamma_{l} \left(\delta \mathbf{u} \right) + \delta \gamma_{nl} \left(\mathbf{u}, \delta \mathbf{u} \right)$$
(14)

The serial developments of the ANM leads to:

$$\begin{cases} \mathbf{u}(a) = \mathbf{u}^{j} + \sum_{p=1}^{N} a^{p} \mathbf{u}_{p} \\ \lambda(a) = \lambda^{j} + \sum_{p=1}^{N} a^{p} \lambda_{p} \end{cases}$$
(15)

■ By introducing (15) into (12), by identifying the order p terms, we obtain the following variational formulation:

$$\int_{\Omega} \delta \gamma^{T} \left(\mathbf{u}^{j}, \delta \mathbf{u} \right) \mathbf{D} \left(\gamma_{l} \left(\mathbf{u}_{p} \right) + 2 \gamma_{nl} \left(\mathbf{u}^{j}, \mathbf{u}_{p} \right) \right) d\Omega + \int_{\Omega} \delta \gamma_{nl}^{T} \left(\mathbf{u}_{p}, \delta \mathbf{u} \right) \mathbf{D} \left(\gamma_{l} \left(\mathbf{u}^{j} \right) + 2 \gamma_{nl} \left(\mathbf{u}^{j}, \mathbf{u}^{j} \right) \right) d\Omega
- \lambda_{p} P_{e} \left(\delta \mathbf{u} \right) + \sum_{r=1}^{p-1} \int_{\Omega} \delta \gamma_{nl}^{T} \left(\mathbf{u}_{p-r}, \delta \mathbf{u} \right) \mathbf{D} \left(\gamma_{l} \left(\mathbf{u}_{r} \right) + 2 \gamma_{nl} \left(\mathbf{u}^{j}, \mathbf{u}_{r} \right) + \sum_{s=1}^{r-1} \gamma_{nl} \left(\mathbf{u}_{r-s}, \mathbf{u}_{s} \right) \right)
+ \int_{\Omega} \delta \gamma^{T} \left(\mathbf{u}^{j}, \delta \mathbf{u} \right) \mathbf{D} \left(\sum_{r=1}^{p-1} \gamma_{nl} \left(\mathbf{u}_{p-r}, \mathbf{u}_{r} \right) \right) d\Omega = 0$$
(16)

■ For the first order, the resulting variational formulation is:

$$\int_{\Omega} \delta \gamma^{T} \left(\mathbf{u}^{j}, \delta \mathbf{u} \right) \mathbf{D} \left[\gamma_{l} \left(\mathbf{u}_{1} \right) + 2 \gamma_{nl} \left(\mathbf{u}^{j}, \mathbf{u}_{1} \right) \right] d\Omega
+ \int_{\Omega} \delta \gamma_{nl}^{T} \left(\mathbf{u}_{1}, \delta \mathbf{u} \right) \mathbf{D} \left[\gamma_{l} \left(\mathbf{u}^{j} \right) + 2 \gamma_{nl} \left(\mathbf{u}^{j}, \mathbf{u}^{j} \right) \right] d\Omega - \lambda_{1} P_{e} \left(\delta \mathbf{u} \right) = 0$$
(17)

■ P2 finite element model is used which leads to the linear systems: (for each branch).

• Order 1:
$$\begin{bmatrix} \mathbf{K}_{T} \end{bmatrix} \hat{\mathbf{U}}_{1} = \mathbf{F}$$

$$\begin{cases} \lambda_{1} = \frac{1}{\sqrt{1 + \hat{\mathbf{U}}_{1}^{T} \hat{\mathbf{U}}_{1}}} \\ \mathbf{U}_{1} = \lambda_{1} \hat{\mathbf{U}}_{1} \end{bmatrix}$$
• Order 2:
$$\begin{bmatrix} \mathbf{K}_{T} \end{bmatrix} \mathbf{U}_{2} = \lambda_{2} \mathbf{F} + \mathbf{F}_{2}^{nl} \quad \left(\mathbf{F}_{2}^{nl} \text{ computed using } \mathbf{U}_{1} \right)$$

$$\begin{cases} \lambda_{2} = -\lambda_{1} \left(\mathbf{U}_{2}^{nl} \right)^{T} \mathbf{U}_{1} \\ \mathbf{U}_{2} = \frac{\lambda_{2}}{\lambda_{1}} \mathbf{U}_{1} + \mathbf{U}_{2}^{nl} \end{cases}$$
with,
$$\mathbf{U}_{2}^{nl} = \begin{bmatrix} \mathbf{K}_{T} \end{bmatrix}^{-1} \mathbf{F}_{2}^{nl}$$

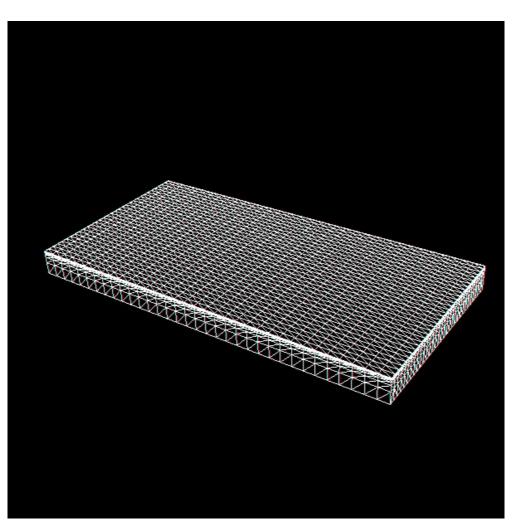
• Order p:
$$\begin{bmatrix} \mathbf{K}_T \end{bmatrix} \mathbf{U}_p = \lambda_p \mathbf{F} + \mathbf{F}_p^{nl} \left(\mathbf{F}_p^{nl} \text{ computed using } \mathbf{U}_1, \mathbf{U}_2, \mathbf{K}, \mathbf{U}_{p-1} \right)$$

$$\begin{cases} \lambda_p = -\lambda_1 \left(\mathbf{U}_p^{nl} \right)^T \mathbf{U}_1 \\ \mathbf{U}_p = \frac{\lambda_p}{\lambda_1} \mathbf{U}_1 + \mathbf{U}_p^{nl} \end{cases}$$
 with,
$$\mathbf{U}_p^{nl} = \begin{bmatrix} \mathbf{K}_T \end{bmatrix}^{-1} \mathbf{F}_p^{nl}$$

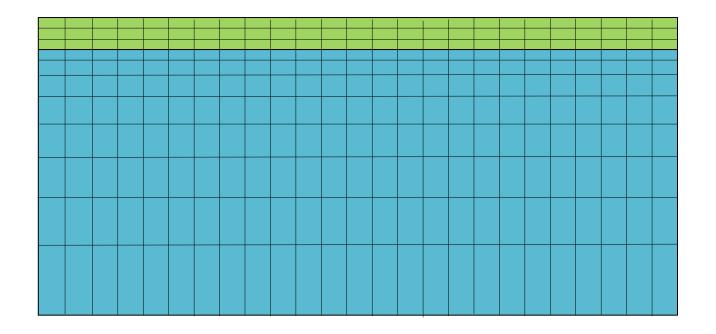
The direction of the path in the branch is determined knowing the tangent of the displacement for the previous branch.

- Main advantages of using FreeFem++:
 - use volume Finite Elements
 - powerful open source mathematical libraries are available
 - allows to compute problems with huge number of degrees of freedom by using parallel computing on HPC
 - allows to check stability of solutions

- 3D Finite Element mesh
- Tetrahedral Finite Elements
- P2 quadratic Lagrangian interpolation
- Number of nodes: 187 005
- CPU time (almost 3 days)
- improvement with parallel programming



- 3D mesh obtained from a 2D mesh
- 2D mesh: uniform mesh in the layer, geometrical rate in the substrate



- Let us present the main features of the FreeFem++ code:
 - macro are used to define the Green Lagrange tensor and its differential:

Definition of the finite element space

```
fespace Vh(Th3D,[P2,P2,P2]);
```

ullet Creation of the tangent matrix $ig[\mathbf{K}_Tig]$ and the second member $ig[\mathbf{F}]$

ullet Computation of ${\bf F}_2^{nl}$

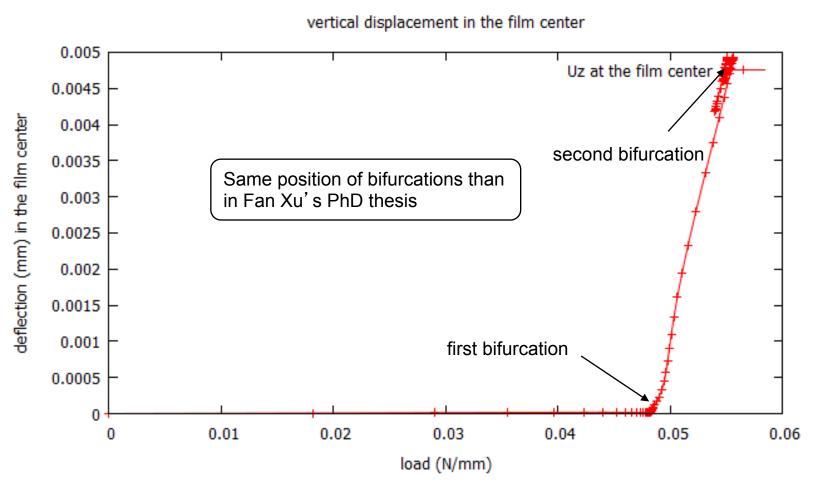
```
 \begin{aligned} & \text{varf PbFnl2} \ ([\text{utmp,vtmp,wtmp}],[\text{uuu,vvv,www}]) = -\inf 3d(\text{Th3D,reg3DL}) \\ & ( \ (d\text{GammaNL}(\text{u[1],v[1],w[1],uuu,vvv,www}))'*(\text{DL*}(\text{GammaL}(\text{u[1],v[1],w[1]})+2*\text{GammaNL}(\text{u[0],v[0],w[0],u[1],v[1],w[1]}))) \\ & + (d\text{Gamma}(\text{u[0],v[0],w[0],uuu,vvv,www}))'*(\text{DL*}(\text{GammaNL}(\text{u[1],v[1],w[1],v[1],w[1],v[1],w[1]}))) \\ & + \text{on}(\text{lsymmid,utmp=0.}) + \text{on}(\text{lrightmid,lrightlmid,vtmp=0.}) + \text{on}(\text{ldownmid,wtmp=0}); \\ & \text{Fnlu[] = PbFnl2}(0,Vh);} \end{aligned}
```

ullet Computation of \mathbf{F}_{n}^{nl} for p>2 : double loop sum is needed

3D implementation in the FreeFem++ environment environment

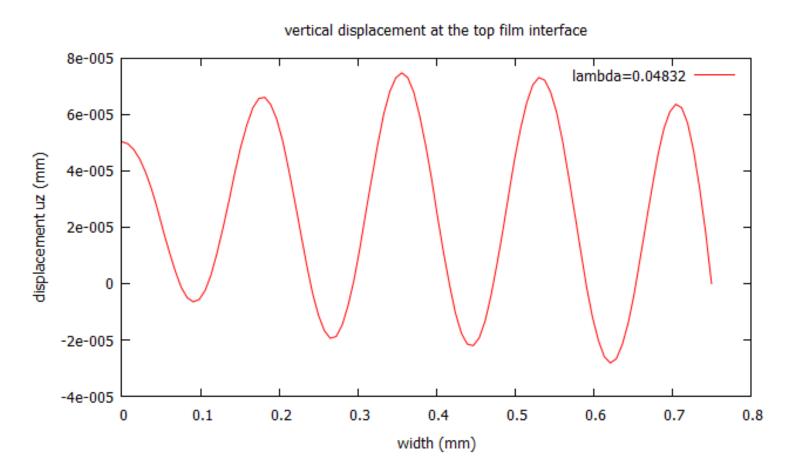
```
for (int ir=1;ir<ipn;ir++)
                  varf\ PbFnla\ ([utmp,vtmp,wtmp],[uuu,vvv,www]) = -\inf 3d(Th3D,reg3DL)\ (\ (dGammaNL(u[ipn-ir],v[ipn-ir],w[ipn-ir],uuu,vvv,www))^*(DL^*(GammaL(u[ir],v[ir],w[ir])^*(DL^*(GammaL(u[ir],v[ir],w[ir])^*(DL^*(GammaL(u[ir],v[ir],w[ir])^*(DL^*(GammaL(u[ir],v[ir],w[ir])^*(DL^*(GammaL(u[ir],v[ir],w[ir])^*(DL^*(GammaL(u[ir],v[ir],w[ir])^*(DL^*(GammaL(u[ir],v[ir],w[ir])^*(DL^*(GammaL(u[ir],v[ir],w[ir])^*(DL^*(GammaL(u[ir],v[ir],w[ir])^*(DL^*(GammaL(u[ir],v[ir],w[ir])^*(DL^*(GammaL(u[ir],v[ir],w[ir])^*(DL^*(GammaL(u[ir],v[ir],w[ir])^*(DL^*(GammaL(u[ir],v[ir],w[ir])^*(DL^*(GammaL(u[ir],v[ir],w[ir])^*(DL^*(GammaL(u[ir],v[ir],w[ir])^*(DL^*(GammaL(u[ir],v[ir],w[ir])^*(DL^*(GammaL(u[ir],v[ir],w[ir])^*(DL^*(GammaL(u[ir],v[ir],w[ir])^*(DL^*(GammaL(u[ir],v[ir],w[ir])^*(DL^*(GammaL(u[ir],v[ir],w[ir])^*(DL^*(GammaL(u[ir],v[ir],w[ir])^*(DL^*(GammaL(u[ir],v[ir],w[ir])^*(DL^*(GammaL(u[ir],v[ir],w[ir])^*(DL^*(GammaL(u[ir],v[ir],w[ir])^*(DL^*(GammaL(u[ir],v[ir],w[ir])^*(DL^*(GammaL(u[ir],v[ir],w[ir])^*(DL^*(GammaL(u[ir],v[ir],w[ir])^*(DL^*(GammaL(u[ir],w[ir],w[ir])^*(DL^*(GammaL(u[ir],w[ir],w[ir])^*(DL^*(GammaL(u[ir],w[ir],w[ir])^*(DL^*(GammaL(u[ir],w[ir],w[ir])^*(DL^*(GammaL(u[ir],w[ir],w[ir])^*(DL^*(GammaL(u[ir],w[ir],w[ir])^*(DL^*(GammaL(u[ir],w[ir],w[ir])^*(DL^*(GammaL(u[ir],w[ir],w[ir])^*(DL^*(u[ir],w[ir])^*(DL^*(u[ir],w[ir])^*(DL^*(u[ir],w[ir])^*(DL^*(u[ir],w[ir])^*(DL^*(u[ir],w[ir])^*(DL^*(u[ir],w[ir])^*(DL^*(u[ir],w[ir])^*(DL^*(u[ir],w[ir])^*(DL^*(u[ir],w[ir])^*(DL^*(u[ir],w[ir])^*(DL^*(u[ir],w[ir])^*(DL^*(u[ir],w[ir])^*(DL^*(u[ir],w[ir])^*(DL^*(u[ir],w[ir])^*(DL^*(u[ir],w[ir])^*(DL^*(u[ir],w[ir])^*(DL^*(u[ir],w[ir])^*(DL^*(u[ir],w[ir])^*(DL^*(u[ir],w[ir])^*(DL^*(u[ir],w[ir])^*(DL^*(u[ir],w[ir])^*(DL^*(u[ir],w[ir])^*(DL^*(u[ir],w[ir])^*(DL^*(u[ir],w[ir])^*(DL^*(u[ir],w[ir])^*(DL^*(u[ir],w[ir])^*(DL^*(u[ir],w[ir])^*(DL^*(u[ir],w[ir])^*(DL^*(u[ir],w[ir])^*(DL^*(u[ir],w[ir])^*(DL^*(u[ir],w[ir])^*(DL^*(u[ir],w[ir])^*(DL^*(u[ir],w[ir])^*(DL^*(u[ir],w[ir])^*(DL^*(u[ir],w[ir])^*(DL^*(u[ir],w[ir])^*(DL^*(u[ir],w[ir])^*(DL^*(u[ir],w[ir])^*(D
+2*GammaNL(u[0],v[0],w[0],u[ir],v[ir],w[ir])));
                   Fnlutmp[] = PbFnla(0,Vh);
                   Fnlu[] = Fnlu[] + Fnlutmp[];
                   for (int is=1;is<ir;is++)
                                    varf PbFnlb([utmp,vtmp,wtmp],[uuu,vvv,www]) = - int3d(Th3D,reg3DL) ( (dGammaNL(u[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-ir],v[ipn-i
is],w[ir-is],u[is],v[is],w[is]))) );
                                    Fnlutmp[] = PbFnlb(0,Vh);
                                    Fnlu[] = Fnlu[] + Fnlutmp[];
                  varf\ PbFnlc\ ([utmp,vtmp],[uuu,vvv,www]) = -int3d(Th3D,reg3DL)\ (\ (dGamma(u[0],v[0],w[0],uuu,vvv,www))'*(DL*(GammaNL(u[ipn-ir],v[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn-ir],w[ipn
ir],u[ir],v[ir],w[ir]))) );
                       Fnlutmp[] = PbFnlc(0,Vh);
                       Fnlu[] = Fnlu[] + Fnlutmp[];
varf PbFnICL ([utmp,vtmp,wtmp],[uuu,vvv,www]) = on(lsymmid,utmp=0.) + on(lrightmid,lrightlmid,vtmp=0.) + on(ldownmid,wtmp=0);
Fnlutmp[] = PbFnlCL(0,Vh);
Fnlu[] = Fnlu[] + Fnlutmp[];
```

Numerical results



Bifurcation curve of the film / substrate under uniaxial compression

Numerical results



Vertical displacement for $\lambda = 0.04832$ at the top film interface

Conclusions

- it has been demonstrated that ANM can be used to study non linear problems
- the FreeFem++ numerical development tool is very efficient to implement the ANM algorithm of non linear mechanical problems
- it has been applied to simulate the non linear mechanical behavior of thin films on soft substrates
- both bifurcation curve and plot of the vertical displacement at the top interface have been shown
- the positions of the two first bifurcations are close to Fan Xu's results

Thank you for your attention