Mutual Admittances of an Infinite Periodic Surface Acoustic Waves Transducer using an Original Coupled FEM/BIE Numerical Model

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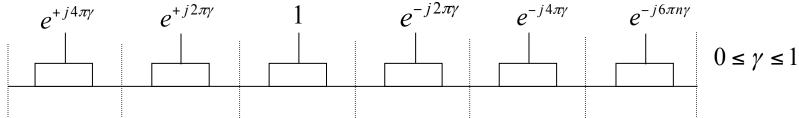
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Introduction

■ The mutual admittances of periodic SAW transducer can be easily computed from the harmonic admittance of the periodic array?

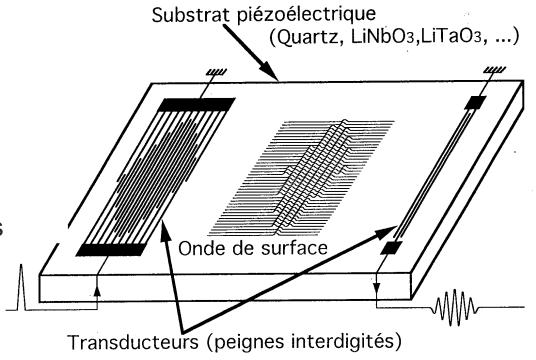


- Various numerical models have been published to compute the harmonic admittance of a periodic SAW transducer (Ballandras et alt. en 2002, Ventura et alt. en 2001, Hashimoto et alt. 2011).
- In 2013, Hecht et alt. published an original numerical model that uses an efficient variational formulation to deal with the periodicity of the geometrical boundary conditions written using the powerful FreeFem++ environment.
- Theoretical basis of the coupled FEM/BIE model and derivation of mutual admittances will be given, numerical simulations are shown

Principles of SAW IDT device

How is built a SAW device

- Piezoelectric substrate
- SAW IDT transducer
- The Surface Acoustic Wave is
 - launched and detected
 - propagating



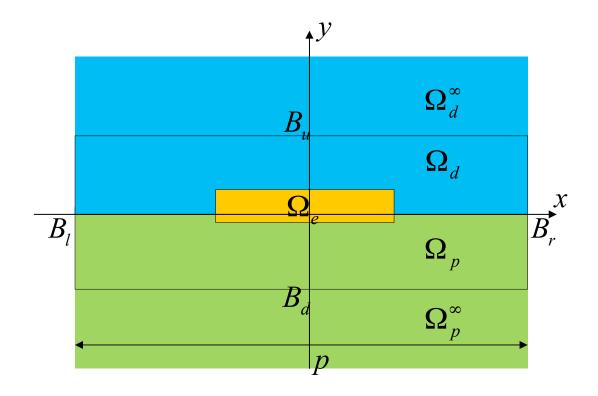


Main dates of piezoelectricity and its applications

- ► 1880 : discovery of piezoelectricity(Pierre et Jacques Curie)
- ▶ 1885 : Rayleigh proved the existence of a surface acoustic wave propagating at a lower velocity than bulk acoustic waves
- ▶ 1915 : first use of piezoelectric effects in acoustic : SONAR
- ▶ 1920 : application to electronics : High quality coefficient Quartz resonators
- ▶ 1950 : fabrication of piezoelectric ceramics
- ▶ 1960 : piezoelectric thin film
- ▶ 1965 : first interdigitated SAW transducers (Hartmann patent)

- The periodic array of electrodes is driven with an harmonic electrical potential $V_n(\gamma) = V_0 e^{-j2\pi n\gamma}$ $\gamma \in [0,1]$
- Using Bloch-Floquet's theorem, it is possible to model only one period of the infinite periodic array of electrodes $\psi(x) = e^{-j2\pi\gamma x/p}\psi_p(x)$
- 2D analysis (very long electrode) : plain strain approximation
- electrical assumption: no dielectric losses in the electrode
- mechanical assumption : the metallic electrode are homogeneous isotropic, elastic materials
- Build a coupled FEM / BIE model which combines the Finite Element Model (finite geometries) and the Boundary Integral Equation (semi infinite domains)

- lacktriangle The finite part of the model is bounded by the boundaries $\left(B_u^{},B_l^{},B_d^{},B_r^{}
 ight)$
- The Green domain Ω_n is the piezoelectric substrate.
- The orange domain Ω_{ρ} is the partially buried metallic electrode.
- The blue domain Ω_d is the dielectric medium (air of vacuum) above the array.



Definitions:

- u mechanical displacement
- ρ mass density, ω pulsation
- ϕ electrical potential
- E electrical field, D electrical displacement field
- T stress tensor, S strain tensor
- C^Eelastic tensor for a constant electrical field
- e piezoelectric coupling tensor, e dielectric tensor for a constant strain

- Differential equations:
 - ightharpoonup The Piezoelectric domain Ω_p and the Elastic domain Ω_e obeys Newton's second law:

$$\nabla \cdot \mathbf{T} = \rho \frac{\partial^2 \mathbf{u}}{\partial t^2}$$

The Piezoelectric domain Ω_p , the Elastic domain Ω_e , and the dielectric domain Ω_d obeys the quasistatic Maxwell's equation:

$$\nabla \cdot \mathbf{D} = 0$$

■ The constitutives equations:

In Ω_p (piezoelectric domain)

$$\begin{cases} \mathbf{T}_{ij} = \mathbf{C}_{ijkl}^{\mathbf{E}} \mathbf{S}_{kl} - \mathbf{e}_{ijk} \mathbf{E}_{k} \\ \mathbf{D}_{i} = \mathbf{e}_{kli} \mathbf{S}_{kl} + \varepsilon_{ik}^{\mathbf{S}} \mathbf{E}_{k} \end{cases}$$

In Ω_d (dielectric domain) $\mathbf{D}_i = \boldsymbol{\varepsilon}_{ik} \mathbf{E}_k$

$$\mathbf{D}_{i} = \varepsilon_{ik} \mathbf{E}_{k}$$

In
$$\Omega_e$$
 (elastic domain)

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$$\Omega_e$$
 (elastic domain)
$$\mathbf{T}_{ij} = (\lambda + \mu) \delta_{ij} \mathbf{S}_{kk} - 2\mu \mathbf{S}_{ij}$$

$$\mathbf{S}_{ij} = \frac{1}{2} \left(\frac{\partial u_i}{\partial x_j} + \frac{\partial u_j}{\partial x_i} \right) \qquad \mathbf{E}_i = -\frac{\partial \Phi}{\partial x_i}$$

- The γ periodic boundary conditions:
 - For the interfaces B_l and $B_r(y>0)$: $\Phi(+p/2,y) = e^{-j2\pi\gamma}\Phi(-p/2,y)$
 - For the interfaces B_l and $B_r(y < 0)$: $\begin{cases} \mathbf{u}(+p/2, y) = e^{-j2\pi\gamma} \mathbf{u}(-p/2, y) \\ \Phi(+p/2, y) = e^{-j2\pi\gamma} \Phi(-p/2, y) \end{cases}$

Variational Formulation

Finds (\mathbf{u}, ϕ) in $V^3(\Omega \cup \Omega) \times V^3(\Omega)$ (satisfying $\phi = 1$ in the electrode) such that for all $(\mathbf{v}, \psi)^p$ in $V^3(\Omega_p \cup \Omega_e) \times V^3(\Omega)$ (satisfying $\phi = 0$ in the electrode)

$$\int_{\Omega_{p} \cup \Omega_{E}} \overline{\mathbf{S}}(\mathbf{v}) \cdot \mathbf{T}(\mathbf{u}) d\Omega - \rho \omega^{2} \int_{\Omega_{p} \cup \Omega_{E}} \overline{\mathbf{v}} \cdot \mathbf{u} d\Omega
- \int_{\Omega_{p} \cup \Omega_{E} \cup \Omega_{D}} \overline{\mathbf{E}}(\psi) \cdot (\mathbf{eS}(u) + \varepsilon \mathbf{E}(\phi)) d\Omega
= \int_{B_{d}} \overline{\mathbf{v}} \cdot (\mathbf{T} \cdot \mathbf{n}) d\Gamma + \int_{B_{u} \cup B_{d}} \overline{\psi}(\mathbf{D}(\phi) \cdot \mathbf{n}) d\Gamma$$

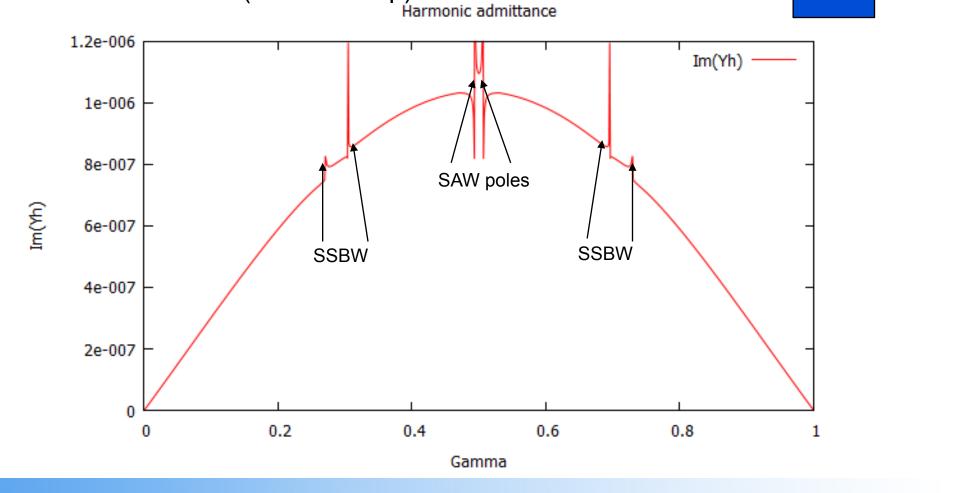
- $V_{\gamma}(\Omega)$ is the mathematical space of $L^{2}(\Omega)$ satisfying γ -harmonic periodic boundary conditions
- It is possible to transform the weak formulation into a weak formulation with only periodic physical fields

Numerical implementation

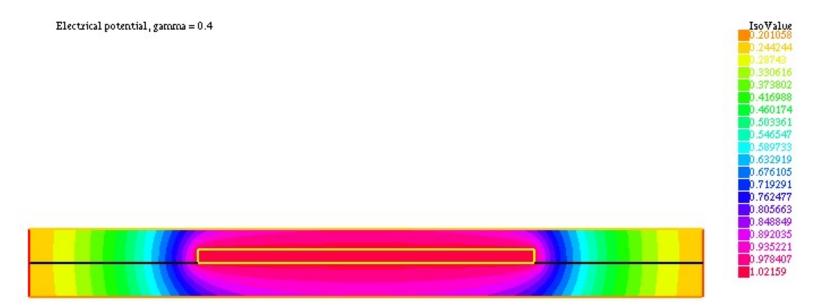
- Numerical implementation in the powerfull FreeFem++ environment developped by F. Hecht, O. Pironneau, A. Le Hyaric, from Laboratoire Jacques Louis Lions, Université Pierre et Marie Curie in France, and O. Ohtsuka from Hiroshima university in Japan : http://www.freefem++.org/ff++/
- In the variational formulation, the volume integrals are easily taken into account with FEM, special care is needed for the surface integrals (BIE part of the model) involving semi-infinite Green's function for the dielectric and the piezoelectric half spaces.
- All the mathematical details are given in:
 F. Hecht, P. Ventura, and, P. Dufilié, "Original coupled FEM/BIE Numerical Model for Analyzing Infinite Periodic Surface Acoustic Wave Transducer", Journal of Computational Physics, Vol. 246, August 2013, pp. 265-274.

Plot of the harmonic admittance

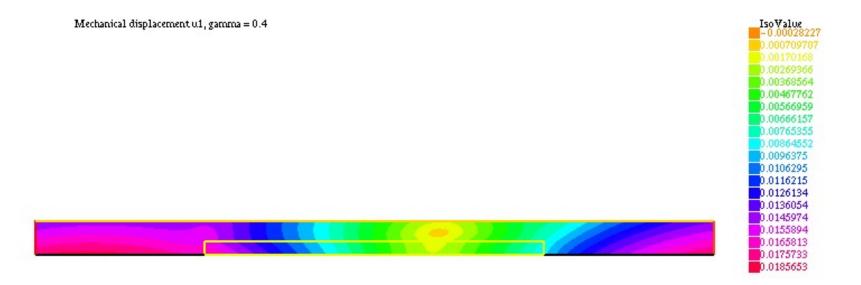
■ Fixed frequency (0.499 MHz). The piezoelectric substrate is Y+39° X propagation Quartz, the period of the transducer is 3.1145 mm, the electrode are assumed rectangular with 2% buried depth and 2% metal thickness (relative to 2p).



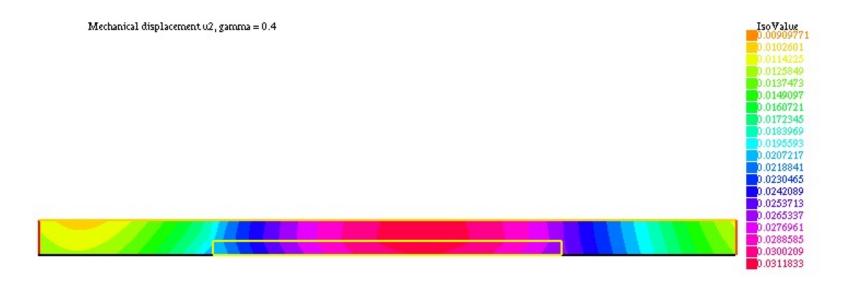
Plot of the electrical field



Plot of u1



Plot of u2



- The mutual admittance Y_m is related to the harmonic admittance $Y(\gamma)$ $Y_m = \int_0^1 Y(\gamma) e^{j2\pi m\gamma} d\gamma$
- because of very fast variation of $Y(\gamma)$ close to the SAW pole, a special mathematical treatment is needed.
- SAW contribution to the mutual admittance:

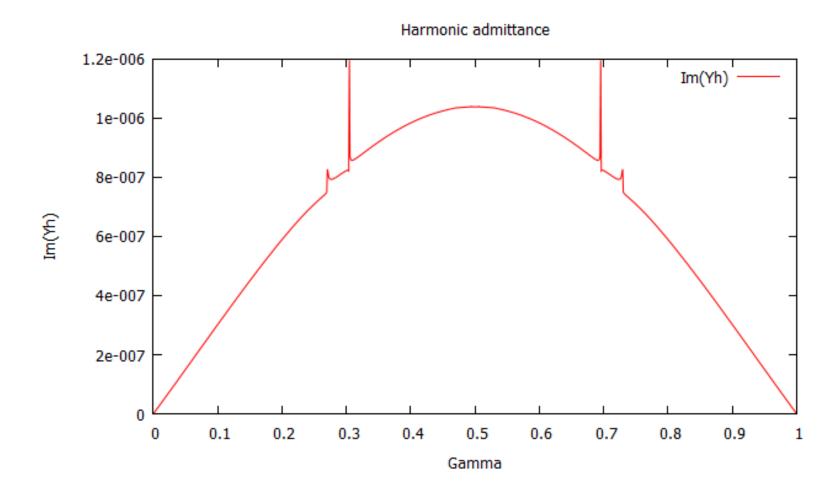
$$\begin{cases} Y_{SC} \exp\left(-j \left| m \right| \phi_{SC}\right) \text{ for } m \neq 0 \\ Y_{SC} \left(1 + j / \tan\left(\phi_{SC}/2\right)\right) \text{ for } m = 0 \quad \text{(to satisfy charge conservation)} \end{cases}$$

SAW contribution to the harmonic admittance

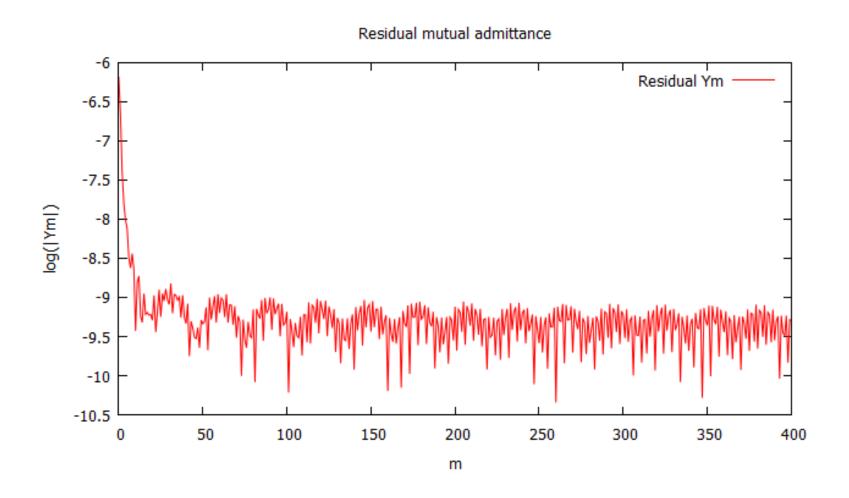
$$Y_{S}(\gamma) = \frac{jY_{SC}}{\tan(\phi_{SC}/2)} \frac{1 - \cos(2\pi\gamma)}{\cos(\phi_{SC}) - \cos(2\pi\gamma)}$$

lacksquare Y_{SC} and ϕ_{SC} are computed using an iterative algorithm

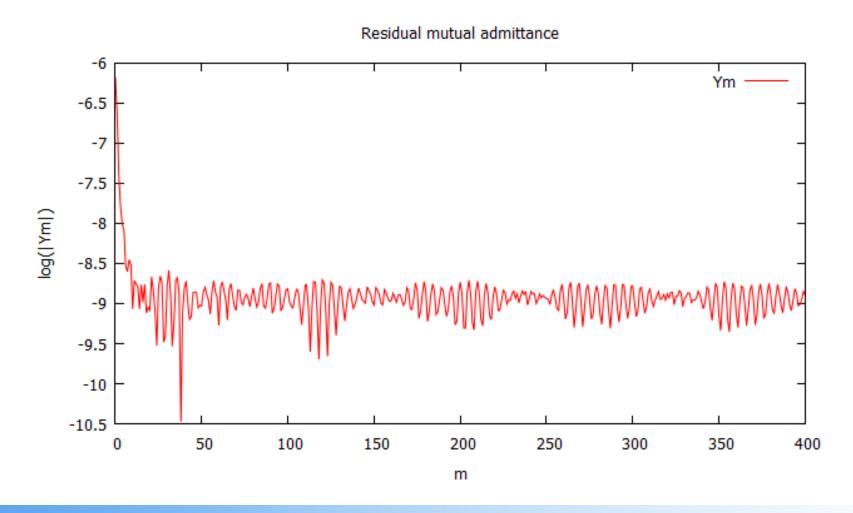
■ residual harmonic admittance



■ residual mutual admittance computed by FFT



 mutual admittance after adding the analytical expression of the SAW contribution



Conclusions

■ A new coupled FEM/BIE numerical model that computes the harmonic admittance of an infinite periodic SAW transducer has been developed in the powerful FreeFem++ environment.

- Using an efficient numerical technique to extract the SAW contribution, we have been able to compute the mutual admittance of the SAW transducer with a very good accuracy.
- This numerical technique can be generalized for various periodic physical problems (acoustics, electromagnetics, ...)