FreeFEM days 2021, 10 Dec.

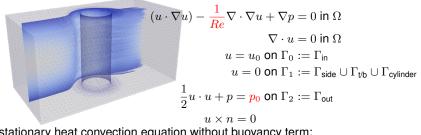
Shape optimization for a heat exchanger in Navier-Stokes flow with dynamic pressure at the outlet

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heat exchanger with cylinder shape in Navier-Stokes flow

boundary $\partial \Omega = \Gamma_{in} \cup \Gamma_{out} \cup \Gamma_{side} \cup \Gamma_{t/b} \cup \Gamma_{cylinder}$ stationary Navier-Stokes equations:



stationary heat convection equation without buoyancy term:

$$\begin{split} (u\cdot\nabla)\theta - \kappa\nabla\cdot\nabla\theta &= 0 \text{ in } \Omega \\ \theta &= 0 \text{ on } \widetilde{\Gamma}_0 := \Gamma_{\text{in}} \\ \theta &= 1 \text{ on } \widetilde{\Gamma}_1 := \Gamma_{\text{t/b}} \cup \Gamma_{\text{cylinder}} \\ \nabla\theta\cdot n &= 0 \text{ on } \widetilde{\Gamma}_2 := \Gamma_{\text{side}} \cup \Gamma_{\text{out}} \end{split}$$

 ϕ : shape deformation $\Omega(\phi) = \phi(\widehat{\Omega})$ from the reference domain $\widehat{\Omega}$ cost function:

$$J(\phi, \{u, p, \theta\}) = -\int_{\widetilde{\Gamma}_1(\phi)} \kappa \nabla \theta \cdot n \to \min$$

weak formulation for Navier-Stokes eqs. with dynamic pressure cnd.

$$V=\{v\in H^1(\Omega)^3\,;\, v=0 \text{ on } \Gamma_0\cup\Gamma_1\,, v\times n=0 \text{ on } \Gamma_2\},\, Q=L^2(\Omega)$$

$$\begin{split} & \int_{\Omega} (\nabla \times u) \times u \cdot v - \frac{1}{2} \int_{\Omega} (u \cdot u) \nabla \cdot v + \frac{1}{Re} \int_{\Omega} \nabla \times u \cdot \nabla \times v \\ & - \int_{\Omega} \nabla \cdot v p - \int_{\Omega} \nabla \cdot u q = - \int_{\Gamma} p_0 \, v \cdot n \quad \forall (v,q) \in V \times Q \end{split}$$

solution of the weak formulation \Rightarrow strong form + boundary condition for velocity satisfying $\nabla \cdot u = 0$, applying integration by parts

$$\begin{split} 0 &= \int_{\Omega} (\nabla \times u) \times u \cdot v + \frac{1}{2} \int_{\Omega} \nabla (u \cdot u) \cdot v - \frac{1}{2} \int_{\partial \Omega} (u \cdot u) v \cdot n \\ &+ \frac{1}{Re} \int_{\Omega} \nabla \times (\nabla \times u) \cdot v + \frac{1}{Re} \int_{\partial \Omega} n \times (\nabla \times u) \cdot v \\ &+ \int_{\Omega} v \cdot \nabla p - \int_{\partial \Omega} v \cdot n p + \int_{\Gamma_2} p_0 \, v \cdot n \\ &= \int_{\Omega} \left((\nabla \times u) \times u + \frac{1}{2} \nabla (u \cdot u) \right) \cdot v - \int_{\Gamma_0 \cup \Gamma_1 \cup \Gamma_2} \left(\frac{1}{2} (u \cdot u) + p - p_0 \right) (v \cdot n) \\ &+ \frac{1}{Re} \int_{\Omega} (\nabla (\overline{\mathbf{V}} \cdot \mathbf{u}) - \nabla \cdot \nabla u + \nabla p) \cdot v - \frac{1}{Re} \int_{\Gamma_0 \cup \Gamma_1 \cup \Gamma_2} (\nabla \times u) \cdot (v \times n) \end{split}$$

NB:
$$(u \cdot \nabla u)u = (\nabla \times u) \times u + \frac{1}{2}\nabla(u \cdot u)$$
 and $(\nabla \times u) \times u = \nabla(\nabla \cdot u) - \nabla \cdot \nabla u$

finite element approximation to Navier-Stokes eqs. with curl-curl form

 $V_h = \{v|_K \in P_1(K)^3 ; v = 0 \text{ on } \Gamma_0 \cup \Gamma_1, v \times n = 0 \text{ on } \Gamma_2\}, Q_h = \{q|_K \in P_1(K)\}$ two difficulties in P1/P1 approximation to curl-curl form

- ▶ P1/P1 approx. does not satisfy the inf-sup cond. $\Leftrightarrow \exists p \neq 0 \ B^T p = 0$

$$\begin{split} &\int_{\Omega} (\nabla \times u) \times u \cdot v - \frac{1}{2} \int_{\Omega} (u \cdot u) \nabla \cdot v + \frac{1}{Re} \int_{\Omega} \nabla \times u \cdot \nabla \times v + \int_{\Omega} \nabla \cdot u \nabla \cdot v \\ &- \int_{\Omega} \nabla \cdot v p + \int_{\Omega} \nabla \cdot u q + \delta \sum_{K} h_{K}^{2} \int_{k} \nabla p \cdot \nabla q = - \int_{\Gamma_{2}} p_{0} \, v \cdot n \quad \forall (v,q) \in V_{h} \times Q_{h} \end{split}$$

stiffness matrix for a linearlized equation

$$\begin{bmatrix} A_1 + A_0 + C & B^T \\ -B & \delta D \end{bmatrix} \quad \begin{cases} \begin{bmatrix} [A_0]_{ij} = \int_{\Omega} \nabla \times \varphi_j \cdot \nabla \times \varphi_i \\ C_{ij} = \int_{\Omega} \nabla \cdot \varphi_j \nabla \cdot \varphi_i \\ B_{ij} = -\int_{\Omega} \nabla \cdot \varphi_j \psi_i \\ D_{ij} = \sum_K h_K^2 \int_K \nabla \psi_j \cdot \nabla \psi_i, \ \delta > 0 \end{cases}$$

 $A_0 = A_0^T$, $C = C^T$, KerC = KerB, Im $A_0 = \text{Ker}B$ stiffness matrix of the Stokes eqs., without unsymmetric (nonlinear) term

$$\begin{bmatrix} A_0 + C & B^T \\ -B & \delta D \end{bmatrix} \text{ is coercive } \Rightarrow \text{additive Schawrz preconditioner in } \text{hpddm}$$

▶ a proof of $(A_0 + C)u = 0 \Rightarrow u = 0$ $\text{Ker}C = \text{Ker}B \ni A_0u = -Cu \in \text{Im}C \Rightarrow A_0u = 0 \land u \in \text{Ker}C = \text{Ker}B$ A_0 is coercive on $\text{Ker}B \Rightarrow u = 0$ C has the same effect as B^TB

weak formulation for thermal convection diffusion

$$\Theta = \{ \theta \in H^1(\Omega) \, ; \, \theta = 0 \text{ on } \widetilde{\Gamma}_0 \cup \widetilde{\Gamma}_1 \}$$

$$\int_{\Omega} (u \cdot \nabla)\theta \xi + \kappa \int_{\Omega} \nabla \theta \cdot \nabla \xi = 0 \quad \forall \theta \in \Theta$$

solution of the weak formulation \Rightarrow strong form + boundary condition for given u

$$\begin{split} \int_{\Omega} \left((u \cdot \nabla) \theta - \kappa \nabla \cdot \nabla \theta \right) \xi + \kappa \int_{\widetilde{\Gamma}_0 \cup \widetilde{\Gamma}_1 \cup \widetilde{\Gamma}_2} \nabla \theta \cdot n \xi &= 0 \quad \forall \xi \in \Theta \\ \Rightarrow \nabla \theta \cdot n &= 0 \text{ on } \widetilde{\Gamma}_2 \end{split}$$

surface integration is converted to domain integration η^* : harmonic extension of $\eta=1$ on $\widetilde{\Gamma}_2$, $\eta=0$ on $\widetilde{\Gamma}_0\cup\widetilde{\Gamma}_1$ θ is the solution of the strong form for given u

$$\begin{split} &-\kappa\int_{\widetilde{\Gamma}_2}\nabla\theta\cdot n\eta\\ &=-\kappa\int_{\widetilde{\Gamma}_2}\nabla\theta\cdot n\eta-\int_{\Omega}\left((u\cdot\nabla)\theta-\nabla\cdot\nabla\theta\right)\eta^*\\ &=-\kappa\int_{\widetilde{\Gamma}_2}\nabla\theta\cdot n\eta-\int_{\Omega}\left(u\cdot\nabla)\theta\eta^*-\kappa\int_{\Omega}\nabla\theta\cdot\nabla\eta^*+\int_{\widetilde{\Gamma}_0\cup\widetilde{\Gamma}_1\cup\widetilde{\Gamma}_2}\kappa\nabla\theta\cdot n\eta\\ &=-\int_{\Omega}(u\cdot\nabla)\theta\eta^*-\kappa\int_{\Omega}\nabla\theta\cdot\nabla\eta^*=-\int_{\Omega}\nabla\theta\cdot u\eta^*-\kappa\int_{\Omega}\nabla\theta\cdot\nabla\eta^* \end{split}$$

cost function and Lagrangean

$$J(\phi, \{u, p, \theta\}) = -\int_{\Gamma_2(\phi)} \kappa \nabla \theta \cdot n$$
$$= -\int_{\Omega(\phi)} \nabla \theta \cdot u \eta^* - \kappa \int_{\Omega(\phi)} \nabla \theta \cdot \nabla \eta^*$$

constraint for the thermal fluid

$$\begin{split} &A(\phi,\{u,p,\theta\},\{v,q,\xi\}) \\ &= \int_{\Omega} (\nabla \times u) \times u \cdot v - \frac{1}{2} \int_{\Omega} (u \cdot u) \nabla \cdot v + \frac{1}{Re} \int_{\Omega} \nabla \times u \cdot \nabla \times v + \int_{\Omega} \nabla \cdot u \nabla \cdot v \\ &- \int_{\Omega} \nabla \cdot v p + \int_{\Omega} \nabla \cdot u q + \delta \sum_{K} h_{K}^{2} \int_{k} \nabla p \cdot \nabla q + \int_{\Gamma_{2}} p_{0} \, v \cdot n \\ &+ \int_{\Omega(\phi)} (\nabla \theta \cdot u) \xi + \kappa \int_{\Omega(\phi)} \nabla \theta \cdot \nabla \xi \end{split}$$

Lagrangean is defined as

$$\mathcal{L}(\phi,\{u,p,\theta\},\{v,q,\xi)\}=J(\phi,\{u,p,\theta\})+A(\phi,\{u,p,\theta\},\{v,q,\xi\})$$
 state problem

 $d\mathcal{L}_{\{v,q,\xi\}}(\phi,\{u,p,\theta\},\{v,q,\xi\})[\{\delta v,\delta q,\delta \xi\}] = 0$

adjoint problem using solution $\{u, p, \theta\}$ of the state problem

$$d\mathcal{L}_{\{u,p,\theta\}}(\phi,\{u,p,\theta\},\{v,q,\xi\})[\{\delta u,\delta p,\delta \theta\}] = 0$$

state and adjoint problems

state problem to find $\{u, p, \theta\}$

$$\begin{split} 0 = & d\mathcal{L}_{\{v,q,\xi\}}(\phi,\{\textbf{\textit{u}},\textbf{\textit{p}},\theta\},\{v,q,\xi\})[\{\delta v,\delta q,\delta\xi\}] = A(\varphi,\{\textbf{\textit{u}},\textbf{\textit{p}},\theta\},\{\delta v,\delta p,\delta\xi\}) \\ = & \int_{\Omega}(\nabla\times\textbf{\textit{u}})\times\textbf{\textit{u}}\cdot\delta v - \frac{1}{2}\int_{\Omega}(\textbf{\textit{u}}\cdot\textbf{\textit{u}})\nabla\cdot\delta v + \frac{1}{Re}\int_{\Omega}\nabla\times\textbf{\textit{u}}\cdot\nabla\times\delta v + \int_{\Omega}\nabla\cdot\textbf{\textit{u}}\nabla\cdot\delta v \\ & - \int_{\Omega}\nabla\cdot\delta v\textbf{\textit{p}} + \int_{\Omega}\nabla\cdot\textbf{\textit{u}}\delta q + \delta\sum_{\mathbf{\textit{v}}}h_{K}^{2}\int_{k}\nabla\textbf{\textit{p}}\cdot\nabla\delta q + \int_{\Gamma_{0}}p_{0}\,\delta v\cdot n \end{split}$$

$$+ \int_{\Omega(\phi)} (\nabla \boldsymbol{\theta} \cdot \boldsymbol{u}) \delta \xi + \kappa \int_{\Omega(\phi)} \nabla \boldsymbol{\theta} \cdot \nabla \delta \xi$$

adjoint problem to find $\{v, q, \xi\}$ using solution $\{u, p, \theta\}$ of the state problem

adjoint problem to find
$$\{v,q,\xi\}$$
 using solution $\{u,p,\theta\}$ of the state proble $0=d\mathcal{L}_{\{u,p,\theta\}}(\varphi,\{u,p,\theta\},\{v,q,\xi\})[\{\delta u,\delta p,\delta \theta\}]$

$$\begin{split} &= \int_{\Omega} (\nabla \times \delta u) \times \boldsymbol{u} \cdot \boldsymbol{v} + \int_{\Omega} (\nabla \times \boldsymbol{u}) \times \delta \boldsymbol{u} \cdot \boldsymbol{v} - \int_{\Omega} (\delta \boldsymbol{u} \cdot \boldsymbol{u}) \nabla \cdot \boldsymbol{v} \\ &+ \frac{1}{Re} \int_{\Omega} \nabla \times \delta \boldsymbol{u} \cdot \nabla \times \boldsymbol{v} + \int_{\Omega} \nabla \cdot \delta \boldsymbol{u} \nabla \cdot \boldsymbol{v} \\ &- \int_{\Omega} \nabla \cdot \boldsymbol{v} \delta \boldsymbol{p} + \int_{\Omega} \nabla \cdot \delta \boldsymbol{u} \boldsymbol{q} + \delta \sum_{K} h_{K}^{2} \int_{k} \nabla \delta \boldsymbol{p} \cdot \nabla \boldsymbol{q} \end{split}$$

$$+ \int_{\Omega(\phi)} (\nabla \delta \theta \cdot u) \frac{\boldsymbol{\xi}}{\boldsymbol{\xi}} + \int_{\Omega(\phi)} (\nabla \theta \cdot \delta u) \frac{\boldsymbol{\xi}}{\boldsymbol{\xi}} + \kappa \int_{\Omega(\phi)} \nabla \delta \theta \cdot \nabla \frac{\boldsymbol{\xi}}{\boldsymbol{\xi}}$$

$$- \int_{\Omega(\phi)} \nabla \theta \cdot \delta u \eta^* - \int_{\Omega(\phi)} \nabla \delta \theta \cdot u \eta^* - \kappa \int_{\Omega(\phi)} \nabla \delta \theta \cdot \nabla \eta^*$$

shape derivative: assuming material derivative vanished

shape derivative (material derivative)

$$u(\phi + \varphi)(x + \varphi(x)) - u(\phi)(x) \simeq u'(\phi)[\varphi](x)$$

shape derivative of gradient with $F(\varphi) = I + \nabla \varphi$

$$\nabla_z u(\phi + \varphi)(x + \varphi(x)) \simeq F^{-T}(\varphi)\nabla_x u(\phi)(x) + (\nabla_x u)'(\phi)[\varphi](x)$$

derivative of domain integral of scalar function u

$$f(\phi, u(\phi)) = \int_{\Omega(\phi)} u(\phi)(x) dx \implies d f_{\phi}(\phi, u)[\varphi] = \int_{\Omega(\phi)} (u' + u(\nabla \cdot \varphi)) dx$$

derivative of domain integral of ∇u

$$f(\phi, \nabla u(\phi)) = \int_{\Omega(\phi)} \nabla u(\phi)(x) dx$$

$$\Rightarrow df_{\phi}(\phi, \nabla u)[\varphi] = \int_{\Omega(\phi)} \left\{ (\nabla u)'[\varphi] - (\nabla \varphi)^T \nabla u + (\nabla \cdot \varphi) \nabla u \right\} dx$$

- ▶ differential operators of 1st order, ∇ , ∇ × to vector valued u are written with tensor of $(\nabla u)^T$ like $\nabla \cdot u = (\nabla u) : I = (\nabla u)^T : I$
- $ightharpoonup
 abla
 ightharpoonup -(
 abla \delta \varphi)^T
 abla and multiplying
 abla \cdot \delta \varphi$ to the integrand

shape derivative of integration of curl-curl form

$$\int_{\Omega(\phi+\varphi)} (\nabla_z \times u(\phi+\varphi)) \cdot (\nabla_z \times v(\phi+\varphi)) - \int_{\Omega(\phi)} (\nabla \times u(\phi)) \cdot (\nabla \times v(\phi))$$

 $\nabla \times u$ is computed by component-wise as

$$\nabla \times u = \begin{bmatrix} (\nabla u)^T : \begin{bmatrix} 0 & 0 & 0 \\ 0 & 0 & 1 \\ 0 & -1 & 0 \end{bmatrix}, \ (\nabla u)^T : \begin{bmatrix} 0 & 0 & -1 \\ 0 & 0 & 0 \\ 1 & 0 & 0 \end{bmatrix}, \ (\nabla u)^T : \begin{bmatrix} 0 & 1 & 0 \\ -1 & 0 & 0 \\ 0 & 0 & 0 \end{bmatrix} \end{bmatrix}^T$$

$$\int_{\Omega(\phi+\varphi)} (\nabla_z u(\phi+\varphi))^T : B_1 - \int_{\Omega(\phi)} (\nabla u(\phi))^T : B_1$$

$$\simeq \int_{\Omega(\phi)} ((F^{-T}(\phi) + F^{-T}'(\phi)[\varphi]) (\nabla u(\phi))^T + (\nabla u)^{T}'(\phi)[\varphi]) : B_1(\omega(\phi) + \omega'(\varphi)) - (\nabla u(\phi))^T : B_1$$

$$= \int_{\Omega(\phi)} ((F - (\psi) + F - (\psi)[\varphi]) (\nabla u(\psi)) + (\nabla u) - (\psi)[\varphi]) \cdot DI(\omega(\psi) + \omega \cdot (\varphi)[\varphi]) \cdot DI(\omega(\psi) + \omega$$

$$\simeq \int_{\Omega(\phi)} (\nabla u(\varphi))^T : B_1(\nabla \cdot \varphi) - (\nabla \varphi)^T (\nabla u(\phi))^T : B_1 + (\nabla u)^{T'}(\phi)[\varphi] : B_1$$

supposing $(\nabla u)^{T'}(\phi) = 0$, and using

$$[(\nabla \varphi)^T (\nabla u)^T : B_1, \ (\nabla \varphi)^T (\nabla u)^T : B_2, \ (\nabla \varphi)^T (\nabla u)^T : B_3]^T = ((\nabla \varphi)^T \nabla) \times u$$

shape derivative is calculated as

$$\int_{\Omega(\phi)} (\nabla \times u) \cdot (\nabla \times v) (\nabla \cdot \varphi) - \left(\left((\nabla \varphi)^T \nabla \right) \times u \right) \cdot (\nabla \times v) - (\nabla \times u) \cdot \left(\left((\nabla \varphi)^T \nabla \right) \times v \right)$$

shape derivative of the Lagrangean of the NS eqs. $d\mathcal{L}_{\phi}(\phi, \{u, p, \theta\}, \{v, q, \xi\})[\varphi]$

$$= -\int_{\Omega(V)} \nabla \theta \cdot u \eta^* (\nabla \cdot \varphi) - (\nabla \varphi)^T \nabla \theta \cdot u \eta^*$$

$$J_{\Omega(\phi)}$$

$$-\kappa \int_{\Omega(\phi)} \nabla \theta \cdot \nabla \eta^* (\nabla \cdot \varphi) - (\nabla \varphi)^T \nabla \theta \cdot \nabla \eta^* - \nabla \theta \cdot (\nabla \varphi)^T \nabla \eta^*$$

+
$$\int_{\Omega(\phi)} ((\nabla \times u) \times u) \cdot v(\nabla \cdot \varphi) + \left(\left(((\nabla \varphi)^T \nabla) \times u \right) \times u \right) \cdot v$$

$$-\frac{1}{2}\int_{O(t)}(u\cdot u)(\nabla\cdot v)(\nabla\cdot \varphi)+(u\cdot u)((\nabla\varphi)^T\nabla)\cdot v)$$

$$= \frac{1}{2} \int_{\Omega(\phi)} (u \cdot u)(\mathbf{v} \cdot \mathbf{v})(\mathbf{v} \cdot \boldsymbol{\varphi}) + (u \cdot u)((\mathbf{v} \cdot \boldsymbol{\varphi}) \cdot \mathbf{v}) \cdot \mathbf{v}$$

$$+ \frac{1}{Re} \int_{\Omega(\phi)} (\nabla \times u \cdot \nabla \times v) (\nabla \cdot \varphi) - ((\nabla \varphi)^T \nabla) \times u \cdot \nabla \times v - \nabla \times u \cdot ((\nabla \varphi)^T \nabla) \times v$$

$$+ \int_{\Omega(\phi)} \nabla \cdot u \nabla \cdot v (\nabla \cdot \varphi) - (\nabla \varphi)^T \nabla \cdot u \nabla \cdot v - \nabla \cdot u (\nabla \varphi)^T \nabla \cdot v$$

$$(\nabla \cdot arphi) - (
abla arphi)^T
abla \cdot arphi$$

$$\varphi$$
) $-\left((\nabla\varphi)^T\nabla\right)\cdot v$

$$(\nabla \varphi)^T \nabla \cdot v p$$

$$0 - \left((\nabla \varphi)^T \nabla \right) \cdot v \, p$$

$$+ \delta \sum_{K} h_{K}^{2} \int_{K(\phi)} \left\{ \nabla p \cdot \nabla q (\nabla \cdot \varphi) - (\nabla \varphi)^{T} \nabla p \cdot q - \nabla p \cdot (\nabla \varphi)^{T} \nabla q \right\}$$

$$+ \kappa \int_{\Omega(\phi)} \nabla \theta \cdot \nabla \xi (\nabla \cdot \varphi) - (\nabla \varphi)^T \nabla \theta \cdot \nabla \xi - \nabla \theta \cdot (\nabla \varphi)^T \nabla \xi$$

$$(\varphi)^T \nabla (x) \times u \times u \cdot v$$

$$u \cdot \nabla \times v - \nabla \times u \cdot \left((\nabla \varphi)^T \nabla \right)$$

$$u(\nabla \varphi)^T \nabla \cdot v$$

$$-\int_{\Omega(\phi)} \nabla \cdot v \, p(\nabla \cdot \varphi) - \left((\nabla \varphi)^T \nabla \right) \cdot v \, p + \int_{\Omega(\phi)} \nabla \cdot u \, q(\nabla \cdot \varphi) - \left((\nabla \varphi)^T \nabla \right) \cdot u \, q$$

$$\left\{
abla q
ight\}$$

$$(\nabla \varphi)^T \nabla \theta \cdot u \mathcal{E}$$

$$+ \int_{\Omega(\phi)} \nabla \theta \cdot u \xi (\nabla \cdot \varphi) - (\nabla \varphi)^T \nabla \theta \cdot u \xi$$

gradient flow solver with inner product by elasticity equations

gradient flow solver with H^1 -norm

$$((\frac{\varphi^{k+1} - \varphi^k}{\Delta \tau}, \varphi)) = -d\mathcal{L}_{\phi} \left(\varphi^k, \{u^k, p^k, \theta^k\}, \{v^k, q^k, \xi^k\}\right) [\varphi] \ \forall \varphi$$
$$((\psi, \varphi)) = d\mathcal{L}_{\phi} \left(\varphi^k, \{u^k, p^k, \theta^k\}, \{v^k, q^k, \xi^k\}\right) [\varphi] \ \forall \varphi$$
$$\varphi^{k+1} = \varphi^k - \Delta \tau \psi$$

 $\{u^k,p^k,\theta^k\}$: solution of the state problem at k-th pseudo time step $\{v^k,q^k,\xi^k\}$: solution of the adjoint problem at k-th pseudo time step elasticity equations to update domain with Lamé constants $\lambda=\mu=1$

$$2\mu \int_{\Omega} D(\psi) : D(\varphi) + \lambda \int_{\Omega} \nabla \cdot \psi \nabla \cdot \varphi = d\mathcal{L}_{\phi} \left(\varphi^{k}, \{u^{k}, p^{k}, \theta^{k}\}, \{v^{k}, q^{k}, \xi^{k}\} \right) [\varphi]$$

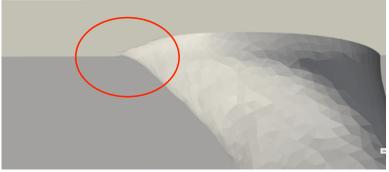
boundary conditions for moving surface of the cylinder whose edges are fixed on top/bottom surfaces within a box

$$\begin{split} \psi &= 0 \text{ on } \Gamma_{\text{in}} \cup \Gamma_{\text{out}} \cup \Gamma_{\text{side}} \\ \psi \cdot n &= 0 \text{ on } \Gamma_{\text{t/b}} \\ D(\psi)n &= 0 \text{ on } \Gamma_{\text{cylinder}} \end{split}$$

 inequality condition to keep cylinder surface within upper/bottom boundary

inequality conditions to keep the upper/bottom surfaces

surface Γ_{cylinder} is updated by $\vec{x} - \Delta \tau \psi(\vec{x})$



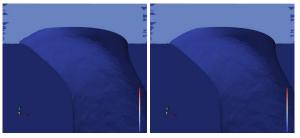
inequality condition to keep cylinder surface within upper/bottom boundary

$$-1 \leq (\vec{x} - \Delta \tau \psi(\vec{x}))|_z \leq 1 \iff \frac{z-1}{\Delta \tau} \leq \psi(\vec{x})|_z \leq \frac{z+1}{\Delta \tau} \text{ on } \Gamma_{\text{cylinder}}$$

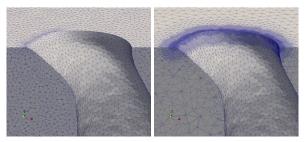
variational inequality can be solved by

- IPOPT : easy to implement but in sequential not in distributed way
- primal-dual active set method : need to be written by FreeFEM script but can be parallelized

re-meshing procedure with mmg3d and Gmsh



FEM node $\Gamma_{
m cylinder}$ is moved to $\Gamma_{
m top}$ with z=1



- mmg3d to coarsen the surface triangulation
- ► Gmsh to generate refined mesh for deformed surface with updated edge

FreeFEM script to change surface label

thanks to Dr. P. Jolivet

```
mesh3 Th:
int[int] ifaces(ThG.nt * 4);
// retrieving the original label of the boundary surface
for (int k = 0; k < Th.nbe; k++) {
 int kk = Th.be(k).Element:
  int ll = Th.be(k).whoinElement;
  ifaces[kk * 4 + 111 = Th.be(k).label;
// set new label
  int k = targetsurfacelement;
  int kk = Th.be(k).Element;
  int ll = Th.be(k).whoinElement;
  int newlabel:
  ifaces[kk * 4 + 11] = newlabel;
Th = change(Th, flabel = ifaces[nuTet * 4 + nuFacel);
```

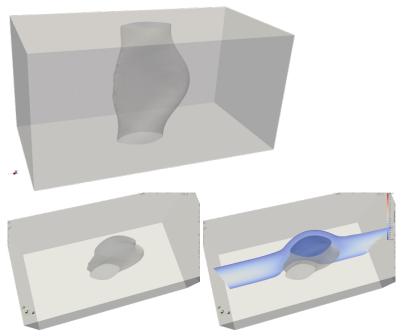
whoinElement returns vertex id, which is the opposite side of the boundary surface and inside of the domain, and also indicates element surface id

technique to write shape gradient operator in FreeFEM

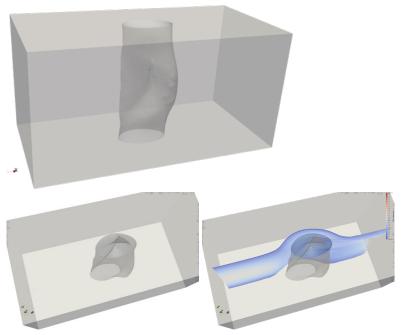
```
for shape derivative of curl-curl form
                       \int_{\Omega} \left( (\nabla \varphi)^T \nabla \right) \times u \cdot \nabla \times v
macro DOT(uu, vv) ((uu)' * (vv)) //'
macro Rot (uu) [(dy(uu[2]) - dz(uu[1])),
                (dz(uu[0]) - dx(uu[2])),
                (dx(uu[1]) - dv(uu[0])) //
macro ShapeTensor(ww) [[dx(ww[0]), dx(ww[1]), dx(ww[2])],
                         [dy(ww[0]), dy(ww[1]), dy(ww[2])],
                         [dz(ww[0]), dz(ww[1]), dz(ww[2])]] //
macro ShapeRot (aa, uu)
[((aa[1][0])*dx(uu[2])+(aa[1][1])*dy(uu[2])+(aa[1][2])*dz(uu[2])
 -(aa[2][0])*dx(uu[1])+(aa[2][1])*dy(uu[1])+(aa[2][2])*dz(uu[1]))
 ((aa[2][0])*dx(uu[0])+(aa[2][1])*dy(uu[0])+(aa[2][2])*dz(uu[0])
 -(aa[0][0])*dx(uu[2])+(aa[0][1])*dy(uu[2])+(aa[0][2])*dz(uu[2]))
 ((aa[0][0])*dx(uu[1])+(aa[0][1])*dy(uu[1])+(aa[0][2])*dz(uu[1])
 -(aa[1][0])*dx(uu[0])+(aa[1][1])*dy(uu[0])+(aa[1][2])*dz(uu[0]))] //
func Pk4 = [P1, P1, P1, P1];
fespace Wh(Th, Pk4);
Wh [uu, uu, uu, pp], [vv, vv, vv, qq]; // state/adjoint sols.
varf SensitivityRHS([w1, w2, w3], [phi1, phi2, phi3]) //
= int3d(Th,qfV=qfV1)( // lowest order
   DOT(ShapeRot(ShapeTensor([phi1, phi2, phi3]), [uu1, uu2, uu3]),
       Rot([vv1, vv2, vv3]))
```

- ▶ Rot macro receives and returns 3-component vector
- possible application of successive multiplication of ShapeTensor
 lowest order numerical guadrature for constant value of integrand in an element

preliminary results: Re=1



 $\label{eq:results: Re = 50} \textbf{preliminary results: } Re = 50$



summary

- weak formulation with curl-curl is used for the Navier-Stokes equations with dynamic pressure condition to perform engineering setting.
- P1/P1 approximation with adding div-div form for the velocity and element-wise grad-grad form for the pressure leads to coercive stiffness matrix for low Reynolds number, which is easily solved by additive Schwarz preconditioner in hpddm.
- shape derivative for domain variations is calculated by using the solutions of state/adjoint problems.
- gradient solver is used to find extremal of the Lagrangean using shape derivative.
- inequality conditions is supposed to keep the deformed domain within a certain box, which is solved by IPOPT.
- mmg3d is used to coarsen the mesh of the cylinder surface and Gmsh is used to regenerate the whole mesh with mesh refinement in the vicinity of the cylinder edge with updating of surface label.

on-going

- 2nd order shape derivative, which is useful for backward Euler method for the gradient flow solver and for Newton iteration
- primal-dual active set method, which can be run on distributed memory environment, to solve variational inequality

references

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