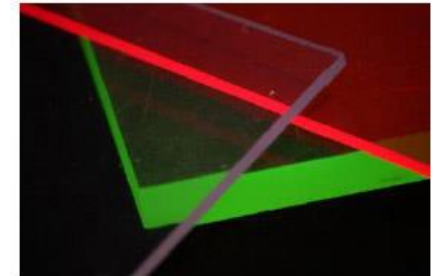
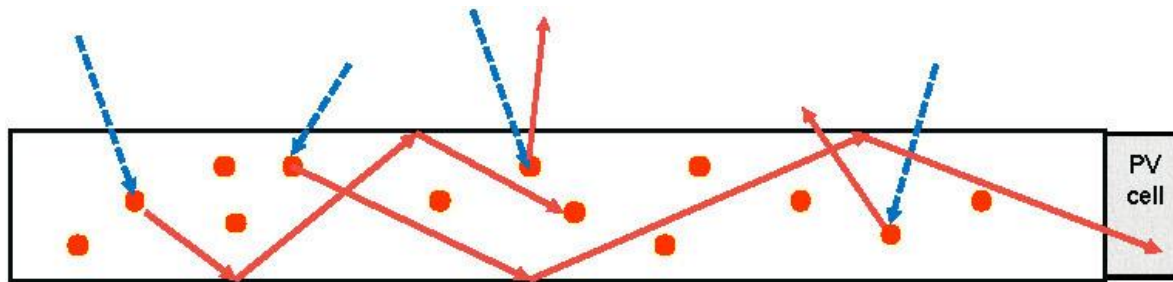
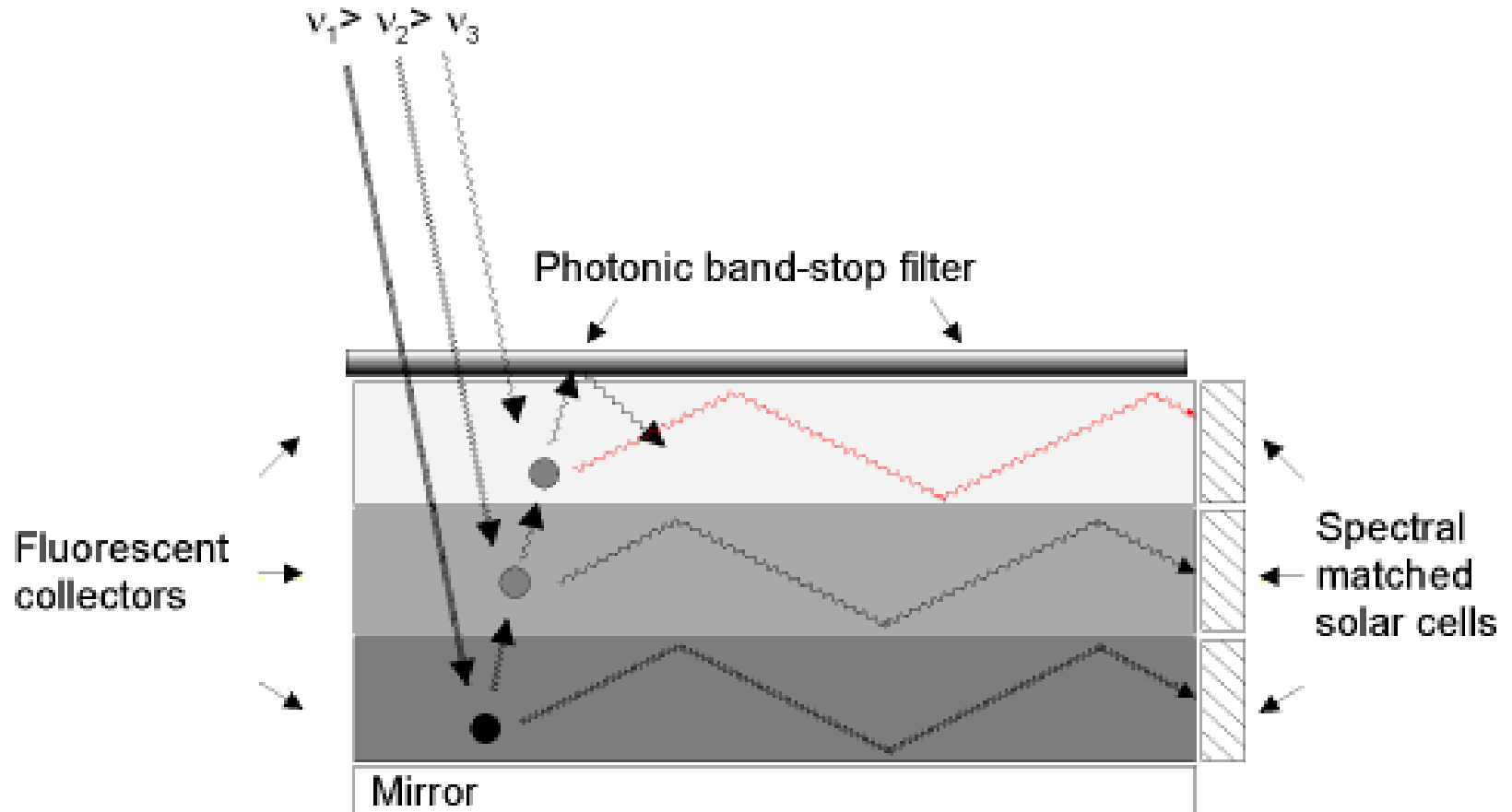


Cool concept of the week

Fluorescent concentrators



Fluorescent concentrators



Goldschmidt FhG ISE



UNIK 4450/9450 – Schedule

- 30/8 Solar cell fundamentals
- 6/9 Solar cell efficiency
- 13/9 Semiconductor theory
- 20/9 Generation
- 27/9 Recombination and lifetime
- 4/10 Silicon
- 11/10 Junctions
- 18/10 Solar cells
- 25/10 Silicon solar cells I (@IFE)
- 1/11 Silicon solar cells II
- 8/11 Light management
- 15/11 Alternative solar cells
- 29/11 Solar modules & systems
- 6/12 Q&A
- Oral exam (Week 50)

Generation



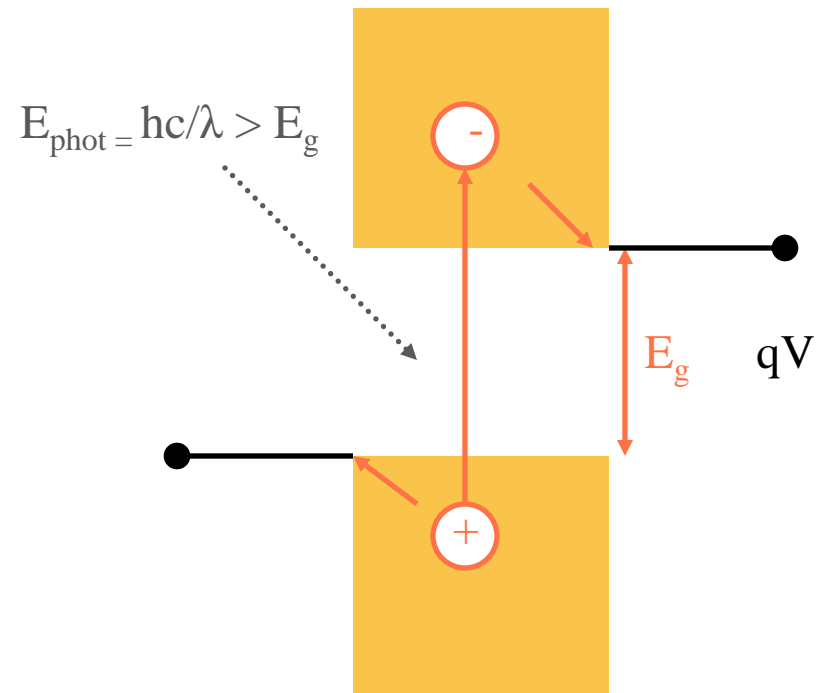
Agenda

1. Where are we?
2. Formulation of the transport problem
3. Generation and recombination – fundamentals
4. Absorption – the macroscopic view
5. Absorption – the microscopic view

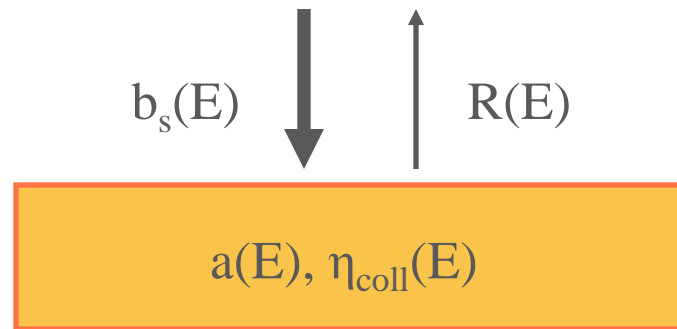
Where are we?

Main principle – solar cell operation

- The **short circuit current density** is directly proportional to the number of absorbed/generated and successfully collected charge carriers
 - Generation is the main topic of this lecture
- A junction is responsible for supplying a driving force for the photocurrent in a solar cell
- We want to be able to collect as many excited h^+ and e^- as possible while maintaining as big a voltage difference as possible between the contacts

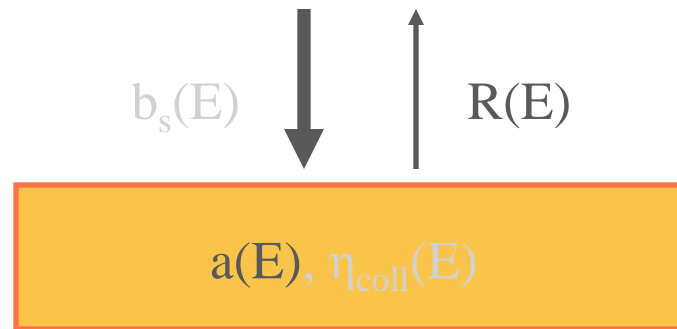


Photogenerated current



$$I_{\text{SC}}(E) = q \cdot A \cdot ([1 - R(E)] \cdot \eta_{\text{coll}}(E) \cdot a(E) \cdot b_s(E))$$

Today



$$I_{\text{SC}}(E) = q \cdot A \cdot ([1 - R(E)] \cdot \eta_{\text{coll}}(E) \cdot a(E) \cdot b_s(E))$$

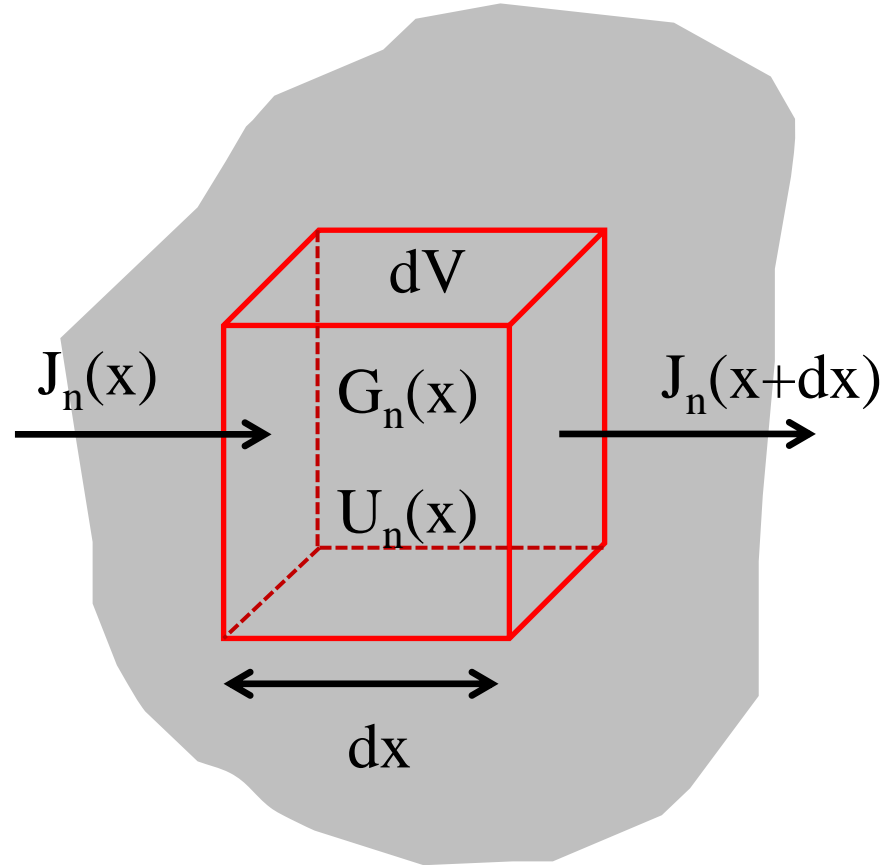
Nomenclature

- $G_n(x)$ generation rate of electrons
- $U_n(x)$ recombination rate of electrons

- $G_p(x)$ generation rate of holes
- $U_p(x)$ recombination rate of holes

Formulation of the transport problem

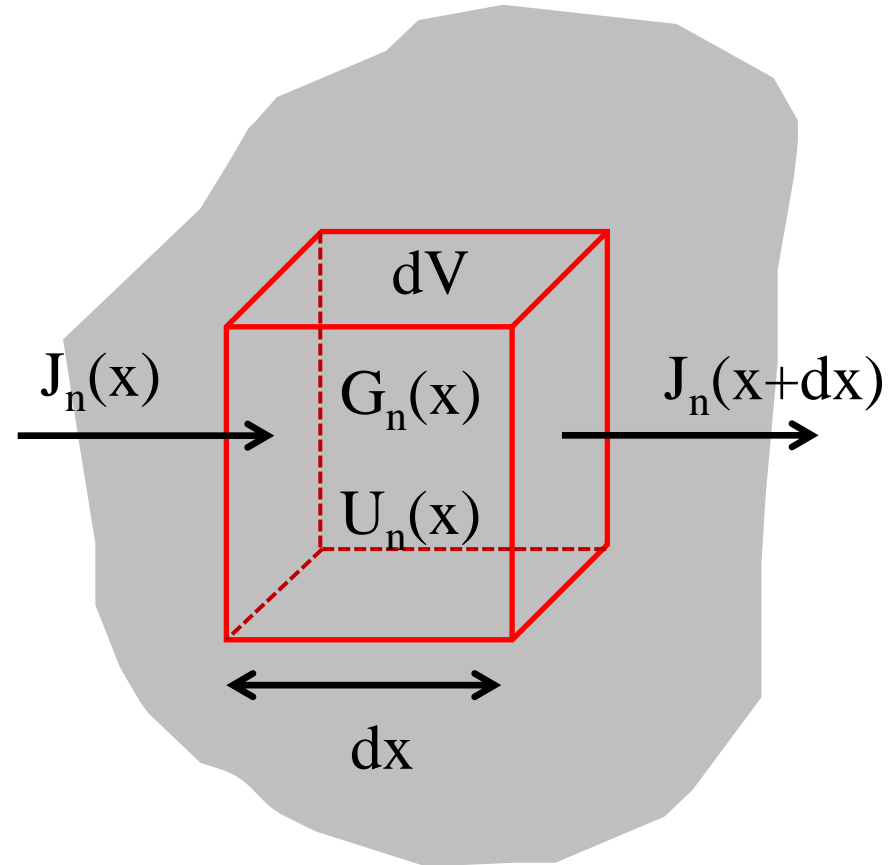
Transport in semiconductors



$$\frac{\delta n}{\delta t} = \frac{1}{q} \cdot \frac{\delta}{\delta x}(J_n(x)) + G_n(x) - U_n(x)$$

Transport in semiconductors

- Solar cell device operation:
 1. The number of charge carriers must be conserved
 2. The electrostatic potential $\phi(x)$ obeys Poisson's equation



Formulation of the transport problem

$$\delta n / \delta t = 1/q \cdot \delta / \delta x (J_n(x)) + G_n(x) - U_n(x)$$

$$\delta p / \delta t = 1/q \cdot \delta / \delta x (J_p(x)) + G_p(x) - U_p(x)$$

$$(\delta / \delta x)^2 \varphi(x) = (q / \epsilon_0 \epsilon_s) (-Q_{\text{fixed}} + n - p)$$

Steady state

$$\delta n / \delta t = 1/q \cdot \delta / \delta x (J_n(x)) + G_n(x) - U_n(x) = 0$$

$$\delta p / \delta t = 1/q \cdot \delta / \delta x (J_p(x)) + G_p(x) - U_p(x) = 0$$

$$(\delta / \delta x)^2 \varphi(x) = (q / \epsilon_0 \epsilon_s) (-Q_{\text{fixed}} + n - p)$$

Formulation of the transport problem

$$\delta n / \delta t = 1/q \cdot \delta / \delta x (J_n(x)) + G_n(x) - U_n(x)$$

$$\delta p / \delta t = 1/q \cdot \delta / \delta x (J_p(x)) + G_p(x) - U_p(x)$$

$$(\delta / \delta x)^2 \varphi(x) = (q / \epsilon_0 \epsilon_s) (-Q_{\text{fixed}} + n - p)$$

Formulation of the transport problem

$$J_n(x) = \mu_n n \cdot \delta / \delta x (E_{Fn}(x))$$

$$J_p(x) = \mu_p p \cdot \delta / \delta x (E_{Fp}(x))$$

Current densities

Total current density:

$$J_{\text{TOT}} = J_n + J_p$$

$$J_{\text{TOT}} = \mu_n n \nabla E_{Fn} + \mu_p p \nabla E_{Fp}$$

Electron current density:

$$J_n = \underbrace{+q D_n \nabla n}_{\text{DIFFUSION}} + \underbrace{\mu_n n [qE - \nabla \chi - kT \nabla \ln(N_c)]}_{\text{DRIFT}}$$

Hole current density:

$$J_p = -q D_p \nabla p + \mu_p p [qE - \nabla \chi - \nabla E_g + kT \nabla \ln(N_c)]$$

- A net current density can be obtained in many ways!

Current densities in a homojunction

Total current density:

$$J_{\text{TOT}} = J_n + J_p$$

$$J_{\text{TOT}} = \mu_n n \nabla E_{Fn} + \mu_p p \nabla E_{Fp}$$

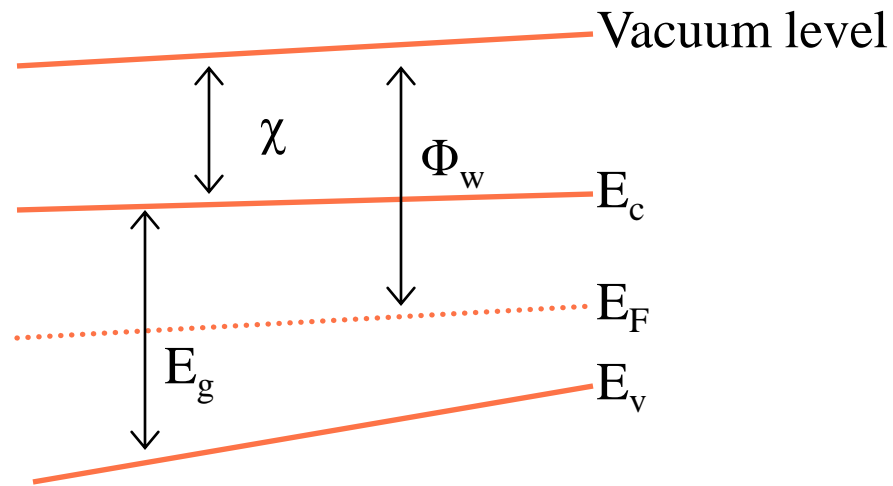
Electron current density:

$$J_n = +q D_n \nabla n + \mu_n n q E$$

Hole current density:

$$J_p = -q D_p \nabla p + \mu_p p q E$$

Driving forces



χ = electron affinity
 E_g = band gap energy
 E_F = Fermi level

E_c = conduction band energy
 E_v = valence band energy
 Φ_w = work function

Formulation of the transport problem

- The conservation equation for electrons then becomes

$$1/q \cdot \delta/\delta x(J_n(x)) + G_n(x) - U_n(x) = 0$$

- We then also recall the Einstein relation

$$\mu_n = qD_n/kT$$

Formulation of the transport problem

- We combine with the expression for $J_n(x)$

$$1/q \cdot \delta/\delta x (q D_n (\delta/\delta x) n + \mu_n n q E) + G_n(x) - U_n(x) = 0$$

- We rearrange terms and get the following equation

$$D_n (\delta^2/\delta x^2) n + \mu_n E (\delta/\delta x) n + \mu_n n (\delta/\delta x) E + G_n(x) - U_n(x) = 0$$

Formulation of the transport problem

- If E is 0 or constant we get

$$D_n(\delta^2/\delta x^2)n + \mu_n E(\delta/\delta x)n + G_n(x) - U_n(x) = 0$$

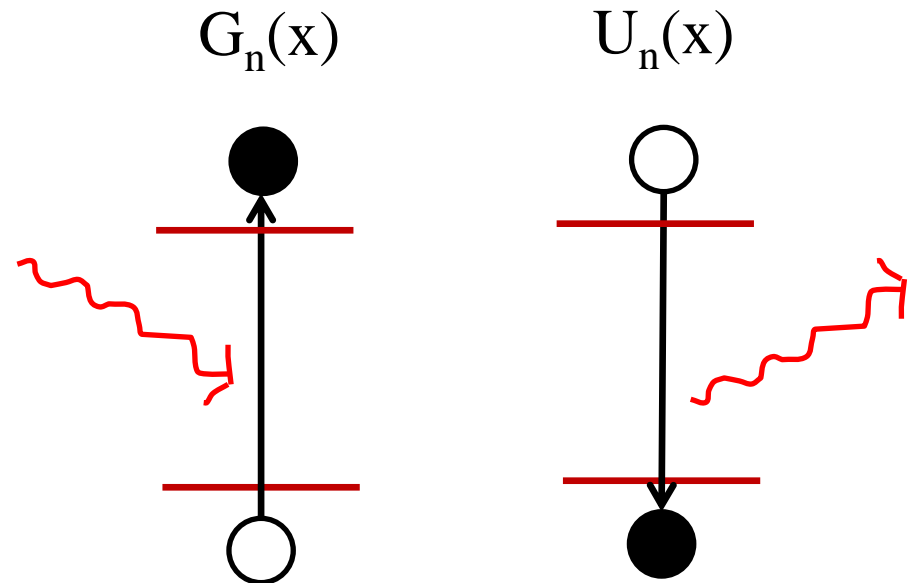
- With a final rearrangement, this finally becomes

$$(\delta^2/\delta x^2)n + (qE/kT)(\delta/\delta x)n + G_n(x)/D_n - U_n(x)/D_n = 0$$

Recombination and generation – basics

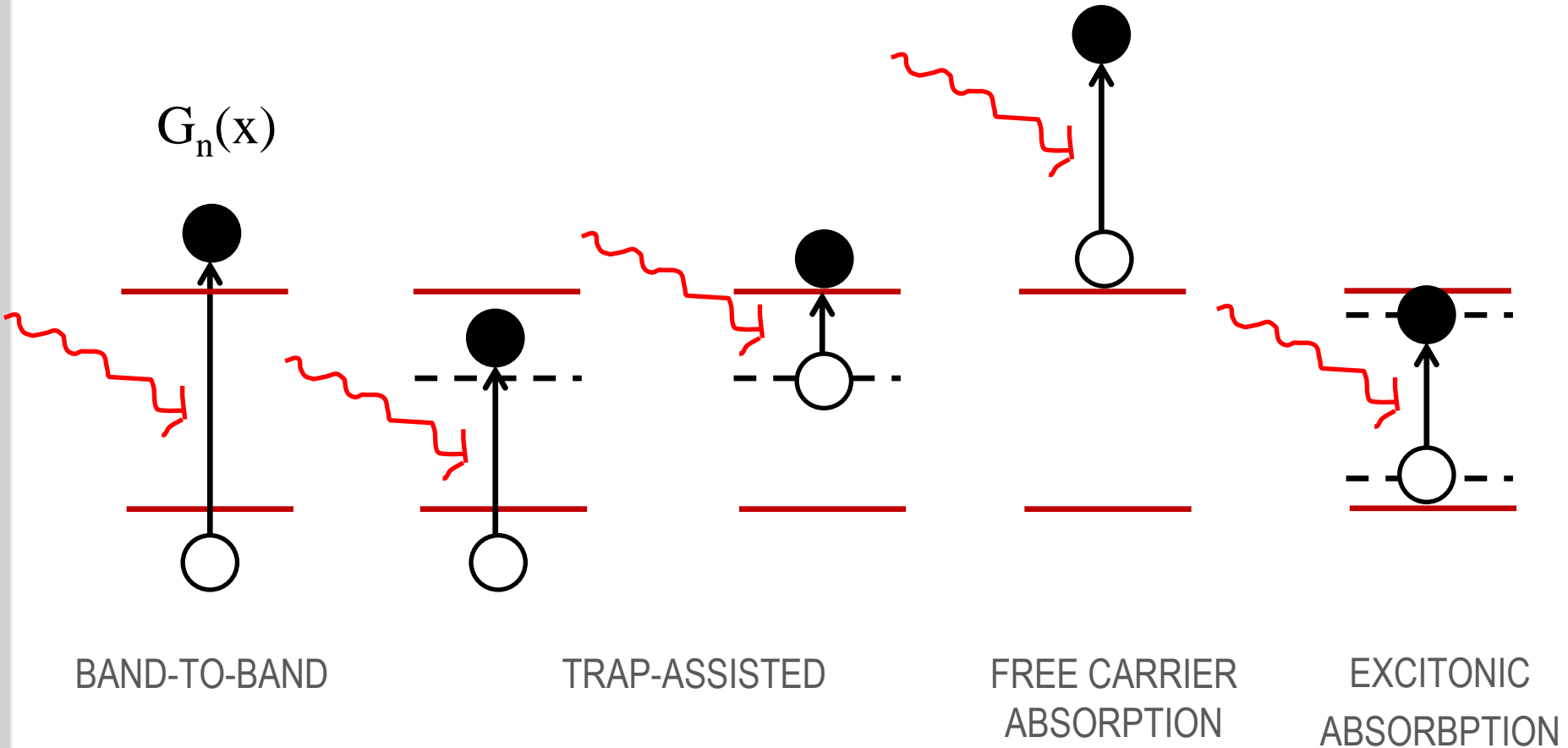
Recombination and generation

- Energy conservation
- Every generation process has an equivalent recombination process



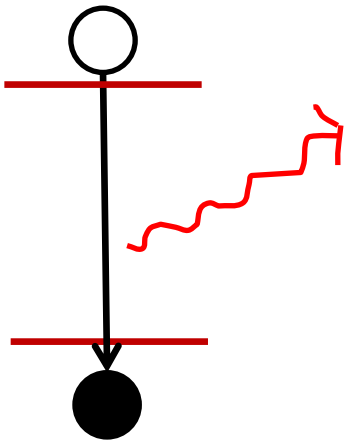
Generation processes

$G_n(x)$

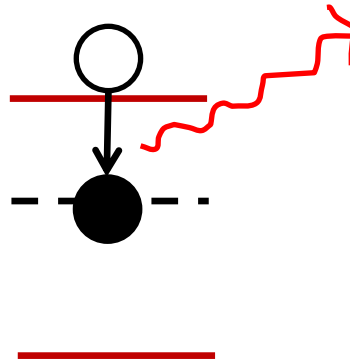


Recombination processes

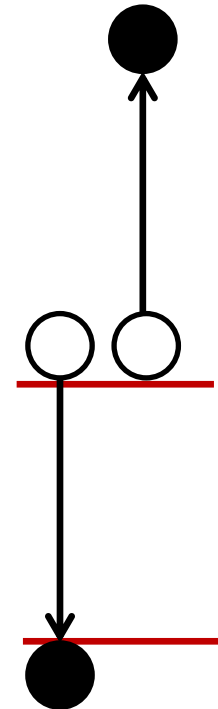
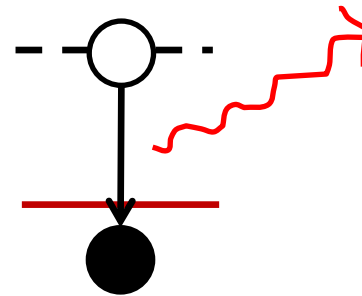
$U_n(x)$



RADIATIVE
(BAND-TO-BAND)



SHOCKLEY-REED-
HALL



AUGER

Thermal generation and recombination

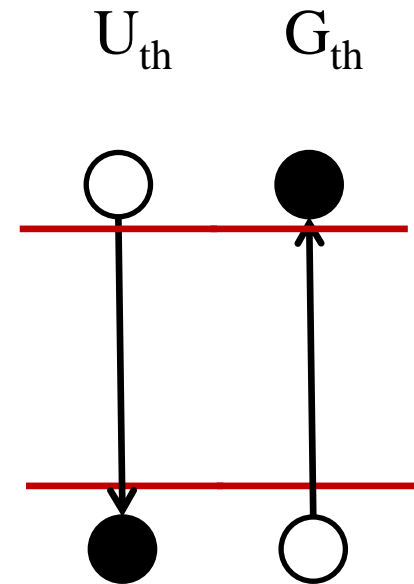
- At $T = 0$:

$$G_{n,th} = U_{n,th} = G_{p,th} = U_{p,th} = 0$$

- At $T > 0$ and in equilibrium:

$$G_{n,th} = U_{n,th}$$

$$G_{p,th} = U_{p,th}$$



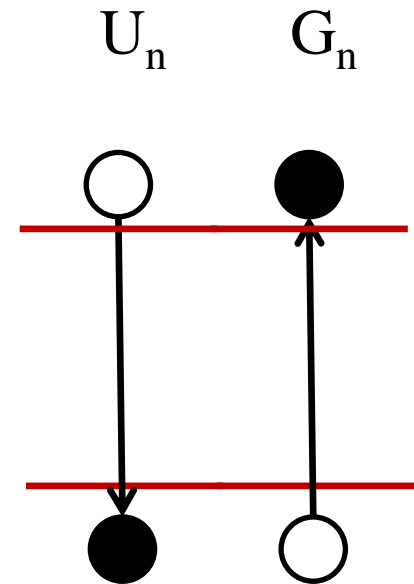
Net generation and recombination

- In the following, we will only be interested in additional generation and recombination
- We define net generation and recombination rates:

$$G_n = G_{n,tot} - G_{n,th}$$

$$U_n = U_{n,tot} - U_{n,th}$$

- ...and similar for holes

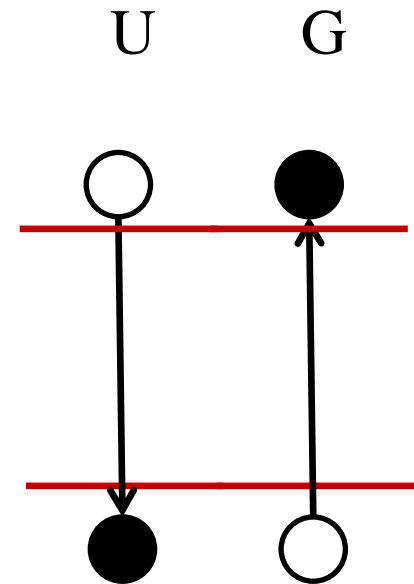


Band to band processes

- In our calculations, we will assume that band-to-band processes dominate.

$$G_n = G_p = G$$

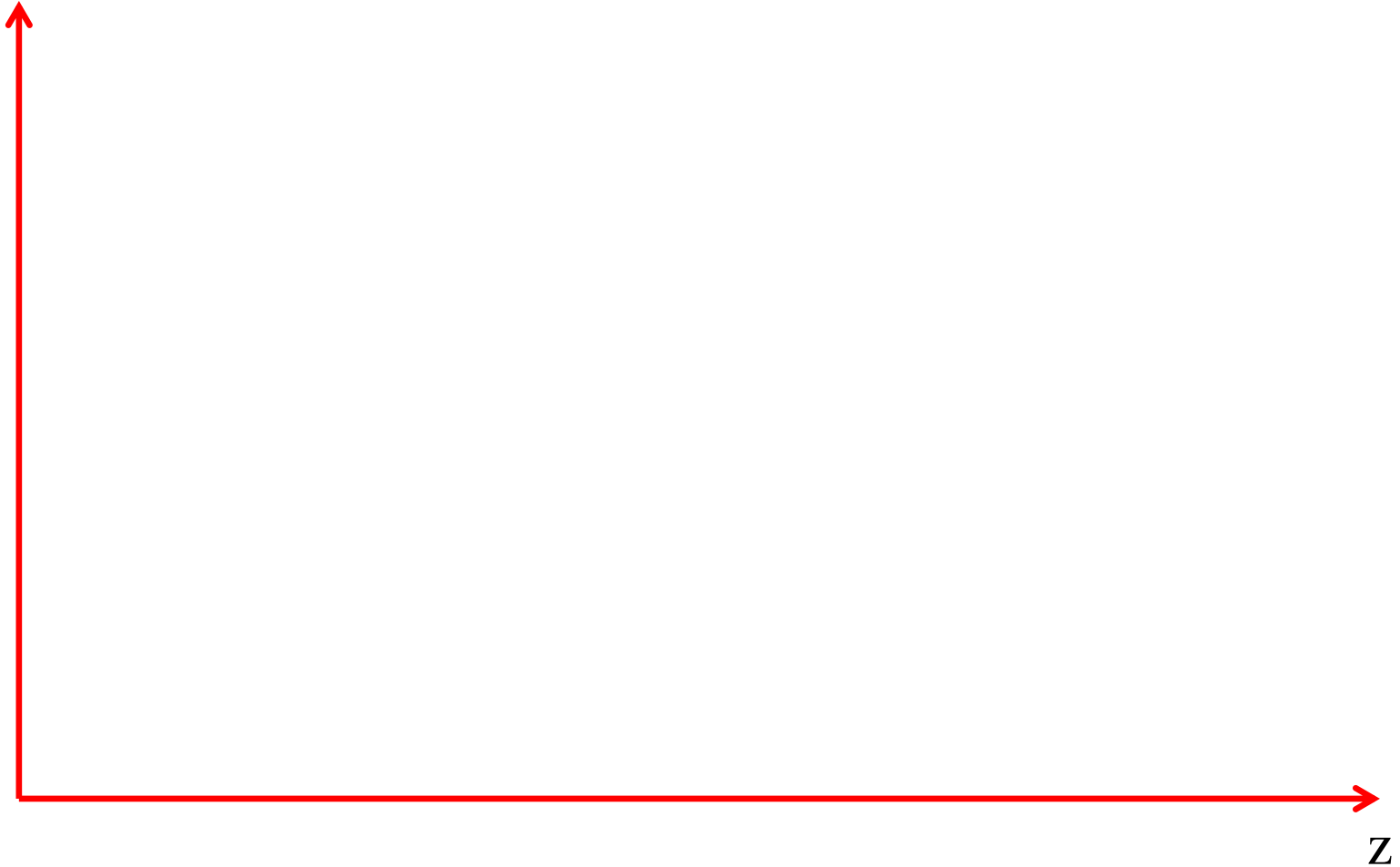
$$U_n = U_p = U$$



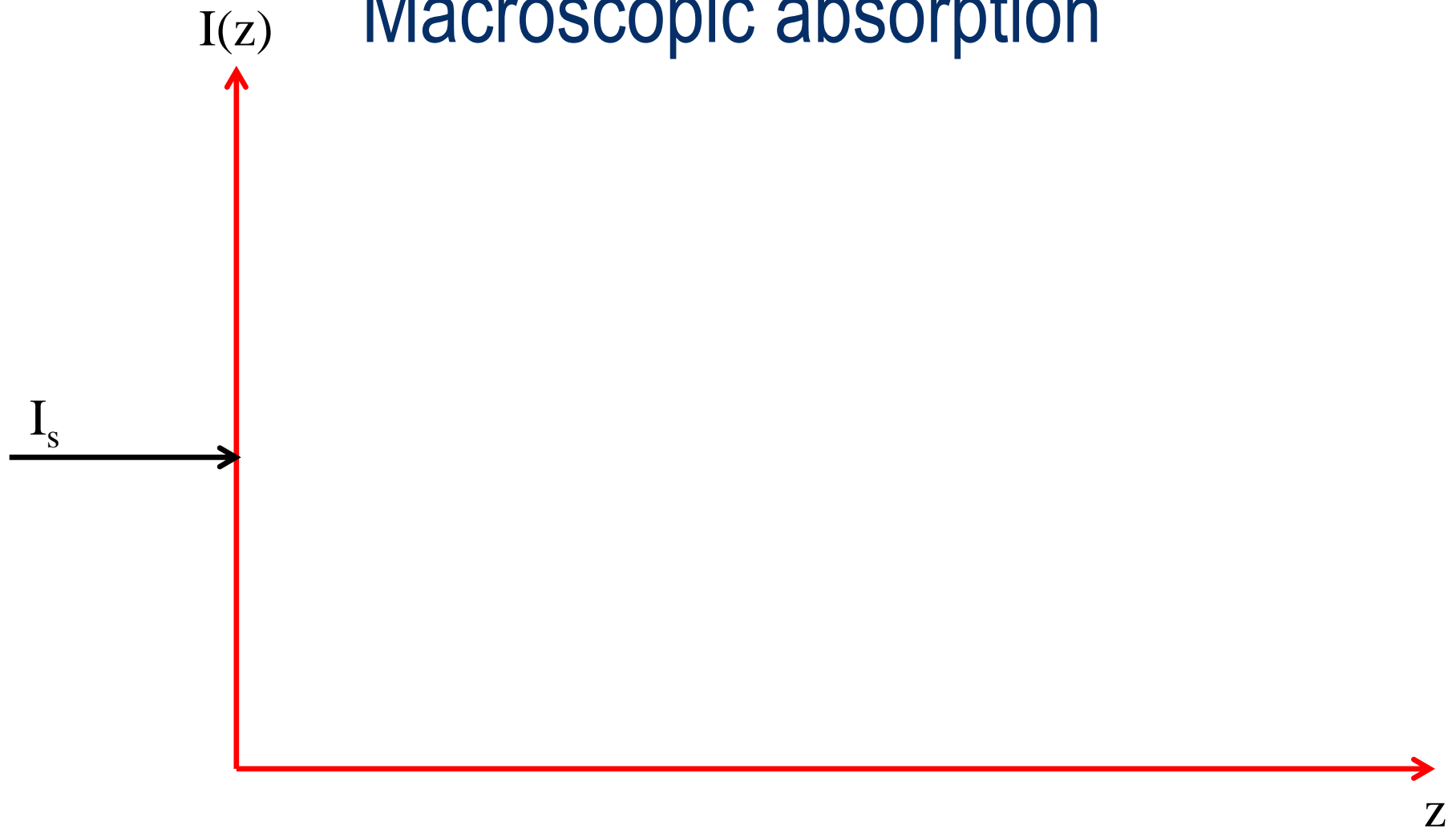
Absorption – the macroscopic view

Macroscopic absorption

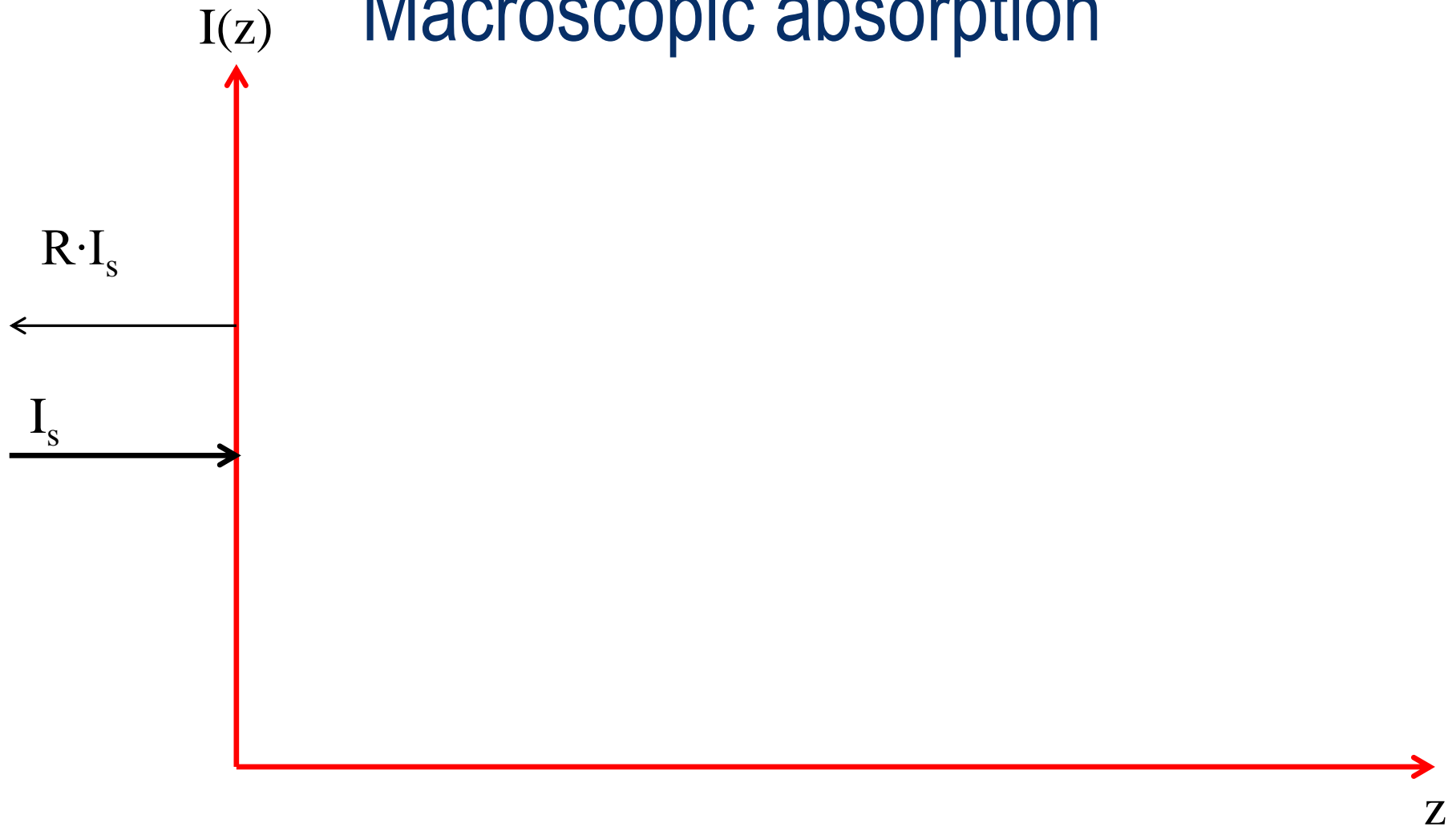
$I(z)$



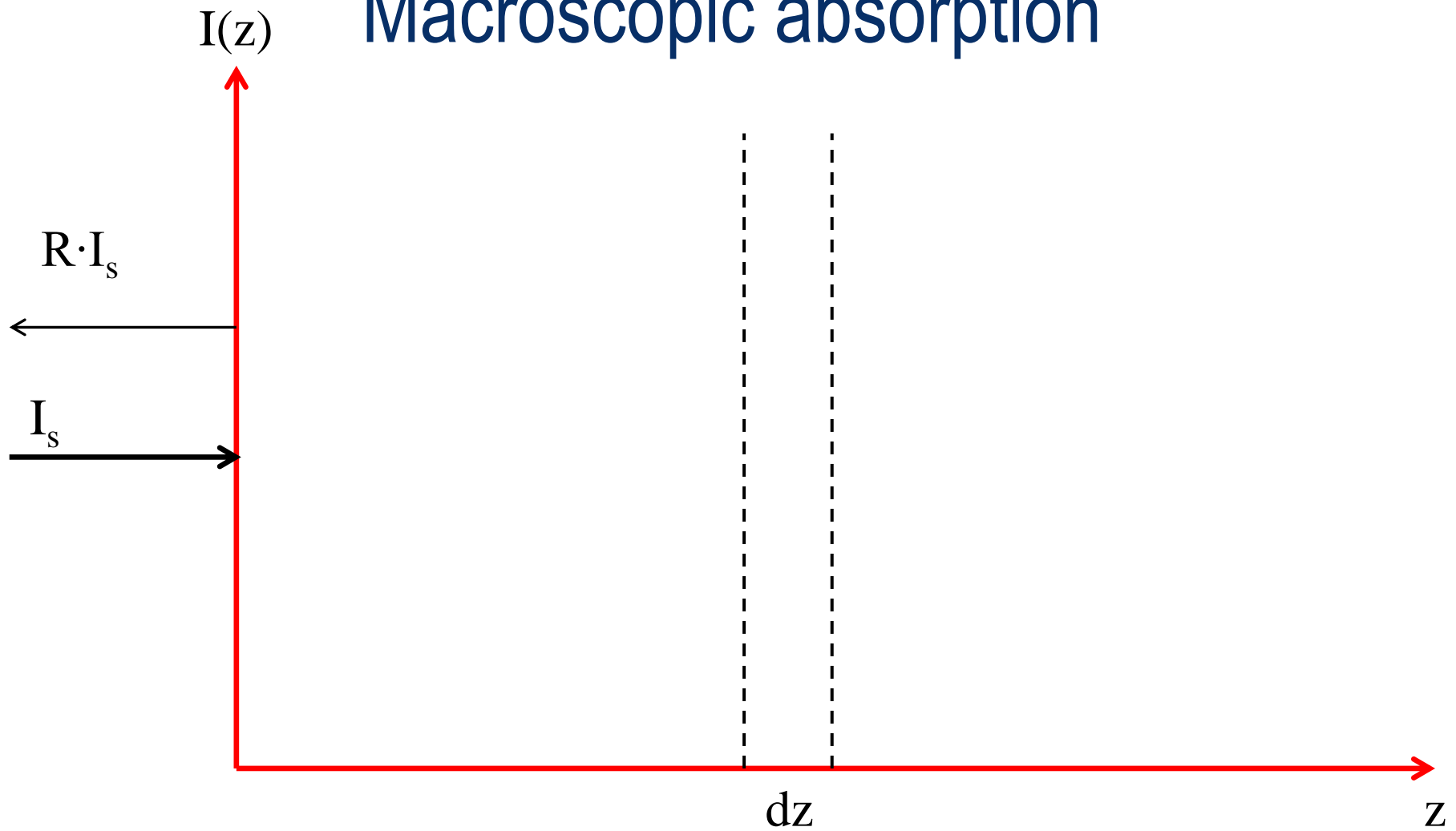
Macroscopic absorption



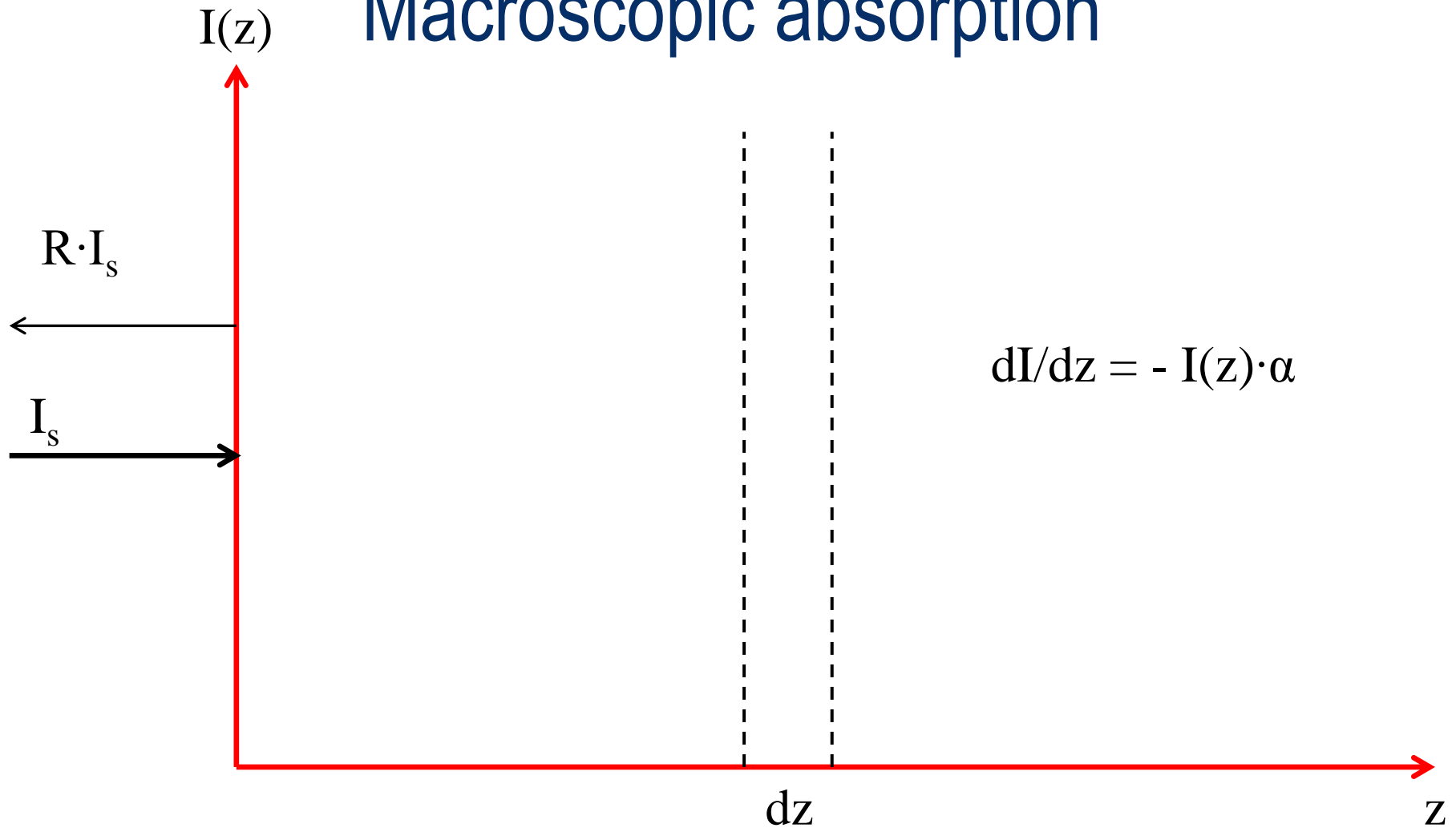
Macroscopic absorption



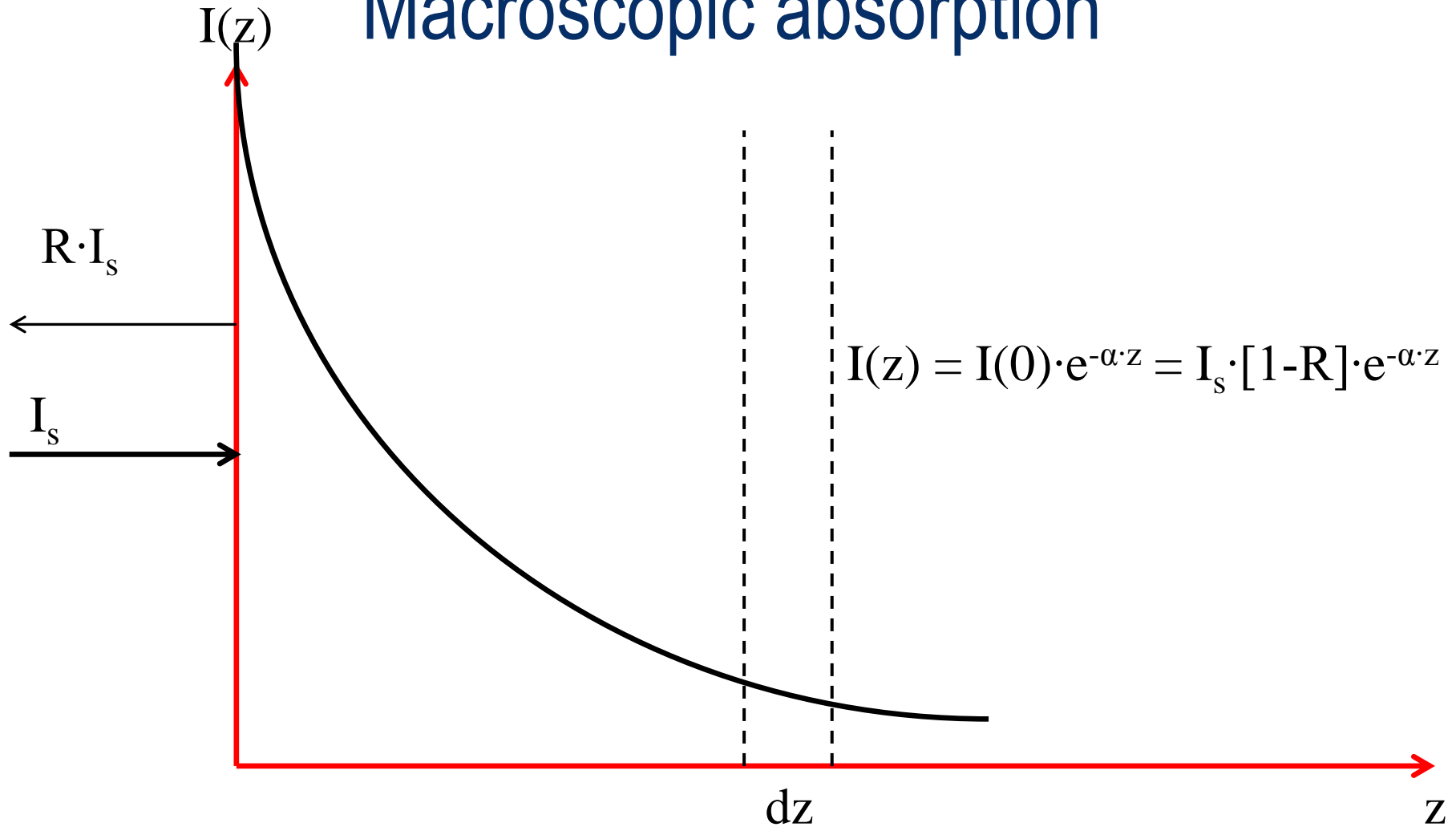
Macroscopic absorption



Macroscopic absorption



Macroscopic absorption



Energy and spatial dependence

- All important parameters are energy dependent

$$I(z,E) = I_s(E) \cdot [1 - R(E)] \cdot e^{-\alpha(E) \cdot z}$$

- Variations in $\alpha(E)$ in space in real solar cells

$$I(z,E) = I_s(E) \cdot [1 - R(E)] \cdot e^{-\int \alpha(E,z') \cdot dz'}$$

Generation current densities

- The rate of carrier generation at x per volume is given by

$$g(E,x) = b_s(E,x) \cdot \alpha[E,x]$$

- With respect to $b_s(E)$, this then becomes

$$g(E,x) = [1-R(E)] \cdot b_s(E) \cdot \alpha[E,x] \cdot e^{-\int \alpha(E,x') \cdot dx'}$$

Generation current densities

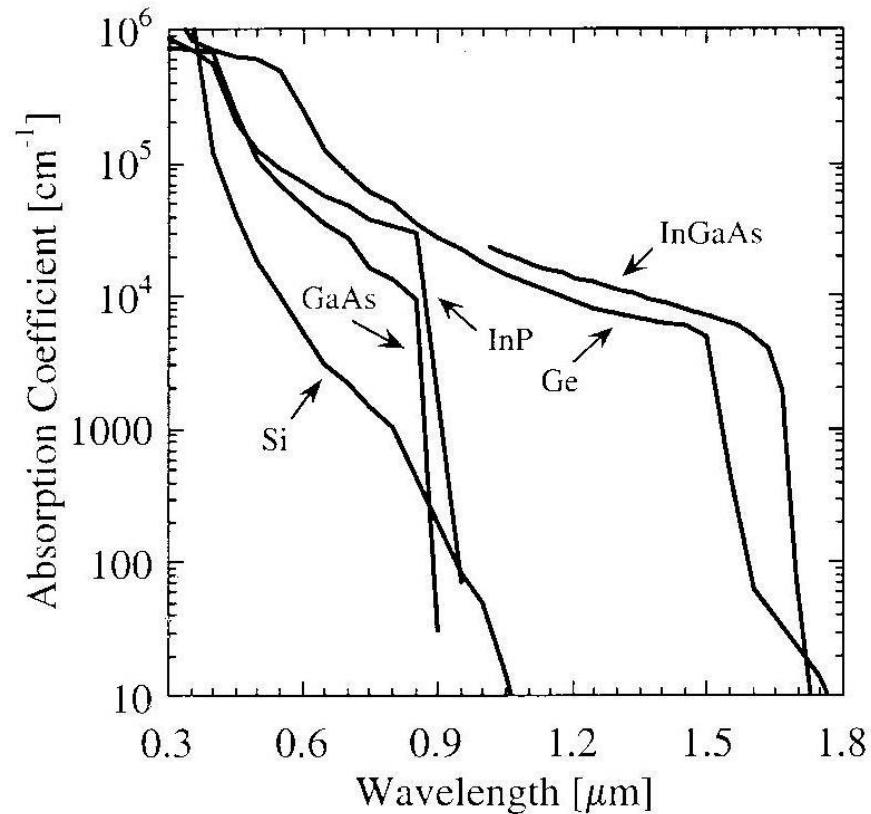
- The total generation current density is given by

$$G(x) = \int g(E, x) dE$$

- This term is then used to express the transport problem

$$D_n(\delta^2/\delta x^2)n + \mu_n E(\delta/\delta x)n + \mu_n n(\delta/\delta x)E + \mathbf{G}_n(\mathbf{x}) - U_n(x) = 0$$

Material dependence



Handbook of Optical Constants of Solids, edited by Edward D. Palik, (1985), Academic Press NY.

Clever student question:

"Why?"

Clever student question (refined):

”Where does the material dependence come from?”

Absorption – the microscopic view

Microscopic absorption

- $G(E, \mathbf{x})$ and $U(E, \mathbf{x})$ can be calculated from first principles

- The way ahead:

- Definition of initial and final states
- Determining transition probabilities
 - Fermi's golden rule
- Determining transition rates
- Determining net transition rates
- Determining total, spectral transition rates

$ i\rangle$	$ f\rangle$
$w_{i \rightarrow f}$	$w_{f \rightarrow i}$

$r_{i \rightarrow f}$	$r_{f \rightarrow i}$
-----------------------	-----------------------

$$r_{if}$$

$$r(E)$$

Our system

$|f\rangle$ ————— E_f

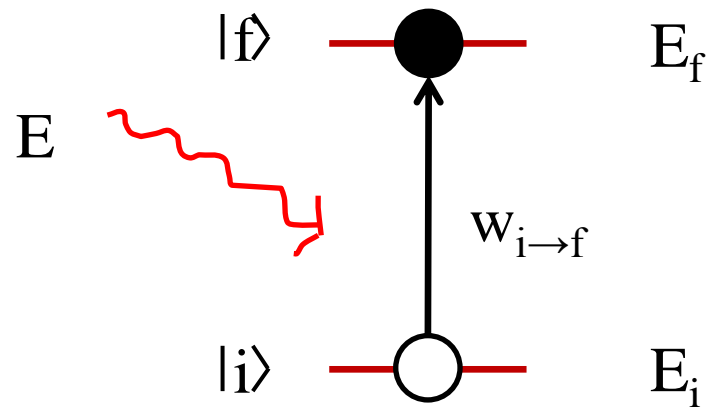
$|i\rangle$ ————— E_i

Our system

$|f\rangle$ ————— E_f

$|i\rangle$ —●— E_i

Our system



Transition probabilities

(Fermi's golden rule)

$$w_{i \rightarrow f} = \frac{2\pi}{\hbar} |\langle f | \mathbf{H} | i \rangle|^2 \cdot \delta(E_f - E_i \pm E)$$

Transition probabilities

(Fermi's golden rule)

$$w_{i \rightarrow f} = \frac{2\pi}{h} |\langle i | \mathbf{H} | f \rangle|^2 \cdot \delta(E_f - E_i \pm E)$$

- Matrix element
 - Quantifies the transition probability
 - Is what we will show how to calculate

- Energy conservation
 - The energies must match precisely

Transition probability \rightarrow Transition rate

$|f\rangle$ ————— E_f f_f

$|i\rangle$ ————— E_i f_i

Transition rates

- The transition rates take the distribution functions into account:

$$r_{i \rightarrow f} = 2\pi/\hbar |\langle i | \mathbf{H} | f \rangle|^2 \cdot \delta(E_f - E_i \pm E) \cdot f_i \cdot (1 - f_f)$$

$$r_{f \rightarrow i} = 2\pi/\hbar |\langle i | \mathbf{H} | f \rangle|^2 \cdot \delta(E_f - E_i \pm E) \cdot f_f \cdot (1 - f_i)$$

Net transition rates

- The net transition rate is the difference between $r_{i \rightarrow f}$ and $r_{f \rightarrow i}$

$$r_{if} = r_{i \rightarrow f} - r_{f \rightarrow i}$$

$$r_{if} = \frac{2\pi}{h} |\langle i | \mathbf{H} | f \rangle|^2 \cdot \delta(E_f - E_i \pm E) \cdot [f_i \cdot (1 - f_f) - f_f \cdot (1 - f_i)]$$

- Which finally becomes for a single pair of states

$$r_{if} = \frac{2\pi}{h} |\langle i | \mathbf{H} | f \rangle|^2 \cdot \delta(E_f - E_i \pm E) \cdot (f_i - f_f)$$

Total spectral transition rates

- To get to the total spectral transition rates, we must sum over all pairs of initial and final states:

$$r(E) = 2\pi/\hbar \cdot \iint |\langle i | \mathbf{H} | f \rangle|^2 \cdot \delta(E_f - E_i \pm E) \cdot (f_i - f_f) \cdot g_i(E_i) \cdot g_f(E_f) dE_i dE_f$$

Generation rate

- The total rate of band to band transitions is finally found by also taking the density of available photons into account:

$$G = \int r(E) \cdot g_{\text{photon}}(E) \cdot dE$$

Microscopic description of absorption

- The net transition rate is given by

$$r(E) = 2\pi/\hbar \cdot \iint |\langle i | \mathbf{H} | f \rangle|^2 \cdot \delta(E_f - E_i \pm E) \cdot (f_i - f_f) \cdot g_i(E_i) \cdot g_f(E_f) dE_i dE_f$$

Microscopic description of absorption

- For an electromagnetic field of strength E_0 , a polarization vector ϵ and angular frequency ω , for wavelengths longer than interatomic distances H is given by the dipole approximation as

$$H = (iqE_0/2m_0\omega) \epsilon \cdot p$$

- Here, p is the quantum mechanical momentum operator

Microscopic description of absorption

- We define $M_{cv} = |\langle v, i | \epsilon \cdot p | c, f \rangle|$
- We then get the following expression:

$$r(E) = (2\pi/\hbar) \cdot (q^2 E_0^2 \hbar^2 / 2m_0 E^2) \cdot \iint M_{cv}^2 \delta(E_c - E_v \pm E) \cdot (f_v - f_c) \cdot g_c \cdot g_v dE_c dE_v$$

Microscopic description of absorption

- The transition rate represents the net rate of photon absorption
- The electromagnetic field transfers energy to the semiconductor at a net rate

$$\delta U_E / \delta t = -E \cdot r(E)$$

- For a plane wave, this corresponds to the cycle-averaged rate of intensity loss, $\delta I / \delta x$

Microscopic description of absorption

- Since

$$\delta I / \delta x = \delta U_E / \delta t$$

- We get:

$$-\alpha I = \delta U_E / \delta t$$

Microscopic description of absorption

- Furthermore

$$I = U_e c / n_s$$

- and

$$U_E = n_s^2 E_0^2 / 8\pi$$

Microscopic description of absorption

- We finally get the following relation between the transition rate and the absorption coefficient:

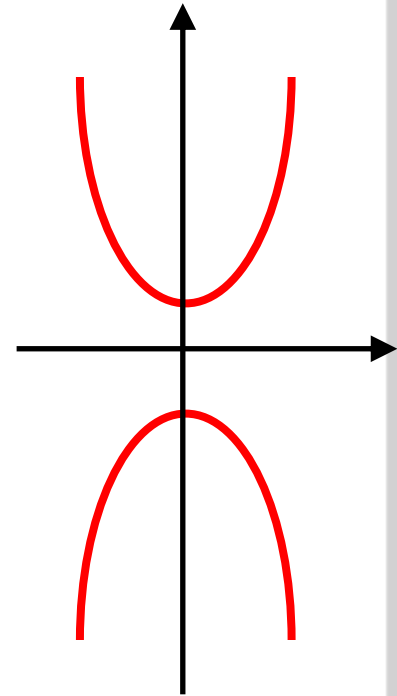
$$\alpha = (n_s E_0 / c U_E) r(E)$$

- This equation can now be used to calculate α for different materials

Direct band semiconductors

- For semiconductors with a direct band, the calculations yield:

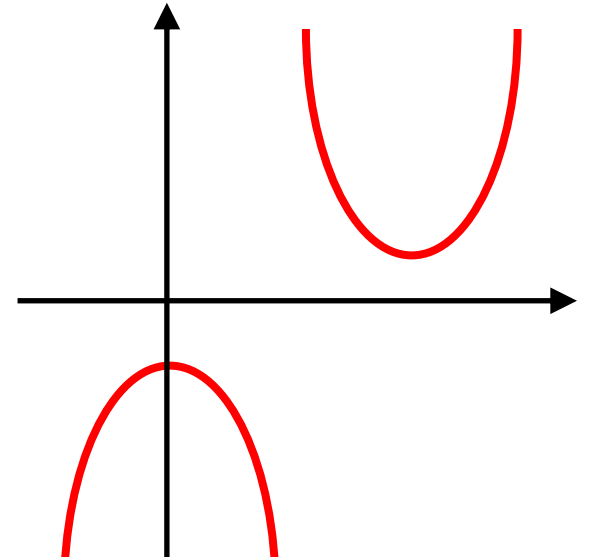
$$\alpha(E) = \alpha_0(E - E_g)^{1/2}$$



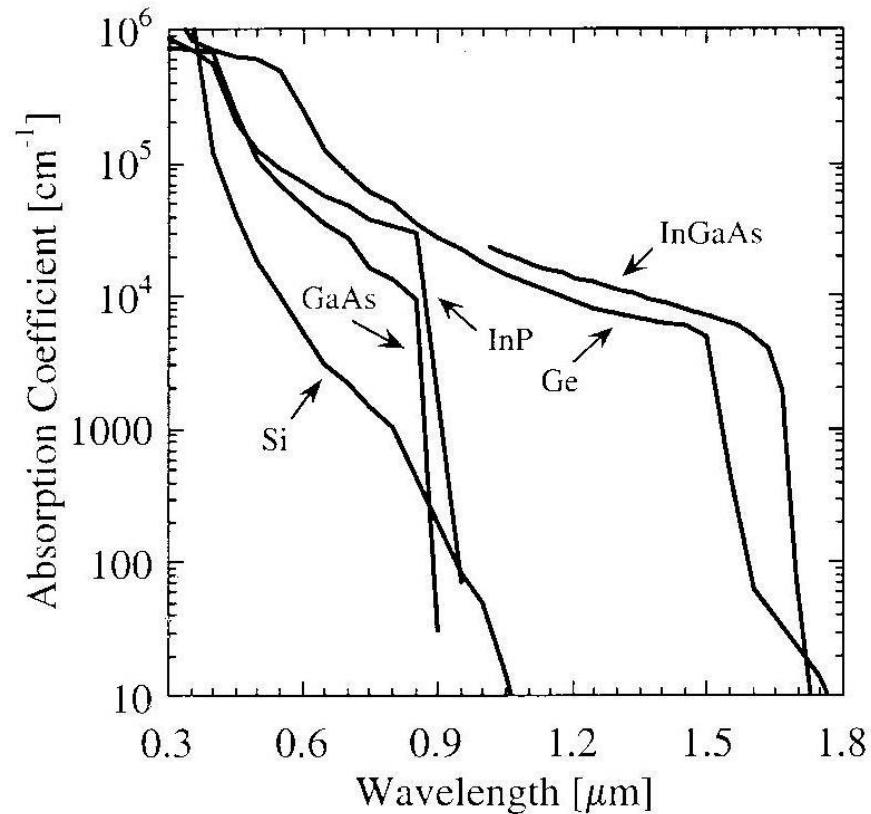
Indirect band semiconductors

- For semiconductors with an indirect band, the calculations yield:

$$\alpha(E) = \alpha_0(E - E_g)^2$$



Material dependence



Handbook of Optical Constants of Solids, edited by Edward D. Palik, (1985), Academic Press NY.