



Modélisation et Simulation des Plasmas

G. Fuhr
Guillaume.Fuhr@univ.amu.fr



Forward Differences

Forward finite-divided-difference formulas:

First Derivative

$$f'(x_i) = \frac{f(x_{i+1}) - f(x_i)}{h}$$

Error

$$O(h)$$

$$f'(x_i) = \frac{-f(x_{i+2}) + 4f(x_{i+1}) - 3f(x_i)}{2h}$$

$$O(h^2)$$

Second Derivative

$$f''(x_i) = \frac{f(x_{i+2}) - 2f(x_{i+1}) + f(x_i)}{h^2}$$

$$O(h)$$

$$f''(x_i) = \frac{-f(x_{i+3}) + 4f(x_{i+2}) - 5f(x_{i+1}) + 2f(x_i)}{h^2}$$

$$O(h^2)$$

Third Derivative

$$f'''(x_i) = \frac{f(x_{i+3}) - 3f(x_{i+2}) + 3f(x_{i+1}) - f(x_i)}{h^3}$$

$$O(h)$$

$$f'''(x_i) = \frac{-3f(x_{i+4}) + 14f(x_{i+3}) - 24f(x_{i+2}) + 18f(x_{i+1}) - 5f(x_i)}{2h^3}$$

$$O(h^2)$$



Backward Differences

Backward finite-divided-difference formulas:

First Derivative

$$f'(x_i) = \frac{f(x_i) - f(x_{i-1})}{h}$$

Error

$$O(h)$$

$$f'(x_i) = \frac{3f(x_i) - 4f(x_{i-1}) + f(x_{i-2}))}{2h}$$

$$O(h^2)$$

Second Derivative

$$f''(x_i) = \frac{f(x_i) - 2f(x_{i-1}) + f(x_{i-2}))}{h^2}$$

$$O(h)$$

$$f''(x_i) = \frac{2f(x_i) - 5f(x_{i-1}) + 4f(x_{i-2}) - f(x_{i-3}))}{h^2}$$

$$O(h^2)$$

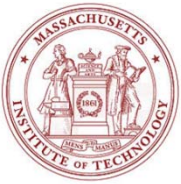
Third Derivative

$$f'''(x_i) = \frac{f(x_i) - 3f(x_{i-1}) + 3f(x_{i-2}) - f(x_{i-3}))}{h^3}$$

$$O(h)$$

$$f'''(x_i) = \frac{5f(x_i) - 18f(x_{i-1}) + 24f(x_{i-2}) - 14f(x_{i-3}) + 3f(x_{i-4}))}{2h^3}$$

$$O(h^2)$$



Centered Differences

Centered finite-divided-difference formulas:

First Derivative

$$f'(x_i) = \frac{f(x_{i+1}) - f(x_{i-1}))}{2h}$$

Error

$O(h)$

$$f'(x_i) = \frac{-f(x_{i+2}) + 8f(x_{i+1}) - 8f(x_{i-1}) + f(x_{i-2}))}{12h}$$

$O(h^2)$

Second Derivative

$$f''(x_i) = \frac{f(x_{i+1}) - 2f(x_i) + f(x_{i-1}))}{h^2}$$

$O(h)$

$$f''(x_i) = \frac{-f(x_{i+2}) + 16f(x_{i+1}) - 30f(x_i) + 16f(x_{i-1}) - f(x_{i-2}))}{12h^2}$$

$O(h^2)$

Third Derivative

$$f'''(x_i) = \frac{f(x_{i+2}) - 2f(x_{i+1}) + 2f(x_{i-1}) - f(x_{i-2}))}{2h^3}$$

$O(h)$

$$f'''(x_i) = \frac{-f(x_{i+3}) + 8f(x_{i+2}) - 13f(x_{i+1}) + 13f(x_{i-1}) - 8f(x_{i-2}) + f(x_{i-3}))}{8h^3}$$

$O(h^2)$



Example: Taylor Table for the Adams-Moulton 3-steps (4 time-nodes) Method

Denoting $h \equiv \Delta t$, $\phi \equiv u$, $\frac{du}{dt} = u' = f(t, u)$ and $u'_n = f(t_n, u^n)$, one obtains for $K = 2$:

$$u^{n+1} - u^n = \sum_{k=-K}^1 \beta_k f(t_{n+k}, u^{n+k}) \Delta t = h \left[\beta_1 f(t_{n+1}, u^{n+1}) + \beta_0 f(t_n, u^n) + \beta_{-1} f(t_{n-1}, u^{n-1}) + \beta_{-2} f(t_{n-2}, u^{n-2}) \right]$$

Taylor Table:

- The first row (Taylor series) + the last 5 rows (Taylor series for each term) must sum to zero
- This can be satisfied up to the 5th column (4th order term)
- Hence, the AM method with 4-time levels is 4th order accurate

	u_n	$h \cdot u'_n$	$h^2 \cdot u''_n$	$h^3 \cdot u'''_n$	$h^4 \cdot u''''_n$
u_{n+1}	1	1	$\frac{1}{2}$	$\frac{1}{6}$	$\frac{1}{24}$
$-u_n$	-1				
$-h\beta_1 u'_{n+1}$		$-\beta_1$	$-\beta_1$	$-\beta_1 \frac{1}{2}$	$-\beta_1 \frac{1}{6}$
$-h\beta_0 u'_n$		$-\beta_0$			
$-h\beta_{-1} u'_{n-1}$		$-\beta_{-1}$	β_{-1}	$-\beta_{-1} \frac{1}{2}$	$\beta_{-1} \frac{1}{6}$
$-h\beta_{-2} u'_{n-2}$		$-(-2)^0 \beta_{-2}$	$-(-2)^1 \beta_{-2}$	$-(-2)^2 \beta_{-2} \frac{1}{2}$	$-(-2)^3 \beta_{-2} \frac{1}{6}$

solving for the β_k 's $\Rightarrow \beta_1 = 9/24$, $\beta_0 = 19/24$, $\beta_{-1} = -5/24$ and $\beta_{-2} = 1/24$



Example of Adams Methods for Time-Integration

Explicit Methods. (Adams-Bashforth, with ABn meaning nth order AB)

$u_{n+1} = u_n + hu'_n$	Euler
$u_{n+1} = u_{n-1} + 2hu'_n$	Leapfrog
$u_{n+1} = u_n + \frac{1}{2}h[3u'_n - u'_{n-1}]$	AB2
$u_{n+1} = u_n + \frac{h}{12}[23u'_n - 16u'_{n-1} + 5u'_{n-2}]$	AB3

Implicit Methods. (Adams-Moulton, with AMn meaning nth order AM)

$u_{n+1} = u_n + hu'_{n+1}$	Implicit Euler
$u_{n+1} = u_n + \frac{1}{2}h[u'_n + u'_{n+1}]$	Trapezoidal (AM2)
$u_{n+1} = \frac{1}{3}[4u_n - u_{n-1} + 2hu'_{n+1}]$	2nd-order Backward
$u_{n+1} = u_n + \frac{h}{12}[5u'_{n+1} + 8u'_n - u'_{n-1}]$	AM3



Practical Time-Integration Methods for CFD

- High-resolution CFD requires large discrete state vector sizes to store the spatial information
- This means that up to two times (one on each side of the current time step) have often been utilized (3 time-nodes): $u^{n+1} - u^n = h \left[\beta_1 f(t_{n+1}, u^{n+1}) + \beta_0 f(t_n, u^n) + \beta_{-1} f(t_{n-1}, u^{n-1}) \right]$
- Rewriting this equations in a way such that differences wrt. the Euler's method are easily seen, one obtains ($\theta = 0$ for explicit schemes):

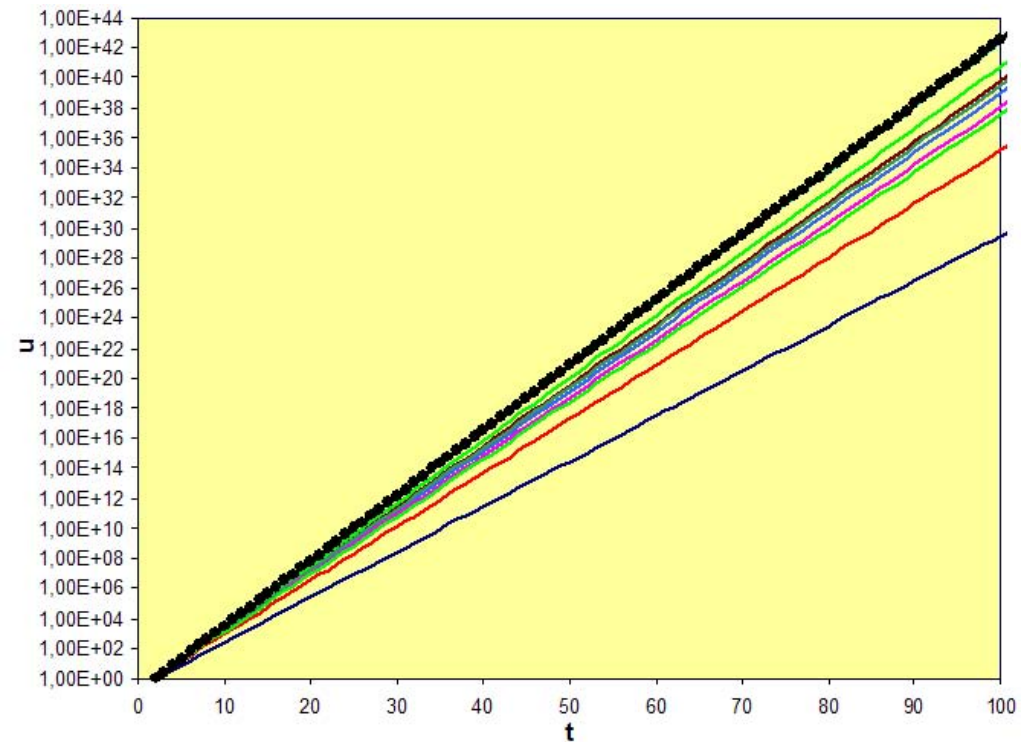
$$(1 + \xi) u^{n+1} = \left[(1 + 2\xi) u^n - \xi u^{n-1} \right] + h \left[\theta f(t_{n+1}, u^{n+1}) + (1 - \theta + \varphi) f(t_n, u^n) - \varphi f(t_{n-1}, u^{n-1}) \right]$$

θ	ξ	φ	Method	Order
0	0	0	Euler	1
1	0	0	Implicit Euler	1
1/2	0	0	Trapezoidal or AM2	2
1	1/2	0	2nd-order Backward	2
3/4	0	-1/4	Adams type	2
1/3	-1/2	-1/3	Lees	2
1/2	-1/2	-1/2	Two-step trapezoidal	2
5/9	-1/6	-2/9	A-contractive	2
0	-1/2	0	Leapfrog	2
0	0	1/2	AB2	2
0	-5/6	-1/3	Most accurate explicit	3
1/3	-1/6	0	Third-order implicit	3
5/12	0	1/12	AM3	3
1/6	-1/2	-1/6	Milne	4



- $$\frac{du}{dt} = u \text{ avec } u(t=0) = 1$$

- ☐ *Euler,*
- ☐ *Adams-Bashforth 1,*
- ☐ *Milne 1,*
- ☐ *Adams-Bashforth 2,*
- ☐ *prédicteur-correcteur et trapèze,*
- ☐ *Milne 3,*
- ☐ *Adams-Bashforth 3,*
- ☐ *Milne 3,*
- ☐ *Runge-Kutta d'ordre 4*



Method	Order	Formula
Forward	1	$\phi^{n+1} = \phi^n + hF(\phi^n)$
Backward	1	$\phi^{n+1} = \phi^n + hF(\phi^{n+1})$
Asselin Leapfrog	1	$\phi^{n+1} = \overline{\phi^{n-1}} + 2hF(\phi^n)$ $\overline{\phi^n} = \phi^n + \gamma(\overline{\phi^{n-1}} - 2\phi^n + \phi^{n+1})$
Leapfrog	2	$\phi^{n+1} = \phi^{n-1} + 2hF(\phi^n)$
Adams– Bashforth	2	$\phi^{n+1} = \phi^n + \frac{h}{2} [3F(\phi^n) - F(\phi^{n-1})]$
Trapezoidal	2	$\phi^{n+1} = \phi^n + \frac{h}{2} [F(\phi^{n+1}) + F(\phi^n)]$
Runge–Kutta	2	$q_1 = hF(\phi^n), \quad \phi_1 = \phi^n + q_1$ $q_2 = hF(\phi_1) - q_1, \quad \phi^{n+1} = \phi_1 + q_2/2$
Magazenkov	2	$\phi^n = \phi^{n-2} + 2hF(\phi^{n-1})$ $\phi^{n+1} = \phi^n + \frac{h}{2} [3F(\phi^n) - F(\phi^{n-1})]$
Leapfrog– Trapezoidal	2	$\phi_1 = \phi^{n-1} + 2hF(\phi^n)$ $\phi^{n+1} = \phi^n + \frac{h}{2} [F(\phi_1) + F(\phi^n)]$
Adams– Bashforth	3	$\phi^{n+1} = \phi^n + \frac{h}{12} [23F(\phi^n) - 16F(\phi^{n-1}) + 5F(\phi^{n-2})]$
Adams– Moulton	3	$\phi^{n+1} = \phi^n + \frac{h}{12} [5F(\phi^{n+1}) + 8F(\phi^n) - F(\phi^{n-1})]$
ABM Predictor– Corrector	3	$\phi_1 = \phi^n + \frac{h}{2} [3F(\phi^n) - F(\phi^{n-1})]$ $\phi^{n+1} = \phi^n + \frac{h}{12} [5F(\phi_1) + 8F(\phi^n) - F(\phi^{n-1})]$
Runge–Kutta	3	$q_1 = hF(\phi^n), \quad \phi_1 = \phi^n + q_1/3$ $q_2 = hF(\phi_1) - 5q_1/9, \quad \phi_2 = \phi_1 + 15q_2/16$ $q_3 = hF(\phi_2) - 153q_2/128, \quad \phi^{n+1} = \phi_2 + 8q_3/15$
Runge–Kutta	4	$q_1 = hF(\phi^n), \quad q_2 = hF(\phi^n + q_1/2)$ $q_3 = hF(\phi^n + q_2/2), \quad q_4 = hF(\phi^n + q_3)$ $\phi^{n+1} = \phi^n + (q_1 + 2q_2 + 2q_3 + q_4)/6$

TABLE 2.1. Summary of methods for the solution of ordinary differential equations. The second- and third-order Runge–Kutta methods are low-storage variants; $h = \Delta t$.

<i>Method</i>	<i>Storage Factor</i>	<i>Efficiency Factor</i>	<i>Amplification Factor</i>	<i>Phase Error</i>	<i>Max s</i>
Forward	2	0	$1 + \frac{s^2}{2}$	$1 - \frac{s^2}{3}$	0
Backward	*	∞	$1 - \frac{s^2}{2}$	$1 - \frac{s^2}{3}$	∞
Asselin Leapfrog	3	< 1	$1 - \frac{\gamma s^2}{2(1-\gamma)}$	$1 + \frac{(1+2\gamma)s^2}{6(1-\gamma)}$	< 1
Leapfrog	2	1	1	$1 + \frac{s^2}{6}$	1
Adams-Bashforth-2	3	0	$1 + \frac{s^4}{4}$	$1 + \frac{5}{12}s^2$	0
Trapezoidal	*	∞	1	$1 - \frac{s^2}{12}$	∞
Runge-Kutta-2	2	0	$1 + \frac{s^4}{8}$	$1 + \frac{s^2}{6}$	0
Magazenkov	3	0.67	$1 - \frac{s^4}{4}$	$1 + \frac{s^2}{6}$	0.67
Leapfrog-Trapezoidal	3	0.71	$1 - \frac{s^4}{4}$	$1 - \frac{s^2}{12}$	1.41
Adams-Bashforth-3	4	0.72	$1 - \frac{3}{8}s^4$	$1 + \frac{289}{720}s^4$	0.72
Adams-Moulton-3	*	0	$1 + \frac{s^4}{24}$	$1 - \frac{11}{720}s^4$	0
ABM Predictor-Corrector-3	4	0.60	$1 - \frac{19}{144}s^4$	$1 + \frac{1243}{8640}s^4$	1.20
Runge-Kutta-3	2	0.58	$1 - \frac{s^4}{24}$	$1 + \frac{s^4}{30}$	1.73
Runge-Kutta-4	4 [†]	0.70	$1 - \frac{s^6}{144}$	$1 - \frac{s^4}{120}$	2.82

[†] A storage factor of 3 may be achieved following the algorithm of Blum (1962).

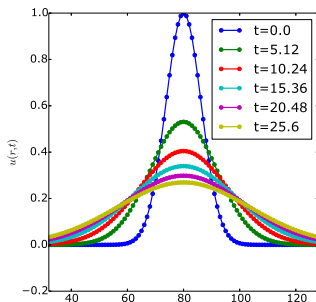
TABLE 2.2. Characteristics of the schemes listed in Table 2.1. The amplification factor and relative phase change are for well-resolved solutions to the oscillation equation, and $s = \kappa \Delta t$. "Max s " is the maximum value of $\kappa \Delta t$ for which the solution is nonamplifying. The storage and efficiency factors are defined in the text. No storage factor is given for implicit schemes.

Diffusion and C.F.L.

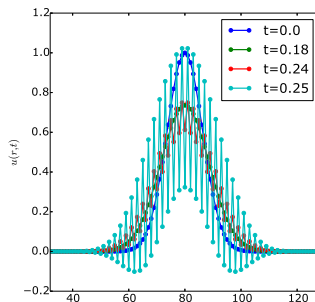
Test Case

$$\begin{aligned}\partial_t u(r, t) &= D \partial_r^2 u(r, t), \\ u(r, t = 0) &= e^{-(r-r_0)^2/\sigma^2}\end{aligned}\tag{10}$$

$n_{CFL} = 0.1$

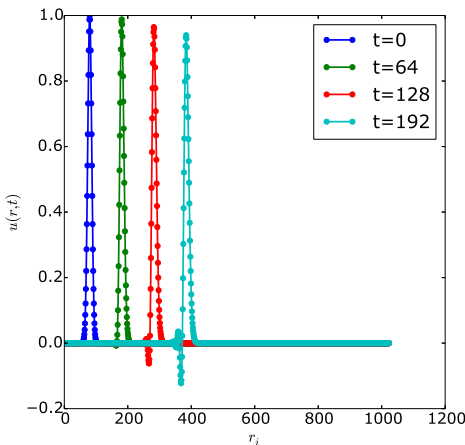


$n_{CFL} = 1.0$



Advection simulation

- Numerical results using $u_0(r) = \exp(-r^2/\sigma^2)$
- C.F.L : $n_{CFL} = 0.4$

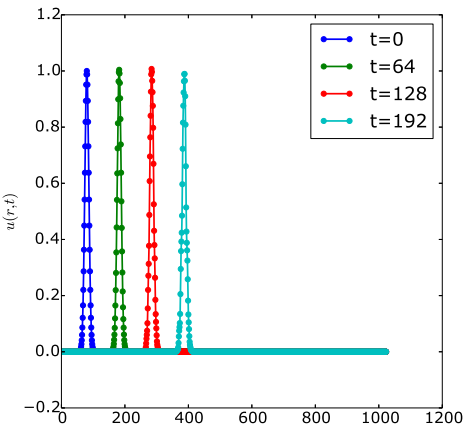


- system not conservative (diffusion)
- numerical oscillations generated
- phenomena disappears when $\Delta r \rightarrow 0$

Advection simulation

Ref. : C. S. Shu

- Numerical results using $u_0(r) = \exp(-r^2/\sigma^2)$
- C.F.L : $n_{CFL} = 0.4$



- 5th order WENO scheme used here
- smooth interpolation locally around discontinuities

$$\partial_t u(r, t) = -\frac{1}{\Delta r} \left(\hat{f}_{r+1/2} - \hat{f}_{r-1/2} \right)$$

$$\hat{f}_{r+1/2} = \hat{f}(u_{r-2}, \dots, u_{r+2})$$

- Pro: parallel implementation