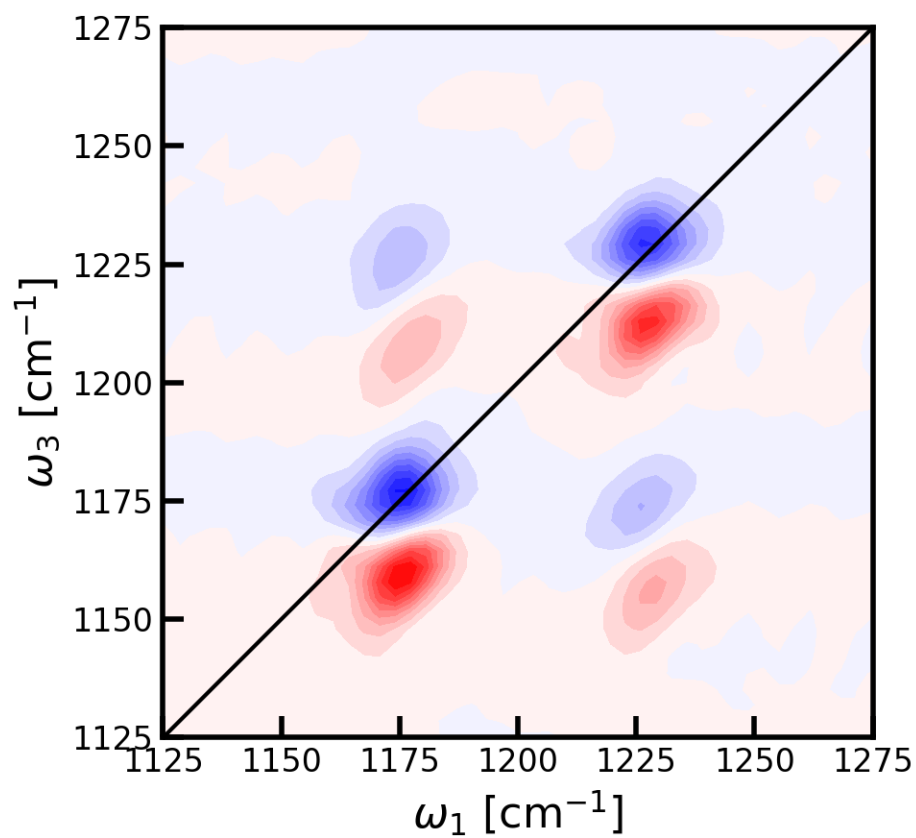


NISE MANUAL

VERSION 3.4

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Contents

Contents	1
1 Introduction	3
2 Installation and options	5
2.1 Building	5
2.2 Installation Trouble Shooting	6
2.3 Parallelization	8
2.3.1 Efficiency considerations	8
2.4 Changelog	9
2.4.1 Version 3.4 in progress	9
2.4.2 Version 3.3	10
2.4.3 Version 3.2	10
2.4.4 Version 3.1	10
3 Program descriptions	11
3.1 Units	11
3.2 Programs	11
3.3 Translate	11
3.3.1 File formats	13
3.3.2 NISE (Full nonadiabatic simulation)	15
3.3.3 Program output	18
3.4 Fourier transform	20
4 Tutorial	21
5 Details on the available techniques	25
5.1 Analyse	25
5.2 Correlation/Autocorralation	26
5.3 Pop (population transfer) ⁺	27
5.4 MCFRET ⁺	27
5.5 Redfield	28
5.6 Dif (diffusion) ⁺	28

5.7	Ani (anisotropy)	29
5.8	DOS ⁺	29
5.9	Absorption	29
5.10	Luminescence	29
5.11	LD (linear dichroism)	30
5.12	CD (circular dichroism) ⁺	30
5.13	Raman	30
5.14	SFG (sum-frequency generation)	31
5.15	2DIR* (two-dimensional infrared)	31
5.16	2DSFG (two-dimensional sum-frequency generation)	31
5.17	2DUVvis* (two-dimensional electronic spectroscopy)	31
5.18	CG_2DES (coarse-grained two-dimensional electronic spectroscopy) ⁺	32
5.19	2DIRRaman*	32
5.20	2DFD (fluorescence detected two-dimensional spectroscopy)	32
6	Information for developers	33
7	Acknowledgement	35

Chapter 1

Introduction

The main developer of the NISE3 code is Thomas la Cour Jansen. Please, cite the appropriate references [1–5] when publishing work using this code. The code allows the calculation of the linear absorption, linear dichroism, sum-frequency generation, two-dimensional spectra (IR, UVvis, and SFG), population transfer, exciton diffusion and integrated anisotropy using the full nonadiabatic semi-classical numerical integration of the Schrödinger equation approach [2] and the sparse matrix optimization approach [4]. The code allows treating spectra of diverse systems involving intra- and intermolecular energy transfer [1, 3, 6–8], non-Gaussian dynamics [9, 10], surfaces [5], and chemical exchange [11]. This manual is not intended as an introduction to two-dimensional spectroscopy. The user is directed to the references including recent reviews [2, 12–15] and books for more information [16–18].

The code is made available as open source on github.com as link [GH1acour/NISE_2017](https://github.com/GH1acour/NISE_2017). The version available is intended to be friendly for developers to contribute improvements. Changes in the code are at own risk and should be reported clearly in publications. Developers wanting to contribute to the official version of the code are suggested to contact the main developer in advance to ensure that the contribution is not conflicting with that of other developers. More information for developers is given in Chapter 6. Feedback on the program and the manual are welcome via e-mail: t.l.c.jansen@rug.nl.

Current developers include: Thomas la Cour Jansen (main developer), Floris Westerman (cmake and MPI implementation), and members of the Computational Spectroscopy group at the University of Groningen.

The code use wavenumbers (cm^{-1}) for frequencies and times are in femtoseconds (fs). The transition dipoles should preferably be given in Debye and transition polarizabilities in \AA^3 . For using the on-the-fly coupling calculations the transition dipoles should be provided in Debye and positions should be provided in Ångström. Users are discouraged from trying to use other units at this may lead to numerical instabilities or unit conversion errors. The final results will scale accordingly.

In this manual keywords used in the NISE input files are highlighted as **Keyword**. Variables defined with these keywords are highlighted as **variable**. Code to be typed in the linux command line or to be used in a script is highlighted as **code**.

Chapter 2

Installation and options

A demonstration for installing NISE can be found on YouTube: <https://www.youtube.com/watch?v=npvV9U0FmDg&t=7s>. NISE has a number of dependencies:

- FFTW3 library, possibly with OpenMP/MPI support. See <http://www.fftw.org/>.
- LAPACK library, often preinstalled. See <http://www.netlib.org/lapack/>.
- CMake v3.10 or higher
- MPI v3 implementation, such as OpenMPI, MPICH (Unix) or MS-MPI (Windows)
- Modern C compiler, implementing a recent OpenMP version
- LaTeX + BibTeX distribution if you want to build the documentation. Not required
- DISLIN, python, MATLAB, or gnuplot, for plotting the results using the included scripts

As some FFTW3 installations do not come with the correct CMake compatibility, it might be useful to use vcpkg as package manager for C++ libraries. This package manager is cross-platform compatible.

2.1 Building

In order to build and compile the NISE software, it is recommended to use the CMake build system. This will try to automatically link the libraries and pick the correct compiler settings.

1. Extract the source code files.
2. Create a `build` directory in the main folder using `mkdir build`.
3. Run `cmake ..` inside this new `build` directory.

4. If `cmake` was successful, run `make` in the same directory to start compilation.
5. All executables should be available in a new `bin` directory in the main folder.

Additional steps to install on MacOS and Windows can be found in section 2.2.

There are several options you can provide to the `cmake` command in order to customize your build:

- `-DCMAKE_BUILD_TYPE`: By default, this is set to `RelWithDebInfo`. Other options include `Debug`, `Release` and `MinSizeRel`. Refer to the CMake documentation for more information
- `-DGENERATE_DOCS`: When not set, CMake will attempt to compile this documentation only when building a Release build. You can override this by setting this variable to `true` or `false`.
- `-DACCURATE_MATHS`: When set, CMake will compile the code to use accurate mathematics implementations (as is default when using a C compiler). When not set, the compiler will use the fast-math option (`-ffast-math`, `/fp:fast` or equivalent) which may yield upto 2x speed-up at the minor cost of numerical accuracy.

After running CMake, the following build targets will have been provided for `make`:

- `all`: Same as not providing a build target, will build all source code for the program, but will skip the documentation and the examples
- `2DFFT`: Will build the 2DFFT executable, used to process results
- `translate`: Will build the translation utility, used to convert between input formats
- `NISE`: Will build the main NISE executable
- `doc`: Will build this documentation from scratch
- `examples`: Will build the code necessary for the examples, used later in this document.

2.2 Installation Trouble Shooting

If the automatic installation procedure outlined above does not work this section contains a few potential solutions. It is of course important first to verify that the libraries specified in the dependencies are available.

If the FFTW libraries are not detected by the `cmake` routine the following `cmake` options may be specified by hand:


```
cmake .. -DFFTW_ROOT=/cm/shared/apps/fftw/openmpi/gcc/64/3.3.8/lib
-DFFTW_LIBRARY=/cm/shared/apps/fftw/openmpi/gcc/64/3.3.8/lib
-DFFTW_INCLUDE_DIRS=/cm/shared/apps/fftw/openmpi/gcc/64/3.3.8/
include
```

The names of the directory locations must then be changed to the actual locations on the given system. The example above is for a fftw library compiled with the gcc compiler.

The program can also be installed on the Mac OSX system. However, the standard compiler does not come with OpenMP support. An OpenMP library must therefore be installed first. This can be done using the homebrew system. Details of the installation of homebrew are given on <http://brew.sh>. Then an OpenMP library as libomp can be installed with `brew install libomp`. If the version of cmake is also not recent enough a suitable cmake version can be installed with homebrew as well. Finally, the programme can be build following the general instructions above for building, but with using the command:

```
/usr/local/bin/cmake .. -DCMAKE_C_COMPILER="clang"
-DOpenMP_C_LIB_NAMES="libomp" -DOpenMP_CXX_LIB_NAMES="libomp"
-DOpenMP_libomp_LIBRARY="/usr/local/lib/libomp.dylib"
-DOpenMP_C_FLAGS="-Xpreprocessor -fopenmp /usr/local/lib/libomp.
dylib
-I/usr/local/opt/include"
```

Here, the location of the installed cmake version and the libomp version may have to be changed to match the location when these were installed by homebrew.

Alternatively macports can be used in a very similar way to homebrew. Install libomp with `sudo port install libomp`. The location of omp.h may be different than expected by CMake, which may be fixed with

```
sudo ln -s /opt/local/include/libomp/omp.h /opt/local/include/omp.h or
sudo port install libomp +top_level.
```

The code can also be installed on a Raspberry Pi 4. This requires the installation of the already discussed packages which can be done with:

```
sudo apt update
sudo apt install -y cmake
sudo apt install libopenblas-dev
sudo apt install libfftw3-dev
sudo apt install libopenmpi-dev
sudo apt install python3-matplotlib
```

When running on this system one needs to keep in mind the limited disk space and memory of the system. However, the NISE program itself easily installs and run on this system.

The program can also be installed with Windows 10 and newer, using the Windows Subsystem for Linux (WSL). This is a feature built into Windows that allows users to run Linux virtual machines. To install WSL, follow the instructions at <https://learn.microsoft.com/en-us/windows/wsl/install>. After that, a Linux terminal can be opened from the Start Menu. The process to install NISE will be identical to the installation process on Raspberry Pi 4.

2.3 Parallelization

NISE is equipped with support for MPI and OpenMP, to provide a tiered parallelization solution. Currently, both the time consuming 2DIR and 2DUVvis techniques support MPI. The linear techniques do not have a parallel implementation as they are generally fast. (Techniques relying on LAPACK including Luminescence may use the MLK OpenMP support for limited speedup.)

For the two-dimensional techniques, it is recommended to understand the implemented approach for parallelization in order to achieve good performance. Each run will calculate a specified number of samples, each for 21 different polarization directions. The calculation time for each polarization direction is determined by the chosen values for `t1max`, `t2max`, and `t3max`.

All polarization directions may efficiently be calculated in parallel using MPI, distributed over all registered tasks (more explanation follows later). As long as you have sufficiently many samples, this will scale very well. In general, it is recommended to have a fraction or multiple of 21 as number of tasks, in order to make sure that no cores are simply waiting around after completing their part of the calculations.

Within each polarization direction, loops over the `t1` coherence times are parallelized using OpenMP. Due to communication overhead and data sharing difficulties, this does not scale as well as the MPI parallelization. If possible, it is recommended to overprovision your cores, i.e. to make the system spawn more threads than there are cores available. The larger the computation per polarization direction (so higher `t2`, `t3`, system size **Singles**), the better this part will scale. It is recommended that the number of OpenMP threads is either small compared to `t1max` or that `t1max+1` is equal to an integer times the number of OpenMP threads.

For example, to run 4 tasks with each 12 threads with OpenMPI, use the following command:

```
mpirun -np 4 -x OMP_NUM_THREADS=12 --bind-to-socket ~/pathToNISE/NISE inputFile
```

For cluster computing refer to the manual for the cluster. Special commands as `srun` may be required for SLURM systems. The MPI implementation require all input files to be located at a disk available to all nodes.

2.3.1 Efficiency considerations

Some considerations and examples to achieve higher performance:

- As OpenMP parallelization uses shared memory, it is necessary to limit each task to one node. If possible, it is recommended to limit a task even to one socket, or in case of more modern chips, 1 NUMA node. However, this might make the runtime for one polarization direction too high, so it might be worthwhile to trade some efficiency for shorter runtimes.

- It is recommended to overprovision your OpenMP threads by a factor of 2-3. So if a NUMA node has 12 cores, you could pin the threads of this task to this NUMA node and tell OpenMP to create 24-36 threads. Thread pinning differs per OS and more details can be found online. Many popular workload schedulers like SLURM, and some MPI implementations, offer this built-in (for example, `--bind-to-socket` in the command above).
- The most efficient workload division also depends on the problem size, for larger problems with fewer samples, the OpenMP scaling is more efficient than for smaller problems with more samples.
- For example, on a machine with 2 12-core Xeon processors (single NUMA node per socket), it is most efficient to run 4 tasks, with 12 threads assigned to each task (overprovisioning). However, for very large problems with only 2 or 3 samples, it might be better to scale down to 2 tasks, each with 24 threads. For smaller problems with many samples, 8 tasks with 6 threads might be better. In general, it is good to do some quick performance tests beforehand.

2.4 Changelog

2.4.1 Version 3.4 in progress

Work by Thomas Jansen, Carleen D. N. van Hengel, Hoang Long Nguyen, Kai Zhong, Vesna Eric, Gijsbert ten Hoven, Stephanie Gonzalez Migoni, Marick Manrho, Ana Cunha, and Kim van Adrichem

- Added the calculation of Redfield transfer matrices
- Added calculation of spectral densities, correlation functions and lineshape functions
- Added CG-2DES (Kai Zhong, Stephanie Gonzalez Migoni)
- Added MCFRET (Hoang Long Nguyen, Kai Zhong, Vesna Eric, Gijsbert ten Hoven, Marick Manrho, and Kim van Adrichem)
- Improved manual for Windows 10 installation. (Hoang Long Nguyen)
- The 2DFFT code was cleaned up
- Raman and 2DIRaman techniques were added (Carleen D. N. van Hengel)
- openMP parallel CD and DOS were implemented
- Diffusion calculations were added
- On the fly transition-dipole and extended-dipole coupling schemes were implemented

- Added project file option to project on substructures - including multi segment options for linear techniques
- Added inhomogeneous and homogeneous apodization functions
- Added automatic check on Singles keyword providing warning to user if it may be incorrect
- Implemented speed-up (x2) of 2D calculations by removing if statements in propagation of double excited states (2DIR by Kim van Adrichem)

2.4.2 Version 3.3

Work by Thomas Jansen

- Extended MPI support to 2DIR
- Implemented true two-state behaviour for the *UVvis techniques
- Added linear dichroism
- Changed timing information for two-dimensional calculations to percentage of full calculation based
- Important change in naming of 2DIR sub techniques to contain IR at the end (GBIR, SEIR, EAIR, etc.).

2.4.3 Version 3.2

Work by Floris Westerman

- Added CMake build system and improved cross-platform compatibility
- Added MPI support to offer significantly better scaling across multiple nodes, instead of just a single one
- Improved code efficiency, around 4x speed-up of main algorithm code of *UVvis techniques

2.4.4 Version 3.1

- Included OpenMP support for two-dimensional calculations

Chapter 3

Program descriptions

This chapter contains the description of the various input parameters, file formats, and output files for the NISE code.

3.1 Units

The code use wavenumbers (cm^{-1}) for frequencies and times are in femtoseconds (fs). Users are discouraged from trying to use other units at this may lead to numerical instabilities or unit conversion errors. The transition dipoles and transition polarizabilities may be given in any desired units. The final results will scale accordingly. However, for using the on-the-fly coupling calculations the transition dipoles should be provided in Debye and positions should be provided in Ångström.

3.2 Programs

The NISE package contains three major programs. These are:

translate [A program that can convert the Hamiltonian and dipole trajectory between different formats.]

NISE [The general code for calculating spectra.]

2DFFT [Do the 2D Fourier transform.]

These programs will be described below.

3.3 Translate

The Hamiltonian **must** be saved in binary format (GROBIN/SKIBIN) for use in the NISE program. The translate program converts between different formats. The program also

allows selecting specific sites in a Hamiltonian file and modification of the fundamental frequencies corresponding to isotope labeling. The input file format is:

InputEnergy [filename]

InputDipole [filename (only needed for GRO/SKI formats)]

InputDipoleX [filename (only needed for MIT format)]

InputDipoleY [filename (only needed for MIT format)]

InputDipoleZ [filename (only needed for MIT format)]

InputAnharm [filename (only needed for SKI format)]

InputOverto [filename (only needed for SKI format)]

InputAlpha [filename (only available for GRO format)]

OutputEnergy [filename]

OutputDipole [filename (only needed for GRO/SKI formats)]

OutputDipoleX [filename (only needed for MIT format)]

OutputDipoleY [filename (only needed for MIT format)]

OutputDipoleZ [filename (only needed for MIT format)]

OutputAnharm [filename (only needed for SKI format)]

OutputOverto [filename (only needed for SKI format)]

OutputAlpha [filename (only available for GRO format)]

Singles [number of oscillators]

Doubles [number of doubly excited states]

Skip [Doubles=Neglect doubly excited states, needed for SKIBIN]

Length [number of timesteps in the trajectory files]

Anharmonicity [A fixed anharmonicity. Only needed for generating GROBIN file from format without double excited states.]

InputFormat [GROBIN/GROASC/MITASC/SKIBIN]

OutputFormat [GROBIN/GROASC/MITASC/SKIBIN]

Modify [Leave this out if you do not wish to modify the Hamiltonian. This keyword requires that the keywords Select, Label, and Shift are also given.]

Select [Number of amide units to include, if selected number is less than the number of units in the original a list of the units to include should be given on the following line (i.e. 0 1 2 3 5 if 6 units are in the original and we want unit 4 excluded) Note that Singles should be the number of residues in the original file. This keyword is only used if the Modify keyword is used.]

Label [On the following line the for each selected unit the isotope labeling is given (0=na-tive, 1=C13, 2=O18) C13 gives a 41 wavenumber frequency shift and O18 gives a 60 wavenumber shift in this implementation. This keyword is only used if the Modify keyword is used.]

Shift [On the following line for each selected unit a frequency shift can be provided. This can be used for modifying proline or terminal frequencies or for isotope labeling. This keyword is only used if the Modify keyword is used.]

3.3.1 File formats

GROBIN is the binary format needed for use with the NISE program. This contains information on the energy of both singly and doubly excited states and allows for fluctuating anharmonicities. GROASC is a text readable format. MITASC and MITTXT are text formats used by the MIT static spectrum code. SKIBIN is a binary format used by the Skinner group. It contain the singly excited states and a separate file for fluctuating anharmonicities and a file for fluctuating overtone transition dipoles. SPECTRON is a text format used by the Mukamel group SPECTRON code. The GROBIN format is the only format that allows storing the doubly excited states Hamiltonian. This will be generated automatically using harmonic rules and the anharmonicity given by the Anharmonicity keyword.

GROBIN and SKIBIN:

The GROBIN and SKIBIN formats are very similar. The energy file has the following format. For each snapshot in the trajectory they contain first an integer typically containing the number of the snapshot. This integer is not used in the calculation, but can be used for control purposes. Then the single excited Hamiltonian is given in floats as a tridiagonal upper matrix. For the GROBIN format the double excited Hamiltonian might be provided afterwards again in a tridiagonal upper matrix.

The dipole file has the following format. For each snapshot in the trajectory they contain first an integer typically containing the number of the snapshot. This integer is not used in the calculation, but can be used for control purposes. Then the x components of the transition dipole matrix are given in floats, followed by the y and z components. For the GROBIN format the transition dipole matrix for the single to double transitions may then be given, again with the x components given first. In Figure 3.1 the format of the file is also illustrated.

Timestep Number (Integer)	Site 1	Site 2	Site ..	Site N
	All x-components (Float)			
Timestep Number (Integer)	Site 1	Site 2	Site ..	Site N
	All y-components (Float)			
Timestep Number (Integer)	Site 1	Site 2	Site ..	Site N
	All z-components (Float)			

Figure 3.1: The file structure for the dipole, and position files. For the extended dipole position file first all x-components of the first set sites are given and then all x-components of the second set of sites.

The anharmonicity file is only given in the SKIBIN format. For each snapshot in the trajectory they contain first an integer typically containing the number of the snapshot. This integer is not used in the calculation, but can be used for control purposes. Then the diagonal anharmonicities are then given in floats.

The overtone file is only given in the SKIBIN format. For each snapshot in the trajectory they contain first an integer typically containing the number of the snapshot. This integer is not used in the calculation, but can be used for control purposes. Then the overtone transition dipoles are then given in floats. Only the same site overtone transition dipoles are given. Transitions like $\langle 1 | \mu | 12 \rangle$ are determined by the harmonic rules.

The transition polarizability file which is not available in the SKIBIN format has the following format. For each snapshot in the trajectory they contain first an integer typically containing the number of the snapshot. This integer is not used in the calculation, but can be used for control purposes. Then the xx components of the transition polarizability matrix are given in floats, followed by the xy, xz, yy, yz, and zz components.

The position files are available in the GROBIN format and used to store either the center positions of the chromophores or the positions of the two sites in each chromophore defining the extended dipole coupling. This file is used when calculating CD spectra, excitation diffusion, or when using the on-the-fly coupling calculations. The structure of the file is as for the dipole files (see Figure 3.1).

GROASC:

The energy file contain one Hamiltonian snapshot for each line. The upper tridiagonal matrix is saved and each number is separated by a space. The transition dipoles are stored in one file. Each line contains one snapshot with the numbers separated by a space. All the x components for one snapshot are stored first followed by the y and z components. Only the single excitation Hamiltonian and the transition dipoles from the ground state to the singly excited state are saved.

MITASC:

The energy file contain one Hamiltonian snapshot for each line. The whole square matrix is saved and each number is separated by a tab. The transition dipoles are stored in one file for each cartesian component. Each line contains one snapshot with the numbers separated by a tab. Only the single excitation Hamiltonian and the transition dipoles from the ground state to the singly excited state are saved.

MITTXT:

The energy file contain one row of the Hamiltonian snapshot for each line. The whole square matrix is saved and each number is separated by a space. Snapshots are separated by an empty line. The transition dipoles are stored in one file for each cartesian component. Each line contains the cartesian coordinates of one transition dipole. The numbers separated by a space. Snapshots are separated by an empty line. Only the single excitation Hamiltonian and the transition dipoles from the ground state to the singly excited state are saved.

SPECTRON:

In the energy file each snapshot is stored after a line containing the word SNAPSHOT and the number of the snapshot (starting from 1). After this the upper tridiagonal matrix of the Hamiltonian snapshot is stored with one row on each line. The numbers are separated by a space. The dipole file also contain a line with the word SNAPSHOT and the number of that snapshot. This is followed by the transition dipoles stored with the x, y and z components for each site on one line separated by a space. Only the single excitation Hamiltonian and the transition dipoles from the ground state to the singly excited state are saved.

3.3.2 NISE (Full nonadiabatic simulation)

The nonadiabatic linear and third-order response can be simulated with the 'NISE' program utilizing the numerical integration of the Schrödinger equation (NISE) approach [1,2]. The sparse and double excited state propagation is described in Ref. 3. The coupling propagation scheme is described in Ref. 4. The parallelization of the code is described in Ref. 19. The program require the Hamiltonian trajectory in binary format. The NISE program use the following input:

Hamiltonianfile [File name]

Dipolefile [File name]

Alphafile [File name] (Transition polarizability file. Only needed for Raman and 2DIRRaman calculations)

Anharmonicfile [File name]

Overtonedipolefile [File name]

HamiltonianType [Full/Coupling/TransitionDipole/ExtendedDipole] (Full is the default)

Couplingfile [File name]

PDBfile [File name]

Length [Number of snapshots in trajectory]

Samplerate [Number of snapshots between ensemble averaging]

Lifetime [The life time in fs included in 2D response functions]

Homogeneous [Homogeneous lifetime for exponential apodization in fs, only included after FFT]

Inhomogeneous [Inhomogeneous lifetime for gaussian apodization in fs, only included after FFT]

Timestep [The length of each timestep in fs]

RunTimes [t1 max] [t2] [t3 max, all in timesteps]

Threshold [The threshold for the sparse matrix approximation, typical value 0.001]

Anharmonicity [0 = anharmonicities from file used, all other values result in the use of a fixed anharmonicity with that value]

Singles [Number of singly excited states]

Propagation [Sparse/Coupling/RK4 default is Sparse] (Coupling recommended for fast calculations)

Couplingcut [Value in cm^{-1} below which the couplings are neglected, default 0, only used in the Coupling propagation scheme. ¹]

Trotter [The number of trotter steps in the Paarmann approximation for double excited states, 5 recommended for the Sparse propagation scheme, for the Coupling propagation scheme this is the general number of trotter steps the recommended value is then 1]

ImaginaryTimesteps [The number of trotter steps in the imaginary time propagation for calculating the effect of the equilibrium density matrix, typically a value of 5 is sufficient. If the value is set to 0 exact matrix diagonalization is used.]

MinFrequencies [minw1] [minw2] [minw3, all in reciprocal cm^{-1}]

MaxFrequencies [maxw1] [maxw2] [maxw3, all in reciprocal cm^{-1}]

¹From the 16/8-2022 version couplings below the cutoff are neglected when reading in the Hamiltonian or when it is constructed on the fly. NOTE that before this the approximation was NOT applied during t_2 for two-dimensional spectroscopic calculations and it was not used for the Analysis of the Hamiltonian.

Technique [Options include: Absorption / DOS / Luminescence / CD / Raman / 2DIR / 2DUVvis / 2DIRraman / Pop / MCFRET / Dif / Ani / Analyse / CG_2DES / FD_CG_2DES , these techniques and additional options are explained in Section 5]

FFT [Number of points on each axis in 2DFFT, if bigger than max times zero padding is used. If left out no zeropadding is used (recommended).]

Format [Matlab/Dislin/Python/Gnuplot Different output formats for the 2D spectral files. For Matlab format rephasing and non-rephasing spectra are not added. The Matlab format refer to a format used by the Tokmakoff group. Dislin and Python are the same and the format used in the tutorials within the package. Gnuplot is suited for gnuplot and an example is given in the tutorial. Details are given in the section on Fourier transforms.]

BeginPoint [The number of the first sample calculated in this run]

EndPoint [The number of the last sample calculated in this run, if this keyword is left out all samples will be included]

Project (Should be followed by either two new lines specifying the sites to project on directly in the file or one line specifying the filename of a file to read the site numbers from. In the first case the first line should read Sites [Number] identifying how many sites are included in the projection and the second line with a list of integer numbers identifying the sites included counting from zero. In the other case the following line should read **Projectfile** [Name], where the file named [Name] contain the number of sites on the first line and then a list of the site numbers. For MCFRET the project file is instead used to specify to which segment each site belongs. In this case the number of sites must be equal to the number of singles.)

Basis [Site / Adiabatic / Average] The options for the basis to use in population transfer calculations. The standard option is the site basis.

PrintLevel [0 / 1 / 2] (Default is 0, the higher number the more detailed (timing) info is given.)

The minimum and maximum frequencies are used for two things. First, the average of minw1 and maxw1 are used to shift the frequencies during the simulation. By shifting all transitions by the average the oscillations of the time-evolution are reduced and the simulations can be performed with longer distance between the time points. Second, the output is reduced to frequencies within the min and the max which reduce file size.

The Lifetime is only applied during the coherence times (t_1 and t_3) according to the formula in Ref. 4. The waiting time dependence can trivially be accounted for by multiplying the whole spectrum with $\exp(-t_2/T_1)$. For apodization with an exponential or gaussian function the Homogeneous and Inhomogeneous keywords are used. These only

change the final spectra and not the response functions. The equations used are

$$R'(t) = R(t) \exp \left(-t/2t_{\text{homo}} - t^2/2t_{\text{inhomo}}^2 \right) \quad (3.1)$$

for the linear techniques,

$$R'(t_1, t_2, t_3) = R(t_1, t_2, t_3) \exp \left(-(t_1 + t_3)/2t_{\text{homo}} - (t_1 - t_3)^2/2t_{\text{inhomo}}^2 \right) \quad (3.2)$$

for rephasing signals and

$$R'(t_1, t_2, t_3) = R(t_1, t_2, t_3) \exp \left(-(t_1 + t_3)/2t_{\text{homo}} - (t_1 + t_3)^2/2t_{\text{inhomo}}^2 \right) \quad (3.3)$$

for non-rephasing signals. As default no apodization is performed. The apodizations are inactive when set to zero.

The HamiltonianType allow the use of different ways to store the Hamiltonian trajectories. The default setting is Full, where the full Hamiltonian is given in the Hamiltonianfile, if Coupling is specified the couplings are assumed to be constant and given in the Couplingfile. The Hamiltonianfile then contain the trajectory of fluctuating diagonal elements. If TransitionDipole is specified the couplings are calculated on the fly using the transition-dipole coupling scheme. The chromophore positions must then be provided in the Positionfile and the transition-dipoles in the Dipolefile. As for the Coupling method the HamiltonianFile then only contains the diagonal elements. The dipoles are then assumed in Debye and Positions in Ångström. If ExtendedDipole is selected the extended-dipole coupling scheme is used to calculate coupling on the fly. The dipole magnitude is then taken from the Dipolefile. The position of the two transition point-charges must be provided in the Positionfile, with the first point corresponding to odd timepoints and the second point corresponding to even timepoints.

The double excitation Hamiltonian is propagated using the Trotter formula scheme [3, 20].

During the calculation the program will create a log file called NISE.log. For linear techniques this file will be updated every time a new sample has been calculated. The update contains timing information and therefore allows for the estimation of when the complete calculation finishes. For 2D techniques the timing information will given directly in the output file tracking the progress percentage.

For exciton diffusion calculations an additional input file containing the site positions is needed. This file should be named Position.bin and have the following format. The first entry should be the size of the box in float (all sides are assumed equally long). This is followed by the x, y, and z coordinates for each site again in floats. After the coordinates of the first snapshot the coordinates for the second follows directly and the coordinates for all snapshots should be stored.

3.3.3 Program output

The linear absorption calculation provides the linear response function in time domain in the file TD_Absorption.dat and the linear dichroism response function in time domain in

the file RLD.dat. The first column is the time in femto seconds. The other columns are the real and imaginary parts of the response function. It is recommended to check that these have decayed to a small value close to zero in the calculated time interval. If this is not the case the system has not lost its coherence within the time determined by t1 max and the spectrum will be too broad. The linear absorption is given in the file Absorption.dat and the linear dichroism in the file LDIR.dat. The first line is the frequency in wavenumbers, the second line is the absorption spectrum, and a third line contain zeroes. For the linear dichroism calculation the z-axis is assumed to be the unique axis.

The 2D calculations (2DIR and 2DUVvis) provides the files R(par/per/cro)(I/II).dat. These files contain the time domain third-order response functions for different polarization directions [13,21]. The first three columns are the times t1, t2, and t2 in femtoseconds. The two last columns are the real and imaginary parts of the response functions. The frequency domain response functions are found using a double Fourier transform (see section 3.4). It is recommended to check that the response function has decayed within the calculated time intervals.

The 2DIRaman calculations provides the files R(par/per/cro).IRaman.(I/II).dat. These files contain the time domain third-order response functions for different polarization directions [13,21]. The first three columns are the times t1, t2, and t2 in femtoseconds. The two last columns are the real and imaginary parts of the response functions. The frequency domain response functions are found using a double Fourier transform (see section 3.4). It is recommended to check that the response function has decayed within the calculated time intervals.

The population transfer calculation provides two files. Pop.dat contain two columns. The first is the time in femtoseconds the second is the probability that if a site was initially excited it is still excited after the given time. This is averaged over all sites. The file PopF.dat contain more detailed information. The first column is the time. The following columns are the probability that if a particular site was excited initially other sites are excited later. The first N columns are the populations of the N sites following an initial excitation of the first site. Then follows the populations after excitation of the second site etc. This file might be a very large file for big systems.

The MCFRET calculation provides, which contains the rate matrix in the file **RateMatrix.dat**. As required for a master equation the sum of the rates out of a segment is (minus the value of the diagonal element) is equal to the sum of the rates of what is transferred to other segments and the sum of each column is zero. It also outputs the quantum corrected rate matrix, QC.RateMatrix.dat, with the Oxtoby quantum correction factor $c = \frac{2}{1 + e^{-\hbar\omega_{ab}/K_B T}}$. Notice the unit used for the rates is ps^{-1} as processes where MCFRET is relevant are typically slow!

For the exciton diffusion the output file is Dif.dat It has three columns. The first is the time in femto seconds. The second column is the mean square displacement of the wave function assuming that one site is initially excited. The third column is the mean square displacement of the center of the wave function assuming that one site is initially excited. The program averages over all sites as initial sites. One should be aware that periodic

boundary conditions are applied in this calculation and the mean square displacements will saturate [3].

For the integrated anisotropy calculation the output file is Ani.dat. It has three columns. The first is the time in femtoseconds. The second is the integrated anisotropy and the last is the orientational correlation function.

The Hamiltonian analysis provides the delocalization length/size according to the definition of Thouless [22] directly in the program standard output.

3.4 Fourier transform

The 2DFFT program use the same input as the NISE program. For 2DIR and 2DUVvis the Fourier transformed response is written in files named Rw(par/per/cro).(I/II).dat, where par, per, and cro denote parallel, perpendicular and cross polarized signals. I and II denote the k_I and k_{II} contributions. When Dislin format is used the 2D correlation spectrum is saved in the files 2D.(par/per/cro).dat and the 'broad pump narrow probe' pump probe signal is stored in PP.(par/per/cro).dat. The first column is ω_1 , the second is ω_3 , the third is the dispersive signal, and the last column is the absorptive signal. Perl scripts connected with Dislin are available for plotting or the user can use his or her favorite plotting program. The Dislin format is also used by the python plotting code included in the tutorial files. The Gunplot output is only differing from the Dislin format in an empty line following each new value of ω_1 . A Gnuplot example file is included in the tutorial. A python plotting script which use the Dislin output format file is included. The cover picture of the manual is generated with this python script. For Matlab format only the Rw files are created. These then contain a matrix with the response. An additional file waxis.dat is created with the values of the frequencies.

If a file called Waitingtimes.dat is present all the waiting times given in the file (in fs) will be calculated simultaneously assuming that the response function files are named RparI_0fs.dat etc. The resulting frequency domain files will be named 2Dpar_0fs.dat etc.

For 2DIRraman the Fourier transformed response is written in files named Rw(par/per/cro)IRraman.(I/II).dat, where par, per, and cro denote parallel, perpendicular and cross polarized signals. I and II denote the k_I and k_{II} contributions. When Dislin format is used the 2D correlation spectrum is saved in the files 2DIRraman.(par/per/cro).(I/II).dat and signal integrated over ω_1 is stored in PPIRraman.(I/II).(par/per/cro).dat. The first column is ω_1 , the second is ω_3 , the third is the dispersive signal, and the last column is the absorptive signal. Note that 2DIRraman rephasing and non-rephasing signals cannot be added directly and are therefore stored separately. The non-rephasing ω_1 frequencies are about twice that of the single excitation transition frequency as first coherence corresponds to a coherence between a double excited state and the ground state.

Chapter 4

Tutorial

Extended tutorials are available on github via <https://github.com/GHlacour/NISE.Tutorials>. A very simple tutorial does come with the program itself and is described below.

A tutorial example exists for two coupled chromophores. To activate this you need to run `make examples` in the `build` directory. The tutorial files are located in the folder `examples/tutorial`. The program `stochastic` creates the input trajectories in GROASC format and input files can be found for converting to GROBIN and for the calculation of linear and two-dimensional spectra. The `stochastic` code allow the user to construct a simple trajectory for two coupled three level system. It is assumed that these levels are coupled to a set of overdamped Brownian oscillators [17]. The user can vary the length of the trajectory, the duration of the time interval between snapshots, the width of the frequency distribution, the correlation time, and the degree of correlation between the two three level system fundamental frequency fluctuations. The energy difference and the average frequency for the two fundamental frequencies can be adjusted as can the coupling between them. The angle between the assumed fixed transition dipoles can be adjusted as well. An example for parameters is given in the script named `run`.

After generating the frequency trajectory with the `stochastic` program one can look at the generated trajectories in `Energy.tex` and `Dipole.tex`. These should be converted to binary input for the NISE program. This is done with the `translate` program with the input given in the file `inpTra`. The command line entry:

```
~/NISE_2017/bin/translate inpTra
```

will accomplish this. (It is assumed that you installed NISE in your home directory. If the programme was installed somewhere else you need to run the programme from the alternative bin directory where it was installed.)

When the binary trajectory exist we can calculate the linear absorption using the input from the file `input1D`.

```
~/NISE_2017/bin/NISE input1D
```

One should plot the output in `TD_Absorption.dat` and `Absorption.dat` to get familiar with this. It is important that the response function has decayed almost completely within the

simulated time to avoid ringing, noise, and other artifacts in the spectrum. If the spectrum is noisy a shorter time can be selected for the apodization function or if possible one can average over more realizations either by making the trajectory longer or decreasing the number of steps between realizations as set by the keyword **Samplerate**. The program should finish in a about 10 seconds. The python script `plot.py` will plot both the response function and the absorption spectrum. Simply type `python plot.py` assuming that python 3 is installed. Alternatively, you can plot the two first columns of each file with your own favourite plotting programme.

The linear dichroism signal can be calculated by changing the **Technique** to **LD** which generate corresponding files with the names TD_LD.dat and LD.dat. In the tutorial this is not so interesting as it is simply minus half of the absorption signal, since the dipoles are in the x, y plane and the unique angle for linear dichroism is defined to be the z-axis.

The two-dimensional infrared spectra are calculated with the input file input2D. This corresponds to the spectrum of two coupled three-level systems in this case, where the anharmonicity is 20 cm^{-1} for both three-level systems as specified in the input2D file. The programme is executed with the command:

```
~/NISE_2017/bin/NISE input2D
```

The percentage wise progress of the calculation can be followed in the output. One should expect that the calculation of one sample point takes about 10 seconds with the given settings. For a test about 100 samples are sufficient, while for nice spectra at least 1000 samples are needed. Typically more samples are needed for two-dimensional spectra than for linear absorption spectra and the more different distinct chromophore environments are present the more samples are needed. The given example should take about 15 minutes to simulate (of course depending on the computer). If you have multiple cores on your computer you can run the MPI parallel version typing:

```
mpirun -n 4 ~/NISE_2017/bin/NISE input2D
```

The number 4 tells mpi that the job should be split in 4. For this small system mpi is likely more efficient than OpenMP on typical computer architectures. The output of the NISE programme is the time domain response functions. To get the frequency domain spectra the **2DFFT** program must be executed as described below. For longer simulations you should use the queue system of your computer cluster. Check the instructions on running on your local computer cluster to do this.

The NISE code was originally build to consider coupled three-level systems, which is what is needed to simulate infrared spectra. If one want to simulate two-dimensional electronic spectra coupled two-level systems usually need to be considered. This can now be achieved with the **2DUVvis** technique. Previously this was achieved by using an anharmonicity that moves the third-level out of the spectral window and sufficiently far away that the coupling between the third-level for each monomer and double excitations involving pairs of sites can be neglected [4, 23].

The 2D simulations can be performed in parallel using openMP. The program does this by distributing the calculations for different values of t1 on different CPUs. When

running parallel it is advantageous to ensure that $([t1max]+1)$ divided by the number of used CPUs is an integer number or slightly smaller than an integer number. To run the code parallel the `OMP_NUM_THREADS` variable has to be set equal to the number of CPUs that one want to use in the submission script (or if not running under a queueing system the variable has to be set in the used linux shell). On most clusters one should further specify to the queueing system how many CPUs are needed when submitting the job. The user is referred to the manual of the cluster used for this kind of information as it varies from system to system. For parts of the code, where the mkl library is used for solving eigen value problems these may be run in parallel as well by setting the `MKL_NUM_THREADS` variable equal to the number of CPUs used for this task. The general experience is that for small problems running parallel is not efficient. The mkl library also does not perform well on large numbers of CPUs and is only recommended for working with large Hamiltonians. For large problems the calculation of polarization contributions and samples can efficiently be spread over different CPUs and/or nodes for efficient massively parallel calculations. 2D simulations for mid-size systems can efficiently be spread over 10s of nodes each with more than 100 CPUs.

The response functions calculated in time domain can be converted to frequency domain spectra using the double Fourier transform code with the same input as for the calculation of the response function.

```
~/NISE_2017/bin/2DFFT input2D
```

The output can be visualized with one of the plotting scripts provided with the tutorial. The recommended method is using the `plot2D.py` python script, which is written to plot the parallel polarization spectra. The line `Data = np.loadtxt('2D.par.dat')` in the beginning of the script can be changed to load and plot one of the other data files (perpendicular spectrum `2D.per.dat` or cross polarization `2D.cro.dat`). The plot at the cover of the manual is the parallel polarization 2DIR spectrum obtained from the tutorial.

The user is encouraged to play with the settings of the run script in the tutorial to get familiar with the code. The simple example can also be useful for getting familiar with the basics of two-dimensional spectroscopy in particular regarding the effects of coupling, correlation, relative angles between coupled chromophores.

The delocalization length [22] can be calculated with the analysis option as given in the example input file `inputAnalyse`. Note that for including the complete trajectory `[t1max]` should be set to 1.

```
~/NISE_2017/bin/NISE inputAnalyse
```

By specifying a frequency range the delocalization of eigenstates within that range is also determined. Furthermore, the density matrix average over all states within the given range is determined along with the density matrix weighted with the square of the transition-dipole for each eigenstate.

A more extensive set of tutorials can be found at https://github.com/GHlacour/NISE_Tutorials. This include examples for the LH2 light-harvesting system, FMO, and

LHCII. This tutorial repository also contains useful scripts for plotting the spectra calculated with this code.

Chapter 5

Details on the available techniques

In this Chapter, more details on special options and the theory behind each technique is given. The techniques marked with a star (*) are implemented with OpenMP and MPI options (see [24]), while techniques implemented with OpenMP options are marked with a plus (+). All other techniques are implemented for a single CPU. Running these techniques on a cluster reserving multiple nodes or CPUs will be waste of time. When running them interactively a warning may be produced, which can in that case be ignored.

5.1 Analyse

The Hamiltonian is analysed and different statistical properties are provided including the average delocalization size of Thouless [22]. This inverse participation ratio (IPR) is expressed as:

$$D_{IPR} = \left\langle \frac{1}{N} \sum_i \left(\sum_j |c_{ij}|^4 \right)^{-1} \right\rangle. \quad (5.1)$$

The Manhattan exciton size is calculated as [25]:

$$D_{MES} = \left\langle \frac{1}{N} \sum_i \left(\sum_j |c_{ij}| \right)^2 \right\rangle. \quad (5.2)$$

In a file named **Analyse.dat** statistics is given for each site including. This include the average energy and the standard deviation. The average coupling strength (signed sum of all couplings of a given site with all other sites) and the standard deviation of this quantity. In a file named **Av_Hamiltonian.dat** the average Hamiltonian is stored as a single snapshot in the GROASC format. Furthermore, the density matrix is calculated for the spectral region defined by the upper and lower bounds of **MinFrequencies** and **MaxFrequencies**. This is done both with and without a weight determined by the contribution to the absorption spectrum as:

$$\rho_{ij} = \sum_k \left\langle c_{ik}^* c_{jk} \Theta(\omega_k - \omega_{min}) \Theta(\omega_{max} - \omega_k) \right\rangle \quad (5.3)$$

Where c_{jk} is the wavefunction coefficient of eigenstate k on site j and ω_k is the wavenumber associated with that eigenstate. The brackets symbolize the ensemble average over the trajectory.

$$\rho_{ij}^{\text{spectral}} = \sum_k \left\langle |\vec{\mu}_k|^2 c_{ik}^* c_{jk} \Theta(\omega_k - \omega_{\min}) \Theta(\omega_{\max} - \omega_k) \right\rangle \quad (5.4)$$

Here $\vec{\mu}_k$ is the transition-dipole vector of eigenstate k . This weighted density matrix, thus, emphasizes the states contributing to the absorption spectrum in the given spectral window. The absolute value density matrix defined as:

$$\rho_{ij}^{\text{ADM}} = \sum_k \left\langle |c_{ik}^* c_{jk}| \Theta(\omega_k - \omega_{\min}) \Theta(\omega_{\max} - \omega_k) \right\rangle, \quad (5.5)$$

reveal information about the delocalization in the system. The participation ration matrix defined as:

$$\rho_{ij}^{\text{PRM}} = \sum_k \left\langle (c_{ik}^* c_{jk})^2 \Theta(\omega_k - \omega_{\min}) \Theta(\omega_{\max} - \omega_k) \right\rangle, \quad (5.6)$$

can also be used to understand the delocalization. These density matrices are stored in the files named `LocalDensityMatrix.dat`, `SpectralDensityMatrix.dat`, `AbsoluteDensityMatrix.dat`, and `ParticipationRatioMatrix.dat`. The former should be identical to the unit matrix if the full region of all eigenstates is included. Segments will be defined based on the values of the absolute value density matrix. When the value between two sites divided by the diagonal value is less than the value set by `Threshold` the two sites will be in the same segment.

5.2 Correlation/Autocorrelation

This technique calculates the frequency correlation functions including cross correlations between all sites. If Autocorrelation is specified instead cross-correlations are skipped, which will speed up the calculations significantly for large systems. The correlation functions are stored up to `t1max` × `Timestep` femto seconds. The first column in the file `CorrelationMatrix.dat` contains the time in femtoseconds, while the following columns contain the correlations functions ordered as in a tridiagonal matrix unless the technique `Autocorrelation` is specified. In that case only the autocorrelation functions are calculated and stored with one column per site. The time-dependent skewness $\langle \omega(t) \omega^2(0) \rangle$, and the time-dependent kurtosis $\langle \omega^2(t) \omega^2(0) \rangle$ are stored in the files `Skewness.dat` and `Kurtosis.dat`, respectively. These files only contain these higher-order correlation functions for individual sites. The spectral density (here defined as the real part of the Fourier transform of the time-correlation function) is calculated for each site and stored in the file `SpectralDensity.dat`. The first column gives the wavenumber and each remaining column gives the spectral density for a site. The lineshape function for each site is calculated and stored in the file `LineshapeMatrix.dat`. The first column gives the time and the following columns pairwise give the real and imaginary part of the lineshape function

for each site using the temperature given in the input. The file `Av_ExpLineshape.dat` contain the the function $\exp(-g(t))$, where $g(t)$ is the average of the site lineshape functions. For a collection of uncoupled chromophores with identical average energy, this will correspond to the response function including thermal correction. The first colum is the time in femto seconds, while the two other columns are the real and imaginary parts.

5.3 Pop (population transfer)⁺

The population transfer is calculated between sites. In general, this is governed by the equation:

$$P_{fi}(t) = \langle |U_{fi}(t, 0)|^2 \rangle \quad (5.7)$$

Here f and i are the final and initial sites. Generally two files are generated. In `Pop.dat` average of the population remaining on the initial site (P_{ii}) is calculated resulting in a value starting at 1 (all popiulation is on the initial state) and decaying to the equilibrium value $1/N$ (equal population on all states). In the `PopF.dat` file a full matrix op all combinations of initial and final states is given. The columns start start with the first initial state and all possible final states continuing to the second possible initial state and all possible final states. This file may be very large for large system sizes! The calculation as default is calculated in the site basis. A different basis can be defined using the keyword **Basis**. The choices `Local`, `Average`, and `Adiabatic` specifies the Local/Site, average exciton, and instantaneous adiabatic basis, respectively.

5.4 MCFRET⁺

The MCFRET part calculates the Excitation Energy Transfer (EET) rate with Multichromophoric Förster Resonance Energy Transfer (MC-FRET) method. [26–28] The rate is calculated based on the equation:

$$k = \frac{2\text{Re}}{\hbar^2} \int_0^\infty dt \text{Tr}[J^{\text{AD}} E^{\text{D}}(t) J^{\text{DA}} I^{\text{A}}(t)], \quad (5.8)$$

with $I^{\text{A}}(t)$ and $E^{\text{D}}(t)$ are the absorption time domain matrix and emission time domain matrix.

The techniques MCFRET-Absorption and MCFRET-Emission can provide the time domain absorption matrix, and time domain emission matrix separately. For MCFRET-Emission the expectation value of the density matrix is required, which can be obtained with MCFRET-Density. The technique of MCFRET-Coupling determines the average coupling between segments over a trajectory. The intermediate results are stored in the files:

`TD_absorption_matrix.dat` for the time domain absorption matrix, `TD_emission_matrix.dat` for the time domain emission matrix, and `CouplingMCFRET.dat` for the average couplings. The technique of MCFRET-Rate can be used to calculate the EET rate with the input

of the time domain absorption matrix, time domain emission matrix, and the coupling to generate the results in steps and introduce possible changes in only some of these steps. The calculated rate matrix is written in the file `RateMatrix.dat` (in units of ps^{-1}). The MCFRET-Analyse technique calculates the expectation value of the segment energies and the quantum corrected rate matrix, which is stored in '`QC_RateMatrix.dat`' (this includes the Oxtoby correction to give a thermal distribution). The intermediate rate response is stored in `RateFile.dat`. It is recommended to validate that the response do decay to zero within the chosen simulation time. The coherence decay matrix is stored in `CoherenceMatrix.dat`. The coherence decay should be faster than the transfer rate for MCFRET to be a suitable approximation for the energy transfer. The MCFRET technique will perform all steps sequentially to calculate the EET rate matrix, but they can be performed sequentially, which may be useful if one for example only change the couplings between segments. The MCFRET keyword, thus, corresponds to running the following techniques sequentially: MCFRET-Density, MCFRET-Absorption, MCFRET-Emission, MCFRET-Coupling, MCFRET-Rate, and MCFRET-Analyse.

5.5 Redfield

This technique calculates the Redfield rate matrix (see Eq. 33 of Ref. 29). It requires the spectral density as calculated with the Autocorrelation technique and the average Hamiltonian, which can be obtained using the Analyse technique. The rate equation is then:

$$k_{ij} = \sum_{n=1} |c_{in}|^2 |c_{jn}|^2 C_n(\omega_j - \omega_i), \quad (5.9)$$

For negative frequency differences (upward energy transfer) the spectral density is suppressed using the Boltzmann factor: $C_n(-\omega) = \exp(-\omega/k_B T) C_n(\omega)$. Here, ω_j and ω_i are the eigen energies of the exciton states i and j . This approximation is applicable in the limit of weak phonon coupling. That is when the values of the spectral density are small compared to the resonance couplings between sites. If segments are specified in the input the excitons will be determined only within these segments and only intra-segment rates are determined, as transfer between weakly coupled segments are likely better described using the MCFRET theory. The calculated rate matrix is stored in the file `RedfieldRates.dat` and the exciton eigenstate information is stored in the `RedfieldRates.dat` file.

5.6 Dif (diffusion)⁺

The mean square displacement of excitons are calculated using the positions provided in the Positions file. Periodic boundary conditions are applied assuming a cube box. The box size must be provided in the first column of the Positions file. Two diffusion properties are calculated. The diffusion of the exciton center is calculated by calculating the center of the excitation given it started on a specific site j : $x_{\text{ex},i}(t) = \langle \sum_j (x_j(t) - x_i(0)) | U_{ji}(t, 0) |^2$

and then determining the mean square displacement as $MSD_{\text{ex}}(t) = \langle \sum_i x_{\text{ex},i}(t)^2 \rangle$. The other measure is based on the probability of starting on one site and being detected on another site $MSD_{\text{site}}(t) = \langle \sum_{ji} (x_j(t) - x_i(0)) |U_{ji}(t, 0)|^2 \rangle$. $MSD_{\text{site}}(t)$ is stored in the second column of the Dif.dat output file and $MSD_{\text{ex}}(t)$ is stored in the third column. The first column contain the time in femtoseconds. The unit for the mean square displacements is the distance unit provided on the position file (Ångström recommended) squared per femtosecond.

5.7 Ani (anisotropy)

Not implemented yet (check NISE_2015)

5.8 DOS⁺

The density of states is calculated using the response function:

$$I(t) = \langle \text{Tr} U(t, 0) \rangle \exp(-t/2T_1). \quad (5.10)$$

Both the real and imaginary parts are stored. The Fourier transform is the frequency domain density of states, which is stored in the file DOS.dat. T_1 is the lifetime, which is often simply used as an appodization function to smoothen the spectrum.

5.9 Absorption

The linear absorption is calculated using the first-order response function [30]

$$I(t) = \sum_{\alpha}^{x,y,z} \langle \mu_{\alpha}(t) U(t, 0) \mu_{\alpha}(0) \rangle \exp(-t/2T_1). \quad (5.11)$$

Both the real and imaginary parts are stored. The Fourier transform is the frequency domain absorption, which is stored in the file Absorption.dat. T_1 is the lifetime, which is often simply used as an appodization function to smoothen the spectrum.

5.10 Luminescence

The luminescence is calculated using the first-order response function

$$I(t) = \sum_{\alpha}^{x,y,z} \langle \frac{1}{Z} \mu_{\alpha}(t) U(t, 0) \exp(-H(0)/k_B T) \mu_{\alpha}(0) \rangle \exp(-t/2T_1). \quad (5.12)$$

Both the real and imaginary parts are stored. The Fourier transform is the frequency domain luminescence, which is stored in the file **Luminescence.dat**. T_1 is the lifetime,

which is often simply used as an appodization function to smoothen the spectrum. The Boltzmann term containing the Hamiltonian at time zero (H) and the temperature (to be specified in the input) ensure the emission from a thermalized population of the excited state ignoring a potential Stoke's shift and effects of vibronic states. The spectrum is normalized with the partition function (Z).

5.11 LD (linear dichroism)

The linear dichroism is calculated identically to the linear absorption except the average of the absorption in the x and y directions is subtracted from the absorption in the z direction. This corresponds to a perfect linear dichroism setup, where the molecules are aligned along the z-axis.

5.12 CD (circular dichroism)⁺

The circular dichroism is calculated using the first-order response function

$$I(t) = \sum_{\alpha} \sum_{nm}^{x,y,z} \langle r_{nm} \mu_{\alpha,n}(t) \times [U(t, 0) \mu_{\alpha,m}(0)] \rangle \exp(-t/2T_1). \quad (5.13)$$

Both the real and imaginary parts are stored. The Fourier transform is the frequency domain absorption, which is stored in the file **Absorption.dat**. T_1 is the lifetime, which is often simply used as an appodization function to smoothen the spectrum. Note that the positions of the chromophores must be provided in the **Positionfile**. The calculation does *not* account for periodic boundary positions and the positions must at all times along the position trajectory be specified accordingly.

5.13 Raman

The Raman response is calculated as [31, 32]

$$I^{VV/VH}(t) = \sum_{a,b,c,d}^{x,y,z} A_{abcd}^{VV/VH} \langle \alpha_{ab}(t) U(t, 0) \alpha_{cd}(0) \rangle \exp(-t/2T_1), \quad (5.14)$$

where $A_{abcd}^{VV/VH}$ is the orientational weighting factor as defined in Ref. 32. Both the all parallel (VV) and perpendicular (VH) components are calculated. The calculation requires that the transition-polarizability (**Alphafile**) file is provided with all six independent components stored. The calculated spectra are stored in the files **Raman_VV.dat** and **Raman_VH.dat**.

5.14 SFG (sum-frequency generation)

The sum-frequency generation signal is calculated using the first-order response function [33]

$$I(t) = \sum_{a,b,c}^{x,y,z} A_{abc}^{SSP/PPP} \langle \alpha_{cb}(t) U(t, 0) \mu_a(0) \rangle \exp(-t/2T_1). \quad (5.15)$$

Both the real and imaginary parts are stored. The Fourier transform is the frequency domain absorption, which is stored in the files `SFG_PPP.dat` and `SFG_SSP.dat`. For PPP A_{abc}^{PPP} is zero except for $a = b = c = 1$, where it is one. For SSP A_{abc}^{SSP} is zero for all components except $a = z$ and $b = c = x, y$, where it is one half. T_1 is the lifetime, which is often simply used as an apodization function to smoothen the spectrum.

5.15 2DIR* (two-dimensional infrared)

This calculates the two-dimensional infrared spectra assuming coupled three level systems. The techniques `GBIR` (ground state bleach), `SEIR` (stimulated emission), and `EAIR` (excited state absorption) provides these contributions separately. Furthermore, the sum of the ground state bleach and the stimulated emission can be calculated with the `noEA` technique keyword. The expressions for the response functions are given in ref. [1].

5.16 2DSFG (two-dimensional sum-frequency generation)

Not implemented yet (check NISE_2015)

5.17 2DUVvis* (two-dimensional electronic spectroscopy)

This calculates the two-dimensional infrared spectra assuming coupled two level systems. The techniques `GBUVvis` (ground state bleach), `SEUVvis` (stimulated emission), and `EAUVvis` (excited state absorption) provides these contributions separately. Furthermore, the sum of the ground state bleach and the stimulated emission can be calculated with the `noEAUVvis` technique keyword.

5.18 CG_2DES (coarse-grained two-dimensional electronic spectroscopy)⁺

This calculates the two-dimensional electronic spectra with the coarse-grained approach. [34] This is especially efficient for large molecular systems that separate the system into different segments and calculate the 2DES based on the doorway-window function. The technique `CG_2DES_doorway` provides the doorway part of the CG-2DES results. The techniques `CG_2DES_window_SE`, `CG_2DES_window_GB`, and `CG_2DES_window_EA` provide these window functions separately. The calculation time is independent of the delay time t_2 . `CG_2DES_combine` can be used to calculate the final spectra using precalculated doorway and window functions.

The response functions governing this technique are provided in ref. [34].

If a file called `Waitingtimes.dat` is present all the waiting times given in the file (in fs) will be calculated simultaneously and the response function files will be named `RparI_0fs.dat` etc.

5.19 2DIRRaman*

This calculates the two-dimensional infrared raman spectra assuming coupled three level systems. The techniques `2DIRraman1` (rephasing), `2DIRraman2` (non-rephasing 1), `2DIRraman3` (non-rephasing 2) provides these contributions separately. The rephasing diagram and the non-rephasing diagrams are stored separately as they cannot be directly added. The expressions for the response functions are given in ref. [35].

5.20 2DFD (fluorescence detected two-dimensional spectroscopy)

The 2DFD spectrum can be calculated in the approximation that all exciton pairs annihilate to produce a single exciton long before fluorescence occurs with the `noEAUVis` technique. The response functions governing this technique are provided in ref. [36].

Chapter 6

Information for developers

Contributions from developers are welcome. The code is available on github: (https://github.com/GHlacour/NISE_2017). It is advisable to contact the main developer before planning contributions to avoid competing contributions. All new techniques should be added in new separate files which are called from the main `NISE.c` code. Implementations should as far as possible make use of existing modules for reading input and data. New propagation schemes should be added in `propagate.c`, Hamiltonians should be added through `read_trajectory.c` generally useful additions should be made in the `NISE_subs.c` module. Projection schemes should be added in `project.c`.

Bugs, feature requests, and other issues can be reported through the Issues option on GitHub or by contacting the main developer directly.

Chapter 7

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Bibliography

- [1] T. L. C. Jansen and J. Knoester. *J. Phys. Chem. B*, **110**:22910–22916, (2006).
- [2] T. L. C. Jansen and J. Knoester. *Acc. Chem. Res.*, **42**(9):1405–1411, (2009).
- [3] T. L. C. Jansen, B. M. Auer, M. Yang and J. L. Skinner. *J. Chem. Phys.*, **132**:224503, (2010).
- [4] C. Liang and T. L. C. Jansen. *J. Chem. Theory Comput.*, **8**:1706–1713, (2012).
- [5] C. Liang, M. Louhivuori, S. J. Marrink, T. L. C. Jansen and J. Knoester. *J. Phys. Chem. Lett.*, **4**:448–452, (2013).
- [6] D. Cringus, T. L. C. Jansen, M. S. Pshenichnikov and D. A. Wiersma. *J. Chem. Phys.*, **127**:084507, (2007).
- [7] T. L. C. Jansen and J. Knoester. *Biophys. J.*, **94**:1818–1825, (2008).
- [8] A. G. Dijkstra, T. L. C. Jansen and J. Knoester. *J. Phys. Chem. A*, **114**:7315–7320, (2010).
- [9] T. L. C. Jansen, D. Cringus and M. S. Pshenichnikov. *J. Phys. Chem. A*, **113**:6260, (2009).
- [10] S. Roy, M. S. Pshenichnikov and T. L. C. Jansen. *J. Phys. Chem. B*, **115**:5431–5440, (2011).
- [11] T. L. C. Jansen and J. Knoester. *J. Chem. Phys.*, **127**:234502, (2007).
- [12] P. Hamm, M. H. Lim and R. M. Hochstrasser. *J. Phys. Chem. B*, **102**:6123–6138, (1998).
- [13] R. M. Hochstrasser. *Chem. Phys.*, **266**(2-3):273–284, (2001).
- [14] M. Cho. *Chem. Rev.*, **108**:1331, (2008).
- [15] S. Mukamel. *Annu. Rev. Phys. Chem.*, **51**:691, (2000).
- [16] M. Cho. *Two-dimensional Optical Spectroscopy*. CRC Press, Boca Raton, 2009.

- [17] S. Mukamel. *Principles of Nonlinear Optical Spectroscopy*. Oxford University Press, New York, 1995.
- [18] P. Hamm and M. T. Zanni. *Concepts and Methods of 2D Infrared Spectroscopy*. Cambridge University Press, Cambridge, 2011.
- [19] Andy S. Sardjan, Floris P. Westerman, Jennifer P. Ogilvie and Thomas L. C. Jansen. *J. Phys. Chem. B*, **124**(42):9420–9427, (October 2020).
- [20] A. Paarmann, T. Hayashi, S. Mukamel and R. J. D. Miller. *J. Chem. Phys.*, **128**:191103, (2008).
- [21] M. T. Zanni, N.-H. Ge, Y. S. Kim and R. M. Hochstrasser. *P. Natl. Acad. Sci. USA*, **98**:11265, (2001).
- [22] D. J. Thouless. *Phys. Rep.*, **13**:93, (1974).
- [23] C. Olbrich, T. L. C. Jansen, J. Liebers, M. Aghtar, J. Strümpfer, K. Schulten, J. Knoester and U. Kleinekathöfer. *J. Phys. Chem. B*, **115**:8609–8621, (2011).
- [24] Andy S. Sardjan, Floris P. Westerman, Jennifer P. Ogilvie and Thomas L. C. Jansen. *J. Chem. Phys. B*, **124**(42):9420–9427, (2020).
- [25] T. L. C. Jansen. *J. Chem. Phys.*, **162**(7):074113, (2025).
- [26] S. Jang, M. D. Newton and R. J. Silbey. *Phys. Rev. Lett.*, **92**:218301, (2004).
- [27] S. Jang, M. D. Newton and R. J. Silbey. *J. Phys. Chem. B*, **111**:6807–6814, (2007).
- [28] Kai Zhong, Hoang Long Nguyen, Thanh Nhut Do, Howe-Siang Tan, Jasper Knoester and Thomas L. C. Jansen. *J. Chem. Phys.*, **158**(6):064103, (February 2023).
- [29] M. Yang and G. R. Fleming. *J. Chem. Phys.*, **113**:2823, (2000).
- [30] Hong-Guang Duan, Amy L. Stevens, Peter Nalbach, Michael Thorwart, Valentyn I. Prokhorenko and R. J. Dwayne Miller. *J. Phys. Chem. B*, **119**(36):12017–12027, (2015).
- [31] H. Torii. *J. Phys. Chem. A*, **106**:3281, (2002).
- [32] L. Shi, S. M. Gruenbaum and J. L. Skinner. *J. Phys. Chem. B*, **116**(47):13821–13830, (November 2012).
- [33] S. Roy, S. M. Gruenbaum and J. L. Skinner. *J. Chem. Phys.*, **141**(18):18C502, (November 2014).
- [34] Kai Zhong, Hoang Long Nguyen, Thanh Nhut Do, Howe-Siang Tan, Jasper Knoester and Thomas L. C. Jansen. *J. Chem. Theory Comput.*, **20**(14):6111–6124, (July 2024).

- [35] Carleen D. N. van Hengel, Kim E. van Adrichem and Thomas L. C. Jansen. *J. Chem. Phys.*, **158**(6):064106, (February 2023).
- [36] Tenzin Kunsel, Vivek Tiwari, Yassel Acosta Matutes, Alastair T. Gardiner, Richard J. Cogdell, Jennifer P. Ogilvie and Thomas L. C. Jansen. *J. Phys. Chem. B*, **123**(2):394–406, (2019).