

The thing with the Golgi apparatus

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Abstract

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Introduction

1.1 QUANTITATIVE WORK ON THE GOLGI SO FAR

1.2 RUSH SYSTEM

1.2.1 MAN II

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Introduction

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Data processing pipeline

In this chapter I present the work done on processing the rush movies. Several preprocessing steps have been undertaken to improve the quality of the fit, and we present all here. Roughly, we can divide the process in four steps:

1. Segmentation and creation of masks
2. Denoising of movies
3. Calculation of spatial and temporal derivatives
4. The actual fitting

Below we describe each step separately.

3.1 STEP 1: SEGMENTATION

The images obtained from the rush experiments often contain multiple cells. Furthermore, we can also segment the image into roughly three different types: 1) the background, where nothing of interest happens. No cells are present here, 2) the cytoplasm, which is the area where we want to fit our model and 3) the Golgi itself, where we do not necessarily want to fit. Unfortunately, no bright field images were available, making segmentation significantly harder, as no clear cell boundary can be observed. Further complicating the story is the large dynamic range of the movies due to the fluorescence concentrating in the Golgi. The following procedures we present have been developed to deal with these problems. Note that they are empirical methods, i.e. there's no theoretical background as to why they *should* work. However, in practice they do and I haven't found any other method which was able to.

3.1.1 VORONOI DIAGRAM

This method is based on a technique called Voronoi tessellation and doesn't depend on any measure of the intensity. It was developed after noting that since the cargo is spread throughout the ER in the first few frames and as the ER is roughly circumnuclear, we can use this to determine the centre of the cell (roughly). Voronoi tessellation then allows us to divide the frame into areas with just one point per area, i.e. one cell per area (theoretically). More precise, given n coordinates, voronoi tessellation divides the given area into n pieces, where every point in a piece is closest to one coordinate. In practice this means for us that each point in a cell area is closest to its the given cell centre. Figure **ref** shows this. Each calculated cell centre is a red point and the lines depict the borders between each voronoi cell. Assuming the cells don't move too much, they don't cross the cells and thus we apply the voronoi diagram calculated in the first few frames to the entire movie.

3.1.2 INTENSITY

For the fitting however we wish to make a slightly better approach than a voronoi diagram. As stated, we can't find the exact delineation of the cell, but looking at the intensity, we can see an 'area' of interest, separating background from the cell. Since the Golgi is quite bright in the last 200 or so frames, we consider only the intensity for the Golgi, while for the cytoplasm we consider both the intensity and its time derivative. Thus we have two analog but different processes. For the Golgi we do the following:

1. Renormalize the concentration C between 0 and 1.
2. Sum all frames. One then obtains an image such as figure **ref**

$$\sum_{frames} C(x, y, t)$$

3. This image is thresholded, either through an otsu threshold or a manual one, until the mask roughly matches what we want. Note that extreme precision isn't required, since we just want the rough area. This results in figure **ref**

For the cytoplasm we follow the same procedure only now we take the log of sum of the product of the intensity and its time derivative:

$$\log \left(\sum_{frames} C(x, y, t) \cdot \partial_t C(x, y, t) \right)$$

We thus obtain a complete mask for the movie as shown in figure **ref**

3.2 STEP 2 - DENOISING

In order to accurately calculate the derivatives and generally improve the quality of fitting, we wish to denoise and smooth the obtained movies. Denoising and

smoothing is a subject about which many books have been written and there are hundreds of approaches. One oft-used technique is to Fourier transform the signal, cutoff all coefficients above a cutoff frequency and retransform back into the real domain. Next, a Savitzky-Golay filter can be used to finally smooth the result. However, a big issue with all these methods is their non locality. Since our movies have different scales, this is a big problem. Furthermore, they often smooth out sharp peaks. After evaluating several methods, I have settled on a relatively new method presented in **ref**.

The so-called WavinPOD method combines two well-known filtering techniques, known as wavelet filtering and Proper Orthogonal Decomposition. Below we explain each separately. Our explanation is adapted from **ref** and **ref**.

WAVELET FILTER

A wavelet filter is not really the appropriate name, as its more of a transform.

PROPER ORTHOGONAL DECOMPOSITION

Proper orthogonal decomposition is a technique similar to what is known as Principal component Analysis in statistics and falls into the general category of model reduction techniques. It's often used in flow problems to extract coherent structures from turbulent flows. Simply put, in POD we wish to express a function as

WAVINPOD

WavinPOD combines these two techniques in the following way. First, we decompose our problem with a POD transformation. This yields a set of temporal and spatial modes. We select the most energetic modes and wavelet filter these, before transforming them back to the real domain. As shown in **ref**, combining these techniques has an advantage over others.

In our case, we select the number of modes to be used by hand (30 in the case of MANII) and apply a 3-level db4 wavelet. We use a slightly higher than necessary level to increase smoothness. In the figure below we show the result for both a pixel in time and one time snapshot. Note that the result is significantly smoother, but that smaller details have been preserved.

3.3 STEP 3 - DERIVATIVES

Taking spatial and temporal derivatives of these images is not an entirely trivial operation due to the discreteness of the system. More specifically, taking numerical derivatives of data is extremely hard to do properly and becomes even harder in the presence of noise. Next to basic finite difference methods, one can for example use a linear-least-squares fitted polynomial, smoothing spline or a so-called tikhonov-regularizer **ref needed**. Each method comes with its strengths and weaknesses, but one particularly nasty thing for our context is that they don't scale well to higher dimensions and quickly become computationally expensive.

Another issue related to discretization is the size of the grid w.r.t. the size of the features. To see this, we plot a 2D-gaussian with $\sigma = 1$ in figure **ref**.

As expected, the derivative is normal to the isolines of the object. Now consider the discretized version of the object. Taking the naive spatial derivative w.r.t. to each direction means only considering a single row or column of and taking the derivative in that direction. Figure **ref** shows the result of this operation. An artifact is clearly visible: instead of a nice uniform derivative, we see a 'cross'. This effect is a cause of the discretization grid being too large for some smaller, often bright, objects.

To remedy this, one can for example artificially upscale the grid, interpolate the values inbetween, and take the derivatives from this grid. This is not ideal however, since the upscaling requires a large amount of memory and is computationally expensive. Another solution which is common in image processing is applying a

kernel operator. The advantage of a kernel operator is that it is extremely computationally cheap, as it involves convolving the original picture with a differentiation kernel. The differentiation kernel is an approximate version of a finite difference scheme. We use and show here the Sobel filter, which is the most commonly used one.

In a simple finite central difference scheme, we set

$$\frac{dx}{dt} \approx \frac{x_{i+1} - x_{i-1}}{2h}$$

where h is the distance between two points. In terms of a kernel operator, this would look like (the h drops out as the distance in terms of pixels is 1):

$$\frac{1}{2} \cdot \begin{bmatrix} 1 & 0 & -1 \end{bmatrix}$$

And applying it by convoluting it to a matrix gives the x-derivative:

$$\partial_x A \approx A * \begin{bmatrix} 1 & 0 & -1 \end{bmatrix}$$

and analogous for the y-direction. However, as we've seen, looking at just a single row introduces cross-like artifacts. To remedy this, we wish to include diagonal pixels as well. However, the distance between the diagonal pixels and the center pixel is not 1 but $\sqrt{2}$ and furthermore we need to decompose it into \hat{x} and \hat{y} , introducing another factor $\sqrt{2}$. Thus, one obtains the classis 3×3 Sobel filter **ref**:

$$\mathbf{G}_x = \frac{1}{8} \cdot \begin{bmatrix} 1 & 0 & -1 \\ 2 & 0 & -2 \\ 1 & 0 & -1 \end{bmatrix} \quad \mathbf{G}_y = \frac{1}{8} \cdot \begin{bmatrix} 1 & 2 & 1 \\ 0 & 0 & 0 \\ -1 & -2 & -1 \end{bmatrix}$$

Although not extremely accurate, the Sobel filter seems to do the tricks for us. Several other versions such as Scharr or Prewitt exist, offering several benefits such as rotational symmetry, but we have not pursued these. They just change the

coefficients. Although we have shown a 3×3 filter here, the filter can take into account higher order schemes such as a 5×5 or 7×7 . The major benefit of the spatial derivatives as a convolution operator is its computational efficiency: convolutional operations are performed parallel and are extremely fast.

For the time derivative, we apply a second order accurate central derivative scheme, while for the spatial derivatives (both first and second order) we apply the 5×5 Sobel filter. We analyze these in the next chapter

3.4 STEP 4 - FITTING

Now that we have gathered all our data we can use it to fit. We use a simple least squares method.

4

Results data analysis

Here we present the results from our analysis on the RUSH experiments. We only show the results of MANII because this is the only thing we studied.

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4.1 ANALYSIS OF TIME DERIVATIVES

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4.2 ANALYSIS OF FIT

4.2.1 DIFFUSION

4.2.2 ADVECTION

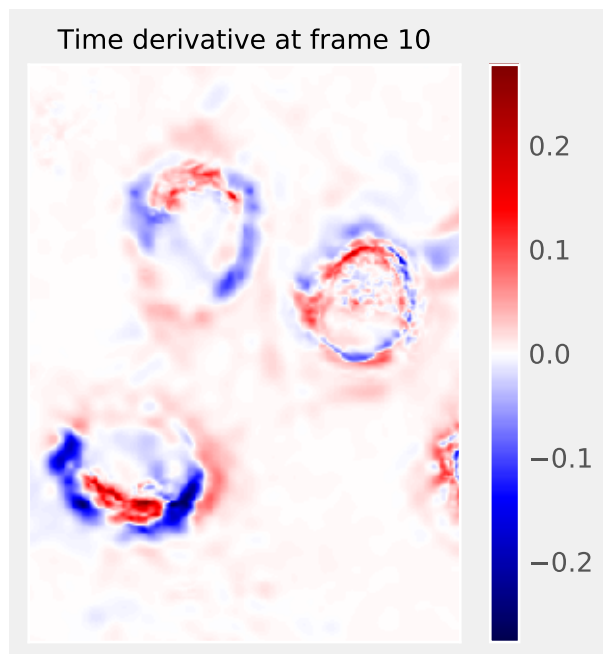


Figure 4.1.1: Wolla zieke subscript.

5

Physics Informed Neural Networks

In previous chapters we showed the difficulty in finding a general fit for a model to spatiotemporal data. Very recently, a set of papers introduced a new technique called Physics Informed neural networks. The set of papers show very powerful and promising results. I've evaluated this technique for use in fitting our model to the RUSH data. Since for many Neural networks are a new technique, this chapter is also a gentle introduction into Neural networks themselves. We have three parts:

- Neural Networks
- Physics Informed Neural networks
- Conclusion

In neural networks we introduce the ideas and some mathematics behind neural

networks in Physics informed neural networks we showcase the concept using a simple toy model and in conclusion we present the evaluation of the technique. Since Neural networks are completely new concept to most, we use a light tone.

5.1 NEURAL NETWORKS

Good source: <http://neuralnetworksanddeeplearning.com/chap2.html>

5.1.1 ARCHITECTURE

The basic of each type of neural network is the same: the neuron. Inspired, but not the same as a biological neural network, the neuron has an input x and transforms it to an output y : $x \rightarrow y$. This is a three step process. The input x is multiplied by a factor a known as a weight and adds a constant known as the bias b . This serves as an input to the *activation function* g of the neuron, so that

$$y = g(ax + b)$$

the role of the activation function is to introduce non-linearity into the system. Many research is ongoing into different activation, but one of the most used is the $\tanh(x)$ function, due to its biological relevance and bounded in and outputs. Several other functions with more favourable properties such as lower computational cost are available as well.

A neuron always has one output, but can have multiple inputs, for example in the case of multiple neurons feeding into another one. One can also have multiple neurons in a single layer. A natural generalization is that a and b become matrices and x and y vectors:

$$\vec{y} = g(A\vec{x} + B)$$

where g is applied element wise. The goal of training a neural network is to find the matrices A and B which make sure that y is closest to the desired outcome.

5.1.2 TRAINING

GRADIENT DESCENT

As the name machine learning indicates, neural networks need to *learn*. When this nomenclature is used, it means that we need to adjust the weights and biases in the network until the output matches the wished-for output. In general this a problem with local minima, but a solution may be found using gradient descent techniques. Consider a function $f(\mathbf{x})$, we which we wish to minimize with respect to \mathbf{x} . One starts with an estimate of \mathbf{x}_0 , which we iteratively refine

$$\mathbf{x}_{n+1} = \mathbf{x}_n - \gamma \nabla f(\mathbf{x}_n)$$

untill the gradient vanishes. This position is where f is minimized w.r.t \mathbf{x} . γ is known as the learning rate and sets the step rate. More advanced techniques can set a variable learning rate etc, but the basis remains similar.

BACK PROPAGATION AND AUTOMATIC DIFFERENTIATION

In the case of neural networks we wish to minimize the cost w.r.t. to each weight w_{jk}^l and bias b_j^l . To do this computationally is not trivial, and most of the NN field uses one algorithm: back propagation. We give a short introduction below. We want to know basically two different things:

$$\frac{\partial C}{\partial w_{jk}^l}, \frac{\partial C}{\partial b_j^l}$$

the change of the cost w.r.t to each weight and bias. Let's define the error neuron j in layer l as $\delta_j^l = \partial C / \partial z_j^l$. We can rewrite this using the chain rule as:

$$\delta_j^l = \sum_k \frac{\partial C}{\partial a_{jk}^l} \frac{\partial a_{jk}^l}{\partial z_j^l}$$

However, the second term is always zero except when $j = k$, so we drop the sum. We also see that the second term is equal to $\sigma'(z_j^l)$. For the last layer L , the first term turns into the derivative of the cost function, finally giving us:

$$\delta_j^L = |a_j^L - y_j| \sigma'(z_j^L)$$

This expression gives us the error in the output layer in terms of its weighted input. This in turn is a function of previous inputs and errors and we thus need to find an expression relating the error in layer l with the error in an layer $l + 1$. Since we have an expression for the error in the last layer, we calculate the errors going down the layers (from L to 0), hence the name backpropagation. Again using the chain rule gives:

$$\delta_j^l = \sum_k \frac{\partial C}{\partial z_{jk}^{l+1}} \frac{\partial z_{jk}^{l+1}}{\partial z_j^l} = \sum_k \delta_k^{l+1} \frac{\partial z_{jk}^{l+1}}{\partial z_j^l}$$

Using the definitions of z_{jk}^l , we finally after substitution obtain:

$$\delta_j^l = \sum_k \delta_k^{l+1} w_{kj}^{l+1} \sigma'(z_j^l)$$

Using these two equations, we can calculate the error in C due to each neuron.

Finally, to calculate the $\partial C / \partial w_{jk}^l$ and

$\partial C / \partial b_j^l$, we relate these to the error, again via the chain rule:

$$\begin{aligned} \frac{\partial C}{\partial b_j^l} \frac{\partial b_j^l}{\partial z_j^l} &= \frac{\partial C}{\partial b_j^l} = \delta_j^l \\ \sum_k \frac{\partial C}{\partial w_{jk}^l} \frac{\partial w_{jk}^l}{\partial z_j^l} &= \delta_j^l \rightarrow \frac{\partial C}{\partial w_{jk}^l} = a_j^{l-1} \delta_j^l \end{aligned}$$

These four equations together make up the the backpropagation algorithm. The complete optimization of the network goes as follows: 1. Complete a forward pass, i.e., calculate the expected outcomes with the current weights and biases. 2.

Calculate the error and back propagate it to obtain the gradients in the weights and biases. 3. Adjust the weights using optimizer (e.g. gradient descent) 4. Return to step 1 and redo the cycle until convergence.

Officially, back propagation is a special case of a technique known as automatic differentiation, which is third type of differentiation, next to numeric and symbolic. It also for machine precision calculation of derivatives by writing it as a chain of simple operations combined with the chain rule, similar to backpropagation. Note that:

$$\delta_j^\circ = \frac{\partial C}{\partial x_j} \frac{\partial x_j}{\partial z_j^\circ}$$

so that:

$$\frac{\partial C}{\partial x_j} = a_j^\circ \delta_j^\circ$$

Thus when learning, we also have access to high precision derivatives with regard to each coordinate!

5.2 PHYSICS INFORMED NEURAL NETWORKS

On the face of things, physics and neural networks seem to satisfy two completely different goals. Whereas physics tries to build (simplified) models, neural networks try to learn a general modelless mapping from the inputs to the outputs. However, recently some approaches have emerged which fuse these seemingly opposite goals in order to do two different things:

1. Use a neural network to simulate/numerically solve equations.
2. Use a neural network to 'fit' a model to spatiotemporal data and even infer a coefficient field.

Initial results have shown very promising: using neural networks to numerically solve models doesn't require any advanced meshing and careful handling of

shocks, whereas the ability to infer coefficient fields from spatiotemporal data hasn't been shown at all to my knowledge.

5.2.1 THE CONCEPT

Consider a set of 1+1D spatiotemporal data, consisting of some property u_i at coordinates (x_i, t_i) . As stated before, a neural network learns the mapping $x, t \rightarrow u$ because it is a universal function approximator through minimizing the mean squared error:

$$MSE = \sum_i (u_i^{in} - u_i^{pred})^2$$

Now assume that we know that $u(x, t)$ is governed by some process which can be modeled as a partial differential equation. We can write (almost) every PDE as:

$$\partial_t u = f(1, u, u_x, u_{xx}, u^2, \dots)$$

where f is some sort of function of u or its spatial derivatives. Now we rewrite it as:

$$g = -\partial_t u + f(1, u, u_x, u_{xx}, u^2, \dots)$$

To satisfy the model, the function g always has to go to zero. The idea behind Physics Informed Neural Networks (PINNs) is that we can include this function as an extra cost to be minimized:

$$cost = \sum_i (u_i^{in} - u_i^{pred})^2 + \sum_i g_i^2 = MSE + PI$$

Since to satisfy the model we need $g \rightarrow 0$, by adding our second 'Physic-informed' term, we effectively penalize solutions not satisfying the

physics we put in: our new term acts as a regularizer. More concretely, this term has two effects:

1. **It prevents overfitting:** Neural Networks can be prone to overfitting. This term prevents that by penalizing variations not described by the physics.
2. **It makes the NN more data-efficient:** we can get good results with not much data.
3. **It allows fitting and prediction:** we can fit and predict based to the terms in f .

The first point is rather technical and interesting for more in-depth, so we'll leave it for now. The second point is very interesting from an experimental point of view. By presupposing some structure in the data in the form of a model, we need significantly less data to get the same result. It's also here that we see the no free lunch theorem: a price has to be paid. By encoding a model into the neural network, we lose the freedom of the neural network to map $x, t \mapsto u$ using any function. For us physicists however, this is actually a blessing but for applications where equationless modeling is useful this is terrible. Note that we also circumvent the issue of numerical derivatives, since we can use the network provided automatic derivatives at machine precision. This combination provides the power of PINNs - a one-two punch consisting of physics regularization and automatic derivatives.

FITTING AND PREDICTION

Point three is however where PINNs really shine. How can the extra term we've included have these effects? To see this, consider the classics physics exercise of calculating the trajectory of a launched object. In this case we know the function g , i.e. $\ddot{y} = -g$. A classical numerical solver would take small steps in time, updating the position and speed of the object each step according to g . A PINN however uses a completely different approach. Given the initial state of the neural

network, it calculates a first trajectory but then keeps adjusting the weights of the network until the cost is minimized, i.e. g goes to zero everywhere. In other words: the Neural network keeps launching the object and adjusting its internal weights until the physics are satisfied at every step of the trajectory. The classical numerical approach goes for one correct try; the neural network just keeps trying until it converges on the correct solution. This approach allows us at the same time to circumvent complex discretization schemes and issues such as solutions blowing up. Such an approach is possible because of the backpropagation algorithm, which allows us to calculate the part of each neuron in the total cost and in which direction to change the weights and biases to minimize the cost.

Interestingly, we can also use this framework to fit models to spatiotemporal data by letting the coefficient of each term be a variable w.r.t to which we minimize too. More concretely, where before the cost was a function of the weights and biases, $cost = f(w_{jk}^l, b_k^l)$, we now let it be a function of the coefficients λ of the PDE as well: $cost = f(w_{jk}^l, b_k^l, \lambda)$. This is shown for all kinds of systems in the papers of [ref]. In theory however, it should also be possible to infer coefficient *fields*, both spatially and temporally varying. We can achieve this by a multi-output Neural network. Instead of outputting just c , we also output the coefficient at that spatiotemporal point, as shown in figure [ref]. We investigate this claim in the next section.

5.2.2 PINNs IN PRACTICE

5.2.2.1 CONSTANT D

5.2.2.2 VARYING D

5.2.2.3 REAL CELL?

5.3 CONCLUSION

5.3.1 WEAK POINTS AND HOW TO IMPROVE

6

Conclusion

Appendix 1: Some extra stuff

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References