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ROUND II Questions & Solutions

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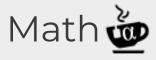
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Art of Problem Solving





DISCUSSION

<u>Discussion</u> <u>eforum.gramoly.org</u> The scoring criteria is as follows: There are no negatives for integer valued problems:

In case of Multiple choice problems, partials will be awarded as follows Marking Scheme: Full marks, +4 if all the correct options are marked. Partial marks, +1 for each correct option provided NO incorrect option is marked. Zero marks, if no option is marked. Negative marks, -1

The raw score of the problem is scaled based on its difficulty as follows: In case of integer valued.

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$$Marks = max(8.5 - ln(N), 4)$$

Where N is the number of people who got the question correct **In case of Multiple choice questions:**

$$Marks = max\left(\left(\frac{8.5 \cdot A}{4} - ln(N)\right), A\right)$$

Where N is the number of people who got the question correct A is the the aggregate marks the person has scored in the question after taking partials into account.

Selection criteria for round 2:

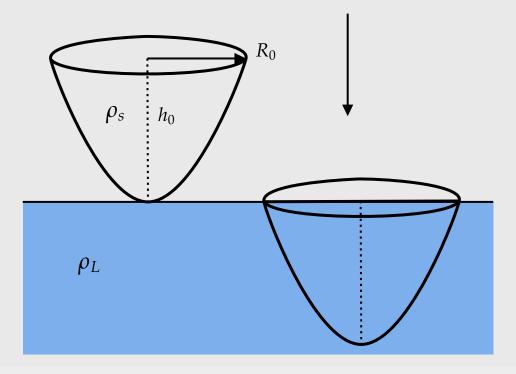
After every participant has been scored the following criteria will be used for selection:

$$\textbf{Final Marks} = \frac{\text{Your marks in round 1}}{\text{Average Marks scored in round 1}} \cdot 10 + \frac{\text{Your marks in round 2}}{\text{Average Marks scored in round 2}} \cdot 90$$

The top 50 from the final mark list will be selected for round 3.

Problem 1

A solid paraboloid of base radius R_0 and height h_0 is having uniform volume mass density ρ_s is inverted and placed just above the surface as shown, of a liquid having volume mass density ρ_L . When paraboloid is released from rest it has been observed that when it becomes completely submerged in the liquid the paraboloid comes to rest again. What is the ratio of $\frac{\rho_L}{\rho_s}$ for such a paraboloid. There exist a uniform gravity of g. Assume that the liquid body is very large so that liquid level is not changing as the paraboloid enters into liquid.



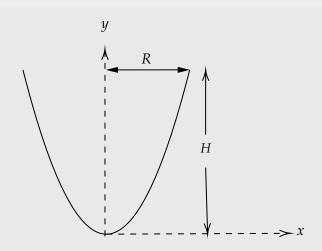
-Proposed by Nitin Sachan

Answer: 3

Net Force on the paraboloid:

$$mg - \rho_L \cdot \frac{1}{2}\pi x^2 y \cdot g = m \cdot \frac{dv}{dt}$$

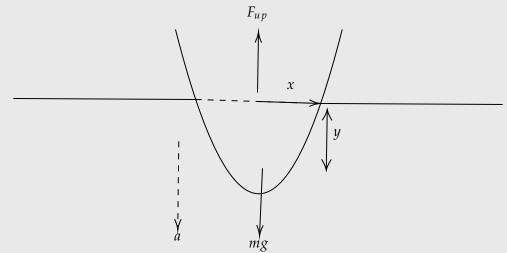
$$v \cdot \frac{dv}{dy} = \frac{dv}{dt} = g - \frac{\rho_L}{\rho_S} \cdot \frac{x^2}{R^2} \cdot \frac{yg}{H} \implies g - \frac{\rho_L}{\rho_S} \cdot \frac{yR^2}{H \cdot R^2} \cdot ygH$$



$$v \cdot \frac{dv}{dy} = g - \frac{g \cdot \rho_L}{H^2 \rho_S} \cdot y^2$$

 $\begin{array}{l} v \cdot \frac{dv}{dy} = g - \frac{g \cdot \rho_L}{H^2 \rho_S} \cdot y^2 \\ \text{Integrating this using } v_i = vf = 0 \text{ and putting appropriate limits we get:} \\ \int_0^0 v \cdot dv = g \cdot \int_0^H dy - \frac{g \rho_L}{H^2 \rho_S} \cdot \int_0^H y^2 dy \end{array}$

$$\int_0^0 v \cdot dv = g \cdot \int_0^H dy - \frac{g\rho_L}{H^2 \rho_c} \cdot \int_0^H y^2 dy$$



$$0 = gH - \frac{g}{3} \frac{\rho_L}{\rho_S} \cdot H$$
Hence,
$$\frac{\rho_L}{\rho_S} = 3$$

Also an energy conservation approach can be applied to shorten the solution.

Solution 1

Problem 2

You have a spectrometer of spectral resolving power $5 \cdot 10^4$. Say n_1 is the maximum value of principal quantum number for which the resulting energy spectra can be distinguished clearly from its neighbours. Also n_0 denotes the minimum value of principal quantum number for which a human eye can see the atomic transition. Find the value of $n_1 - n_0$. We are only working in the Balmer transitions of hydrogen atom.

Value of Rydberg constant is $R = 1 \cdot 10^7$ and maximum wavelength eye can see is 700nm

Proposed by Abhiram Cherukupalli

Answer: 69

$$\frac{1}{\lambda} = R(\frac{1}{4} - \frac{1}{n^2}), \ n \ge 3$$
$$\frac{\Delta \lambda}{\Delta n} \approx \frac{d\lambda}{dn}$$

Note that though this is an approximate method, it can be verified by desmos, and the answers are close.

Now using this logic we can differentiate () to get:

$$\frac{d\lambda}{dn} = -\frac{2R}{n^3}\lambda^2 \approx \frac{\Delta\lambda}{\Delta n}$$

The beautiful step here is to take $\Delta n = 1$, as it has to be distinguished from its neighbours. This gives us:

$$\frac{\Delta\lambda}{\lambda} \approx \frac{2R}{n_1^3}\lambda$$

The spectral power

$$\frac{\lambda}{\Delta \lambda} = \frac{{n_1}^3}{2} (\frac{1}{4} - \frac{1}{n^2}) < 5 \cdot 10^4$$

$$\implies n_1 < 73.69 \implies \boxed{n_1 = 73}$$

Calculation of n_0 is straight-forward:-

$$n_0 = \sqrt{\frac{4R\lambda}{R\lambda - 4}} = 3.055$$

Hence $n_0 = 4$ and $n_1 - n_0 = 69$

Note that we have to use the value of R mentioned in the question not the exact value also note that though the human eye data is close the value is very easy to calculate by hand as we get $\frac{4R\lambda}{R\lambda-4} > 9$ without need of much calculation as $R\lambda = 7$.

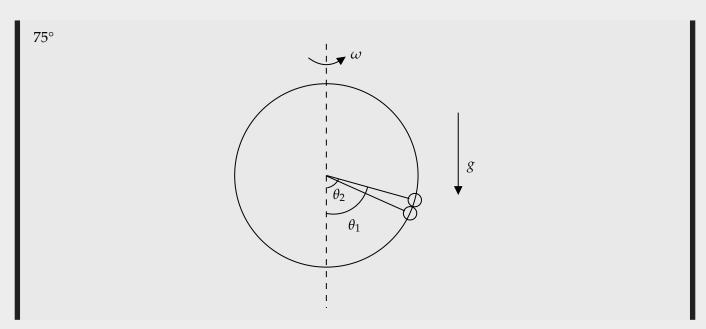
Solution 2

Problem 3

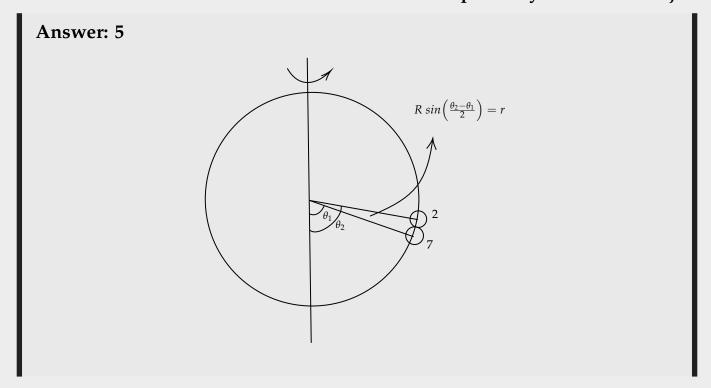
We have a loop of radii 80 cm with 2 identical beads of mass 9 kg and radius 6.97 cm inserted in it, free to move without friction. The loop starts rotating along the diameter with an angular velocity

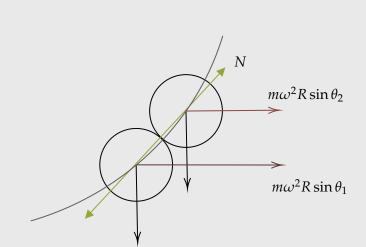
 $\omega = 5 \text{ rad/s}$. If the angle lower bead makes θ_1 and upper bead makes θ_2 then report $\left\lfloor \frac{\theta_1 + 3\theta_2}{50} \right\rfloor$. You may assume that both beads are displaced to the same side of loop.

Use $\arcsin\left(\frac{6.97}{80}\right)\approx 5^{\circ}$, $\cos 5^{\circ}\approx 0.9962$, $\cos 10^{\circ}\approx 0.9848$, $\arccos(0.5058)\approx 60^{\circ}$ and $\arccos(0.2529)\approx 10^{\circ}$

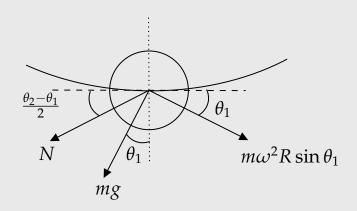


-Proposed by Atharva Mahajan



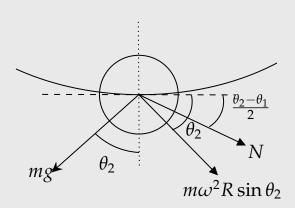


For Ball 1:



$$N\cos\frac{\theta_1-\theta_2}{2}=m\omega^2R\sin\theta_1\cos\theta_1-mg\sin\theta_1$$

For Ball 2:



$$N\cos\frac{\theta_1-\theta_2}{2}=mg\sin\theta_2-m\omega^2R\sin\theta_2\cos\theta_2$$

 $m\omega^2 R \sin \theta_1 \cos \theta_1 - mg \sin \theta_1 = mg \sin \theta_2 - m\omega^2 R \sin \theta_2 \cos \theta_2$

After solving, we get $\theta_1 = 55^{\circ}$ and $\theta_2 = 65^{\circ}$

So,
$$\left| \frac{\theta_1 + 3\theta_2}{50} \right| = 5$$

Note that we can also directly find the position of centre of mass of the two balls to find the angles easily.

Solution 3

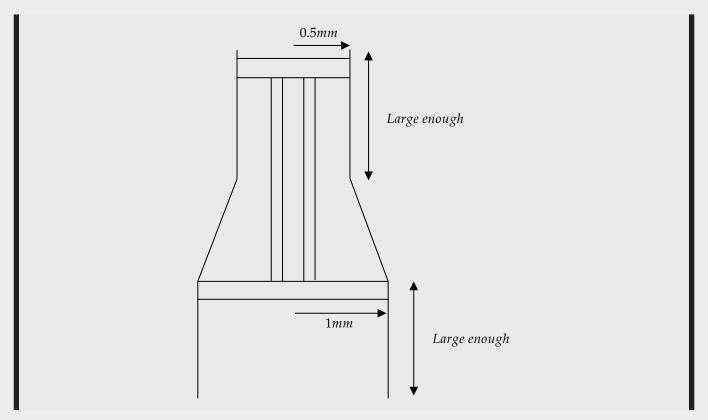
Problem 4

The bouncing container

There is a container containing two circular discs (of radii $0.5 \, mm$ and $1 \, mm$ respectively) which are connected by two light rigid rods of negligible area. The space between them is filled by a unknown ideal liquid, This container is dropped vertically on a table from a height 20cm, with which it inelastically collides. Assume that the container loses all its kinetic energy into rest. It remains at rest for a time t seconds and then it amazingly shoots back up to the same height

Report value of $\lfloor 10 \cdot t \rfloor$ in seconds (round off to nearest integer).

Neglect the hydro-static pressure of the liquid, assume pressure everywhere outside the space between the two discs to be ambient pressure $10^5 Pa$. Neglect all types of frictional forces i.e the discs are free to move up and down. The mass of the disc liquid system is 15 g. Assume the container is large enough so that the bottom piston don't fall out and also the top piston does not reach the conical middle. **Assume the conical Portion of the container is small**



-Proposed by Abhiram Cherukupalli

Answer: 7

The assumption here was to take conical section was of negligible height.

When the impact happens the container comes to rest, but the pistons continue moving, since upper piston has lower area the pistons moving down will cause the volume between them to increase. As the liquid is in compressible, a vacuum will form between the pistons and the force due to pressure difference will provide the necessary impulse for bouncing. Say ΔA is the difference in area of the pistons and M_{total} is the total mass of the system then:

The upward force acting on the water + disc system is:-

$$F = \Delta A P_0 - M_{total} g$$

Impulse equation

$$F \cdot t = 2M_{total} \sqrt{2gh}$$

$$t = \frac{2M_{total}\sqrt{2gh}}{\Delta AP_0 - M_{total}g} \approx 0.7$$

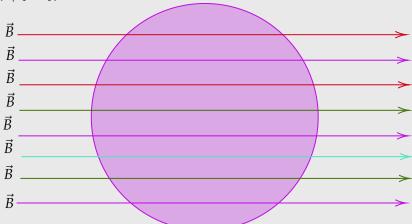
Solution 4

Problem 5

There is a superconducting ring,made up of perfectly conducting wire of mean radius r=2m and cross-sectional area A and of mass 10kg which floats in a gravity free space in a region of uniform magnetic field parallel to its induction vector $\vec{B} = 2T$. Inductance L of the ring is so small that inertia of free electrons cannot be neglected in the current building process. The free electron density in the conductor is n, mass of an electron is m_e and modulus of charge on an electron is e. He found that if the ring is turned slightly about a diameter and released, it will execute simple harmonic motion. If calculated time period is in form of

$$\frac{\sqrt{a\left(L+\frac{b\pi m_e}{ne^2A}\right)}}{C}$$

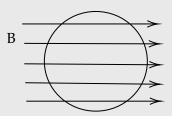
Find the value of a + b - c.



-Proposed by AKIII

Answer: 7

Concept: Given Inductance L of super conducting ring is so small that inductance due to free electrons cannot be ignored.



We'll use the concept of Kinetic inductance.

Here, L' is the kinetic inductance.

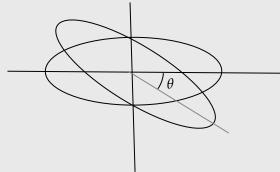
Kinetic energy of electron $=\frac{1}{2}L'i^2 \implies \frac{1}{2}Mv_e^2 = \frac{1}{2}L'i^2$ Now, we know, $i=neAv_d \implies v_d = \frac{i}{neA}$ Also, by definition, $M=nm_eA\cdot 2\pi R$

Hence, from our previous equation, we can write

$$\frac{1}{2}nm_eA \cdot 2\pi R \frac{i^2}{n^2e^2A^2} = \frac{1}{2}L'i^2 \implies \boxed{L' = \frac{2\pi Rm_e}{ne^2A}}$$

So, $L_{\text{net}} = L + L'$ Now the ring is rotated by a small angle, we can calculate the torque τ due to the magnetic field :

$$\tau = \pi R^2 \cdot iB \cos \theta = \frac{B^2 \sin \theta \cos \theta (\pi R^2)^2}{L_{net}}$$



Since its a super conducting ring, flux is constant thus, $L_{net}i = B \sin \theta \pi R^2$ Hence for small angle θ , S.H.M. :

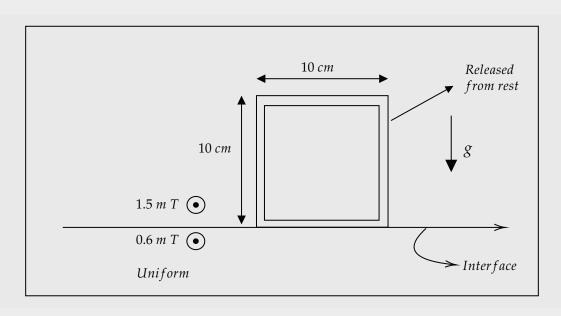
$$T = 2\pi \sqrt{\frac{\frac{mR^2}{2}L_{net}}{B^2(\pi R^2)^2}} = \frac{2}{BR} \sqrt{\frac{mL_{net}}{2}} = \frac{\sqrt{2mL_{net}}}{BR}$$

Solution 5

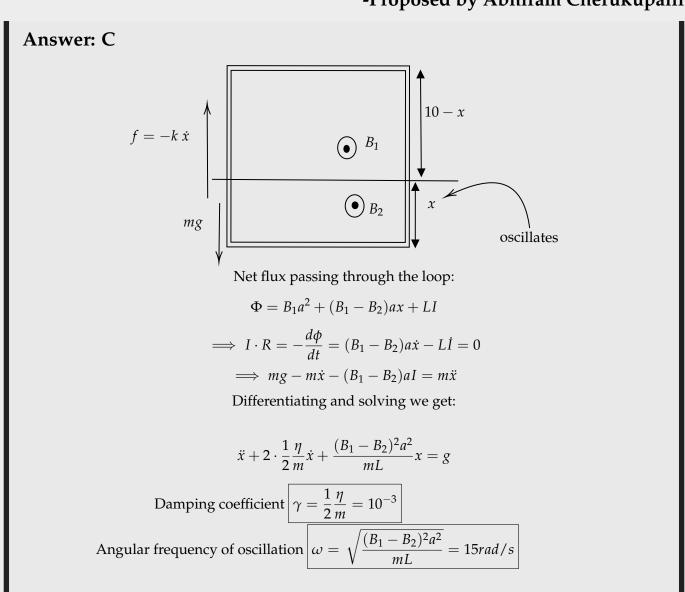
Problem 6

A square loop of negligible resistance, side length 10 *cm*, mass 0.9*g*, total self inductance $4 \cdot 10^{-8} H$ is left in a vertical plane where where its bottom end touches the interface of two different uniform horizontal magnetic fields ($B_{upper} > B_{lower}$) where $B_{upper} = 1.5mT \hat{k}$ and $B_{lower} = 0.6mT \hat{k}$. There is also a drag force, f = -kv ($k = 1.8 \cdot 10^{-6}$) acting on it. This square loop happens to undergo a damped simple harmonic motion, mark the correct options:

- A The damping coefficient of the resulting oscillations γ is $1.5 \cdot 10^{-3} \ Ns/m$
- B The angular frequency omega of the oscillations ω is 10 rad/s
- C The maximum magnetic field of the lower region for SHM to ensue is $9 \cdot 10^{-4} T$
- D The vertical position $z(t) \approx 0.1 \left(1 \frac{e^{0.0015t}}{\cos \phi} \left(\cos \left(5t \phi\right)\right)\right)$, $\phi = \arctan\left(\frac{\gamma}{\sqrt{\omega^2 \gamma^2}}\right)$



-Proposed by Abhiram Cherukupalli



Vertical position as a function of time
$$\boxed{z(t) = 4.44 \left(1 - \frac{e^{0.001t}}{\cos\phi}(\cos(15t - \phi))\right) cm}$$
 where $\phi = \arctan(\frac{\gamma}{\sqrt{\omega^2 - \gamma^2}})$ For SHM to ensue:

$$\frac{g}{\omega^2} < a \implies |B_1 - B_2| > 0.6mT \implies \boxed{B_2 < 0.9mT} \text{ or } B_2 > 2.1mT$$

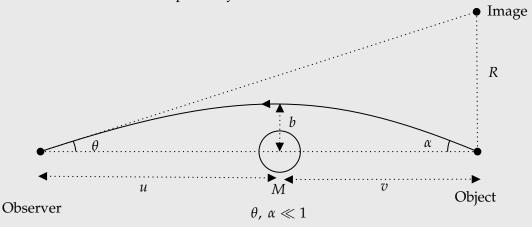
But since $B_{upper} > B_{lower}$ we don't consider the second possibility as it is the best possible option.

Solution 6

Problem 7

In this problem, we will be analysing the formation of image of ring around galaxies using Corpuscluar Theory and Universal Law of Gravitation. (Note: Do not consider any aspects of Relativity while solving this problem, i.e, solve it using only Classical Mechanics)

Consider light as a collection of particles of mass m moving with a very high speed c ($c^2 >> \frac{GM}{b}$, where b is the impact parameter). These particles will deviate when they come under the interaction of a star as shown in figure. Consider a simple setup with one star acting as a lens. Let the distance of object and observer be *v* and *u* respectively. Mark the correct statements:



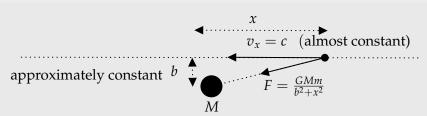
- **A)** The radius of ring observed will be $\sqrt{\frac{2GMv(u+v)}{c^2u}}$
- **B)** The maximum radius of star which can create such a image is $\sqrt{\frac{4GMvu}{c^2(u+v)}}$
- C) For v >> u, the value of angle subtended by ring at observer will be $\sqrt{\frac{4GM}{uc^2}}$
- **D)** The value of $\alpha + \theta$ is $\frac{2GM}{hc^2}$

-Proposed by Prerak

Answers: (A), (D)

First, let's find the deviation angle for impact parameter "b"





Net impulse in y-direction, can be calculated as:

$$\int_{-\infty}^{\infty} \frac{b}{\sqrt{x^2 + b^2}} \frac{GMm}{x^2 + b^2} dt$$

Here, x = ct

$$\implies$$
 Net impulse in y-direction $=\int_{-\infty}^{\infty} \frac{GMmb}{(c^2t^2+b^2)^{3/2}} = \boxed{\frac{2GMm}{bc}}$

So angle of deviation $\approx \left| \frac{2GM}{bc^2} \right|$

Now, R

$$\alpha, \beta \ll 1$$

By exterior angle property, it's clear that

$$\implies \frac{2GM}{bc^2} = \frac{b}{u} + \frac{b}{v} \implies b^2 = \frac{2GM}{c^2} \frac{uv}{u+v} = \theta^2 u^2$$

Hence, we can write θ as

$$\theta = \sqrt{\frac{2GMv}{c^2u(u+v)}}$$

For calculating *R*, we can use basic trigonometry.

$$R = (u + v) \tan \theta \approx \theta (v + u)$$

On substituting the value of θ obtained earlier

$$R \approx \theta(v+u) = \sqrt{\frac{2GMv(u+v)}{uc^2}}$$

Solution 7

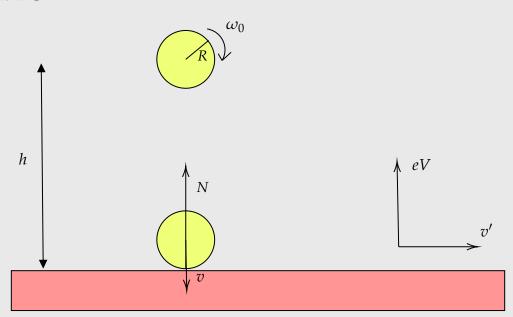
Problem 8

A ball of (R = 10m) from height h = 100m on a surface having coefficient of kinetic friction as $\mu_k = 0.5$. After the first collision with the surface, it rebounds (coefficient of restitution = 0.5) and also moves a maximal horizontal distance of x = 1.29m. Assuming the ball to be a uniform solid sphere, choose the correct options:

- (A) It'll slip on the surface for the complete duration of impact.
- (B) It won't slip for complete duration of impact.
- (C) Maximum angular velocity being given is 0.1 rad/s.
- (D) Maximum angular velocity being given is 0.2 rad/s.

-Proposed by Atharv Shivram Mahajan





$$\int Ndt = mv(1+e)$$

Now first lets say it slips for all the time during collision.

$$\mu \int Ndt = mv'$$

 \implies Time of flight before 2nd collision = $\frac{2eV}{g}$

$$\implies$$
 Range = $x = \frac{mu \int Ndt}{m} \cdot \frac{2eV}{g}$

 \implies Range = $x = \frac{mu \int Ndt}{m} \cdot \frac{2eV}{g}$ On substituting with values and simplifying, we get: x = 100m but x = 1.29m is given. Hence, it can't slip for all the time.

$$\therefore \mu.R. \int_0^t Ndt = I\left(\omega_0 - \frac{v'}{R}\right)$$
and
$$\mu \int_0^t N.dt = mv'$$

Using these two equations, we get:

$$R = \frac{I}{mv'} \left(\omega_0 - \frac{v'}{R} \right)$$

So,

$$\omega_0 = v' \left(\frac{mR^2}{IR} + \frac{1}{R} \right)$$

$$\omega_0 = \frac{7v'}{2R}$$

Also:

$$v' = \frac{xg}{2eV} = \frac{1.29 \times g}{\sqrt{2gh}}$$

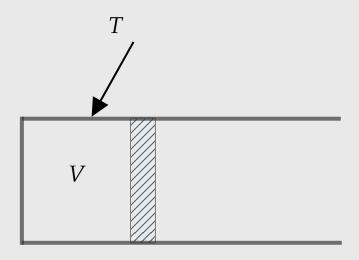
Using this we get:

$$\omega = 0.1 rad/s$$

Solution 8

Problem 9

Photon gas on a carnot engine ride



A container with a piston shown in the figure, has a photon gas characterised by state equations of

Internal energy U(T, V) and pressure P(T) which are

$$U(T,V) = bVT^4$$

$$P(T) = \frac{1}{3}bT^4$$

where b is a dimensioned constant, V is the volume of system, T is its temperature.

Let's take the gas through a carnot cycle, starting with V=0 and temperature of the enclosure T_H and sequentially take it through

Process 1: Isothermal expansion at temperature T_H to a volume V_1 .

Process 2: Adiabatic expansion to a volume V_2 .

Process 3: Isothermal compression to zero volume at temperature T_c .

Process 4: Heating back the enclosure from T_c to T_H at zero volume.

Mark the correct option(s):

(a) The enclosure acts as perfect reflector for photons during process-2.

(b) The change in entropy of the photon gas during process 3 is $-\frac{1}{3}bT_c^3V_2$.

(c) The process equation for process-2 is $PV^{4/3} = \text{constant}$.

(d) Ratio of number of photons in the enclosure at same volume $\frac{V_1}{2}$ but during process-1 to process-3 is $\frac{T_H^4}{T_4^4}$.

-Proposed by Janardanudu Thallaparthi

Answer: AC

Process 1-3 is isothermal

$$U = bVT^4 \to dU = bT^4dV = 3PdV$$

$$P = \frac{b}{3}T^4 \to dQ = dU + dW = 4PdV$$

So,
$$W=\frac{1}{3}bT^4\Delta V$$
 and $Q=\frac{4}{3}bT^4\Delta V$ Also, $S=\frac{4}{3}bVT^3$ so option B is wrong.

For process 2, which is adiabatic "S" should not change $\to VT^3 = \text{constant} \to PV^{\frac{4}{3}} = \text{constant}$, also number of photons remain constant as walls cannot emit or absorb photons during adiabatic process. So walls are considered perfect reflectors in process 2

Also, in adiabatic process, Number of photons N and entropy S are constant

$$\rightarrow N = \alpha S$$
, Where *N* is a constant

$$\rightarrow \frac{N_1}{N_2} = \frac{S_1}{S_2} = \frac{\frac{4}{3}bV_1T_1^3}{\frac{4}{3}bV_2T_2^3} = \frac{T_1^3}{T_2^3}$$
 as same volume

Solution 9

Problem 10

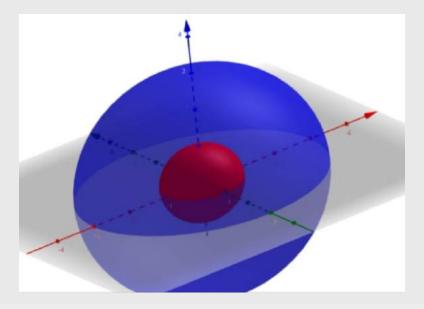
A sphere of radius $r_1 = 1m$ is placed concentrically in a thin shell of radius $r_2 = 3m$ with emmisivities e_1 and e_2 respectively. The inner sphere generates a constant power P. Assume that none of the surfaces allow transmission. Also assume that space between sphere and outer shell is a vacuum.T is steady state temperature of inner sphere.Then Choose the correct option(s):

A.If
$$e_1 = 0.30$$
 and $e_2 = 0.50$, $11P = 12\sigma\pi T^4$

B.If
$$e_1 = 0.30$$
 and $e_2 = 0.50$, $91P = 3\sigma\pi T^4$

C.If
$$e_1 = 0.20$$
 and $e_2 = 0.40$, $49P = 36\sigma\pi T^4$

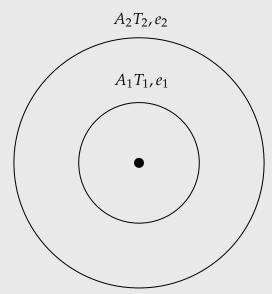
D.If
$$e_1 = 0.20$$
 and $e_2 = 0.40$, $92P = 3\sigma\pi T^4$



-Proposed by AKIII

Answer: (A), (C)

Solution-1:



As seen from outside, net power P should be coming out. Let T be the steady state temperature of the inner shell and T_2 be the steady state temperature of the outer shell.

$$\therefore \sigma e_2 A_2 T_2^4 = P$$

Let total power absorbed by the outer shell be = k. For steady state of the outer shell:

$$k = 2\sigma e_2 A_2 T_2^4$$
$$k = 2P$$

If the outer shell absorbs power, then it must have received $\frac{k}{e_2}$ power. The radiations that outer shell receives are from some of its own radiations that fall on itself, coming from the inner sphere.

∴It has reflected $\frac{k}{e_2}(1-e_2)$ Power.

Out of this,

 $\frac{k}{e_2}(1-e_2).\frac{A_1}{A_2}e_1$ is absorbed by the inner shell.

So, for the inner shell:

$$P\frac{k}{e_2}(1-e_2).\frac{A_1}{A_2}e_1 + P\frac{A_1}{A_2}e_1 = \sigma A_1 e_1 T^4$$

($P\frac{A_1}{A_2}e_1$ is the direct radiation from the outer shell.)

$$\sigma A_1 e_1 T^4 = P \frac{A_1}{A_2} e_1 \frac{2 - e_2}{e_2} + P$$

$$\sigma A_1 e_1 T^4 = P\left(\frac{A_1}{A_2} e_1 \frac{2 - e_2}{e_2} + 1\right)$$

Solution-2:

X=Power falling on inner sphere

Y=Power coming out of the inner sphere(radiated+reflected)

$$\implies Y - X = \sigma$$
$$Y = X(1 - e_1) + e_1 A_1 \sigma T^4$$

where:

 $X(1 - e_1)$ is the reflected part and $e_1A_1\sigma T^4$ is the radiated part. Considering the equation of the outer shell:

$$P\left(\frac{A_1}{A_2}\right) + P\left(\left(1 - \frac{A_1}{A_2}\right)(1 - e_2)\frac{A_1}{A_2}\cdots\right) + P = Ye_2 + Y(1 - e_2)\left(1 - \frac{A_1}{A_2}e_2 + \cdots\right)$$

$$Ye_2 = P\left(2\frac{A_1}{A_2} + e_2 - \frac{A_1}{A_2}e_2\right)$$

$$Y = P\left(\frac{A_1}{A_2}\left(\frac{2 - e_1}{e_2}\right) + 1\right)$$

Also,

$$Ye_1 + P(1 - e_1) = e_1 A_1 \sigma T^4$$

So,

$$P\left(\frac{A_1}{A_2}\left(\frac{2-e_2}{e_2}\right)e_1 + e_1 + 1 - e_1\right) = e_1 A_1 \sigma T^4$$

$$\sigma A_1 e_1 T^4 = P\left(\frac{A_1}{A_2}e_1 \frac{2-e_2}{e_2} + 1\right)$$

 \therefore Energy flow from the inner sphere to the outer sphere is not considered in X and Y, we must account it separately in ().

Solution 10

Problem 11

Consider the region $y \ge 0$ in the cartesian plane filled by a certain material whose refractive index is a function of the vertical distance (y) from the x axis.

$$n(y) = ay \qquad (y > 0)$$

Where $n(0) = n_0$ and a is a constant with dimension $[L]^{-1}$ A ray of light is incident to this region at the origin making an angle θ_0 with the vertical. If

$$y(x) = K \frac{n_0 \sin \theta_0 [\exp\{(2ax)/(n_0 \sin \theta_0)\} + 1]}{a \exp[(ax)/(n_0 \sin \theta_0)]}$$

Report the value of $(2K+1)^{(2K+1)}$ as the answer. If the setup is impossible, enter 0

-Proposed by Muhammed Yaseen Nivas

Answer: 0

As we can see $n(0) = n_0$, but $n(0^+) \to 0$. So, Total internal refraction occurs and the setup is impossible, hence answer is 0.

But in case we assume the ray passes through to the second medium, this is how we solve it: The angle θ , that the tangent to the graph of a function y(x) makes with the y axis clearly satisfies

$$\tan \theta = \frac{1}{y'(x)}$$

Trigonometric manipulations leads us to

$$\sin\theta = \frac{1}{\sqrt{1 + (y')^2}}$$

Now, coming to physics; to solve this problem, we use Snell's law to write that

$$n(y)\sin\theta = c$$

Where *c* is determined by the initial conditions $n(0) = n_0$, $\theta(y) = \theta_0$. Plugging them in,

$$\frac{ay}{\sqrt{1 + (y')^2}} = c \iff \sqrt{\frac{a^2 y^2 - c^2}{c^2}} = \frac{dy}{dx}$$
$$\int_0^x \frac{a dx}{c} = \int_0^y \frac{dy}{\sqrt{y^2 - c^2/a^2}}$$
$$\frac{ax}{c} = \frac{1}{2} \ln \left| \frac{\sqrt{y^2 - c^2/a^2} + y}{\sqrt{y^2 - c^2/a^2} - y} \right|$$

Two cases, either the argument of the logarithm is positive, in which case we can drop the absolute values. We will consider this first.

$$\exp\left(\frac{2ax}{c}\right) = \frac{\sqrt{y^2 - c^2/a^2} + y}{\sqrt{y^2 - c^2/a^2} - y}$$

$$\exp\left(\frac{2ax}{c}\right) \left(\sqrt{y^2 - c^2/a^2} - y\right) = \sqrt{y^2 - c^2/a^2} + y$$

$$(\sqrt{y^2 - c^2/a^2}) (\exp(2ax/c) - 1) = y(\exp(2ax/c) + 1)$$

$$y^2 - c^2/a^2 = y^2\alpha^2$$

Where,

$$\alpha(x) \equiv \frac{\exp(2ax/c) + 1}{\exp(2ax/c) - 1}$$
$$y^2(1 - \alpha^2) = c^2/a^2$$
$$y = \frac{c}{a}\sqrt{\frac{1}{1 - \alpha^2}}$$

$$y(x) = \frac{n_0 \sin \theta_0}{a} \sqrt{\frac{1}{1 - \alpha^2}}$$

This function is undefined, so we have to consider the other case.

$$\exp\left(\frac{2ax}{c}\right) = \frac{\sqrt{y^2 - c^2/a^2} + y}{-\sqrt{y^2 - c^2/a^2} + y}$$

$$\exp\left(\frac{2ax}{c}\right) (y - \sqrt{y^2 - c^2/a^2}) = \sqrt{y^2 - c^2/a^2} + y$$

$$y(\exp\left(\frac{2ax}{c}\right) - 1) = \sqrt{y^2 - c^2/a^2} (\exp\left(\frac{2ax}{c}\right) + 1)$$

$$y^2\beta^2 = y^2 - c^2/a^2$$

$$y^2(1 - \beta^2) = c^2/a^2$$

Where
$$\beta = \frac{\exp(2ax/c)-1}{\exp(2ax/c+1)}$$
.

$$y = \frac{c}{a} \sqrt{\frac{1}{1 - (\frac{\exp(2ax/c) - 1}{\exp(2ax/c) + 1})^2}}$$

$$y = \frac{c}{a} \sqrt{\frac{(\exp(2ax/c) + 1)^2}{(\exp(2ax/c) + 1)^2 - (\exp(2ax/c) - 1)^2}}$$

$$y = \frac{c}{a} \sqrt{\frac{(\exp(2ax/c) + 1)^2}{4 \exp(2ax/c)}}$$

$$y = \frac{n_0 \sin \theta_0 (\exp(2ax/n_0 \sin \theta_0) + 1)}{2a \exp(ax/n_0 \sin \theta_0)}$$

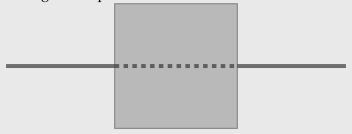
So in case we did assume ray to pass through this would be the trajectory equation.

Solution 11

Problem 12

A cube of side 2*a* and specific gravity *s* floats in a calm water. "Water line section" is the cross-section of cube intersecting free surface of water.

If the cube is given a small angular displacement about centre of water line section and released then,

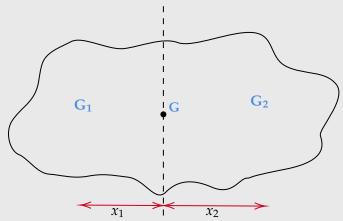


- (a) Volume of cube submerged is increased due to small angular displacement.
- (b) Volume of the cube submerged remains same even after small angular displacement.
- (c) For the stability of rotational equilibrium of cube, the condition required is $4s^2 4s + 1 > 0$
- (d) For the stability of rotational equilibrium of cube, the condition required is $6s^2 6s + 1 > 0$

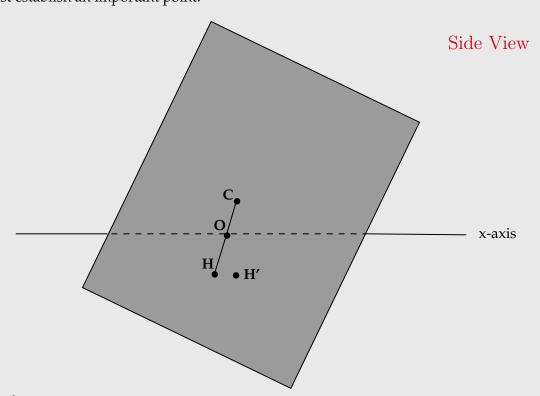
-Proposed by Janardanudu Thallaparthi

Answer: BD

Volume submerged remains same for any arbitrary shaped water line cross section. Top view of this section is shown here:



For small angular displacements, G_1G moves up by x_1 , θ and GG_2 moves down by $x_2\theta$, For equilibrium initially $A_1x_1 = A_2x_2$ where A_1 and A_2 are areas left and right of dashed horizontal line. So, $A_1x_1\theta = A_2x_2\theta$ represents equal volumes by Pappus' theorem. Hence, volume submerged remains unchanged. So, buoyant forces remains unchanged. Let's first establish an important point.



H and H' are centeres of buoyancy before and after the deflection. When a vertical line is drawn through H' and it meets the line HC at M, this point M' is called M eacher. For stability (using torque equation), we need MM > HC i.e. Meta center should lie above M, the center of mass.

Volume element = $(\delta A)(\delta x)$, its depth upon small rotation = $x\theta$ (see side-view). The vertical displacement of center of buoyancy is of order θ^2 and to first order of θ , neglected.

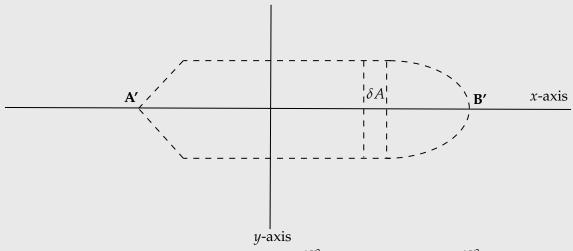
Consider Volume moments about *y*-axis. Let $\bar{x'}$ and \bar{x} be final and initial *x*-positions of center of buoyancy. We have,

$$V\bar{x'} = V\bar{x} - \int_0^{A'} (x\theta)dAx + \int_0^{B'} x\theta(dA)x$$

$$V\underbrace{(\bar{x}' - \bar{x})}_{HH'} = \theta \int_{A'}^{B'} x^2 dA \quad \text{(related to radius of gyration)}$$

$$V(HH') = \theta(AK^2)$$

Top View



 $V(HM\sin\theta)=\theta(AK^2)$ As $\sin\theta\approx\theta$, we get $HM=\frac{AK^2}{V}$. For stability we get $\frac{AK^2}{V}>HC$ So, for cube of side 2a, relative density s $A=4a^2$, $K^2=\frac{a^2}{3}$, $V=8a^3s$ HC=a(1-s), So for a stability,

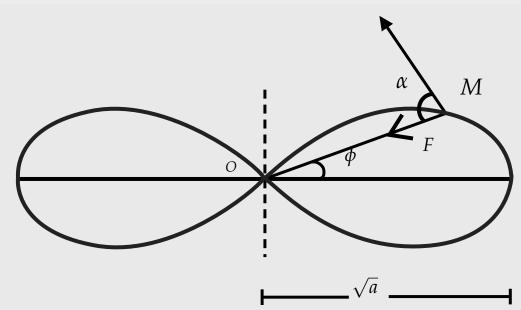
$$\frac{a}{6s} > a(1-s)$$

. Unstable if $6s^2 - 6s + 1 < 0$

Solution 12

Problem 13

A particle of mass m, subjected to a central force F describes the following trajectory : $r^2 = a \cos 2\phi$, where a is a positive constant, r is distance of point from centre of force. At initial moment, $r = r_0$, speed $v = V_0$ and makes an angle α with straight line connecting the point with centre of force!



If the expression for force is given as

$$F = -\frac{KmV_0^2 r_0^2 a^2 \sin^2 \alpha}{r^7}$$

Find the value of *K*

-Proposed by Harshit Gupta

Answer: 3

The magnitude of central force:

$$F(r) = m(\ddot{r} - r\dot{\theta}^2)$$

Conservation of angular momentum gives us:

$$r^2\dot{\theta} = \text{constant} = L$$

r can be converted into terms of $u = \frac{1}{r}$

$$\frac{\mathrm{d}u}{\mathrm{d}\theta} = \frac{\mathrm{d}}{\mathrm{d}t} \left(\frac{1}{r}\right) \frac{\mathrm{d}t}{\mathrm{d}\theta} = -\frac{\dot{r}}{r^2 \dot{\theta}} = -\frac{\dot{r}}{L}$$

$$\frac{\mathrm{d}^2 u}{\mathrm{d}\theta^2} = -\frac{1}{L} \frac{\mathrm{d}\dot{r}}{\mathrm{d}t} \frac{\mathrm{d}t}{\mathrm{d}\theta} = -\frac{\ddot{r}}{L\dot{\theta}} = -\frac{\ddot{r}}{L^2 u^2}$$

Solving all the 3 equations we get the famous Binet Equation:

$$F(r) = -\frac{L^2}{mr^2} \left[\frac{d^2}{d\phi^2} \left(\frac{1}{r} \right) + \frac{1}{r} \right]$$

Now substituting
$$r = \sqrt{a \cos 2\phi}$$
, $L = mv_o r_o \sin \alpha$ and differentiating wrt to ϕ we get:
$$F = -\frac{3ma^2 r_o v_o \sin^2 \alpha}{r^3 \cos^2 \phi} = -\frac{3ma^2 r_o^2 v_o^2 \sin^2 \alpha}{r^7} \text{ hence } \boxed{k=3}$$

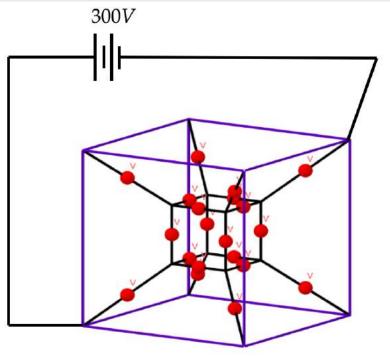
Solution 13

Problem 14

We have a tesseract cubic style arrangement with voltmeters and resistances as shown in the figure. There are identical voltmeters to each side of smaller cube and side connecting to larger cube as shown in fig. below. All sides have equal resistance R, and across 2 ends of bigger cube connected with a battery of 300V. Calculate

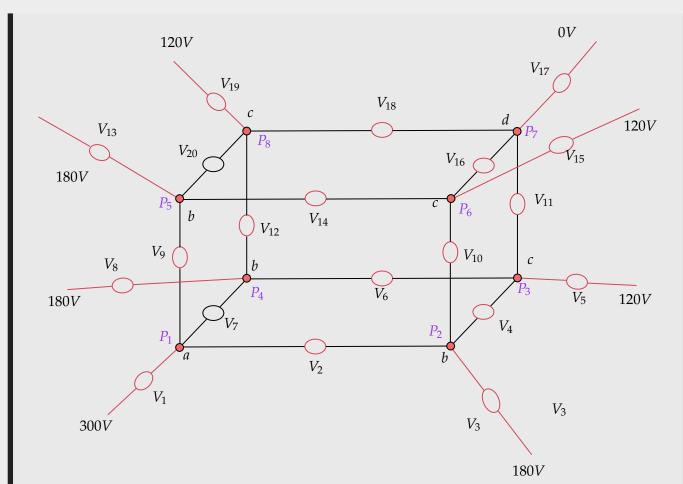
$$\frac{7}{99}\sum_{i=0}^n V_i,$$

where V_i is reading of i^{th} voltmeter



-Proposed by AKIII

Answer: 40



Applying KCL at P1:

$$\frac{a-300}{R} + 3\left(\frac{a-b}{R}\right) = 0$$
$$\Rightarrow 4a = 300 + 3b$$

Applying KCL at P7:

$$\frac{d}{R} + 3\left(\frac{d-c}{R}\right) = 0$$
$$\Rightarrow 4d = 3c$$

Applying KCL at P2:

$$\frac{b-180}{R} + 2\left(\frac{b-c}{R}\right) + \frac{b-a}{R} = 0$$
$$\Rightarrow 4b = a + 2c + 180$$

Applying KCL at P3:

$$\frac{c-d}{R} + 2\left(\frac{c-b}{R}\right) + \frac{c-120}{R} = 0$$

$$\Rightarrow 4c = d + 2b + 120$$

$$\Rightarrow \frac{13}{4}c = 2b + 120 \qquad \text{...from (2)}$$

So,

$$26b = 16c + 2090$$

$$26c = 16b + 960$$

By solving,

$$a = \frac{1380}{7}V$$
, $b = \frac{1140}{7}V$, $c = \frac{960}{7}V$, $d = \frac{720}{7}V$

$$\rightarrow V_1 = 300 - a$$
, $V_2 = V_7 = V_9 = a - b$, $V_3 = V_8 = V_{13} = b - 180$

$$\rightarrow V_4 = V_{10} = V_6 = V_{12} = V_{14} = V_2 = b - c$$

$$\to V_5 = V_{15} = V_{19} = c - 120$$

$$\rightarrow V_{11} = V_{16} = V_{18} = c - d, V_{17} = d$$

$$\Rightarrow \frac{7}{99} \sum_{i=1}^{20} V_i = \frac{7}{99} (300 - a + 3a - 3b + 3b - 540 + 6b - 6c + 3c - 360 + 3c - 3d + d)$$

$$7 \quad (3960)$$

$$= \frac{7}{99} \left(\frac{3960}{7} \right) = 40$$

Solution 14

Problem 15

Power radiated by a point charge:

$$P = k\mu_0^d q^e \zeta^f c^g$$

where k is a dimensionless constant. Here, μ_0 is the permeability of free space, q is the charge, ζ is the acceleration, and c is the speed of light in vacuum.

Consider a fixed positive charge Q situated at fixed point in space (vacuum). At some instant, a Li^{+3} ion (stable enough throughout the problem) is fired from infinity directly towards the charge Q with a velocity v_0 . The ion makes a shortest distance of approach to the fixed charge x_0 ($x_0 > 0$). Atharva (using his supernatural powers) extracts energy radiated away by Li^{+3} ion during its whole journey and with that energy, excites an H atom from ground state to third excited state. Ignore relativistic effects (i.e. $v_0 << c$), and radiative losses on the motion of the particle. Neglect all kind of particle physics interactions.

If $\frac{mv_0^5}{Q} = \psi^4 \times 10^{25}$, where m is the mass of ion, report $[\psi]$ where [.] is the greatest integer function. Take $k = \frac{1}{6\pi}$ in the expression of power radiated by a point charge. You may find this integral helpful:

$$\int_{\xi}^{\infty} \frac{1}{x^4 \sqrt{\frac{1}{\xi} - \frac{1}{x}}} dx = \frac{16}{15\xi^{5/2}}$$

-Proposed by Tarpan Ghosh

Answer: 2

First, we calculate the exact expression for power radiated by a point when accelerated. Using dimensional analysis, it can be shown:

$$P = k \frac{\mu_0 q^2 \zeta^2}{c}$$

This implies, the rate of work done is

$$\frac{dW}{dt} = k \frac{\mu_0 q^2 \zeta^2}{c}$$

Let the charge of Li^{+3} be q and using Coulomb's law,

$$\zeta = \frac{Qq}{4\pi\epsilon_0 mx^2} = \frac{\Omega}{x^2}$$

where $\Omega = \frac{Qq}{4\pi\epsilon_0 m}$

Therefore, the total energy radiated is (since total energy emitted is double of the energy emitted in a single cycle)

$$W = 2 \cdot \frac{k\Omega^2 \mu_0 q^2}{c} \int_{x_0}^{\infty} \frac{1}{v \cdot x^4} dx$$

Now applying conservation of energy (no radiative losses as given in question)

$$v^2 = v_0^2 - \frac{2\Omega}{r}$$

 $x_0 = \frac{2\Omega}{v_0^2}$ Therefore,

$$W = 2 \cdot \frac{k\Omega^2 \mu_0 q^2}{c} \int_{x_0}^{\infty} \frac{1}{v \cdot x^4} dx = 2 \cdot \frac{k\Omega^2 \mu_0 q^2}{c} \int_{x_0}^{\infty} \frac{1}{x^4 \sqrt{v_0 - \frac{2\Omega}{x}}} dx = \frac{2k\Omega^2 \mu_0 q^2}{c\sqrt{2\Omega}} \int_{x_0}^{\infty} \frac{1}{x^4 \sqrt{\frac{1}{x_0} - \frac{1}{x}}} dx$$

Using the integral given in the problem, and further simplifications, we get total energy dissipated

$$E = \frac{8qm}{45c^3Q}v_0^5$$

Since Atharva is using this energy to excite the ground state single electron in *H* atom, we can write

$$\frac{8em}{15c^3Q}v_0^5 = 13.6Ze\left(\frac{1}{n_1^2} - \frac{1}{n_2^2}\right)$$

Substituting all the values, we $\frac{mv_0^5}{Q}=64.5468\times 10^{25}$. By comparing, we get $\psi=2.834449208$. Hence, $[\psi]=2$

Solution 15

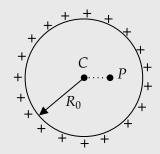
Problem 16

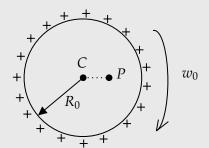
A uniformly charged ring of radius R_o and total charge on the ring is Q_o . There is a point P by a very small distance of x ($x << R_o$) from the centre C in the plane of the ring as shown , The magnitude of electric field at point P due to the stationary ring is E_o . Now the ring is rotated with constant angular

velocity of w_o about an axis passing through its centre and perpendicular to the plane of the ring, The magnitude of magnetic field at point P due to the rotating ring is B_o . Then the ratio of E_o/B_o at point P is of the form

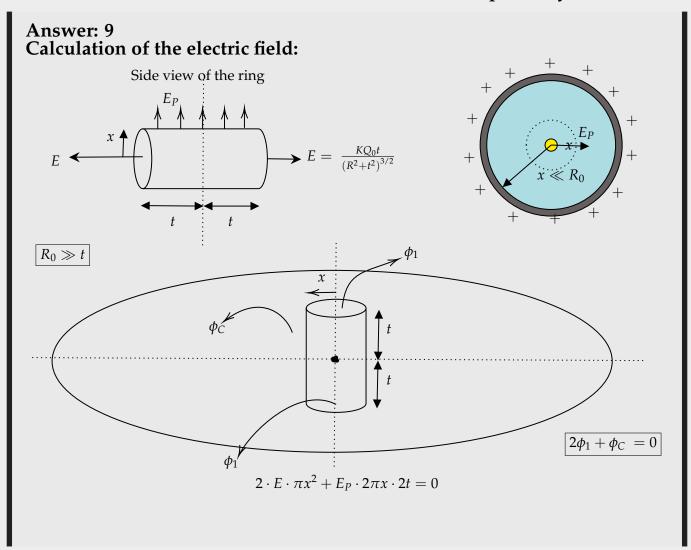
$$\frac{E_o}{B_o} = \frac{\alpha x c^2}{\omega \left[\beta R^2 + \gamma x^2\right]}$$

Where α , β and γ are Integers and c is speed of light. Compute $(\alpha + \beta + \gamma)$.





-Proposed by Nitin Sachan



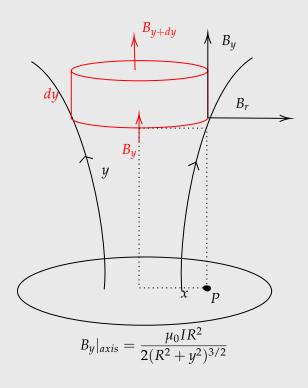
$$E_P = -\frac{E \cdot x}{2t} = \frac{KQ_0t}{(R^2 + t^2)^{3/2}} \cdot \frac{x}{2t}$$

$$E_P = \frac{KQ_0x}{2R^3}$$

Direction towards the center of the ring.

Calculation of Magnetic Field:

$$I = \frac{QW}{2\pi}$$



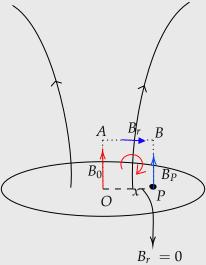
Near the axis :- $x \ll R$

To find B_r : Using Gauss Theorem

$$\begin{split} B_{y+dy} \cdot \pi x^2 - B_y \cdot \pi x^2 + B_r \cdot 2\pi x \cdot dy &= 0 \\ B_r &= \frac{-x}{2} \left(\frac{B_{y+dy} - B_y}{dy} \right) = \frac{-x}{2} \cdot \frac{dB}{dy} \\ B_r &= \frac{-x}{2} \cdot \frac{\mu_0 I R^2}{2} \cdot \frac{d}{dy} [R^2 + y^2]^{-3/2} \\ B_r &= -\frac{\mu_0 I R^2 x}{4} \cdot \frac{-3}{2} \cdot 2y \cdot [R^2 + y^2]^{-5/2} \\ B_r &= \frac{3\mu_0 I R^2 xy}{4(R^2 + y^2)^{5/3}} \end{split}$$

 $h \rightarrow 0 \Rightarrow$ Point *A* and *B* are very close to axis

$$B_0 = \frac{\mu_0 I}{2R}$$



Using Amperes Law:

$$\underbrace{\int_0^A \vec{B} \cdot \vec{dl}}_{+B_0h} + \int_A^B \vec{B_r} \cdot \vec{dl} + \underbrace{\int_B^P \vec{B_p} \cdot \vec{dl}}_{-B_ph} + \underbrace{\int_0^P \vec{B} \cdot \vec{dl}}_{0} = 0$$

$$\int_A^B \vec{B_r} \cdot \vec{dl} = \int_0^x \frac{3\mu_0 I R^2 x h}{4(R^2 + y^2)^{5/3}} \cdot dx = \frac{3\mu_0 I R^2 x^2 h}{8(R^2 + y^2)^{5/3}} = \frac{3\mu_0 I R^2 x^2 h}{8R^5} = \frac{3\mu_0 I x^2 h}{8R^3}$$

 $R \gg h$

$$\Rightarrow +B_0 h - B_p h + \frac{3\mu_0 I x^2 h}{8R^3} = 0$$

$$B_p = B_0 + \frac{3\mu_0 I x^2}{8R^3}$$

$$B_p = \frac{\mu_0 I}{2R} \left[1 + \frac{3x^2}{4R^2} \right]$$

$$B_p = \frac{\mu_0}{2R} \cdot \frac{QW}{2\pi} \left(1 + \frac{3x^2}{4R^2} \right) = \frac{\mu_0 QW}{4\pi R} \left(1 + \frac{3x^2}{4R^2} \right)$$

$$E_p = \frac{kQ_0 x}{2R^3} = \frac{1}{4\pi\epsilon_0} \cdot \frac{Q_0 x}{2R^3} = \frac{Q_0 x}{8\pi\epsilon_0 R^3}$$

$$\frac{E_p}{B_p} = \frac{2x^2}{w[4R^2 + 3x^2]}$$

Solution 16