

The Computation of Discontinuous Solutions of Linear Hyperbolic Equations*

MICHAEL S. MOCK

Rutgers University

AND

PETER D. LAX

Courant Institute

Let L be a linear hyperbolic operator in any number of variables with C^∞ coefficients. As is well known, a solution u of $Lu=0$ which has C^∞ initial data is C^∞ for all time. Let L_h be a difference approximation to L that is stable, and accurate of order ν . Denote by U the solution of $L_h U=0$ whose initial values agree with those of u on the lattice points, h denoting the mesh width of the lattice. According to the basic theory of difference approximations, at all times t at which U is available, and in any fixed range $0 \leq t \leq T$,

$$(1) \quad |u(t) - U(t)| = O(h^\nu).$$

Consider piecewise C^∞ initial data whose discontinuities occur across C^∞ surfaces. It is known that solutions u of $Lu=0$ with such initial data are themselves piecewise C^∞ , and their discontinuities occur across characteristic surfaces issuing from the discontinuity surface of the initial data. What happens when such a solution is approximated by a solution of $L_h U=0$? Does U differ from u by $O(h^\nu)$ in those regions where u is C^∞ , or have the discontinuities hopelessly polluted the approximate solutions even at smooth regions between discontinuities? In a recent paper [2], Majda and Osher have shown that, for second order accurate schemes applied to hyperbolic equations in one space variables and with constant coefficients, U differs from u in smooth regions by $O(h^2)$, provided that at lattice points of discontinuity the initial value of U is taken as the average of the values of u on the two sides of the discontinuity. In this note we show that, for a scheme of any order ν , the *moments* of U approximate those of u with accuracy $O(h^\nu)$, provided that the initial data of U are *pre-processed* appropriately near the discontinuities; this is true for equations in any number of variables, and with variable coefficients. We show further that we can, by *post-processing* the approximate solution, recover the exact solution, as well as its derivatives,

* This work was supported in part by the U.S. Department of Energy Contract EY-76-C-02-3077*000 at the Courant Mathematics and Computing Laboratory, New York University. Reproduction in whole or in part is permitted for any purpose of the United States Government.

with an error $O(h^{\nu-\delta})$, δ as small as we wish, at all points, no matter how close to the discontinuity. The idea of comparing the moments of a discontinuous solution with the moments of its approximation comes from [3]; there it is used also in studying discontinuous solutions of nonlinear conservation laws.

We need the following quadrature result, which goes back to the 18-th century, see [1]:

LEMMA 1. *Let f be any C^∞ function on R_+ , with bounded support. Given any positive integer ν , there exists a quadrature formula accurate of order ν of the form*

$$(2) \quad \int_0^\infty f(x) dx = h \sum_0^\infty w_j f(jh) + O(h^\nu),$$

where the weights w_j depend on ν , but

$$(3) \quad w_j = 1 \quad \text{for } j \geq \nu.$$

EXAMPLE. For $\nu = 4$ we have

$$w_0 = \frac{3}{8}, \quad w_1 = \frac{7}{6}, \quad w_2 = \frac{23}{24}, \quad w_j = 1 \quad \text{for } j > 2.$$

Using (2) twice we get this

COROLLARY. *Let f be a piecewise C^∞ function, the discontinuity occurring at $x = 0$. Let ν be any integer; then*

$$(2)' \quad \int_{-\infty}^\infty f(x) dx = h \sum w'_j f(jh) + O(h^\nu),$$

where

$$(4) \quad w'_0 = 2w_0, \quad w'_j = w_{|j|} \quad \text{for } j \neq 0,$$

and

$$f(0) = \frac{f(0-) + f(0+)}{2}.$$

Similar formulas can be obtained in case the discontinuity of f occurs between the mesh points, say at $x = h\theta$, $|\theta| < 1$, and for the integrals of discontinuous functions of several variables whose discontinuities occur along smooth surfaces.

Let L be a first-order hyperbolic matrix operator

$$(5) \quad L = \partial_t + \sum A_j \partial_j + B, \quad \partial_j = \frac{\partial}{\partial x_j};$$

the coefficients A_j and B are C^∞ functions of x and t ; for simplicity we take them, and all solutions, to be real.

Denote the L_2 scalar product with respect to the x variables by

$$(6) \quad (u, v) = \int u(x) \cdot v(x) dx.$$

Denote by L^* the adjoint of L ; suppose u and v satisfy

$$(7) \quad Lu = 0, \quad L^*v = 0,$$

and one of them vanishes for $|x|$ large. Then, by Green's formula in the slab $0 \leq t \leq T$,

$$(8) \quad (u(T), v(T)) = (u(0), v(0)).$$

Let L_h be a two-level, forward difference approximation:

$$(9) \quad L_h = D_t^+ + \sum S_j T^j, \quad T^j \text{ translation by } jh,$$

Let U be a lattice function that satisfies $L_h U = 0$, i.e.,

$$(10) \quad U_k^{n+1} = \sum C_j U_{k-j}^n, \quad C_j = C_j^n(k).$$

Let V be another lattice function; multiplying (10) by V_k^{n+1} and summing gives

$$(11) \quad \begin{aligned} \sum_k U_k^{n+1} \cdot V_k^{n+1} &= \sum_{k,j} C_j U_{k-j}^n \cdot V_k^{n+1} \\ &= \sum_k U_k^n \cdot \sum_j C_j^*(k+j) V_{k+j}^{n+1}. \end{aligned}$$

Suppose V satisfies the two-level backward adjoint equation

$$L_h^* V = 0$$

defined to be

$$(10)^* \quad V_k^n = \sum C_j^*(k+j) V_{k+j}^{n+1}.$$

Then (11) can be rewritten as

$$(12) \quad (U^{n+1}, V^{n+1})_h = (U^n, V^n)_h,$$

where $(\cdot, \cdot)_h$ denotes the lattice scalar product

$$(13) \quad (U, V)_h = h \sum U_k \cdot V_k.$$

Note that, for u, v in C_0^∞ ,

$$(14) \quad (u, v)_h = (u, v) + O(h^\nu)$$

for any ν .

We draw two conclusions from (12):

(a) If L_h approximates L to order ν , then L_h^* approximates L^* also to order ν .

(b) For all N ,

$$(15) \quad (U^N, V^N) = (U^0, V^0).$$

Proof: (a) L_h approximates L to order ν if

$$(16) \quad U^1 = u(h) + O(h^{\nu+1}),$$

where $u(t)$ is any C^∞ solution of $Lu = 0$ and U the solution $L_h U = 0$ with the same initial data as u :

$$(17) \quad U^0 = u(0).$$

Let $v(t)$ be any C^∞ solution of $L^*v = 0$, V a solution of $L_h^*V = 0$ which equals v at $t = h$:

$$(18) \quad V^1 = v(h).$$

Then by (10)* we can determine V^0 :

$$(19) \quad V^0 = v(0) + h^{p+1}E + O(h^{p+2}),$$

where p is the order to which L_h^* approximates L^* , and E is the leading truncation error. It follows from (16) and (18) using (14), that

$$(20) \quad (U^1, V^1)_h = (u(h), v(h)) + O(h^{\nu+1});$$

on the other hand, by (17) and (19),

$$(21) \quad (U^0, V^0)_h = (u(0), v(0)) + h^{p+1}(u(0), E(0)) + O(h^{p+2}).$$

By (12) the left sides of (20) and (21) are equal, and by (8) the first terms on the right are equal; from this we conclude that $p = \nu$.

(b) Relation (15) follows by summing (12) from 0 to $N-1$. It follows that if L_h is stable, so is L_h^* .

We turn now to discontinuous solutions; for the sake of simplicity we take the number of space dimensions to be 1. Suppose u is a piecewise C^∞ solution of $Lu=0$ whose initial data contain a single discontinuity at, say, $x=0$. Let U be a solution of $L_h U=0$ whose initial data are related to those of u as follows:

$$(22) \quad \begin{aligned} U^0(jh) &= w'_j u(jh, 0), & j \neq 0, \\ U^0(0) &= w'_0(u(0-, 0) + u(0+, 0))/2. \end{aligned}$$

where w'_j are the weights (4) entering formula (2)'. Let $\phi(x)$ be an arbitrary C_0^∞ function; denote by v the solution of

$$(23) \quad L^* v = 0, \quad v(x, T) = \phi(x).$$

Let V be the solution of

$$(23)_h \quad L_h^* V = 0, \quad V^N(jh) = \phi(jh),$$

where N is a time step corresponding to $t=T$. Since v is C^∞ and since L_h^* is accurate of order ν , it follows that if L_h^* is stable, then

$$(24) \quad V^0(jh) = v(jh, 0) + O(h^\nu).$$

We take, in the quadrature formula (2)', $f = u(x, 0) v(x, 0)$; this is a piecewise C^∞ function, with a discontinuity at $x=0$, therefore

$$(25) \quad (u(0), v(0)) = \int f(x) dx = h \sum w'_j u(jh, 0) v(jh, 0) + O(h^\nu).$$

Using the definition (22) of $U_h^0(jh)$, the error estimate (24) and the definition (13) of scalar products for lattice functions, we see from (25) that

$$(26) \quad (u(0), v(0)) = (U^0, V^0)_h + O(h^\nu).$$

Now using (8) and (15) we deduce from (26) that

$$(u(T), v(T)) = (U^N, V^N)_h + O(h^\nu).$$

According to (23) and $(23)_h$, both $v(T)$ and V^N were chosen to be ϕ ; therefore the relation above can be rewritten as

$$(27) \quad (u(T), \phi) = (U^N, \phi)_h + O(h^\nu).$$

We summarize:

THEOREM 1. *Let L_h be a stable two-level difference operator that approximates L to order ν . Let u be a solution of $Lu=0$ whose initial data are piecewise C^∞ with a discontinuity at $x=0$. Let U be the solution of $L_h U=0$ whose initial data are related to those of U by formula (22). Then at any later time T the moments of u and U with any C_0^∞ function ϕ differ by $O(h^\nu)$.*

We show now how to use the weak error estimate (27) to deduce pointwise estimates. To this end we need to know the dependence of the error term in (27) on ϕ ; since the error term comes from (24), it is of order of the size of the $(\nu+1)$ -st derivative of v . Thus

$$(27)' \quad (u(T), \phi) = (U^N, \phi)_h + ch^\nu,$$

where

$$(28) \quad c = O(|\phi|_{\nu+1}),$$

the bars denoting the maximum norm of $\partial_x^{\nu+1}\phi$; for more space variables more derivatives of ϕ are needed in (28). Let $s(x)$ be an auxiliary function whose support is contained in the interval $(-1, 1)$, satisfying

$$(29) \quad \int s(x) dx = 1, \quad \int x^l s(x) dx = 0, \quad l = 1, \dots, p-1,$$

p an arbitrary integer. We set

$$(30) \quad \phi(x) = \frac{1}{\varepsilon} s\left(\frac{x-y}{\varepsilon}\right).$$

For any function g that is C^∞ in $(y-\varepsilon, y+\varepsilon)$

$$g(y) = \int g(x)\phi(x) dx + O(\varepsilon^p).$$

So if the interval $(y - \varepsilon, y + \varepsilon)$ is free of discontinuities of u at time T ,

$$\int u(x, T) \phi(x) dx = u(y, T) + O(\varepsilon^p).$$

Comparing this with (27), gives

$$(31) \quad u(y, T) = (U^N, \phi)_h + O(|\phi|_{\nu+1}) h^\nu + O(\varepsilon^p).$$

From the definition of ϕ in (30) we see that

$$|\phi|_{\nu+1} = O\left(\frac{1}{\varepsilon^{\nu+2}}\right);$$

thus

$$(31)' \quad u(y, T) = (U^N, \phi)_h + O\left(\frac{h^\nu}{\varepsilon^{\nu+2}}\right) + O(\varepsilon^p).$$

We choose ε so that the two error terms are of the same order:

$$(32) \quad \frac{h^\nu}{\varepsilon^{\nu+2}} = \varepsilon^p, \quad \varepsilon = h^{\nu/(\nu+2+p)}.$$

With this choice of ε ,

$$(33) \quad u(y, T) = (U^N, \phi)_h + O(h^{\nu p/(\nu+p+2)}).$$

Taking p large enough we get

THEOREM 2. *Choose ϕ of the form (30), (29); formula (33) recovers $u(y, T)$ with accuracy as close to order ν as we wish if p is taken large enough.*

The same technique yields this

COROLLARY. *Any derivative of u can be recovered with an accuracy as close to order ν as desired.*

The points (y, T) at which u and its derivatives can be recovered by formula (33) and its analogue are subject to the restriction that the interval $(y - \varepsilon, y + \varepsilon)$ be free of discontinuities of u at time T . If, however, we choose the support of s to lie in $(0, 1)$ or $(-1, 0)$, this can be replaced by the one-sided restriction that u be free of discontinuities at time T in either $(y - \varepsilon, y)$ or $(y, y + \varepsilon)$. This allows one to get accurate pointwise estimates right up to the discontinuity.

Bibliography

- [1] Davis, P., and Rabinowitz, P., *Methods of Numerical Integration*, Academic Press, New York, 1974.
- [2] Majda, A., and Osher, S., *Propagation of error into regions of smoothness for accurate difference approximations to hyperbolic equations*, Comm. Pure Appl. Math., Vol. 30, 1977, pp. 671-706.
- [3] Mock, M. S., *A difference scheme employing fourth order viscosity to enforce an entropy inequality*, Proc. Bat-Sheva Seminar on Finite Elements for Nonlinear Problems, Tel-Aviv University, 1977, to appear.

Received December, 1977.