

Continuum damage models *Phase-field method* (Variational approach)

G. Anciaux

Civil Engineering, Materials Science, EPFL



Outline

- Introduction
- Continuum approach: local versus non local continuum damage model
- Phase-field approach
- Dynamic crack branching

Continuum damage mechanics

Definition of a continuous damage variable

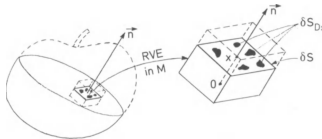
with:

- δS RVE section area
- δS_d area intersecting micro-cracks and/or micro-cavities
- isotropic damage

$$d = \max_{\text{RVE planes}} \frac{\delta S_d}{\delta S}$$

Then:

- $d = 0 \Rightarrow$ Undamaged RVE
- $d = 1 \Rightarrow$ Fully broken RVE in 2 parts



L. Kachanov. *On Time to Rupture in Creep Conditions* (in Russian). Izvestia Akademii Nauk SSSR, Otdelenie Tekhnich. 8,26-31. (1958)

J. Lemaitre. *A Course on Damage Mechanics*. (Springer Berlin Heidelberg, 1996).

Continuum damage mechanics

Helmholtz free energy

$$\psi(\epsilon, D) = \frac{1}{2\rho} [(1-D)\epsilon : C : \epsilon]$$

Constitutive law

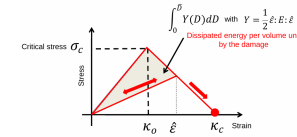
$$\sigma = \rho \frac{\partial \psi}{\partial \epsilon} = (1-D)C : \epsilon$$

Volume energy density

$$\bar{Y} = \rho \frac{\partial \psi}{\partial D} = -\frac{1}{2} \epsilon : C : \epsilon$$

Strain energy density release rate

$$Y(D) = -\bar{Y}(D) = \frac{1}{2} \epsilon : C : \epsilon$$



Dissipated energy

$$E^{\text{dissipated}} = \int Y(D)dD$$

J. Lemaitre. *A Course on Damage Mechanics*. (Springer Berlin Heidelberg, 1996).

Damage evolution

- A set of constraints (comparable to yielding) are necessary:

Positive dissipation constraint

$$Y(D)\dot{D} \geq 0$$

Increasing damage constraint

$$\dot{D} \geq 0$$

Threshold constraint (example)

$$F(Y, d) = Y - Y_d - Sd \leq 0$$

- with Y_d and S material parameters to be fitted to obtain the right dissipation

These constraints allow to compute the damage evolution (similar to plastic flow evolution)

Example: if the constraint is violated ($F > 0$), projecting on the constraint surface leads to:

$$F(Y, d^{n+1}) = F(Y, d^n) + \frac{\partial F}{\partial d} \Delta d = 0$$

which bring the evolution of d :

$$d^{n+1} = \frac{-Y + Y_d}{S}$$

For the constraint the $d \leq 1$ we can change it to:

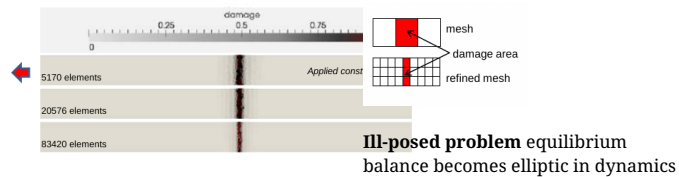
$$d^{n+1} = \min\left(\frac{-Y + Y_d}{S}, 1\right)$$

Remark: F can also be defined by means of an equivalent strain ϵ^e instead of Y , leading to other formulations

J.-J. Marigo. *Formulation d'une loi d'endommagement d'un matériau élastique..* Comptes rendus des séances de l'Académie des sciences. **2**,1309-1312. (1981)

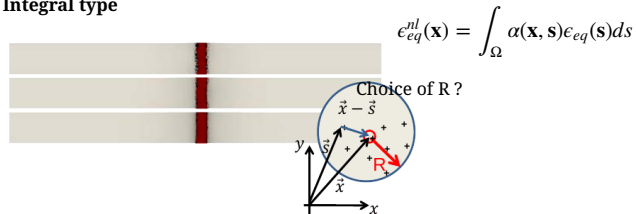
g.-p.-u. family=Vree, W. Brekelmans, g.-p.-u. family=Gils. *Comparison of Nonlocal Approaches in Continuum Damage Mechanics.* Computers & Structures. **55**(4),581-588. (1995)

Local/Non-Local Continuum damage mechanics



- Several methods to avoid localization

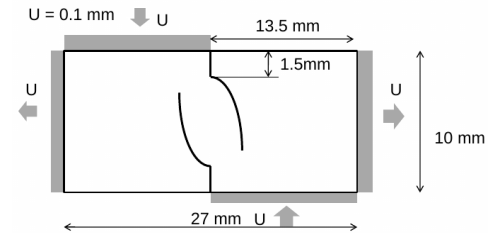
For instance: Gradient, Delay, Integral type



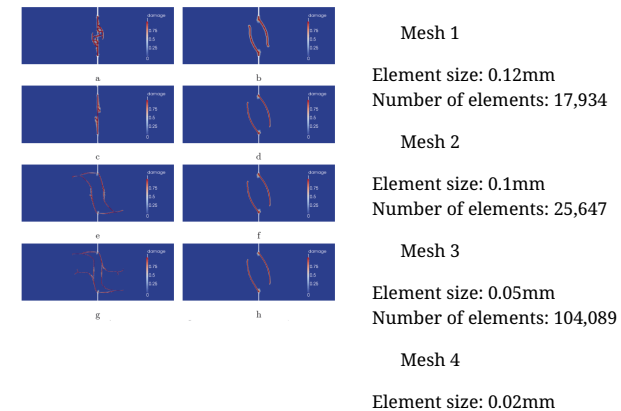
G. Pijaudier-Cabot, Z. Bazant. *Nonlocal Damage Theory.* Journal of Engineering Mechanics. **113**(10),1512-1533. (1987)

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Mesh convergence: Nooru Mohamed test



Mesh convergence: Nooru-Mohamed test



Variational approach to fracture: Phase-field

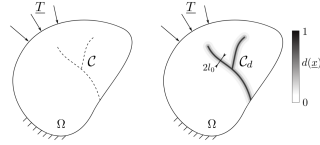
Idea

- The crack discontinuous topology is regularized by a **continuous** phase-field function
 - This function is the damage $d \in [0, 1]$ of the material
 - Introduction of a regularization length scale l_0

Variational approach to fracture: Phase-field

Better suited for

- complex crack paths: dynamic crack branching, instabilities
- crack propagation in heterogeneous media
- multiphysics coupling



With a discontinuous representation

- Γ : the crack path
- G_c : the critical energy release rate
- ψ : the Helmholtz free energy
- $\Psi_0 = \rho\psi$: the energy density

The total energy becomes

$$E(u, \Gamma) = \underbrace{\int_{\Omega \setminus \Gamma} \Psi_0(\epsilon) d\Omega}_{\text{elastic energy}} + \underbrace{G_c \int_{\Gamma} ds}_{\text{dissipated energy}}$$

Regularization leads to

- d : damage phase field (strong link with **damage gradient models**)
- l_0 : characteristic regularization length
- k : residual stiffness at full failure

$$E(u, d) = \int_{\Omega} [(1-d)^2 + k] \Psi_0(\epsilon) d\Omega + \frac{G_c}{2l_0} \int_{\Omega} (d^2 + l_0^2 \|\nabla d\|^2) d\Omega$$

Remarks:

The information of the crack path is **now contained in the phase field d**

Dissipated energy density $\frac{G_c}{2l_0} (d^2 + l_0^2 \|\nabla d\|^2)$

The problem becomes

find u, d minimizers of $E(u, d)$ with the constraint $\Delta d > 0$

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Equilibrium: Euler-Lagrange equations

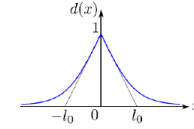
Constitutive law

$$\begin{aligned} \sigma &= [(1-d)^2 + k] \frac{\partial \Psi_0}{\partial \epsilon} \\ &= [(1-d)^2 + k] C : \epsilon \end{aligned}$$

Damage equation

$$0 = 2(1-d)\Psi_0 + \frac{G_c}{l_0} (d + l_0^2 \Delta d)$$

- This equation leads to spreading of damage



Resolution algorithm: alternate minimization

1. At fixed u , solve for d (constrained optimization for irreversibility)
2. At fixed d , solve for u : elastodynamics problem with degraded stiffness

\Rightarrow regularization of the crack surface with a phase-field

Compression/traction separation

$$E(u, d) = \int_{\Omega} \{ [(1-d)^2 + k] \Psi_0^+ + \Psi_0^- \} d\Omega + \frac{G_c}{2l_0} \int_{\Omega} (d^2 + l_0^2 \|\nabla d\|^2) d\Omega$$

with the separation of **compression/traction** strains:

$$\Psi_0^{\pm} = \frac{1}{2} \lambda \langle \text{tr}(\epsilon) \rangle_{\pm}^2 + \mu \text{tr}(\epsilon_{\pm}^2)$$

which then brings:

$$\begin{aligned} \sigma &= [(1-d)^2 + k] \frac{\partial \Psi_0^+}{\partial \epsilon} + \frac{\partial \Psi_0^-}{\partial \epsilon} \\ 0 &= 2(1-d)\Psi_0^+ + \frac{G_c}{l_0} (d + l_0^2 \Delta d) \end{aligned}$$

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Irreversibility of damage using history

$$\begin{aligned} 0 &= 2(1-d)\mathcal{H}(x) + \frac{G_c}{l} (d + l_0^2 \Delta d) \\ \mathcal{H} &= \max_t \Psi_0^+(\epsilon, t) \end{aligned}$$

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Variational approach to fracture:

Phase-field

Choice of dissipation function

$$\frac{G_c}{2l_0} (d^2 + l_0^2 ||\nabla d||^2)$$

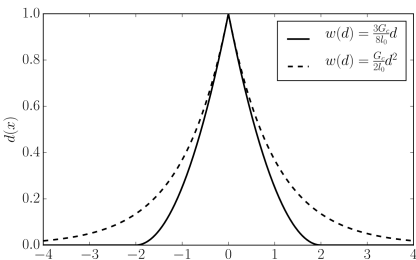
most widely used model in phase-field literature

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$$\frac{3G_c}{8l_0} (d + l_0^2 ||\nabla d||^2)$$

Faster decay of d

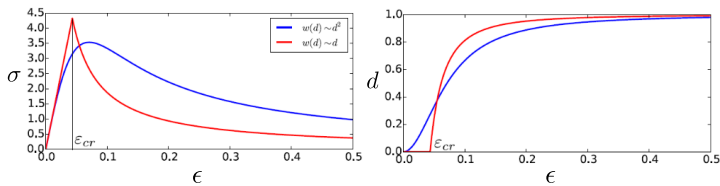
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Comparison of dissipation models

- Un-physical spreading of the phase field

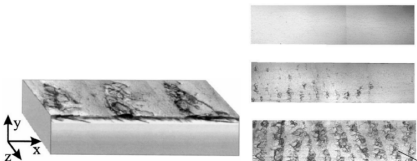


Remark: presence/absence of a purely elastic phase

Introduction: some experimental facts on dynamic fracture

Prediction of a limit crack velocity: c_R (mode I), c_S (mode III) never attained in experiments, rarely exceed $0.4 - 0.6c_R$ explained by **crack tip instabilities** [Fineberg et al.]:

- microbranching ($\sim 0.4c_R$): small ($1 - 100\mu m$ in PMMA) short-lived micro-cracks, highly localized in z direction:

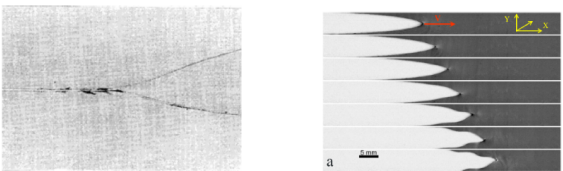


- mirror, mist, hackle patterns

Introduction: some experimental facts on dynamic fracture

Prediction of a limit crack velocity: c_R (mode I), c_S (mode III) never attained in experiments, rarely exceed $0.4 - 0.6c_R$ explained by **crack tip instabilities** [Fineberg et al.]:

- increase of microbranch width \Rightarrow macroscopic branching



- microbranching can be suppressed in thin samples or strongly anisotropic materials \Rightarrow oscillatory instability at $0.9 c_R$

Variational approach to fracture: Phase-field

elastic strain energy density:

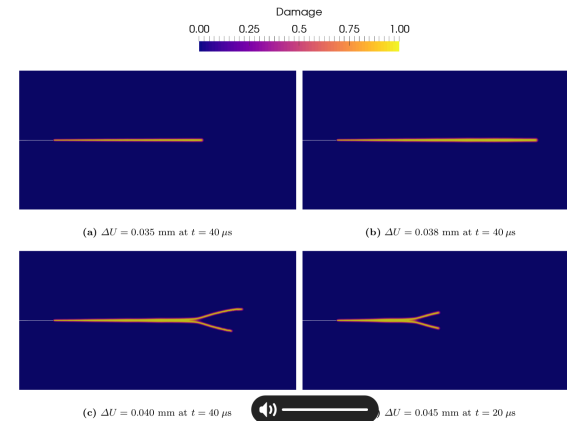
$$\Psi(\epsilon, d) = (1 - d)^2 \Psi_0^+ + \Psi_0^-$$

fracture energy density:

$$\frac{3G_c}{8l_0} (d + l_0^2 \|\nabla d\|^2)$$

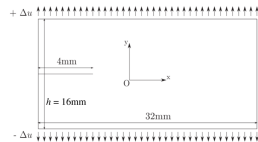
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Variational approach to fracture: Damage zone thickening

Crack branching: Pre-strained plate



Prestrained PMMA plate, fixed boundaries [Zhou, 1996]

$$E = 3.09 \text{ GPa}$$

$$\nu = 0.35$$

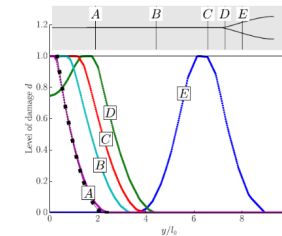
$$\rho = 1180 \text{ kg/m}^3$$

$$G_c = 300 \text{ J/m}^2$$

$$c_R = 906 \text{ m/s}$$

- $\Gamma = \frac{dE}{da} = 2E(\Delta U)^2/h$
 \Rightarrow crack should accelerate to c_R
- transition from straight propagation to branched patterns
- apparent toughness increases with loading/crack velocity

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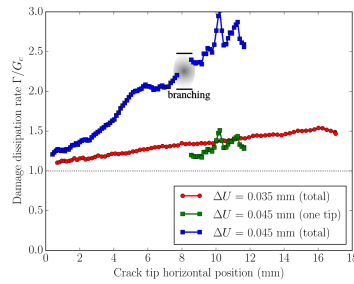


- progressive thickening of the damaged band before branching
- branching viewed as a progressive transition from a widening crack to two crack tips screening each other
- branching angle seems to depend on geometry

Variational approach to fracture: Apparent fracture energy

Damage dissipation rate $\Gamma = \frac{dE_{rac}}{da}$ interpreted as the apparent fracture energy

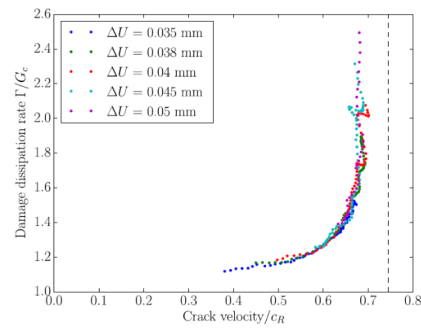
Variational approach to fracture: Dynamic crack branching



suggests a critical value of $\Gamma \approx 2G_c$ associated to branching

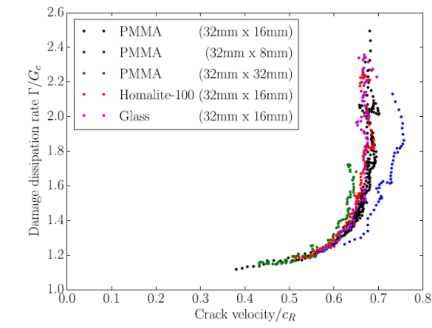
Variational approach to fracture: Velocity toughening

during propagation and before macroscopic branching

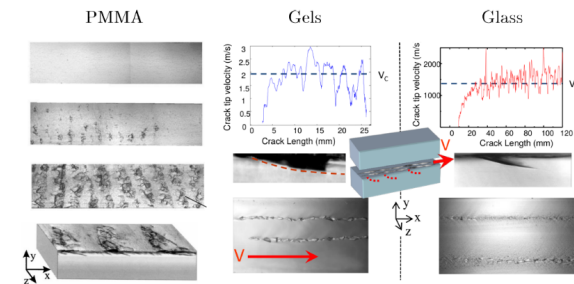


existence of a well-defined $\Gamma(v)$ relationship associated to a velocity-toughening mechanism

Variational approach to fracture: Velocity toughening



But branching is a 3D instability: Fineberg et al.

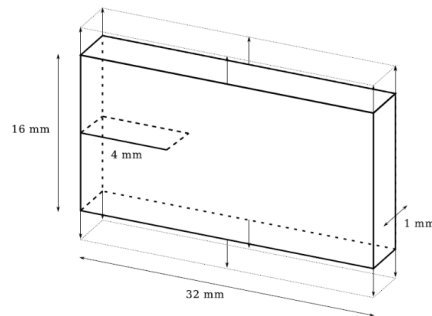


- increasing branch length with loading/velocity
- z-localization: much more localized for Gels/Glass than for PMMA
- x-periodicity: 10 – 100 μm in PMMA, from nm to mm in glass
- microbranching **suppressed for thin samples**

Variational approach to fracture: From 2D to 3D

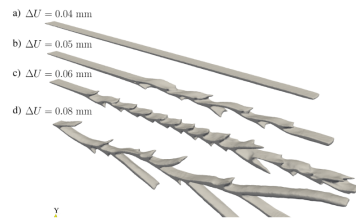
Same setup, same material parameters, now: $L \times H \times W$ plate

starting with $W = 1\text{mm}$, $l_0 = 0.004\text{mm}$



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Variational approach to fracture: Effect of loading

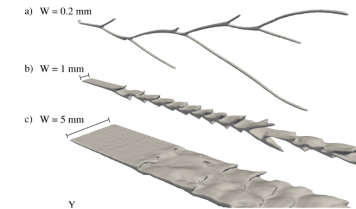


- single straight crack → microscopic branches → macroscopic branches
- nice quasi-periodic regime at intermediate loading:

- less z-invariance at smaller loading consistent with experiments

Variational approach to fracture: Effect of thickness

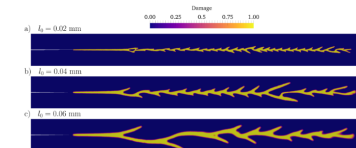
same loading $\Delta U = 0.06\text{mm}$



- microbranching clearly suppressed for small width
- increasing localization with increasing width
- no periodicity for large width:

Variational approach to fracture: Influence of l_0

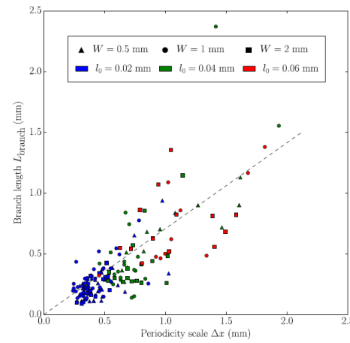
$\Delta U = 0.06\text{mm}$, $W = 1\text{mm}$



- strong effect on Δx
- initiation occurs at roughly the same time
- total dissipated energy almost identical ($\pm 2\%$)
- no microbranching when $l_0 \approx 0.1\text{mm}$
- other plate width suggest $W_{crit} \approx 10l_0$ for this loading

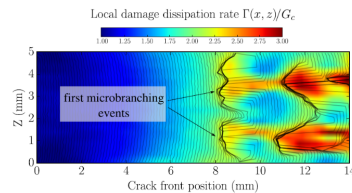
Variational approach to fracture: Emergent geometrical features

Correlation between microbranch length L_{branch} and Δx



Variational approach to fracture: Transition to branching

$W = 5\text{mm}$: localized microbranching events



- first events when $\Gamma \geq 2G_c$ **locally**
- velocity overshoot ahead of the first event initiates the second one
- after that, complex dynamics....

Conclusions

- Variational approach to dynamic fracture shows great prospects (comes at a cost)
- Mesh independency
- Reproduces many experimental features
- Many many open questions... (in dynamics)

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