

Mat: Stokesov izrek:

$$\oint_{\partial S} \vec{E} \cdot d\vec{s} = \int_S (\nabla \times \vec{E}) \cdot d\vec{s}$$

$$\oint_{\partial S} \vec{E} \cdot d\vec{s} = -U[\vec{r}(t_b) = \vec{r}(t_a), \vec{r}(t_a)] = -\text{napetost po (poljubni) razločnem poti}$$

$$\nabla \times \vec{E} = 0 \Rightarrow \oint_{\partial S} \vec{E} \cdot d\vec{s} = 0 \quad \checkmark$$

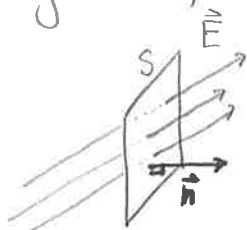
Komentar: omejena veljavnost izreka o el. napetosti

- mirujoči naboji
- brez (časovno spreminjajočega se) magnetnega polja

Dopolnitev: (Faradayev) zakon o el. napetosti (indukcijski zakon)

③ Električni pretok

Primer: homogeno \vec{E} , ravna ploskev



El. pretok - intuitivno: število silnic, ki prebadajo ploskev

Formalno: $\vec{S} = \vec{n} S$; \vec{n} normiran vektor, \perp na S

$$\Phi_E \equiv \epsilon_0 \vec{E} \cdot \vec{S} = \epsilon_0 \vec{E} \cdot \vec{n} S$$

Splošna definicija (nelnogeno polje, katera ploskev):

$$\Phi_E \equiv \epsilon_0 \int_S \vec{E} \cdot d\vec{S} = \epsilon_0 \int_S \vec{E} \cdot \vec{n} dS$$

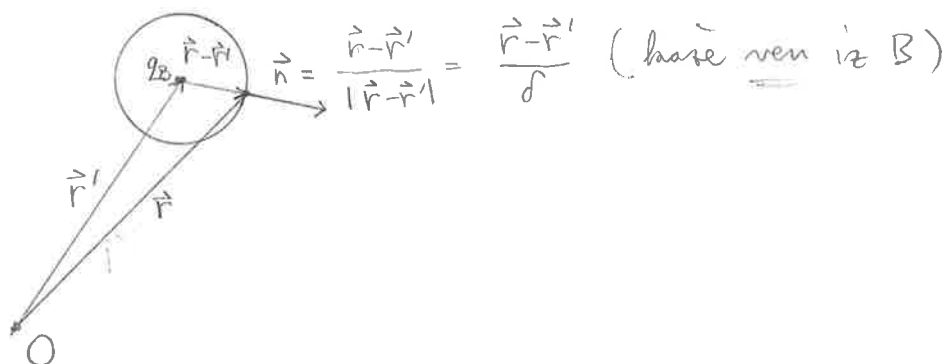
$$\boxed{[\Phi_E] = \frac{As}{Vm} \frac{V}{m} m^2 = As}$$

Primer: el. polje skozi sfero z nabojem q vredišču

$B(\vec{r}', \sigma)$: krogla središčem v \vec{r}' in s polmerom σ

∂B : sfera (radikalna plošev) ; $S = 4\pi\sigma^2$

2D projekcija:



$$\vec{E}(\vec{r}) = \frac{q_B}{4\pi\epsilon_0} \frac{(\vec{r}-\vec{r}')}{|\vec{r}-\vec{r}'|^3} = \frac{q}{4\pi\epsilon_0} \frac{(\vec{r}-\vec{r}')}{\sigma^3}$$

$$\Rightarrow \boxed{\Phi_E(\partial B)} = \epsilon_0 \oint_{\partial B} \frac{q_B}{4\pi\epsilon_0} \frac{(\vec{r}-\vec{r}')}{|\vec{r}-\vec{r}'|^3} \cdot d\vec{S} = \oint \frac{q_B}{4\pi}$$

$$= \frac{q_B}{4\pi\sigma^2} \oint_{\partial B} dS$$

$$= \frac{q_B}{4\pi\sigma^2} S$$

$$= \frac{q_B}{4\pi\sigma^2} 4\pi\sigma^2$$

$$= \boxed{q_B} \text{ zobjeti naboj}$$

Splošna rešica (ižel o Φ_E): polulna radikalna plošev, polulna poravnatelj $S(\vec{r})$ (mirujoči naboj)

Mat : $\boxed{\operatorname{div} \vec{E} \equiv \nabla \cdot \vec{E} = \frac{\partial}{\partial x} E_x + \frac{\partial}{\partial y} E_y + \frac{\partial}{\partial z} E_z}$

Primer: $\vec{E} = \frac{q}{4\pi\epsilon_0} \frac{(\vec{r}-\vec{r}')}{|\vec{r}-\vec{r}'|^3} \Rightarrow E_x = \frac{q}{4\pi\epsilon_0} \frac{(x-x')}{[(x-x')^2 + (y-y')^2 + (z-z')^2]^{3/2}}$

$$E_y = \dots$$

$$E_z = \dots$$

Za $\vec{r} \neq \vec{r}'$: $\frac{\partial}{\partial x} E_x \propto \frac{1}{[\]^{3/2}} - (x-x') \left(-\frac{3}{2}\right) \frac{1}{[\]^{5/2}} \cdot 2(x-x')$

$$= \frac{1}{[\]^{3/2}} - \frac{3(x-x')^2}{[\]^{5/2}}$$

$$\frac{\partial}{\partial y} E_y \propto \frac{1}{[\]^{3/2}} - \frac{3(y-y')^2}{[\]^{5/2}}$$

$$\frac{\partial}{\partial z} E_z \propto \frac{1}{[\]^{3/2}} - \frac{3(z-z')^2}{[\]^{5/2}}$$

$$\Rightarrow \boxed{\nabla \cdot \vec{E} \propto \frac{3}{[\]^{3/2}} - \frac{3[(x-x')^2 + (y-y')^2 + (z-z')^2]}{[\]^{5/2}}}$$

$$= \frac{3}{[\]^{3/2}} - \frac{3}{[\]^{3/2}} = 0$$

Mat : V območje v 3D

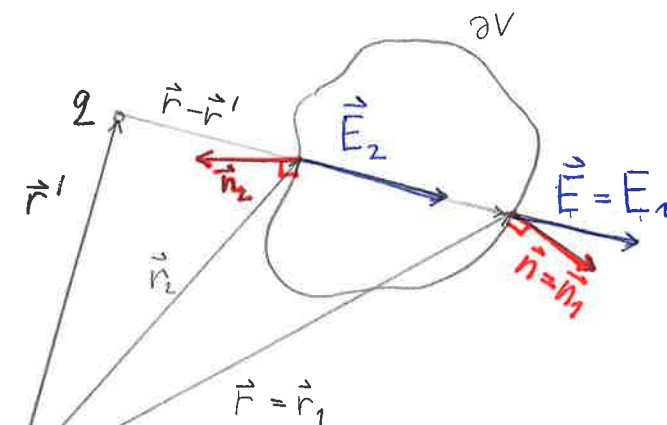
∂V : meja V (zadružena ploskev)

Težava (Gauss, Ostrogradsky):

$$\oint_{\partial V} \vec{E} \cdot d\vec{S} = \int_V \nabla \cdot \vec{E} dV$$

$$\left(\int_{\partial \Omega} \omega = \int_{\Omega} d\omega \right)$$

Primer: Φ_E skozi ∂V , ki ne vsebuje nobenega naboja



$$\forall \vec{r} \in V: \vec{r} \neq \vec{r}': \vec{r} \cdot \vec{E} = 0$$

$$\begin{aligned} \Rightarrow \Phi_E &= \epsilon_0 \oint_{\partial V} \vec{E} \cdot d\vec{S} \\ &= \epsilon_0 \int_V \nabla \cdot \vec{E} dV \\ &= 0 \end{aligned}$$

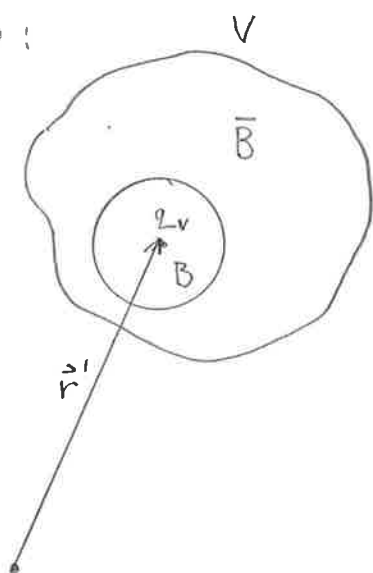
Vsaka silnica, ki vstopi ($\vec{n}_2 \cdot \vec{E}_2 < 0$), tudi izstopi ($\vec{n}_1 \cdot \vec{E}_1 > 0$).

Težava o Φ_E za točko naboj:

$$\Phi_E(\partial V) = \epsilon_0 \oint_{\partial V} \vec{E} \cdot d\vec{S} = q_V \rightarrow \text{zabijati naboj v } V$$

Res: • $q_V = 0$: glej zgornji primer

• $q_V \neq 0$:



$$V; B(\vec{r}', \delta) \subset V; \bar{B} = V \setminus B \quad (B \cap \bar{B} = \emptyset, B \cup \bar{B} = V)$$

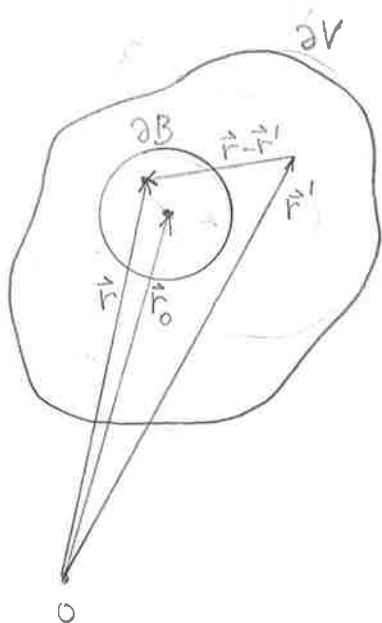
$$q_B = q_V, q_{\bar{B}} = 0 (\nabla \cdot \vec{E} = 0)$$

$$\begin{aligned} \Rightarrow \epsilon_0 \oint_{\partial V} \vec{E} \cdot d\vec{S} &= \epsilon_0 \int_V \nabla \cdot \vec{E} dV \\ &= \epsilon_0 \int_B \nabla \cdot \vec{E} dV + \epsilon_0 \int_{\bar{B}} \nabla \cdot \vec{E} dV \\ &= \epsilon_0 \int_B \nabla \cdot \vec{E} dV = \dots \end{aligned}$$

$$\begin{aligned} \dots &= \epsilon_0 \int_{\partial B} \vec{E} \cdot d\vec{S} \\ &= q_B \\ &= \boxed{q_V} \end{aligned}$$

Komentar: velja za poljubno območje, tudi za kroglo $B(\vec{r}_0, \sigma)$, ki ni centrirana v \vec{r}'

Splošen izrek o Φ_E



Ω : nabito telo

∂V : ploščer, skozi katere gledamo Φ_E

$\rho(\vec{r}')$: gostota naboja v okolici \vec{r}'

$$q_V = \int_V \rho(\vec{r}') dV'$$

$$\text{Izrek: } \boxed{\Phi_E(\partial V) = \epsilon_0 \oint_{\partial V} \vec{E} \cdot d\vec{S} = q_V}$$

Res: $B(\vec{r}_0, \sigma)$: krogla s središčem v \vec{r}_0 in s polmerom σ

$$\left. \begin{aligned} B \subset V \\ \bar{B} \equiv V \setminus B \end{aligned} \right\} B \cup \bar{B} = V, B \cap \bar{B} = \emptyset$$

← Coulombov zakon (+ načelo superpozicije)

$$\begin{aligned} \vec{r} \in B; \quad \vec{E}(\vec{r}) &= \int_{\Omega} \frac{\rho(\vec{r}')}{4\pi\epsilon_0} \frac{(\vec{r}-\vec{r}')}{|\vec{r}-\vec{r}'|^3} dV' = \underbrace{\int_B \frac{\rho(\vec{r}')}{4\pi\epsilon_0} \frac{(\vec{r}-\vec{r}')}{|\vec{r}-\vec{r}'|^3} dV'}_{\vec{E}_B} + \underbrace{\int_{\bar{B}} \frac{\rho(\vec{r}')}{4\pi\epsilon_0} \frac{(\vec{r}-\vec{r}')}{|\vec{r}-\vec{r}'|^3} dV'}_{\vec{E}_{\bar{B}}} \\ &= \vec{E}_B + \vec{E}_{\bar{B}} \end{aligned}$$

$$\Phi_E(\partial B) = \epsilon_0 \oint_{\partial B} \vec{E}(\vec{r}) dV = \epsilon_0 \int_B \nabla \cdot \vec{E} dV = \epsilon_0 \int_B \nabla \cdot (\vec{E}_B + \vec{E}_{\bar{B}}) dV =$$

$$= \epsilon_0 \int_B \nabla \cdot \vec{E}_B dV + \epsilon_0 \int_B \nabla \cdot \vec{E}_{\bar{B}} dV$$

$$\vec{E}_{\bar{B}} = \int_{\bar{B}} \frac{\rho(\vec{r}')}{4\pi\epsilon_0} \frac{(\vec{r}-\vec{r}')}{|\vec{r}-\vec{r}'|^3} dV' ; \vec{r}' \in \bar{B} \Rightarrow \vec{r}' \notin B \Rightarrow \vec{r}' \neq \vec{r}$$

$$\Rightarrow \nabla \cdot \vec{E}_{\bar{B}} = \nabla \cdot \int_B \frac{\rho(\vec{r}')}{4\pi\epsilon_0} \frac{(\vec{r}-\vec{r}')}{|\vec{r}-\vec{r}'|^3} dV'$$

$$= \int_B \frac{\rho(\vec{r}')}{4\pi\epsilon_0} \underbrace{\nabla \cdot \left\{ \frac{(\vec{r}-\vec{r}')}{|\vec{r}-\vec{r}'|^3} \right\}}_{\substack{\text{"} \\ 0 (\vec{r} \neq \vec{r}')}} dV'$$

$$= 0$$

$$\Rightarrow \epsilon_0 \int_B \nabla \cdot \vec{E}_B dV = 0 \quad (\text{malaj izven } B \text{ ne prispeva k el. polju znotaj } \partial B, \text{ velja splošno za poljubno območje } V \text{ z mejno plosčino } \partial V)$$

$$\Rightarrow \Phi_E(\partial B) = \epsilon_0 \int_B \nabla \cdot \vec{E}_B dV$$

$$= \epsilon_0 \int_B \left\{ \int_B \frac{\rho(\vec{r}')}{4\pi\epsilon_0} \frac{(\vec{r}-\vec{r}')}{|\vec{r}-\vec{r}'|^3} dV' \right\} dV$$

$$= \int_B \left\{ \int_B \frac{\rho(\vec{r}')}{4\pi} \nabla \cdot \left\{ \frac{(\vec{r}-\vec{r}')}{|\vec{r}-\vec{r}'|^3} \right\} dV' \right\} dV$$

$$= \int_B \rho(\vec{r}') \left\{ \int_B \frac{1}{4\pi} \nabla \cdot \left\{ \frac{(\vec{r}-\vec{r}')}{|\vec{r}-\vec{r}'|^3} \right\} dV' \right\} dV'$$

$$\int_B \frac{1}{4\pi} \nabla \cdot \left\{ \frac{(\vec{r}-\vec{r}')}{|\vec{r}-\vec{r}'|^3} \right\} dV = \int_{\partial B} \frac{1}{4\pi} \frac{(\vec{r}-\vec{r}')}{|\vec{r}-\vec{r}'|^3} dV$$

$$= \frac{1}{2} \oint_{\partial B} \frac{q}{4\pi} \frac{(\vec{r}-\vec{r}')}{|\vec{r}-\vec{r}'|^3} dS$$

$$= \frac{1}{2} \cdot q$$

$$= 1$$

električni polje skozi sfero z nabojem q v centru (glej zgoraj!)

$$\Rightarrow \Phi_E(\partial B) = \varepsilon_0 \int_B \nabla \cdot \vec{E} \, dV = \int_B \rho(\vec{r}) \, dV$$

$$\varepsilon_0 \int_B \nabla \cdot \vec{E} \, dV = \varepsilon_0 \nabla \cdot \vec{E}(\vec{r}_c) V_B ; \vec{r}_c \in B$$

$$\int_B \rho(\vec{r}') \, dV' = \rho(\vec{r}_c') V_B ; \vec{r}_c' \in B$$

$$\Rightarrow \varepsilon_0 \nabla \cdot \vec{E}(\vec{r}_c) = \rho(\vec{r}_c') ; \forall \delta$$

$$\delta \rightarrow 0 \Rightarrow \vec{r}_c, \vec{r}_c' \rightarrow \vec{r}_0$$

$$\Rightarrow \boxed{\varepsilon_0 \nabla \cdot \vec{E}(\vec{r}_0) = \rho(\vec{r}_0) ; \forall \vec{r}_0 \in V_0}$$

$$\boxed{\varepsilon_0 \nabla \cdot \vec{E}(\vec{r}) = \rho(\vec{r})} \text{ ker } \sigma \Phi_E \text{ v } \underline{\text{diferencialni obliki}}.$$

$$\Rightarrow \Phi_E(\partial V) = \varepsilon_0 \oint_{\partial V} \vec{E} \cdot d\vec{S} = \varepsilon_0 \int_V \nabla \cdot \vec{E} \, dV = \int_V \rho(\vec{r}) \, dV = q_V \text{ (dokazan tudi v integralni obliki)} \quad \checkmark$$

- Komentar: izpel temelji na Coulombovem zakonu, ki (striktno) velja za mirujoče naboje. Eksperimenti pa kažejo na to, da trditev (tako v integralni kot tudi diferencialni obliki) velja splošno - tudi za premikajoče se naboje \Rightarrow izpel o električnem polju \rightarrow zakon o električnem polju.

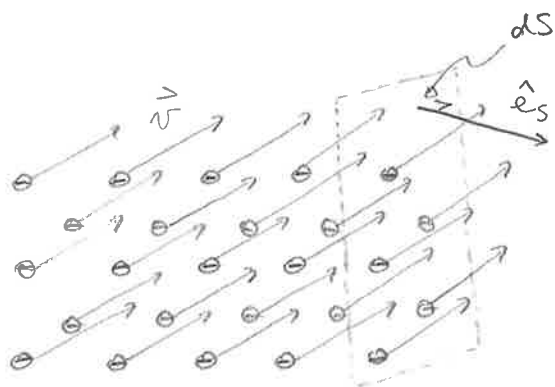
III. ELEKTRIČNI TOK

① gostota električnega toka, električni tok, kontinuitetna enačba

Analogija (konceptualna):

- Električni pretok Φ_E : pretakanje (silnic) el. polja skozi ploskev
- Električni tok I : pretakanje el. nabojev skozi ploskev

Prosti (nevezani) naboji:



- $n = \frac{N}{V}$; številčna gostota prostih nabojev (> 0)

- q_1 (v našem primeru $q_1 < 0$)

$\Rightarrow \rho_e = n q_1$: gostota naboja; $[\rho_e] = \frac{As}{m^3}$

- \vec{v} : (povprečna) hitrost nabojev (upr. pod vplivom el. sile zunanjega el. polja)

Def: $\boxed{\vec{j}_e \equiv \rho_e \vec{v} = q_1 n \vec{v}}$

vektor gostote el. toka

↔ Analogija: $\epsilon_0 \vec{E}$ = vektor, gostote el. polja

$\boxed{[\vec{j}_e] = \frac{As}{m^3} \frac{m}{s} = \frac{A}{m^2}}$

$\vec{v} = v \hat{e}_r$, kjer je $v = |\vec{v}| > 0$

$\Rightarrow \vec{j}_e = q_1 n v \hat{e}_r$

$\Rightarrow \boxed{j_e = |\vec{j}_e| = |q_1| n v} > 0$

$\Rightarrow \boxed{\vec{j}_e = |q_1| n v \frac{q_1}{|q_1|} \hat{e}_r = j \hat{e}_j}$, kjer je $\boxed{\hat{e}_j \equiv \frac{q_1}{|q_1|} \hat{e}_r = \begin{cases} +\hat{e}_r; q_1 > 0 \\ -\hat{e}_r; q_1 < 0 \end{cases}}$

Plošter $d\vec{S} = \hat{e}_s ds$:

Površina plošterve dS "dovolj" majhna, da je

a) plošter \sim rama

b) \vec{j}_e v neposredni okolici \sim konst. (homogeno vektorsko polje)

Def: $\boxed{dI \equiv \vec{j}_e \cdot d\vec{S}}$

El. tok skozi $d\vec{S}$

$$\boxed{[dI] = \frac{A}{m^2} m^2 = A}$$

! Analogija: $d\Phi_E \equiv \epsilon_0 \vec{E} \cdot d\vec{S} =$ električni pretok skozi $d\vec{S}$

Splošno: kriva plošter, nehomogeno polje \vec{j}_e :

Def: $\boxed{I \equiv \int_S \vec{j}_e \cdot d\vec{S}}$

! Analogija: $\Phi_E \equiv \int_S \epsilon_0 \vec{E} \cdot d\vec{S}$

Primer: homogeno polje \vec{j}_e , rama plošter \vec{S}

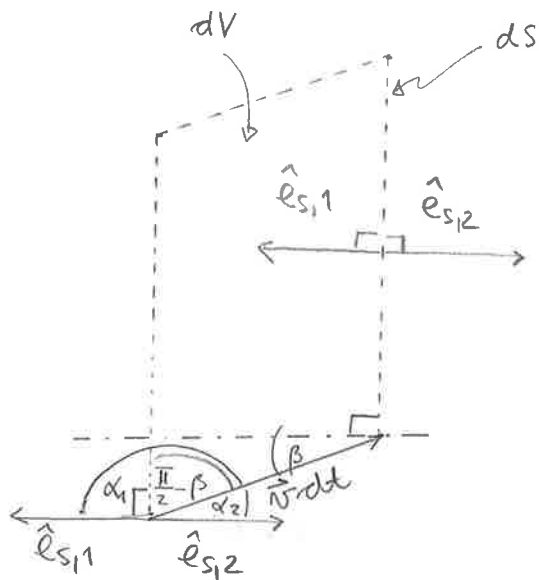
$$\Rightarrow \boxed{I = \int_S \vec{j}_e \cdot d\vec{S} = \int_S j_e \cos \phi dS = j_e \cos \phi \int_S dS = j_e \cos \phi S = \vec{j}_e \cdot \vec{S}}$$

$\phi: \angle \vec{j}_e, \vec{S} = \text{konst.}$
 $j_e = \text{konst.}$

Primer: enako, kot v prejšnjem primeru, in $\hat{e}_s = \pm \hat{e}_j$ (\vec{S} pravokotna na \vec{j})

$$\Rightarrow I = \begin{cases} +j_e S (>0) & ; \hat{e}_s = +\hat{e}_j \\ -j_e S (<0) & ; \hat{e}_s = -\hat{e}_j \end{cases}$$

Interpretacija el. toka



• β : vpadni kot majhni (ločni) kot med \vec{v} in \perp na dS ; $\cos\beta \geq 0$

• $dV = dS v dt \cos(\beta)$ prostornina, iz katere v času $[0, dt]$ vsi majhni prečja (deono) plošči dS

$$\Rightarrow dN = n dV = n v \cos\beta dS dt$$

$$\Rightarrow \boxed{d^3q = q_1 \cdot dN = q_1 n v \cos\beta dS dt}$$

= majh, ki se v časovnem intervalu dolžine dt prečja kot plošči dS v smeri \vec{v}

$$\Rightarrow \boxed{d^2\left(\frac{\partial q}{\partial t}\right) = q_1 n v \cos\beta dS}$$

Kako se to izrazi z $dI = \vec{j}_e \cdot d\vec{S}$?

α : \hat{e}_r, \hat{e}_s

$$(1) \alpha_1 = \frac{\pi}{2} + \frac{\pi}{2} - \beta = \pi - \beta$$

$$a) q_1 > 0 \Rightarrow \hat{e}_j = \hat{e}_r \Rightarrow \boxed{\vec{j}_e \cdot d\vec{S} = j_e dS \hat{e}_j \cdot \hat{e}_s = \overset{q_1}{|q_1|} n v dS \overset{-\cos\beta}{\cos(\pi - \beta)}} \\ = -q_1 n v \cos\beta dS = -\boxed{d^2\left(\frac{\partial q}{\partial t}\right)}$$

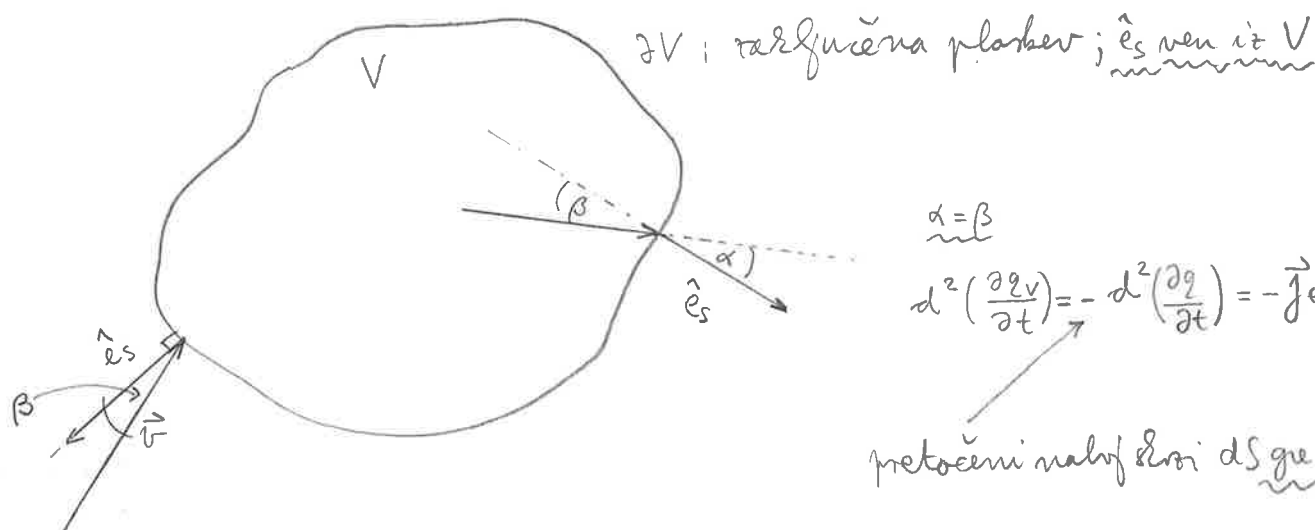
$$b) q_1 < 0 \Rightarrow \hat{e}_j = -\hat{e}_r \Rightarrow \boxed{\vec{j}_e \cdot d\vec{S} = j_e dS \hat{e}_j \cdot \hat{e}_s = -j_e dS \hat{e}_r \cdot \hat{e}_s = -\overset{-q_1}{|q_1|} n v dS \overset{-\cos\beta}{\cos(\pi - \beta)}} \\ = -q_1 n v \cos\beta dS = -\boxed{d^2\left(\frac{\partial q}{\partial t}\right)}$$

$$(2) \alpha_2 = \beta$$

$$a) q_1 > 0 \Rightarrow \hat{e}_j = \hat{e}_r \Rightarrow \boxed{\vec{j}_e \cdot d\vec{S} = \overset{q_1}{|q_1|} n v dS \cos\beta = q_1 n v \cos\beta dS = \boxed{d^2\left(\frac{\partial q}{\partial t}\right)}}$$

$$b) q_1 < 0 \Rightarrow \hat{e}_j = -\hat{e}_r \Rightarrow \boxed{\vec{j}_e \cdot d\vec{S} = - \overset{-q_1}{|q_1|} n v \cos \beta dS = q_1 n v \cos \beta dS = d^2 \left(\frac{\partial q}{\partial t} \right)}$$

$$\Rightarrow dI = \vec{j}_e \cdot d\vec{S} = \pm d^2 \left(\frac{\partial q}{\partial t} \right); \pm \text{ nabyj, ki se na enoto časa pretoci skozi } dS \text{ (interpretacija } dI)$$



$$\alpha = \pi - \beta:$$

$$d^2 \left(\frac{\partial q_V}{\partial t} \right): \text{sprememba naboja v } V \text{ na enoto časa}$$

$$d^2 \left(\frac{\partial q_V}{\partial t} \right) = + d^2 \left(\frac{\partial q_V}{\partial t} \right) = - \vec{j}_e \cdot d\vec{S}$$

↑ pretoceni nabyj skozi dS pride v V

$$\text{V obeh primerih: } \boxed{d^2 \left(\frac{\partial q_V}{\partial t} \right) = - \vec{j}_e \cdot d\vec{S}}$$

$$\Rightarrow \boxed{\left. \frac{\partial q_V}{\partial t} \right|_{\partial V} - \oint_{\partial V} \vec{j}_e \cdot d\vec{S}}$$

Sprememba naboja v V na enoto časa zaradi pretakanja naboja skozi ∂V

Zakon o ohranitvi (električnega) naboja - kontinuitetna enačba

Fizikalni zakon: (neto) električnega naboja ne moremo ne ustvariti ne izničiti.

Primeri:

$$\gamma \rightarrow e^+ e^-$$

$$q=0 \quad \underbrace{+q_0 \quad -q_0}_{q=0}$$

$$n \rightarrow p + e^- + \bar{\nu}_e$$

$$q=0 \quad \underbrace{+q_0 \quad -q_0 \quad 0}_{q=0}$$

$$p \rightarrow n + e^+ + \nu_e$$

$$\underbrace{+q_0 \quad 0 \quad +q_0 \quad 0}_{+q_0}$$

\Rightarrow Če hočemo v nekem območju V (neto) naboj povečati (zmanjšati), ga moramo (skoz ∂V) pripeljati iz okolice (odpeljati v okolico). Oz. neto naboj v V se (na enoto časa) spremeni le toliko, kolikor ga (na enoto časa) skoz ∂V pripeljemo iz okolice (ali odpeljemo v okolico).

Formalno:

$$q_V = \int_V \rho_e(\vec{r}) dV$$

$$\Rightarrow \boxed{\frac{\partial q_V}{\partial t} = \frac{\partial}{\partial t} \int_V \rho_e(\vec{r}) dV = \int_V \frac{\partial \rho_e}{\partial t} dV}$$

$$\boxed{\frac{\partial q_V}{\partial t} = \frac{\partial q_V}{\partial t} \Big|_{\partial V}} \quad \text{Zakon o ohranitvi naboja}$$