

# CS 498QC Final Project - Literature Review

## “On the Hardness of Detecting Macroscopic Superpositions”

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# Superposition and Mixture

## Definition

*Equal superpositions*  $|\psi\rangle$  and  $|\phi\rangle$  of orthogonal states  $|x\rangle$  and  $|y\rangle$ :

$$|\psi\rangle = \frac{1}{\sqrt{2}} |x\rangle + \frac{1}{\sqrt{2}} |y\rangle, \quad |\phi\rangle = \frac{1}{\sqrt{2}} |x\rangle - \frac{1}{\sqrt{2}} |y\rangle$$

$$|x\rangle = \frac{1}{\sqrt{2}} |\psi\rangle + \frac{1}{\sqrt{2}} |\phi\rangle, \quad |y\rangle = \frac{1}{\sqrt{2}} |\psi\rangle - \frac{1}{\sqrt{2}} |\phi\rangle$$

## Definition

An *equal mixture* of  $|x\rangle$  and  $|y\rangle$ , or equivalently,  $|\psi\rangle$  and  $|\phi\rangle$ :

$$\rho_{\text{mixed}} = \frac{1}{2} |x\rangle \langle x| + \frac{1}{2} |y\rangle \langle y| = \frac{1}{2} |\psi\rangle \langle \psi| + \frac{1}{2} |\phi\rangle \langle \phi|$$

# Distinguishing

## Definition

A **distinguisher** for  $|\psi\rangle$  and  $|\phi\rangle$  with bias  $\Delta$  is a circuit  $U$  that accepts  $|\psi\rangle$  with probability  $p$  and  $|\phi\rangle$  with probability  $p - \Delta$ .

Specifically, we measure the first qubit after applying  $U$  and accept the state if the outcome is  $|0\rangle$ .

$$\langle 0 | \text{tr}_{aux} \left( U |\psi\rangle \langle \psi| U^\dagger \right) | 0 \rangle = p$$

$$\langle 0 | \text{tr}_{aux} \left( U |\phi\rangle \langle \phi| U^\dagger \right) | 0 \rangle = p - \Delta$$

# Distinguishing

## Lemma

*A distinguisher  $U$  for  $|\psi\rangle$  and  $|\phi\rangle$  also distinguishes them from an equal mixture of  $|x\rangle$  and  $|y\rangle$  with bias  $\Delta/2$ .*

## Proof:

$$\begin{aligned} & \langle 0 | \text{tr}_{aux} \left( U \rho_{mixed} U^\dagger \right) | 0 \rangle \\ &= \langle 0 | \text{tr}_{aux} \left( U \left( \frac{1}{2} |\psi\rangle \langle \psi| + \frac{1}{2} |\phi\rangle \langle \phi| \right) U^\dagger \right) | 0 \rangle \\ &= \frac{1}{2} \langle 0 | \text{tr}_{aux} \left( U |\psi\rangle \langle \psi| U^\dagger \right) | 0 \rangle + \frac{1}{2} \langle 0 | \text{tr}_{aux} \left( U |\phi\rangle \langle \phi| U^\dagger \right) | 0 \rangle \\ &= \frac{|a|^2 + |c|^2}{2} = p - \frac{\Delta}{2} \end{aligned}$$

□

# Swapping

## Definition

A **swapper** of  $|x\rangle$  and  $|y\rangle$  with error  $\varepsilon$  is a unitary  $U$  that satisfies

$$\text{“fidelity” of } U := \frac{|\langle y|U|x\rangle| + |\langle x|U|y\rangle|}{2} = 1 - \varepsilon.$$

When  $\varepsilon = 0$ , either  $U$  or  $-U$  is a perfect swapper.

## Definition

The **swap complexity**  $\mathcal{S}_\varepsilon(|x\rangle, |y\rangle)$  is the minimum number of gates needed for a swapper of  $|x\rangle$  and  $|y\rangle$  with error at most  $\varepsilon$ .

# The Hardness of Detecting (Macroscopic) Superpositions

## Theorem (Distinguishability $\leftrightarrow$ Swappability)

*The circuit complexity of distinguishing  $|\psi\rangle$  and  $|\phi\rangle$  from an equal mixture of  $|x\rangle$  and  $|y\rangle$  with bias  $\Delta$  is the same as the circuit complexity of swapping  $|x\rangle$  and  $|y\rangle$  with fidelity  $2\Delta$ , i.e.,  $S_{1-2\Delta}(|x\rangle, |y\rangle)$ .*

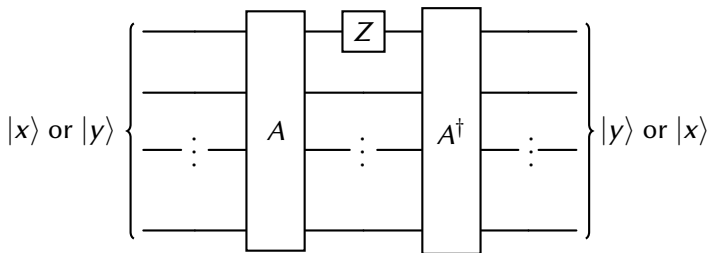
## Corollary

*When building an effective swapper for two orthogonal macroscopic states  $|x\rangle$  and  $|y\rangle$  is infeasible, it would also be infeasible to distinguish an equal superposition of  $|x\rangle$  and  $|y\rangle$  from an equal mixture of  $|x\rangle$  and  $|y\rangle$ , which means it's technologically impossible to prepare or even observe such superpositions.*

# Distinguishability $\rightarrow$ Swappability

## Theorem

*If we can construct a distinguisher  $A$  for  $|\psi\rangle$  and  $|\phi\rangle$  with bias  $\Delta$ , then with black-box access to  $A$  and  $A^\dagger$  and one additional gate, we can construct a swapper for  $|x\rangle$  and  $|y\rangle$  with fidelity  $\Delta$ .*



# Distinguishability $\rightarrow$ Swappability: Perfect Case

Suppose we have a perfect distinguisher  $A$ :

$$A|\psi\rangle = |0\rangle |g_0(\psi)\rangle, \quad A|\phi\rangle = |1\rangle |h_1(\phi)\rangle$$

Then

$$A|x\rangle = A\left(\frac{|\psi\rangle + |\phi\rangle}{\sqrt{2}}\right) = \frac{1}{\sqrt{2}}(|0\rangle |g_0(\psi)\rangle + |1\rangle |h_1(\phi)\rangle)$$

$$A|y\rangle = A\left(\frac{|\psi\rangle - |\phi\rangle}{\sqrt{2}}\right) = \frac{1}{\sqrt{2}}(|0\rangle |g_0(\psi)\rangle - |1\rangle |h_1(\phi)\rangle)$$

$$Z_0 A|x\rangle = \frac{1}{\sqrt{2}}(|0\rangle |g_0(\psi)\rangle - |1\rangle |h_1(\phi)\rangle) = A|y\rangle$$

$$Z_0 A|y\rangle = \frac{1}{\sqrt{2}}(|0\rangle |g_0(\psi)\rangle + |1\rangle |h_1(\phi)\rangle) = A|x\rangle$$

Circuit  $U = A^\dagger Z_0 A$  is a perfect swapper:

$$U|x\rangle = A^\dagger Z_0 A|x\rangle = |y\rangle, \quad U|y\rangle = A^\dagger Z_0 A|y\rangle = |x\rangle.$$



# Distinguishability $\rightarrow$ Swappability

**Proof:** Suppose we have a distinguisher  $A$  for  $|\psi\rangle$  and  $|\phi\rangle$ , such that

$$A|\psi\rangle = a|0\rangle|g_0(\psi)\rangle + b|1\rangle|g_1(\psi)\rangle$$

$$A|\phi\rangle = c|0\rangle|h_0(\phi)\rangle + d|1\rangle|h_1(\phi)\rangle$$

where

$$|a|^2 = p, \quad |c|^2 = p - \Delta,$$

then

$$\langle y|U|x\rangle = \langle y|A^\dagger Z_0 A|x\rangle = -\Delta + \frac{1}{2}[ac(C_0 - C_0^*) + bd(C_1 - C_1^*)]$$

$$\langle x|U|y\rangle = \langle x|A^\dagger Z_0 A|y\rangle = -\Delta + \frac{1}{2}[ac(C_0^* - C_0) + bd(C_1^* - C_1)]$$

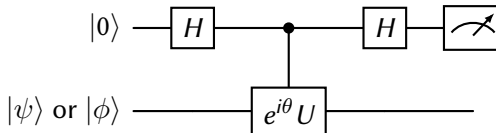
$$\frac{|\langle y|U|x\rangle + \langle x|U|y\rangle|}{2} = \Delta$$



# Swappability $\rightarrow$ Distinguishability

## Theorem

*If we can construct a swapper for  $|x\rangle$  and  $|y\rangle$  with fidelity  $\Delta$ , then with black-box access to a controlled  $e^{i\theta} U$  gate and two additional gates, we can construct a distinguisher for  $|\psi\rangle$  and  $|\phi\rangle$  with bias  $\Delta$ .*



# Swappability $\rightarrow$ Distinguishability: Perfect Case ( $\theta = 0$ )

**If the input state is  $|\psi\rangle$ :** Before controlled  $U$ , the state is

$$|+\rangle |\psi\rangle = \frac{1}{2} (|0\rangle |x\rangle + |0\rangle |y\rangle + |1\rangle |x\rangle + |1\rangle |y\rangle)$$

After controlled  $U$ , the state is

$$\frac{1}{2} (|0\rangle |x\rangle + |0\rangle |y\rangle + |1\rangle |y\rangle + |1\rangle |x\rangle) = |+\rangle |\psi\rangle$$

**If the input state is  $|\phi\rangle$ :** Before controlled  $U$ , the state is

$$|+\rangle |\phi\rangle = \frac{1}{2} (|0\rangle |x\rangle - |0\rangle |y\rangle + |1\rangle |x\rangle - |1\rangle |y\rangle)$$

After controlled  $U$ , the state is

$$\frac{1}{2} (|0\rangle |x\rangle - |0\rangle |y\rangle + |1\rangle |y\rangle - |1\rangle |x\rangle) = |-\rangle |\phi\rangle$$

# Swappability $\rightarrow$ Distinguishability

**Proof:** Suppose  $U$  satisfies

$$U|x\rangle = a|y\rangle + c|x\rangle + f|w\rangle$$

$$U|y\rangle = b|x\rangle + d|y\rangle + g|z\rangle$$

where  $|a + b|/2 = \Delta$ . The probabilities of accepting  $|\psi\rangle$  and  $|\phi\rangle$  are

$$\Pr_{\psi}(|+\rangle) = \frac{1}{2} + \frac{1}{4}\Re\left(e^{i\theta}(a + b + c + d)\right)$$

$$\Pr_{\phi}(|+\rangle) = \frac{1}{2} + \frac{1}{4}\Re\left(e^{i\theta}(-a - b + c + d)\right).$$

The bias is

$$\Pr_{\psi}(|+\rangle) - \Pr_{\phi}(|+\rangle) = \frac{\Re(e^{i\theta}(a + b))}{2} \leq \frac{|a + b|}{2} = \Delta$$

The equality is achieved when we set  $\theta = -\arg(a + b)$ . □

## Theorem

*For any  $0 \leq \Delta \leq 1$ , there exist a pair of orthogonal  $n$ -qubit states  $|x\rangle$  and  $|y\rangle$  and a swapper  $U$  of size  $O(1)$  that swaps them with fidelity  $\Delta$ , but to build a swapper with fidelity  $\Delta + \omega(0.5^{n/3}\sqrt{\log n})$  requires  $\omega(2^{n/3})$  gates.*

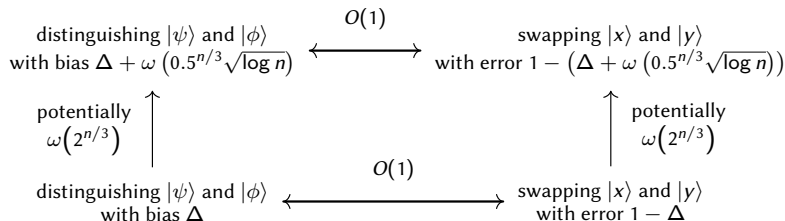
## Corollary

*For any  $0 \leq \Delta \leq 1$ , there exist a pair of  $n$ -qubit states  $|\psi\rangle$  and  $|\phi\rangle$  and a distinguisher  $A$  of size  $O(1)$  that distinguishes them with bias  $\Delta$ , but to build a distinguisher with bias  $\Delta + \omega(0.5^{n/3}\sqrt{\log n})$  requires  $\omega(2^{n/3})$  gates.*

# Summary

For any two orthogonal states  $|x\rangle$  and  $|y\rangle$  and

$$|\psi\rangle = \frac{1}{\sqrt{2}}|x\rangle + \frac{1}{\sqrt{2}}|y\rangle, \quad |\phi\rangle = \frac{1}{\sqrt{2}}|x\rangle - \frac{1}{\sqrt{2}}|y\rangle$$



# Tightness: Background

## Lemma

*For random  $n$ -qubit states  $|x\rangle$  and  $|y\rangle$ ,  $\Pr(|\langle x | y \rangle| \geq \varepsilon) \leq e^{-\varepsilon^2(2^n-1)}$ .*

The probability is *doubly* exponentially small in  $n$ .

## Lemma

*For random  $n$ -qubit states  $|x\rangle$  and  $|y\rangle$ , the probability that we can construct a circuit  $U$  such that  $\langle y | U | x \rangle \geq \varepsilon$ , for some  $\varepsilon$  of order  $O(0.5^{n/3} \sqrt{\log n})$ , using only  $O(2^{n/3})$  gates from a universal gate set of size  $n^{O(1)}$ , is  $0.5^{O(2^{n/3} \log n)}$ .*

The probability is still *doubly* exponentially small in  $n$ .

# Tightness: Proof (1/3)

If we randomly sample 8 states  $\{\psi_0, \psi_1, \dots, \psi_7\}$  from  $\mathbb{C}^{2^{n-3}}$ , they will be pairwise *almost* orthogonal with high probability. We could make them orthogonal by adding 3 index qubits, i.e.

$$|\bar{k}\rangle = |k\rangle \otimes |\psi_k\rangle, k \in \{0, 1\}^3.$$

Based on  $\{|\bar{k}\rangle\}$  we can construct a orthonormal basis of  $\mathbb{C}^{2^n}$ .



## Tightness: Proof (2/3)

For the following two orthogonal states:

$$\begin{aligned}|x\rangle &= \sqrt{a-b} \left( \frac{1}{2} \sum_{k=0}^3 |\bar{k}\rangle \right) + \sqrt{b} \left( \frac{|\bar{4}\rangle + |\bar{5}\rangle}{\sqrt{2}} \right) + \sqrt{c} |\bar{6}\rangle \\ |y\rangle &= \sqrt{a-b} \left( \frac{1}{2} \sum_{k=0}^3 i^k |\bar{k}\rangle \right) + \sqrt{b} \left( \frac{|\bar{4}\rangle - |\bar{5}\rangle}{\sqrt{2}} \right) + \sqrt{c} |\bar{7}\rangle,\end{aligned}$$

the following swapper of size  $O(1)$  achieves fidelity  $(a+b)/2$ :

$$U = \left( \sum_{k=0}^3 i^k |k\rangle \langle k| + |4\rangle \langle 4| - |5\rangle \langle 5| + |6\rangle \langle 6| + |7\rangle \langle 7| \right) \otimes \mathbb{I}_{2^{n-3}}$$

## Tightness: Proof (3/3)

Consider any swapper  $U$  of size  $O(2^{n/3})$  and its representation in a basis where  $\{|\bar{k}\rangle\}$  are the first 8 basis vectors. Let  $V$  be identical to  $U$  in the first  $8 \times 8$  entries and 0 everywhere else.  $V$  achieves the same fidelity as  $U$ :

$$a' = \langle y | U | x \rangle = \langle y | V | x \rangle, \quad b' = \langle x | U | y \rangle = \langle x | V | y \rangle$$

For any  $i \neq j$ , with an overwhelming probability,

$$|\langle \bar{i} | V | \bar{j} \rangle| = O(0.5^{n/3} \log n),$$

which means that

$$V = \sum_{k=0}^7 \beta_k e^{i\theta_k} |\bar{k}\rangle \langle \bar{k}| \pm O(0.5^{n/3} \log n), \text{ where } 0 \leq \beta_k \leq 1.$$

Therefore,

$$|a' + b'| \leq |a + b| + O(0.5^{n/3} \log n)$$