

# MATH 4347 Homework 5

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## 7.7

The general form of the full Fourier series is

$$\phi(x) = \frac{1}{2}A_0 + \sum_{n=1}^{\infty} A_n \cos nx + B_n \sin nx$$

where

$$A_n = \frac{1}{\pi} \int_{-\pi}^{\pi} \phi(x) \cos nx dx = \frac{1}{\pi} \int_{-\pi}^0 -\sin x \cos nx dx + \frac{1}{\pi} \int_0^{\pi} \sin x \cos nx dx = \frac{2((-1)^n + 1)}{\pi(1 - n^2)}$$

$$B_n = \frac{1}{\pi} \int_{-\pi}^{\pi} \phi(x) \sin nx dx = \frac{1}{\pi} \int_{-\pi}^0 -\sin x \cos nx dx + \frac{1}{\pi} \int_0^{\pi} \sin x \cos nx dx = 0$$

Therefore the series is

$$\begin{aligned} |\sin x| &= \frac{2}{\pi} + \sum_{n=1}^{\infty} \frac{2((-1)^n + 1)}{\pi(1 - n^2)} \cos nx = \frac{2}{\pi} + \sum_{\text{even}} \frac{4}{\pi(1 - n^2)} \cos nx \\ &= \frac{2}{\pi} + \sum_{n=1}^{\infty} \frac{4}{\pi(1 - 4n^2)} \cos 2nx \end{aligned}$$

Since the series converges pointwise, at  $x = 0$

$$|\sin 0| = \frac{2}{\pi} + \sum_{n=1}^{\infty} \frac{4}{\pi(1 - 4n^2)} \cos 0 \Rightarrow \boxed{\sum_{n=1}^{\infty} \frac{1}{4n^2 - 1} = \frac{1}{2}}$$

At  $x = \pi/2$

$$|\sin \frac{\pi}{2}| = \frac{2}{\pi} + \sum_{n=1}^{\infty} \frac{4}{\pi(1 - 4n^2)} \cos n\pi$$

$$1 = \frac{2}{\pi} + \sum_{n=1}^{\infty} \frac{4(-1)^n}{\pi(1 - 4n^2)}$$

$$\boxed{\sum_{n=1}^{\infty} \frac{(-1)^n}{4n^2 - 1} = \frac{1}{2} - \frac{\pi}{4}}$$

**b**

Let  $g(x)$  be a even 2-periodic function, whose value in interval  $[-1, 1]$  is  $g(x) = x^2$ . The Fourier coefficients are

$$A_0 = \int_{-1}^1 x^2 = \frac{2}{3}, \quad A_n = \int_{-1}^1 x^2 \cos n\pi x = \frac{4(-1)^n}{n^2\pi^2}, \quad B_n = \int_{-1}^1 x^2 \sin n\pi x = 0$$

The full series is

$$x^2 = \frac{1}{3} + \sum_{n=1}^{\infty} \frac{4(-1)^n}{n^2\pi^2} \cos n\pi x$$

Evaluated at  $x = 1$

$$1 = \frac{1}{3} + \sum_{n=1}^{\infty} \frac{4}{n^2\pi^2} \Rightarrow \sum_{n=1}^{\infty} \frac{1}{n^2} = \frac{\pi^2}{6}$$

Evaluated at  $x = 0$

$$\sum_{n=1}^{\infty} \frac{(-1)^n}{n^2} = -\frac{\pi^2}{12}$$

## 8.2

If  $u(x) \equiv c$ , then  $\max_{\bar{U}} u = \max_{\partial U} u \equiv c$ .

Otherwise, by the strong form of the maximum principle,

$$\begin{aligned} \forall x \in U : u(x) < \max_{\partial U} u &\Rightarrow \max_U u < \max_{\partial U} u \\ \max_{\bar{U}} u = \max\{\max_{\partial U} u, \max_U u\} &= \max_{\partial U} u \end{aligned}$$

### 5.2.11

The complex form of full Fourier series is

$$\phi(x) = \sum_{n=-\infty}^{\infty} c_n e^{in\pi x/l} = e^x$$

where the coefficients are

$$\begin{aligned} c_n &= \frac{1}{2l} \int_{-l}^l e^x e^{-in\pi x/l} dx \\ &= \frac{1}{2l} \int_{-l}^l e^{(1-in\pi/l)x} dx \\ &= \frac{1}{2(l-in\pi)} e^{(1-in\pi/l)x} \Big|_{-l}^l \\ &= \frac{e^{l-in\pi} - e^{-l-in\pi}}{2(l-in\pi)} \\ &= (-1)^n \frac{e^l - e^{-l}}{2(l-in\pi)} \end{aligned}$$

So the series is

$$e^x = \sum_{n=-\infty}^{\infty} (-1)^n \frac{e^l - e^{-l}}{2(l - in\pi)} e^{in\pi x/l} = \boxed{\sum_{n=-\infty}^{\infty} (-1)^n \sinh l \frac{l + in\pi}{l^2 + n^2\pi^2} e^{in\pi x/l}}$$

And since  $e^{i\theta} = \cos \theta + i \sin \theta$ ,

$$\begin{aligned} e^x &= \frac{e^l - e^{-l}}{2l} + \sum_{n=-1}^{-\infty} (-1)^n \frac{e^l - e^{-l}}{2(l - in\pi)} \cos \frac{n\pi x}{l} + \sum_{n=1}^{\infty} (-1)^n \frac{e^l - e^{-l}}{2(l - in\pi)} \cos \frac{n\pi x}{l} \\ &\quad + \sum_{n=-1}^{-\infty} (-1)^n \frac{ie^l - ie^{-l}}{2(l - in\pi)} \sin \frac{n\pi x}{l} + \sum_{n=1}^{\infty} (-1)^n \frac{ie^l - ie^{-l}}{2(l - in\pi)} \sin \frac{n\pi x}{l} \\ &= \frac{e^l - e^{-l}}{2l} + \sum_{n=1}^{\infty} (-1)^n \frac{e^l - e^{-l}}{2(l + in\pi)} \cos \frac{n\pi x}{l} + \sum_{n=1}^{\infty} (-1)^n \frac{e^l - e^{-l}}{2(l - in\pi)} \cos \frac{n\pi x}{l} \\ &\quad + \sum_{n=1}^{\infty} (-1)^n \frac{ie^{-l} - ie^l}{2(l + in\pi)} \sin \frac{n\pi x}{l} + \sum_{n=1}^{\infty} (-1)^n \frac{ie^l - ie^{-l}}{2(l - in\pi)} \sin \frac{n\pi x}{l} \\ &= \frac{e^l - e^{-l}}{2l} + \sum_{n=1}^{\infty} (-1)^n \frac{(e^l - e^{-l})l}{l^2 + n^2\pi^2} \cos \frac{n\pi x}{l} + \sum_{n=1}^{\infty} (-1)^n \frac{n\pi(e^{-l} - e^l)}{l^2 + n^2\pi^2} \sin \frac{n\pi x}{l} \\ &= \boxed{\frac{\sinh l}{l} + 2 \sinh l \sum_{n=1}^{\infty} \frac{(-1)^n}{l^2 + n^2\pi^2} \left[ l \cos \frac{n\pi x}{l} - n\pi \sin \frac{n\pi x}{l} \right]} \end{aligned}$$

### 5.3.4

**a**

Let  $v = u - U$ , then all derivatives of  $v$  is the same as  $u$ , and  $v(0, t) = u(0, t) - U = 0$ . Separation of variables gives

$$\frac{X''}{X} = \frac{T'}{kT} = \lambda, \quad v = X(x)T(t)$$

$\lambda = 0$  under the boundary condition gives only the trivial solution, and for  $\lambda > 0$  the boundary conditions cannot be both satisfied, therefore  $\lambda < 0$ . Let  $\lambda = -\beta^2$ , then

$$X = A \cos \beta x + B \sin \beta x$$

Boundary condition at 0 implies  $A = 0$ , and boundary condition at  $x = l$  implies

$$\cos \beta l = 0 \Rightarrow \beta_n = \left(n - \frac{1}{2}\right) \frac{\pi}{l} = (2n - 1) \frac{\pi}{2l}$$

$$X_n = \sin \beta_n x$$

The corresponding  $T_n$  are

$$T' = -k\beta_n^2 T \Rightarrow T = e^{-k\beta_n^2 t}$$

So the general solution is

$$\begin{aligned}
v &= \sum_{n=1}^{\infty} A_n e^{-k\beta_n^2 t} \sin \beta_n x, \quad \beta_n = (2n-1)\frac{\pi}{2l} \\
v(x, 0) &= \sum_{n=1}^{\infty} A_n \sin \beta_n x = u(x, 0) - U = -U \\
\Rightarrow A_n &= \frac{2}{l} \int_0^l -U \sin \beta_n x = -\frac{2U}{l\beta_n} = -\frac{4U}{(2n-1)\pi} \\
\Rightarrow v &= -\sum_{n=1}^{\infty} \frac{4U}{(2n-1)\pi} e^{-k\pi^2(2n-1)^2 t/4l^2} \sin \frac{(2n-1)\pi}{2l} x \\
u = v + U &= \boxed{U - \sum_{n=1}^{\infty} \frac{4U}{(2n-1)\pi} e^{-k\pi^2(2n-1)^2 t/4l^2} \sin \frac{(2n-1)\pi}{2l} x}
\end{aligned}$$

**b**

Let  $a_n$  be the  $n$ -th term of the series. Apply the ratio test:

$$\begin{aligned}
\lim_{n \rightarrow \infty} \left| \frac{a_{n+1}}{a_n} \right| &= \lim_{n \rightarrow \infty} \left| \frac{2n-1}{2n+1} \frac{\sin \beta_{n+1} x}{\sin \beta_n x} e^{-2k\pi^2 n t/l^2} \right| \\
&= \lim_{n \rightarrow \infty} \frac{2n-1}{2n+1} e^{-2k\pi^2 n t/l^2} \left| \frac{\sin \beta_{n+1} x}{\sin \beta_n x} \right| \\
&\leq \lim_{n \rightarrow \infty} \frac{2n-1}{2n+1} e^{-2k\pi^2 n t/l^2} = 0
\end{aligned}$$

Therefore, the series converges.

**c**

The error is smaller than the second term, which is

$$-\frac{4U}{\pi} e^{-k\pi^2 t/4l^2}$$

In order for the error to be within  $\epsilon$ , it suffices to let

$$\left| \frac{4U}{\pi} e^{-k\pi^2 t/4l^2} \right| = \frac{4|U|}{\pi} e^{-k\pi^2 t/4l^2} < \epsilon$$

$$\boxed{t > -\frac{4l^2}{k\pi^2} \log \frac{\epsilon\pi}{4|U|}}$$

### 5.3.10

**a**

Proof by induction:

Base case: To show that  $(Z_2, Z_1) = 0$ , it's sufficient to show that  $(Y_2, X_1) = 0$ :

$$\begin{aligned}(Y_2, X_1) &= (X_2, X_1) - \left(X_2, \frac{X_1}{\|X_1\|}\right) \left(\frac{X_1}{\|X_1\|}, X_1\right) \\ &= (X_2, X_1) - (X_2, X_1) \frac{(X_1, X_1)}{\|X_1\|^2} \\ &= (X_2, X_1) - (X_2, X_1) = 0\end{aligned}$$

Inductive step: Suppose  $\forall n, m \leq k, n \neq m : (Z_n, Z_m) = 0$ . Then in order to prove that  $\forall n, m \leq k+1, n \neq m : (Z_n, Z_m) = 0$ , we only need to prove that  $(Z_{k+1}, Z_n) = 0$  for all  $n \leq k$ , because all other cases are already proven in the  $k$ -th step. And to that end, it's sufficient to show that  $(Y_{k+1}, Z_n) = 0$ . Since the vector in  $\{Z_n : n \leq k\}$  are orthonormal by inductive hypothesis,

$$\begin{aligned}(Y_{k+1}, Z_n) &= (X_{k+1}, Z_n) - (X_{k+1}, Z_1)(Z_1, Z_n) - \cdots - (X_{k+1}, Z_n)(Z_n, Z_n) - \cdots \\ &= (X_{k+1}, Z_n) - (X_{k+1}, Z_n)(Z_n, Z_n) \\ &= (X_{k+1}, Z_n) - (X_{k+1}, Z_n) = 0\end{aligned}$$

By induction, orthogonality holds for all  $k$ .

**b**

Let  $X_1 = \cos x + \cos 2x$ , and  $X_2 = 3 \cos x - 4 \cos 2x$ , then

$$(X_1, X_1) = (\cos x, \cos x) + (\cos 2x, \cos 2x) = \pi, \quad \|X_1\| = \sqrt{\pi}$$

$$Z_1 = \frac{X_1}{\|X_1\|} = \boxed{\frac{1}{\sqrt{\pi}}(\cos x + \cos 2x)}$$

$$(X_2, Z_1) = \frac{1}{\sqrt{\pi}} \left[ 3(\cos x, \cos x) - 4(\cos 2x, \cos 2x) \right] = -\frac{\sqrt{\pi}}{2}$$

$$\begin{aligned}Y_2 &= X_2 - (X_2, Z_1)Z_1 = 3 \cos x - 4 \cos 2x - \left(-\frac{\sqrt{\pi}}{2}\right) \frac{1}{\sqrt{\pi}}(\cos x + \cos 2x) \\ &= 3 \cos x - 4 \cos 2x + \frac{1}{2}(\cos x + \cos 2x) = \frac{7}{2} \cos x - \frac{7}{2} \cos 2x\end{aligned}$$

$$(Y_2, Y_2) = \frac{49}{4} \frac{\pi}{2} + \frac{49}{4} \frac{\pi}{2} = \frac{49\pi}{4}, \quad \|Y_2\| = \frac{7\sqrt{\pi}}{2}$$

$$Z_2 = \frac{Y_2}{\|Y_2\|} = \boxed{\frac{1}{\sqrt{\pi}}(\cos x - \cos 2x)}$$

### 5.4.7

**a**

The general formula is

$$\phi(x) = \frac{1}{2}A_0 + \sum_{n=1}^{\infty} A_n \cos n\pi x + B_n \sin n\pi x$$

where

$$A_n = \int_{-1}^1 \phi(x) \cos n\pi x dx = \int_{-1}^0 (-1-x) \cos n\pi x dx + \int_0^1 (1-x) \cos n\pi x dx = 0$$

$$B_n = \int_{-1}^1 \phi(x) \sin n\pi x dx = \int_{-1}^0 (-1-x) \sin n\pi x dx + \int_0^1 (1-x) \sin n\pi x dx = \frac{2}{n\pi}$$

Therefore the full series is

$$\phi(x) = \sum_{n=1}^{\infty} \frac{2}{n\pi} \sin n\pi x$$

**b**

The first three non-zero terms are

$$\frac{2}{\pi} \sin \pi x + \frac{1}{\pi} \sin 2\pi x + \frac{2}{3\pi} \sin 3\pi x$$

**c**

Obviously,

$$\|\phi(x)\|^2 = \int_{-1}^1 \phi^2(x) dx < \infty$$

Therefore, the series converges in the mean square sense.

**d**

$$\phi'(x) = -1, \text{ for } x \in (-1, 0) \cup (0, 1)$$

Since  $\phi$  and  $\phi'$  are both piecewise continuous, the series converges pointwise. [At  $x = 0$ , it converges to  $\frac{1}{2}(-1 + 1) = 0$ .]

**e**

$\phi(x)$  has a discontinuity at 0, so  $\phi \notin C^2[-1, 1]$ , and therefore the series does not converge uniformly.

### 5.5.4

**a**

Separation of variables gives

$$-\frac{X''}{X} = -\frac{T'}{kT} = \lambda$$

For  $\lambda = 0$ ,

$$\begin{aligned} X'' = 0 &\Rightarrow X = A + Bx \\ T' = 0 &\Rightarrow T = \text{const.} \end{aligned}$$

For  $\lambda > 0$ ,

$$\begin{aligned} -X'' = \beta^2 X &\Rightarrow X = C \cos \beta x + D \sin \beta x \\ T' = -k\beta^2 T &\Rightarrow T = e^{-k\beta^2 t} \end{aligned}$$

The general solution is

$$u = A + Bx + \sum_{n=1}^{\infty} e^{-k\beta_n^2 t} [C_n \cos \beta_n x + D_n \sin \beta_n x]$$

**b**

As  $t \rightarrow \infty$ , each term in the sum converges to zero, therefore  $\lim u = A + Bx$ .

**c**

From the boundary condition

$$\begin{aligned} u_x(0, t) = u_x(l, t) &= \frac{u(l, t) - u(0, t)}{l} \\ X'(0)T(t) = X'(l)T(t) &= \frac{X(l)T(t) - X(0)T(t)}{l} \\ X'(0) = X'(l) &= \frac{X(l) - X(0)}{l} \end{aligned}$$

Green's first identity in one dimension is

$$vu' \Big|_0^l = \int_0^l v'u'dx + \int_0^l vu''dx$$

Let  $v = u = X$ , then

$$\begin{aligned} LHS &= XX' \Big|_0^l \\ &= X(l)X'(l) - X(0)X'(0) \\ &= X(l)\frac{X(l) - X(0)}{l} - X(0)\frac{X(l) - X(0)}{l} \\ &= \frac{[X(l) - X(0)]^2}{l} \end{aligned}$$

$$\begin{aligned}
RHS &= \int_0^l (X')^2 dx + \int_0^l X X'' dx \\
&= \int_0^l (X')^2 dx - \int_0^l \lambda X^2 dx
\end{aligned}$$

Multiply both sides by  $l$  and swap both sides,

$$l \int_0^l (X')^2 dx - l \int_0^l \lambda X^2 dx = [X(l) - X(0)]^2$$

If  $\lambda$  is negative, then the second integral is negative, which means that

$$l \int_0^l (X')^2 dx < [X(l) - X(0)]^2$$

This contradicts the inequality in Ex. 3. Therefore, there cannot be negative eigenvalues.

**d**

First we can verify that the boundary condition is symmetric:

$$\begin{aligned}
X'_1(l)X_2(l) - X'_2(l)X_1(l) &= X_2(l) \frac{X_1(l) - X_1(0)}{l} - X_1(l) \frac{X_2(l) - X_2(0)}{l} \\
&= \frac{X_1(l)X_2(0) - X_1(0)X_2(l)}{l}
\end{aligned}$$

$$\begin{aligned}
X'_1(0)X_2(0) - X'_2(0)X_1(0) &= X_2(0) \frac{X_1(l) - X_1(0)}{l} - X_1(0) \frac{X_2(l) - X_2(0)}{l} \\
&= \frac{X_1(l)X_2(0) - X_1(0)X_2(l)}{l}
\end{aligned}$$

$$X'_1X_2 - X'_2X_1 \Big|_0^l = 0$$

Therefore eigenfunctions associated with different eigenvalues are orthogonal. At  $t = 0$

$$\phi(x) = u(x, 0) = A + Bx + \sum_{n=1}^{\infty} C \cos \beta_n x + D \sin \beta_n x$$

$$(\phi, 1) = (A + Bx, 1)$$

$$\Rightarrow \int_0^l \phi(x) dx = \int_0^l A + Bx dx = \left[ Ax + \frac{B}{2} x^2 \right]_0^l = lA + \frac{l^2}{2} B$$

$$(\phi, x) = (A + Bx, x)$$

$$\Rightarrow \int_0^l \phi(x) x dx = \int_0^l Ax + Bx^2 dx = \left[ \frac{A}{2} x^2 + \frac{B}{3} x^3 \right]_0^l = \frac{l^2}{2} A + \frac{l^3}{3} B$$

$$\boxed{A = \frac{4}{l} \int_0^l \phi(x) dx - \frac{6}{l^2} \int_0^l \phi(x) x dx}$$

$$\boxed{B = \frac{12}{l^3} \int_0^l \phi(x) x dx - \frac{6}{l^2} \int_0^l \phi(x) dx}$$



### 6.1.2

The equation is  $\Delta u = k^2 u$ . Expressed in spherical coordinates, dropping zero terms:

$$\frac{1}{r} \partial_r^2 (ru) = k^2 u$$

Let  $v = ur$ . Multiply the equation by  $r$

$$v''(r) = k^2 v(r) \Rightarrow v = Ae^{kr} + Be^{-kr}$$

$$u = \frac{v}{r} = \boxed{\frac{1}{r} [Ae^{kr} + Be^{-kr}]}$$

### 6.1.9

**a**

The heat equation is  $u_t = \Delta u$ . In steady state,  $u_t = \Delta u = 0$ . Since  $u$  only depends on  $r$ ,

$$\frac{1}{r} \frac{\partial^2}{\partial r^2} (ru) = 0$$

Let  $v(r) = ru(r)$ , then  $v''(r) = 0 \Rightarrow v(r) = Ar + B$ ,

$$u(r) = \frac{v}{r} = A + \frac{B}{r}$$

$$u'(r) = -B \frac{1}{r^2}$$

Plug in the boundary conditions:

$$\begin{cases} u(1) = A + B = 100 \\ u'(2) = -\frac{B}{4} = -\gamma \end{cases} \Rightarrow \begin{cases} A = 100 - 4\gamma \\ B = 4\gamma \end{cases}$$

Therefore, the temperature distribution is

$$\boxed{u(r) = 100 - 4\gamma + \frac{4\gamma}{r}}$$

**b**

$$u'(r) = -\frac{4\gamma}{r^2} < 0$$

Temperature decreases radially, therefore the highest temperature is obtained at  $r = 1$ , which is  $\boxed{100}$ , and the lowest temperature is obtained at  $r = 2$ , which is  $\boxed{100 - 2\gamma}$

**c**

Yes.  $\gamma = 40$ .