## PHYS 7125 Homework 5

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The total proper time along a curve  $\gamma$  is

$$\tau_{total} = \int_{\gamma} d\tau = \int_{\gamma} \sqrt{-ds^2} = \int_{\gamma} \sqrt{-\left(\frac{2M}{r}-1\right)dt^2 + \left(\frac{2M}{r}-1\right)^{-1}dr^2 - r^2d\Omega^2}$$

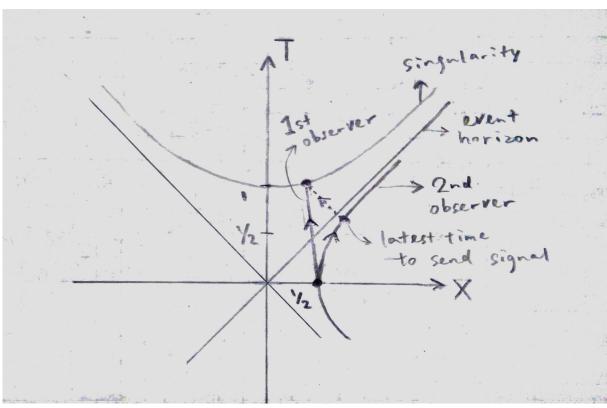
Inside the event horizon, r < 2M, the first and third term in the square root are negative, therefore

$$\tau_{total} < \int_{0}^{2M} \sqrt{\left(\frac{2M}{r} - 1\right)^{-1}} dr$$

$$= \left[ -\sqrt{r(2M - r)} + 2M \cot^{-1} \left(\sqrt{\frac{2M}{r} - 1}\right) \right]_{0}^{2M} = \pi M$$

 $\mathbf{2}$ 

 $\mathbf{a}$ 



## b

Yes, because the observer is massive. (If not, the light signal in (c) could never reach the falling observer, because the two worldlines would be parallel straight lines in a Kruskal diagram.)

 $\mathbf{c}$ 

At constant r = R,

$$T = \frac{1}{2}\sinh\left(\frac{t}{4GM}\right), \quad X = \frac{1}{2}\cosh\left(\frac{t}{4GM}\right), \quad 4X^2 - 4T^2 = 1$$

At t = 0 on this worldline,

$$T = \frac{1}{2}\sinh 0 = 0, \quad X = \frac{1}{2}\cosh 0 = \frac{1}{2}$$

At the singularity,

$$T = \cosh\left(\frac{t}{4GM}\right), \quad X = \sinh\left(\frac{t}{4GM}\right), \quad T^2 - X^2 = 1$$

A timelike straight line passing through (1/2,0) can be expressed as

$$X = kT + \frac{1}{2}$$

where -1 < k < 1. When the first observer reaches singularity,

$$(1-k^2)T^2 - kT - \frac{5}{4} = 0 \quad \Rightarrow \quad \left| T_s = \frac{k + \sqrt{5 - 4k^2}}{2(1-k^2)}, \quad X_s = \frac{k^2 + k\sqrt{5 - 4k^2}}{2(1-k^2)} + \frac{1}{2} \right|$$

To reach this critical point, the photon emitted by the second observer must follow the straight line

$$X = -T + T_{\circ} + X_{\circ}$$

The photon's worldline in the past intersects the second observer's worldline at

$$AT^2 + 4(T_s + X_s)^2 - 8(T_s + X_s)T - AT^2 = 1$$

$$\Rightarrow T = \frac{1}{2}\sinh\left(\frac{t}{4GM}\right) = \frac{4(T_s + X_s)^2 - 1}{8(T_s + X_s)}, \quad \boxed{t = 4GM\sinh^{-1}\left(\frac{4(T_s + X_s)^2 - 1}{4(T_s + X_s)}\right)}$$

which is the latest Schwarzschild time at which the second observer could send the signal. If the worldline of the first observer is vertical in the Kruskal diagram, i.e. k = 0, then  $t \approx 4.70GM$ 

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A perfect fluid is incompressible, therefore  $\nabla_{\mu}\rho = \partial_{\mu}\rho = 0$ . From the continuity equation,

$$\nabla_{\mu}(\rho u^{\mu}) = \rho \nabla_{\mu} u^{\mu} = 0 \quad \Rightarrow \quad \nabla_{\mu} u^{\mu} = 0$$

Since the covariant divergence of the stress-energy tensor must vanish,

$$\nabla_{\mu}T^{\mu}{}_{\nu} = (\nabla_{\mu}p + \nabla_{\mu}\rho)u^{\mu}u_{\nu} + (p+\rho)(\nabla_{\mu}u^{\mu})u_{\nu} + (p+\rho)u^{\mu}(\nabla_{\mu}u_{\nu}) + \delta^{\mu}_{\nu}\nabla_{\mu}p$$

$$= u_{\nu}(\nabla_{\underline{u}}p) + (p+\rho)\nabla_{\underline{u}}u_{\nu} + \nabla_{\nu}p$$

$$= 0$$

The free index  $\nu$  can be dropped, yielding

$$(p+\rho)\nabla_u u = -\nabla p - u\nabla_u p$$