

# MATH 4347 Homework 1

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## 1.4

$$\begin{aligned}u_t &= a'(t)e^{2x} + b'(t)e^x + c'(t) \\u_{xx} &= 4a(t)e^{2x} + b(t)e^x\end{aligned}$$

Compare the two derivatives, we get

$$\begin{cases} a' = 4a \\ b' = b \\ c' = 0 \end{cases} \Rightarrow \begin{cases} a(t) = C_1 e^{4t} \\ b(t) = C_2 e^t \\ c(t) = C_3 \end{cases}$$

## 1.7

The characteristic lines are  $x - 4t = m$ , so  $u = f(x - 4t)$ . From the initial condition

$$u(x, 0) = f(x) = 0, \text{ where } x > 0$$

From the boundary condition

$$u(0, t) = f(-4t) = te^{-t} \text{ where } t > 0$$

Let  $x = -4t$ , then  $t = -\frac{1}{4}x$ .

$$f(x) = -\frac{1}{4}xe^{\frac{1}{4}x} \text{ where } x < 0$$

Combining above results

$$f(x) = \begin{cases} -\frac{1}{4}xe^{\frac{1}{4}x} & \text{for } x < 0 \\ 0, & \text{for } x > 0 \end{cases}$$

$$u = f(x - 4t) = \begin{cases} -\frac{1}{4}(x - 4t)e^{\frac{1}{4}(x - 4t)} & \text{for } 0 < x < 4t \\ 0, & \text{for } x > 4t \end{cases}$$

If the PDE changed to  $u_t - 4u_x = 0$ , then  $u = f(x + 4t)$  and the boundary condition implies that

$$f(4t) = te^{-t} \text{ where } t > 0$$

This is a contradiction because the initial condition tells us that  $f(x) = 0$  for positive arguments, and therefore the equation has no solution.

## 2.5

**a**

Let  $u(x, t) = e^{i\xi x + \sigma t}$ , then  $u_{tt} = \sigma^2 u$ , and  $u_{xxx} = \xi^3 u$ . So the PDE becomes

$$\sigma^2 = -\xi^4 \Rightarrow \sigma = \pm i\xi^2$$

$$u(x, t) = e^{i\xi x \pm i\xi^2 t} = e^{i\xi(x \pm \xi t)}$$

The wave speed depends on  $\xi$ , so the wave is dispersive. The reason for this dependence is that the derivatives in space and time are of different orders, so the resulting dispersion relation  $\sigma(\xi)$  is not linear, so there are powers of  $\xi$  that cannot be factored out. By contrast, the space and time derivatives in wave equation are both of the second order, therefore  $\sigma = \pm c\xi$ , and consequently the wave speed does not depend on wave number.

**b**

$$u_t = \sigma u$$

$$u_x = i\xi u$$

$$u_{xxt} = -\xi^2 \sigma u$$

Dividing the original PDE by  $u$

$$\sigma + ic\xi - \beta\xi^2\sigma = 0 \Rightarrow \sigma = \frac{ic\xi}{\beta\xi^2 - 1}$$

So the solution is

$$u = e^{i\xi x + \frac{ic\xi}{\beta\xi^2 - 1}t} = e^{i\xi(x - \frac{c}{1 - \beta\xi^2}t)}$$

which are traveling waves with speed  $\frac{c}{1 - \beta\xi^2}$ . Compared to the solution of linearized KdV equation

$$u(x, t) = e^{i\xi(x - (c - \beta\xi^2)t)}$$

which has wave speed  $c - \beta\xi^2$ , the wave speed of the solutions of BBM equation is monotonically increasing w.r.t.  $\xi^2$  when  $\xi^2 < \frac{1}{\beta}$  and  $\xi^2 > \frac{1}{\beta}$ , however the wave speed of the solutions of KdV equation is monotonically decreasing w.r.t.  $\xi^2$

## 2.6

**a**

Intuitively, the velocity should be higher when there are less traffic, so  $v(u)$  should be monotonically decreasing.

**b**

Since  $v(u)$  is monotonically decreasing and both  $u$  and  $v$  need to be nonnegative to have physical meanings,  $v_{max} = v(u_{min}) = v(0)$  and  $v(u_{max}) = v_{min} = 0$

**c**

From (b),  $Q(0) = 0v(0) = 0$ , and  $Q(u_{max}) = u_{max}v(u_{max}) = u_{max}0 = 0$ . If  $Q$  doesn't attain maximum at some point in the interval  $(0, u_{max})$ , then  $Q(0) = Q(u_{max}) = 0$  must be the maximum, and for all  $u$  in the interval,  $Q(u)$  must be strictly less than 0 (otherwise these points will also maximum), which is a contradiction because we required that  $u$  and  $v$  both be non-negative. Therefore,  $Q(u)$  must have a maximum in the interval.

**d**

Such function can be constructed. Consider

$$Q(u) = u \left( -\frac{1}{4}u^3 + 2u^2 - \frac{11}{2}u + 6 \right) = -\frac{1}{4}u^4 + 2u^3 - \frac{11}{2}u^2 + 6u$$

whose derivative is

$$Q'(u) = -(u-1)(u-2)(u-3)$$

The function has maximum at 1 and 3, and we can check that

$$v(u) = -\frac{1}{4}u^3 + 2u^2 - \frac{11}{2}u + 6$$

is positive and monotonically decreasing in interval  $(0, 4)$ .

## 3.2

First solve for the characteristic curves:

$$\frac{dx}{dt} = \frac{1}{1+t^2} \Rightarrow x = \arctan t + m$$

So  $u(x, t) = f(m) = f(x - \arctan t)$ . Plug in the initial condition:

$$u(x, 0) = f(x - \arctan 0) = f(x) = \sin x$$

Therefore,

$$u(x, t) = f(x - \arctan t) = \sin(x - \arctan t)$$

### 3.3

Treat the first two terms as the directional derivative of  $u$  in  $(1, 1)$  direction. Consider the characteristic lines  $x = t + k$ , or  $(s, s + k)$  where the equation becomes an ODE:

$$\frac{d}{ds}u + 3u = e^{3s+2k}$$

Solving use integrating factor  $\phi(s) = e^{\int 3ds} = e^{3s}$

$$\begin{aligned} u(s) &= \frac{1}{e^{3s}} \int e^{3s} e^{3s+2k} ds = \frac{1}{6} e^{3s+2k} + A(k) e^{-3s} \\ u(x, t) &= \frac{1}{6} e^{2x+t} + A(x-t) e^{-3t} \end{aligned}$$

Plugin the initial condition to solve for  $A(x)$ :

$$u(x, 0) = \frac{1}{6} e^{2x} + A(x) = x \Rightarrow A(x) = x - \frac{1}{6} e^{2x}$$

Thus the solution is

$$\begin{aligned} u(x, t) &= \frac{1}{6} e^{2x+t} + \left( (x-t) - \frac{1}{6} e^{2(x-t)} \right) e^{-3t} \\ &= \frac{1}{6} e^{2x+t} - \frac{1}{6} e^{2x-5t} + x e^{-3t} - t e^{-3t} \end{aligned}$$

### 3.8

**a**

the LHS can be treated as a total derivative w.r.t to  $t$  on the characteristic curve

$$\frac{dx}{dt} = t \Rightarrow x = \frac{t^2}{2} + k$$

On the characteristic curve  $(s, \frac{s^2}{2} + k)$  the PDE becomes ODE

$$\frac{du}{ds} = u^2 \Rightarrow u = -\frac{1}{s + A(k)} = -\frac{1}{t + A(x - t^2/2)}$$

Plug in the initial condition:

$$u(x, 0) = -\frac{1}{A(x)} = \frac{1}{1+x^2} \Rightarrow A(x) = -1 - x^2$$

Therefore the solution is

$$u(x, t) = -\frac{1}{t - 1 - (x - t^2/2)^2} = \frac{1}{(x - t^2/2)^2 - t + 1}$$

**b**

Given  $t$ ,  $u(x, t)$  attains maximum at critical point

$$\frac{\partial u}{\partial x} = -u^2 \cdot 2(x - \frac{t^2}{2}) = 0 \Rightarrow x_c = \frac{t^2}{2}$$

$$\max_x u(x, t) = u(\frac{t^2}{2}, t) = \frac{1}{1-t}$$

which blows up as  $t \rightarrow 1^-$

### 1.3. #2

Apply Newton's second law to a segment of the chain from  $x = a$  to  $x = b$ :

$$F(b) \sin \theta - F(a) \sin \theta = \int_a^b \rho u_{tt} \frac{1}{\cos \theta} dx$$

$$\int_b^{L_x} \rho g \frac{\sin \theta}{\cos \theta} dx - \int_a^{L_x} \rho g \frac{\sin \theta}{\cos \theta} dx = - \int_a^b \rho g \frac{\sin \theta}{\cos \theta} dx = \int_a^b \rho u_{tt} \frac{1}{\cos \theta} dx$$

Since this must be true for arbitrary  $a, b$ , the integrands must be equal.

$$-\rho g \frac{\sin \theta}{\cos \theta} = \rho u_{tt} \frac{1}{\cos \theta}$$

$$u_{tt} = -g \sin \theta$$

where

$$\sin \theta = \frac{u}{\sqrt{u^2 + x^2}} = \frac{u}{x} + \mathcal{O}(u^2) \approx \frac{u}{x}$$

when  $u$  is small. Therefore the chain satisfies PDE

$$xu_{tt} + gu = 0$$

### 1.5 #2

**a**

No. Suppose there are two different functions  $u_1$  and  $u_2$  that satisfy the equations. Now consider their difference  $v = u_1 - u_2$ .

$$\begin{cases} u_1''(x) + u_1'(x) = f(x) \\ u_2''(x) + u_2'(x) = f(x) \end{cases} \Rightarrow v''(x) + v'(x) = 0 \Rightarrow v(x) = Ae^{-x} + B$$

and the boundary conditions:

$$v'(0) = v(0) = \frac{1}{2}[v'(l) + v(l)]$$

$$\Rightarrow -A = A + B = \frac{1}{2}B$$

$$\Rightarrow B = -2A, \quad v(x) = Ae^{-x} - 2A$$

Thus if we can find any function  $u_0$  that satisfies the problem, then the family of functions

$$u = u_0 + Ce^{-t} - 2C$$

are all valid solutions.

**b**

Rearranging the second equation and applying the Fundamental Theorem of Calculus:

$$\begin{aligned} u'(0) + u(0) &= u'(l) + u(l) \Rightarrow \left[ u'(x) + u(x) \right]_0^l = 0 \\ \Rightarrow \int_0^l u''(x) + u'(x) dx &= \int_0^l f(x) dx = 0 \end{aligned}$$

Therefore in order for the problem to have a solution, the integral of  $f$  from 0 to  $l$  must be zero.