

Trigonometric functions

sin(α + β) = sin α cos β + cos α sin β
cos(α + β) = cos α cos β − sin α sin β
tan(α + β) = $\frac{\tan \alpha + \tan \beta}{1 - \tan \alpha \tan \beta}$
sin(2α) = 2 sin α cos α; tan(2α) = $\frac{2 \tan \alpha}{1 - \tan^2 \alpha}$
cos(2α) = cos^2 α − sin^2 α =
= 2 cos^2 α − 1 = 1 − 2 sin^2 α
sin α + sin β = 2 sin $\frac{\alpha + \beta}{2}$ cos $\frac{\alpha - \beta}{2}$

Hyperbolic functions

sinh(x + y) = sinh x cosh y + cosh x sinh y
cosh(x + y) = cosh x cosh y + sinh x sinh y
tanh(x + y) = $\frac{\tanh x + \tanh y}{1 + \tanh x \tanh y}$

Areas

triangle: $\sqrt{p(p-a)(p-b)(p-c)}$

Combinatorics

$D_{n,k} = \frac{n!}{(n-k)!}$

$P_n^{(m_1,m_2,\dots)} = \frac{n!}{m_1!m_2!...}$

$C_{n,k} = \binom{n}{k} = \frac{n!}{k!(n-k)!}$

$C'_{n,k} = \binom{n+k-1}{k}$

Miscellaneous

$A.B\overline{C} = \frac{ABC-AB}{9 \times C} \frac{0 \times B}{0}$
 $\sqrt{a \pm \sqrt{b}} = \sqrt{\frac{a + \sqrt{a^2 - b}}{2}} \pm \sqrt{\frac{a - \sqrt{a^2 - b}}{2}}$
 $\sum_{k=0}^n a^k = \frac{1-a^{n+1}}{1-a}$
 $\sum_{k=0}^n k a^k = \frac{a}{(1-a)^2} (1 + na^{n+1} - (n+1)a^n)$
 $\sum_{x=1}^n x^3 = \left(\sum_{x=1}^n x\right)^2 = \frac{1}{4}n^2(n+1)^2$
 $\sum_{x=1}^n x^2 = \frac{1}{6}n(n+1)(2n+1)$
 $\frac{d}{dx} \int_0^x g(x,y)dy = \int_0^x \frac{\partial g}{\partial x}(x,y)dy + g(x,x)$
 $\sqrt{z} = \sqrt{\frac{|z| + \text{Re } z}{2}} \pm i \sqrt{\frac{|z| - \text{Re } z}{2}}$

Derivatives

$(a^x)' = a^x \ln a$
 $\tan' x = 1 + \tan^2 x$
 $\cot' x = -1 - \cot^2 x$
 $\operatorname{atan}' x = -\operatorname{acot}' x = \frac{1}{1+x^2}$
 $\operatorname{asin}' x = -\operatorname{acos}' x = \frac{1}{\sqrt{1-x^2}}$
 $\log'_a x = \frac{1}{x \ln a}$
 $\cosh' x = \sinh x$
 $\tanh' x = 1 - \tanh^2 x$
 $\operatorname{atanh}' x = \operatorname{acoth}' x = \frac{1}{1-x^2}$

Integrals

$\int x^a = \frac{x^{a+1}}{a+1}$
 $\int a^x = \frac{a^x}{\ln a}$
 $\int \frac{1}{x} = \ln |x|$
 $\int \tan x = -\ln |\cos x|$
 $\int \cot x = \ln |\sin x|$
 $\int \frac{1}{\sin x} = \ln \left| \tan \frac{x}{2} \right|$
 $\int \frac{1}{\cos x} = \ln \left| \tan \left(\frac{x}{2} + \frac{\pi}{4} \right) \right|$
 $\int \ln x = x(\ln x - 1)$
 $\int \tanh x = \ln \cosh x$
 $\int \coth x = \ln |\sinh x|$
 $\int \frac{1}{\sqrt{a^2 - x^2}} = \operatorname{asin} \frac{x}{a}$
 $\int \frac{1}{a^2 + x^2} = \frac{1}{a} \operatorname{atan} \frac{x}{a}$
 $\int xy = x \int y - \int (\dot{x} \int y)$

Differential equations

$\dot{x} + \dot{a}x = b : x = e^{-a} \left(\int b e^a + c_1 \right)$

Taylor

$\sin x = x - \frac{x^3}{3!} + \frac{x^5}{5!} - \frac{x^7}{7!} + \dots$
 $\cos x = 1 - \frac{x^2}{2} + \frac{x^4}{4!} - \frac{x^6}{6!} + \dots$
 $\tan x = x + \frac{x^3}{3} + \frac{2}{15}x^5 + \frac{17}{315}x^7 + O(x^9)$
 $\frac{1}{\sin x} = \frac{1}{x} + \frac{x}{6} + \frac{7}{360}x^3 + \frac{31}{15120}x^5 + O(x^7)$
 $\frac{1}{\cos x} = 1 + \frac{x^2}{2} + \frac{5}{24}x^4 + \frac{61}{720}x^6 + \frac{277}{8064}x^8 + O(x^{10})$
 $\frac{1}{\tan x} = \frac{1}{x} - \frac{x}{3} - \frac{x^3}{45} - \frac{2}{945}x^5 + O(x^7)$
 $\operatorname{asin} x = x + \frac{x^3}{6} + \frac{3}{40}x^5 + \frac{5}{112}x^7 + O(x^9)$

$\cos \alpha + \cos \beta = 2 \cos \frac{\alpha + \beta}{2} \cos \frac{\alpha - \beta}{2}$
 $\cos \alpha - \cos \beta = -2 \sin \frac{\alpha + \beta}{2} \sin \frac{\alpha - \beta}{2}$
 $2 \sin \alpha \sin \beta = \cos(\alpha - \beta) - \cos(\alpha + \beta)$
 $2 \sin \alpha \cos \beta = \sin(\alpha + \beta) + \sin(\alpha - \beta)$
 $2 \cos \alpha \cos \beta = \cos(\alpha + \beta) + \cos(\alpha - \beta)$
 $\sin \frac{\alpha}{2} = \pm \sqrt{\frac{1 - \cos \alpha}{2}} \quad \cos \frac{\alpha}{2} = \pm \sqrt{\frac{1 + \cos \alpha}{2}}$
 $\tan \frac{\alpha}{2} = \frac{\sin \alpha}{1 + \cos \alpha} = \frac{1 - \cos \alpha}{\sin \alpha} = \pm \sqrt{\frac{1 - \cos \alpha}{1 + \cos \alpha}}$

$\left(\frac{\sinh x}{\cosh x} \right) = \frac{1}{2} \left(\frac{e^x - e^{-x}}{e^x + e^{-x}} \right)$
 $\cosh^2 x - \sinh^2 x = 1$
 $\cosh^2 x = \frac{1}{1 - \tanh^2 x}$

$\sin^2 \alpha = \frac{\tan^2 \alpha}{1 + \tan^2 \alpha} \quad \sin \alpha = \frac{2 \tan \frac{\alpha}{2}}{1 + \tan^2 \frac{\alpha}{2}}$
 $\cos^2 \alpha = \frac{1}{1 + \tan^2 \alpha} \quad \cos \alpha = \frac{1 - \tan^2 \frac{\alpha}{2}}{1 + \tan^2 \frac{\alpha}{2}}$
 $a \sin x + b \cos x =$
 $= \operatorname{sgn}(a) \sqrt{a^2 + b^2} \sin \left(x + \operatorname{atan} \frac{b}{a} \right)$
 $= \operatorname{sgn}(b) \sqrt{a^2 + b^2} \cos \left(x - \operatorname{atan} \frac{a}{b} \right)$
 $\operatorname{acos} x + \operatorname{asin} x = \frac{\pi}{2}$

$\sinh(ix) = i \sin x; \cosh(ix) = \cos x$
 $\left(\frac{\operatorname{asinh} x}{\operatorname{acosh} x} \right) = \log \left(x + \sqrt{x^2 + \left(\frac{1}{-1} \right)} \right)$
 $\operatorname{atanh} x = \frac{1}{2} \log \frac{1+x}{1-x}$

quad: $\sqrt{(p-a)(p-b)(p-c)(p-d) - abcd \cos^2 \frac{\alpha + \gamma}{2}}$

Pick: $A = \left(I + \frac{B}{2} - 1 \right) A_{\text{check}}$

$\langle \operatorname{Re}(ae^{-i\omega t}) \operatorname{Re}(be^{-i\omega t}) \rangle = \frac{1}{2} \operatorname{Re}(ab^*)$

$\operatorname{erf}(x) = \frac{2}{\sqrt{\pi}} \int_0^x e^{-t^2} dt$

$f^{(n)}(z_0) = \frac{n!}{2\pi i} \oint \frac{f(z)dz}{(z-z_0)^{n+1}}$

$f(z) = \sum_{k=-\infty}^{\infty} \left(\frac{1}{2\pi i} \oint \frac{f(z')dz'}{(z'-z_0)^{k+1}} \right) (z-z_0)^k$

$\operatorname{sinc} x = \frac{\sin x}{x}$

$\zeta(s) = \sum_{n=1}^{\infty} \frac{1}{n^s}; \operatorname{Li}_s(z) = \sum_{n=1}^{\infty} \frac{z^n}{n^s}$

$\operatorname{asinh}' x = \frac{1}{\sqrt{x^2+1}}$

$\operatorname{acosh}' x = \frac{1}{\sqrt{x^2-1}}$

$(f^{-1})' = \frac{1}{f'(f^{-1})}$

$\left(\frac{1}{x} \right)' = -\frac{\dot{x}}{x^2}$

$\left(\frac{x}{y} \right)' = \frac{\dot{x}y - x\dot{y}}{y^2}$

$(x^y)' = x^y (\dot{y} \ln x + y \frac{\dot{x}}{x})$

$\frac{\partial(x,y)}{\partial(u,v)} := \frac{\partial x}{\partial u} \frac{\partial y}{\partial v} - \frac{\partial x}{\partial v} \frac{\partial y}{\partial u}$

$\frac{\partial(x,y)}{\partial(u,y)} = \frac{\partial x}{\partial u} \Big|_y = -\frac{\partial x}{\partial y} \Big|_u \frac{\partial y}{\partial u} \Big|_x$

$\frac{\partial x}{\partial y} \Big|_u \frac{\partial y}{\partial u} \Big|_x \frac{\partial u}{\partial x} \Big|_y = -1$

$\frac{\partial x}{\partial u} \Big|_y = \frac{\partial x}{\partial u} \Big|_v - \frac{\partial x}{\partial y} \Big|_u \frac{\partial y}{\partial u} \Big|_v$

$\frac{\partial x}{\partial u} \Big|_v = \frac{\partial x}{\partial y} \Big|_v \frac{\partial y}{\partial u} \Big|_v$

$\int_{-\infty}^{\infty} e^{-x^2} = \sqrt{\pi}$

$\int dx e^{-\frac{1}{2} x^\top A x + b^\top x} =$

$= \frac{1}{\sqrt{\det \frac{A}{2\pi}}} e^{\frac{1}{2} b^\top A^{-1} b}$

$\int_{-\infty}^{\infty} e^{i\omega t} dt = 2\pi \delta(\omega)$

$\int \frac{du}{(1+u^2)^{3/2}} = \frac{u}{\sqrt{1+u^2}}$

$\int d\Omega_d = \frac{d\pi^{d/2}}{\Gamma(\frac{d}{2}+1)}$

$\int_0^\infty \frac{x^3}{e^x - 1} = \frac{\pi^4}{15}$

$\int_0^\infty \frac{x^{n-1}}{e^x - 1} = \Gamma(n) \zeta(n)$

$x\ddot{x} = k\dot{x}^2 : x = c_2^{-1-k} \sqrt{(1-k)t + c_1}$

$\dot{x} + ax^2 = b : x = \sqrt{\frac{b}{a}} \tanh \left(\sqrt{ab}(c_1 + t) \right)$

$\operatorname{atan} x = x - \frac{x^3}{3} + \frac{x^5}{5} - \frac{x^7}{7} + \dots$

$\frac{1}{1-x} = 1 + x + x^2 + x^3 + \dots$

$e^x = 1 + x + \frac{x^2}{2} + \frac{x^3}{3!} + \dots$

$\ln(1+x) = x - \frac{x^2}{2} + \frac{x^3}{3} - \frac{x^4}{4} + \dots$

$\sinh x = x + \frac{x^3}{3!} + \frac{x^5}{5!} + \frac{x^7}{7!} + \dots$

$\cosh x = 1 + \frac{x^2}{2} + \frac{x^4}{4!} + \frac{x^6}{6!} + \dots$

$\tanh x = x - \frac{x^3}{3} + \frac{2}{15}x^5 - \frac{17}{315}x^7 + O(x^9)$

$\frac{1}{\sinh x} = \frac{1}{x} - \frac{x}{6} + \frac{7}{360}x^3 - \frac{31}{15120}x^5 + O(x^7)$

$\frac{1}{\cosh x} = 1 - \frac{x^2}{2} + \frac{5}{24}x^4 - \frac{61}{720}x^6 + \frac{277}{8064}x^8 + O(x^{10})$

$\frac{1}{\tanh x} = \frac{1}{x} + \frac{x}{3} - \frac{x^3}{45} + \frac{2}{945}x^5 + O(x^7)$

$\operatorname{asinh} x = x - \frac{x^3}{6} + \frac{3}{40}x^5 - \frac{5}{112}x^7 + O(x^9)$

$(1+x)^n = 1 + nx + \frac{n(n-1)}{2}x^2 + O(x^3)$

$(1+x)^x = 1 + x^2 - \frac{x^3}{2} + \frac{5}{6}x^4 - \frac{3}{4}x^5 + O(x^6)$

$x! = 1 - \gamma x + \left(\frac{\gamma^2}{2} + \frac{\pi^2}{12} \right) x^2 + O(x^3)$

$J_0(x) = 1 - \frac{x^2}{4} + \frac{x^4}{64} - \frac{x^6}{2304} + \frac{x^8}{147456} + O(x^{10})$

$J_1(x) = \frac{x}{2} - \frac{x^3}{16} + \frac{x^5}{384} - \frac{x^7}{18432} + O(x^9)$

Fourier Fourier: $c_n = \frac{2}{T} \int_0^T f(t) \cos(n \frac{t}{T}) dt$ $\mathcal{F}[f](\omega) = \hat{f}(\omega) = \int dt e^{i\omega t} f(t)$ $f, g \in L^2 : (\hat{f}, \hat{g}) = 2\pi(f, g)$ $\mathcal{F}[\frac{\sin t}{t}] = \sqrt{\frac{\pi}{2}} \chi_{[-1;1]}(\omega)$ $t^{k \leq n} f(t) \in L^1 : \mathcal{F}[t^n f(t)] = (-i)^n \hat{f}^{(n)}$	$f^{(k \leq n)} \in L^1 : \mathcal{F}[f^{(n)}] = (-i\omega)^n \hat{f}$ $\mathcal{F}^2 f = 2\pi f(-t); (\omega \hat{f})' = -\mathcal{F}[t f']$ $f \star g = g \star f; \hat{f} \star \hat{g} = 2\pi \mathcal{F}[f g]$ $f \in L^1, g \in L^p : \mathcal{F}[f \star g] = \hat{f} \hat{g}$ $f \star g(x) = \int f(x-y) g(y) dy$ $(f \star g)' = f' \star g = f \star g'$	$f(x + \Delta) \star g = f \star g(x + \Delta)$ $f \in L^1, g \in L^p \Rightarrow f \star g \in L^p$ $f, g \in L^2 : f \star g = \frac{1}{2\pi} \int \hat{f} \hat{g} e^{-i\omega t} d\omega$ $\ f\ = 1 : \Delta\omega \Delta t \geq \frac{1}{2}$ $\Delta\omega \Delta t = \frac{1}{2} : f(t) = g(t; \bar{t}, \Delta t) e^{-i\bar{\omega} t}$																																																
Distributions $\mathcal{D} := \{f \in C^\infty \mid \exists K \text{ compact} : f(\mathcal{C} K) = 0\}$ $\mathcal{S} := \{f \in C^\infty \mid x^n f^{(k)} \leq A_{nk}\} \supset \mathcal{D}$ $\langle 1, f \rangle := \int f; \langle gT, f \rangle := \langle T, gf \rangle$ $T \in \mathcal{S}' : \langle \mathcal{F}T, f \rangle := \langle T, \mathcal{F}f \rangle$ $\langle T', f \rangle := -\langle T, f' \rangle; \langle \delta, f \rangle := f(0)$	$\langle T \otimes S, \phi \rangle := \langle T(x), \langle S(y), \phi(x+y) \rangle \rangle$ $\langle T \star S, \phi \rangle := \langle T \otimes S, \phi(x+y) \rangle$ $\mathcal{F}1 = 2\pi\delta(\omega); \mathcal{F} \operatorname{sgn} = 2i\mathcal{P} \frac{1}{\omega}$ $\mathcal{F}\theta = i\mathcal{P} \frac{1}{\omega} + \pi\delta(\omega)$ $\lim_{\varepsilon \rightarrow 0} \frac{1}{x+i\varepsilon} = \mathcal{P} \frac{1}{x} - i\pi\delta(x)$ $x^n T = 0 \Rightarrow T = \sum_{k=0}^{n-1} a_k \delta^{(k)}$	$xT = S \Rightarrow T = S/x + k\delta$ $T, S \in \mathcal{D}' : T \otimes S = S \otimes T$ $\sum_{n=0}^\infty e^{inx} = \mathcal{P} \frac{1}{1-e^{ix}} + \pi \sum_{n=-\infty}^\infty \delta(x-2n\pi)$ $\delta^{(n)} \star f = f^{(n)}$ $\delta(g(x)) = \sum_{x_i \in g^{-1}(0)} \frac{\delta(x-x_i)}{ g'(x_i) }$																																																
Bessel functions sol. of $x^2 f'' + x f' + (x^2 - \alpha^2) f = 0$ $\alpha = \text{“order”}$ $J_\alpha = \text{“first kind, normal”}$ $\alpha \in \mathbb{Z}_0 \vee \alpha > 0 : J_\alpha(0) = 0$ $J_0(0) = 1; \text{ otherwise } J_\alpha(0) = \infty$	$J_0(z) = \frac{1}{2\pi} \int_0^{2\pi} d\theta e^{z \cos \theta}; J_1(z) = J_0'(z)$ $\alpha \notin \mathbb{Z} : J_\alpha, J_{-\alpha} \text{ indep.}$ $\alpha \in \mathbb{Z} : J_{-\alpha} = (-1)^\alpha J_\alpha$ $Y_\alpha = \text{“second kind, normal” (also } N_\alpha)$ $\alpha \notin \mathbb{Z} : Y_\alpha = \frac{\cos(\alpha\pi) J_\alpha - J_{-\alpha}}{\sin(\alpha\pi)}$ $\alpha \in \mathbb{Z} : Y_\alpha = \lim_{\alpha' \rightarrow \alpha} Y_{\alpha'}$	$\alpha \in \mathbb{Z} : Y_\alpha, J_\alpha \text{ indep.}$ $\alpha \in \mathbb{Z} : Y_{-\alpha} = (-1)^\alpha Y_\alpha$ $\frac{2\alpha}{x} J_\alpha(x) = J_{\alpha-1}(x) + J_{\alpha+1}(x)$ $2J_\alpha'(x) = J_{\alpha-1}(x) - J_{\alpha+1}(x)$ $\int_0^1 x J_\alpha(x u_{\alpha,m}) J_\alpha(x u_{\alpha,n}) = \frac{\delta_{mn}}{2} J_{\alpha+1}^2(u_{\alpha,m})$ $u_{\alpha,n} = n\text{th. zero of } J_\alpha$ $Z_k(z) = \text{comb. of } e^{\pm i k z}$																																																
Cylindrical harmonics $V(\rho, \phi, z) = \sum_{n=0}^\infty \int dk A_{nk} P_{nk}(\rho) \Phi_n(\phi) Z_k(z)$	$P_{nk}(\rho) = \text{comb. of } J_n(k\rho), Y_n(k\rho)$ $\Phi_n(\phi) = \text{comb. of } e^{\pm i n \phi}$																																																	
Spherical harmonics $\frac{1}{ \vec{r}-\vec{r}' } = \sum_{l=0}^\infty \frac{\min(r,r')^l}{\max(r,r')^{l+1}} P_l(\frac{\vec{r}\vec{r}'}{rr'}); \sum_m Y_{lm} ^2 = \frac{2l+1}{4\pi}$ $P_l(x) = \frac{1}{2^l l!} \frac{d^l}{dx^l} (x^2-1)^l; f = \sum_{l=0}^\infty c_l P_l : c_l = \frac{2l+1}{2} \int_{-1}^1 f P_l$ $P_l(1) = 1; (P_n, P_m) = \frac{2\delta_{nm}}{2n+1}; (Y_{lm}, Y_{l'm'}) = \delta_{ll'} \delta_{mm'}$ $P_0 = 1; P_1 = x; P_2 = \frac{3x^2-1}{2}; Y_{00} = \frac{1}{\sqrt{4\pi}}; Y_{10} = \sqrt{\frac{3}{4\pi}} \cos \theta$	$Y_{11} = -\sqrt{\frac{3}{8\pi}} \sin \theta e^{i\varphi}; Y_{20} = \sqrt{\frac{5}{16\pi}} (3 \cos^2 \theta - 1)$ $Y_{21} = -\sqrt{\frac{15}{8\pi}} \sin \theta \cos \theta e^{i\varphi}; Y_{22} = \sqrt{\frac{15}{32\pi}} \sin^2 \theta e^{2i\varphi}$ $P_{lm}(x) = \frac{(-1)^m}{2^l l!} (1-x^2)^{\frac{m}{2}} \frac{d^{l+m}}{dx^{l+m}} (x^2-1)^l, 0 \leq m \leq l$ $Y_{lm}(\theta, \varphi) = \sqrt{\frac{2l+1}{4\pi} \frac{(l-m)!}{(l+m)!}} e^{im\varphi} P_{lm}(\cos \theta); Y_{l,-m} = (-1)^m Y_{lm}^*$ $P_l(\frac{\vec{r}\vec{r}'}{rr'}) = \frac{4\pi}{2l+1} \sum_{m=-l}^l Y_{lm}^*(\theta', \varphi') Y_{lm}(\theta, \varphi)$																																																	
Inequalities $ a - b \leq a+b \leq a + b $ $x > -1 : 1 + nx \leq (1+x)^n$	$\frac{ a^n-b^n }{ a-b <1} \leq n(1+ b)^{n-1}$ $\sqrt[p]{\sum (a_i+b_i)^p} \leq \sqrt[p]{\sum a_i^p} + \sqrt[p]{\sum b_i^p}$ $\sum a_i b_i \leq (\sum a_i^p)^{\frac{1}{p}} (\sum b_i^{\frac{p}{p-1}})^{\frac{p-1}{p}}$	$x^p y^q \leq (\frac{px+qy}{p+q})^{p+q}$ $\sqrt[p]{\frac{1}{n} \sum a_i^{p \leq q}} \leq \sqrt[q]{\frac{1}{n} \sum a_i^q}$ $\sum (\frac{a_1+\dots+a_i}{i})^p \leq (\frac{p}{p-1})^p \sum a_i^p$ $x \geq 0, \ddot{x} \leq M : \dot{x} \leq \sqrt{2Mx}$ $\frac{1}{1+x} < \ln(1+\frac{1}{x}) < \frac{1}{x}$																																																
Linear algebra $\dim(U+V) = \dim U + \dim V - \dim(U \cap V)$	$\dim V = \dim \ell(V) + \dim(V \cap \ker \ell)$																																																	
Symbols <table><tr><td>A</td><td>B</td><td>Γ</td><td>Δ</td><td>E</td><td>Z</td><td>H</td><td>Θ</td><td>I</td><td>K</td><td>Λ</td><td>M</td></tr><tr><td>α</td><td>β</td><td>γ</td><td>δ</td><td>ϵ/ε</td><td>ζ</td><td>η</td><td>θ/ϑ</td><td>ι</td><td>κ</td><td>λ</td><td>μ</td></tr></table>	A	B	Γ	Δ	E	Z	H	Θ	I	K	Λ	M	α	β	γ	δ	ϵ/ε	ζ	η	θ/ϑ	ι	κ	λ	μ	<table><tr><td>N</td><td>Ξ</td><td>O</td><td>Π</td><td>P</td><td>Σ</td><td>T</td><td>Υ</td><td>Φ</td><td>X</td><td>Ψ</td><td>Ω</td></tr><tr><td>ν</td><td>ξ</td><td>o</td><td>π/ϖ</td><td>ρ/ϱ</td><td>σ/ς</td><td>τ</td><td>v</td><td>ϕ/φ</td><td>χ</td><td>ψ</td><td>ω</td></tr></table>	N	Ξ	O	Π	P	Σ	T	Υ	Φ	X	Ψ	Ω	ν	ξ	o	π/ϖ	ρ/ϱ	σ/ς	τ	v	ϕ/φ	χ	ψ	ω	
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ν	ξ	o	π/ϖ	ρ/ϱ	σ/ς	τ	v	ϕ/φ	χ	ψ	ω																																							
Constants, units $\pi = 3.142$ $e = 2.718$ $\gamma = 5.772 \cdot 10^{-1}$ $G = 6.674 \cdot 10^{-11} \frac{\text{m}^3}{\text{kg s}^2}$ $R = 8.314 \frac{\text{J}}{\text{mol K}}$ $R = 8.206 \cdot 10^{-2} \frac{1 \text{atm}}{\text{mol K}}$	$N_A = 6.022 \cdot 10^{23} \frac{1}{\text{mol}}$ $k_B = 1.381 \cdot 10^{-23} \frac{\text{J}}{\text{K}}$ $k_B = 8.617 \cdot 10^{-5} \frac{\text{eV}}{\text{K}}$ $c = 2.998 \cdot 10^8 \frac{\text{m}}{\text{s}}$ $q_e = 1.602 \cdot 10^{-19} \text{A s}$ $m_e = 9.109 \cdot 10^{-31} \text{kg}$ $m_p = 1.673 \cdot 10^{-27} \text{kg}$	$m_n = 1.675 \cdot 10^{-27} \text{kg}$ $m_e = 5.110 \cdot 10^{-1} \text{MeV}$ $m_p = 9.383 \cdot 10^2 \text{MeV}$ $m_n = 9.396 \cdot 10^2 \text{MeV}$ $m_n - m_p = 1.293 \text{MeV}$ $\text{amu} = 1.661 \cdot 10^{-27} \text{kg}$ $h = 6.626 \cdot 10^{-34} \text{J s}$	$h = 4.136 \cdot 10^{-15} \text{eV s}$ $\varepsilon_0 = 8.854 \cdot 10^{-12} \frac{\text{C}^2}{\text{N m}^2}$ $\frac{1}{4\pi\varepsilon_0} = 8.988 \cdot 10^9 \frac{\text{N m}^2}{\text{C}^2}$ $\mu_0 = 1.257 \cdot 10^{-6} \frac{\text{N}}{\text{A}^2}$ $\mu_B = 9.274 \cdot 10^{-24} \text{A m}^2$ $\mu_B = 5.788 \cdot 10^{-5} \frac{\text{eV}}{\text{T}}$	$\alpha = 7.297 \cdot 10^{-3}$ $\text{barn} = 1 \cdot 10^{-28} \text{m}^2$ $\text{cd}_{555 \text{nm}} = 1.464 \cdot 10^{-3} \frac{\text{W}}{\text{sr}}$ $r_B = 5.292 \cdot 10^{-11} \text{m}$ $\text{Rydberg} = 1.361 \cdot 10^1 \text{eV}$ $r_e = 2.818 \cdot 10^{-15} \text{m}$ $\text{Debye} = 3.336 \cdot 10^{-30} \text{C m}$																																														
Vectors $\varepsilon_{ijk} = \begin{cases} 0 & i = j \vee j = k \vee k = i \\ 1 & i+1 \equiv j \wedge j+1 \equiv k \\ -1 & i \equiv j+1 \wedge j \equiv k+1 \end{cases}$ $\varepsilon_{ijk} \varepsilon_{ilm} = \delta_{jl} \delta_{km} - \delta_{jm} \delta_{kl}$ $\vec{a} \times \vec{b} = \varepsilon_{ijk} a_j b_k \hat{e}_i; (\vec{a} \otimes \vec{b})_{ij} = a_i b_j$	$(\vec{a} \times \vec{b}) \vec{c} = (\vec{c} \times \vec{a}) \vec{b}$ $(\vec{a} \times \vec{b}) \times \vec{c} = -(\vec{b} \vec{c}) \vec{a} + (\vec{a} \vec{c}) \vec{b}$ $(\vec{a} \times \vec{b})(\vec{c} \times \vec{d}) = (\vec{a} \vec{c})(\vec{b} \vec{d}) - (\vec{a} \vec{d})(\vec{b} \vec{c})$ $ \vec{u} \times \vec{v} ^2 = u^2 v^2 - (\vec{u} \vec{v})^2$ $\vec{\nabla} = (\frac{\partial}{\partial x}, \frac{\partial}{\partial y}, \frac{\partial}{\partial z}); \square = \frac{\partial^2}{\partial t^2} - \nabla^2$	$\vec{\nabla} V = \frac{\partial V}{\partial \rho} \hat{\rho} + \frac{1}{\rho} \frac{\partial V}{\partial \phi} \hat{\phi} + \frac{\partial V}{\partial z} \hat{z}$ $\vec{\nabla} \vec{v} = \frac{1}{\rho} \frac{\partial(\rho v_\rho)}{\partial \rho} + \frac{1}{\rho} \frac{\partial v_\phi}{\partial \phi} + \frac{\partial v_z}{\partial z}$ $\vec{\nabla} \times \vec{v} = (\frac{1}{\rho} \frac{\partial v_z}{\partial \phi} - \frac{\partial v_\phi}{\partial z}) \hat{\rho} +$ $+(\frac{\partial v_\rho}{\partial z} - \frac{\partial v_z}{\partial \rho}) \hat{\phi} + \frac{1}{\rho} (\frac{\partial(\rho v_\phi)}{\partial \rho} - \frac{\partial v_\rho}{\partial \phi}) \hat{z}$ $\nabla^2 V = \frac{1}{\rho} \frac{\partial}{\partial \rho} (\rho \frac{\partial V}{\partial \rho}) + \frac{1}{\rho^2} \frac{\partial^2 V}{\partial \phi^2} + \frac{\partial^2 V}{\partial z^2}$																																																

$$\begin{aligned} \vec{\nabla} V &= \frac{\partial V}{\partial r} \hat{r} + \frac{1}{r} \frac{\partial V}{\partial \theta} \hat{\theta} + \frac{1}{r \sin \theta} \frac{\partial V}{\partial \varphi} \hat{\varphi} & \frac{1}{r^2} \frac{\partial}{\partial r} (r^2 \frac{\partial V}{\partial r}) &= \frac{1}{r} \frac{\partial^2}{\partial r^2} (rV) = \frac{2}{r} \frac{\partial V}{\partial r} + \frac{\partial^2 V}{\partial r^2} & \frac{1}{2} \vec{\nabla} v^2 &= (\vec{v} \cdot \vec{\nabla}) \vec{v} + \vec{v} \times (\vec{\nabla} \times \vec{v}) \\ \vec{\nabla} \vec{v} &= \frac{1}{r^2} \frac{\partial (r^2 v_r)}{\partial r} + \frac{1}{r \sin \theta} \frac{\partial (v_\theta \sin \theta)}{\partial \theta} + \frac{1}{r \sin \theta} \frac{\partial v_\varphi}{\partial \varphi} & \vec{\nabla} (\vec{\nabla} \times \vec{v}) &= \vec{\nabla} \times \vec{\nabla} V = 0 & \int \vec{\nabla} \vec{v} d^3 x &= \oint \vec{v} d\vec{S}; \int (\vec{\nabla} \times \vec{v}) d\vec{S} = \oint \vec{v} d\vec{l} \\ \vec{\nabla} \times \vec{v} &= \frac{1}{r \sin \theta} \left(\frac{\partial (v_\varphi \sin \theta)}{\partial \theta} - \frac{\partial v_\theta}{\partial \varphi} \right) \hat{r} + & \vec{\nabla} (f \vec{v}) &= (\vec{\nabla} f) \vec{v} + f \vec{\nabla} \vec{v} & \int (f \nabla^2 g - g \nabla^2 f) d^3 x &= \oint_S (f \frac{\partial g}{\partial n} - g \frac{\partial f}{\partial n}) dS \\ &+ \frac{1}{r} \left(\frac{1}{\sin \theta} \frac{\partial v_r}{\partial \varphi} - \frac{\partial (rv_\varphi)}{\partial r} \right) \hat{\theta} + \frac{1}{r} \left(\frac{\partial (rv_\theta)}{\partial r} - \frac{\partial v_r}{\partial \theta} \right) \hat{\varphi} & \vec{\nabla} \times (f \vec{v}) &= \vec{\nabla} f \times \vec{v} + f \vec{\nabla} \times \vec{v} & \oint \vec{v} \times d\vec{S} &= - \int (\vec{\nabla} \times \vec{v}) d^3 x \\ \nabla^2 V &= \frac{\partial}{\partial r} (r^2 \frac{\partial V}{\partial r}) + \frac{\partial}{\partial r^2} (\sin \theta \frac{\partial V}{\partial \theta}) + \frac{\partial^2 V}{r^2 \sin^2 \theta} & \vec{\nabla} \times (\vec{\nabla} \times \vec{v}) &= -\nabla^2 \vec{v} + \vec{\nabla} (\vec{\nabla} \cdot \vec{v}) & \delta(\vec{r} - \vec{r}_0) &= \frac{\delta(r-r_0) \delta(\theta-\theta_0) \delta(\varphi-\varphi_0)}{r_0^2 \sin \theta_0} \\ & & \vec{\nabla} (\vec{v} \times \vec{w}) &= \vec{w} (\vec{\nabla} \times \vec{v}) - \vec{v} (\vec{\nabla} \times \vec{w}) & \nabla^2 \frac{1}{|\vec{r}-\vec{r}_0|} &= -4\pi \delta(\vec{r} - \vec{r}_0) \\ & & \vec{\nabla} \times (\vec{v} \times \vec{w}) &= (\vec{\nabla} \cdot \vec{w} + \vec{w} \cdot \vec{\nabla}) \vec{v} - (\vec{\nabla} \cdot \vec{v} + \vec{v} \cdot \vec{\nabla}) \vec{w} \end{aligned}$$

Statistics

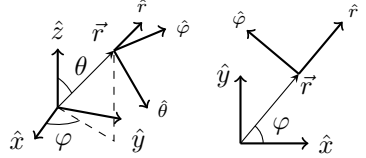
$$\begin{aligned} P(E \cap E_1) &= P(E_1) \cdot P(E|E_1) & g(\vec{x}; \vec{\mu}, V) &= \frac{e^{-\frac{1}{2}(\vec{x}-\vec{\mu})^\top V^{-1}(\vec{x}-\vec{\mu})}}{\sqrt{\det(2\pi V)}} & n \geq 35 : S(x; n) &\approx g(x; 0, 1) \\ \Delta x_{\text{hist}} &\approx \frac{x_{\text{max}} - x_{\text{min}}}{\sqrt{N}} & \text{FWHM}_g &= 2\sigma \sqrt{2 \ln 2} & c(x; a) &= \frac{a}{\pi} \frac{1}{a^2 + x^2} \\ P(x \leq k) &= F(k) = \int_{-\infty}^k p(x) & z &= \frac{x-\mu}{\sigma}; \mu, \sigma[z] = 0, 1 & \sigma_{xy} &= E[xy] - \mu_x \mu_y \leq \sigma_x \sigma_y \\ \text{median} &= F^{-1}(\frac{1}{2}) & \chi^2 &= \sum_{i=1}^n z_i^2; \wp := p[\chi^2] & \rho_{xy} &= \frac{\sigma_{xy}}{\sigma_x \sigma_y}, |\rho_{xy}| \leq 1 \\ E[f(x)] &= \int_{-\infty}^{\infty} f(x) p(x) & \wp(x; n) &= \frac{2^{-\frac{n}{2}}}{\Gamma(\frac{n}{2})} x^{\frac{n}{2}-1} e^{-\frac{x}{2}} & \mu_{f(x)} &\approx f(\mu_x) \\ \mu &= E[x] = \int_{-\infty}^{\infty} x p(x) & \mu_\wp = n, \sigma_\wp^2 &= 2n & \sigma_{fg} &\approx \sigma_{x_i x_j} \frac{\partial f}{\partial x_i} \Big|_{\mu_{x_i}} \frac{\partial g}{\partial x_j} \Big|_{\mu_{x_j}} \\ \alpha_n &= E[x^n] & n \geq 30 : \wp(x; n) &\approx g(x; n, \sqrt{2n}) & \mu &\approx m = \frac{1}{n} \sum_{i=1}^n x_i \\ M_n &= E[(x - \mu)^n] & n \geq 8 : p[\sqrt{2\chi^2}] &\approx g(\sqrt{2n-1}, 1) & s^2 &\approx s^2 = \frac{1}{n-1} \sum_{i=1}^n (x_i - m)^2 \\ \sigma^2 &= M_2 = E[x^2] - \mu^2 & S(x; n) &= \frac{\Gamma(\frac{n+1}{2})}{\sqrt{n\pi} \Gamma(\frac{n}{2})} (1 + \frac{x^2}{n})^{-\frac{n+1}{2}} & s_m^2 &= \frac{s^2}{n} \\ \text{FWHM} &\approx 2\sigma & \mu_S = 0, \sigma_S^2 &= \frac{n}{n-2} & p[\frac{m-\mu}{s_m}] &= S(n) \\ \gamma_1 &= \frac{M_3}{\sigma^3}, \gamma_2 = \frac{M_4}{\sigma^4} & p[z \sqrt{\frac{n}{\chi^2}}] &= S(n) & H_0 \text{ sign.lev.} &= \int_{\text{reject}} p(S|H_0) dS \\ \phi[y](t) &= E[e^{ity}] & & & H_0 \text{ pow.vs. } H_1 &= \int_{\text{reject}} p(S|H_1) dS \end{aligned}$$

Least squares

$$\begin{aligned} f(x) &= mx + q, \quad f(x) = a, & \sigma_{\hat{m}, \hat{q}}^2 &= \frac{\sum \frac{1}{\Delta y^2}}{\sum \frac{1}{\Delta y^2} \cdot \sum \frac{x^2}{\Delta y^2} - (\sum \frac{x}{\Delta y^2})^2} & \hat{b} &= \frac{\sum \frac{xy^2}{\Delta y^2}}{\sum \frac{x^2}{\Delta y^2}}, \sigma_b^2 = \frac{1}{\sum \frac{x^2}{\Delta y^2}} \\ f(x) &= bx, \quad f(x; \theta) = \theta_i h_i(x) & \hat{a} &= \frac{\sum \frac{y}{\Delta y^2}}{\sum \frac{1}{\Delta y^2}}, \sigma_a^2 = \frac{1}{\sum \frac{1}{\Delta y^2}} & H_{ij} &:= h_j(x_i) \\ \hat{m} &= \frac{\sum \frac{1}{\Delta y^2} \cdot \sum \frac{xy}{\Delta y^2} - \sum \frac{x}{\Delta y^2} \cdot \sum \frac{y}{\Delta y^2}}{\sum \frac{1}{\Delta y^2} \cdot \sum \frac{x^2}{\Delta y^2} - (\sum \frac{x}{\Delta y^2})^2} & \hat{\mathbf{a}} &= (\sum V_{\mathbf{y}}^{-1})^{-1} (\sum V_{\mathbf{y}}^{-1} \mathbf{y}) & \hat{\theta} &= (H^\top V_{\mathbf{y}}^{-1} H)^{-1} H^\top V_{\mathbf{y}}^{-1} \mathbf{y} \\ & & V_{\hat{\mathbf{a}}} &= (\sum V_{\mathbf{y}}^{-1})^{-1} & V_{\hat{\theta}} &= (H^\top V_{\mathbf{y}}^{-1} H)^{-1} \\ \sigma_{\hat{q}}^2 &= \frac{\sum \frac{x^2}{\Delta y^2}}{\sum \frac{1}{\Delta y^2} \cdot \sum \frac{x^2}{\Delta y^2} - (\sum \frac{x}{\Delta y^2})^2} & & & & \end{aligned}$$

Kinematics

$$\begin{aligned} \frac{1}{R} &= \left| \frac{v_x a_y - v_y a_x}{v^3} \right| & \theta \equiv \frac{\pi}{2} \rightarrow \ddot{\vec{r}} &= (\ddot{r} - r\dot{\varphi}^2) \hat{r} + (r\ddot{\varphi} + 2\dot{r}\dot{\varphi}) \hat{\varphi} & \vec{A} &= \ddot{\vec{r}} + \vec{A}_T + \vec{\omega} \times (\vec{\omega} \times \vec{r}) + \dot{\vec{\omega}} \times \vec{r} + 2\vec{\omega} \times \dot{\vec{r}} \\ \vec{\omega} &= \dot{\varphi} \cos \theta \hat{r} - \dot{\varphi} \sin \theta \hat{\theta} + \dot{\theta} \hat{\varphi} & \dot{\vec{r}} &= \dot{r} \hat{r} + r \dot{\theta} \hat{\theta} + r \dot{\varphi} \sin \theta \hat{\varphi} & & \\ \dot{\vec{w}} &= \frac{d(\vec{w}\hat{r})}{dt} \hat{r} + \frac{d(\vec{w}\hat{\theta})}{dt} \hat{\theta} + \frac{d(\vec{w}\hat{\varphi})}{dt} \hat{\varphi} + \vec{\omega} \times \vec{w} & \ddot{\vec{r}} \hat{r} &= \ddot{r} - r\dot{\theta}^2 - r\dot{\varphi}^2 \sin^2 \theta & & \\ \theta \equiv \frac{\pi}{2} \rightarrow \dot{\vec{r}} &= \dot{r} \hat{r} + r \dot{\varphi} \hat{\varphi} & \ddot{\vec{r}} \hat{\theta} &= r\ddot{\theta} + 2\dot{r}\dot{\theta} - r\dot{\varphi}^2 \sin \theta \cos \theta & & \\ & & \ddot{\vec{r}} \hat{\varphi} &= r\ddot{\varphi} \sin \theta + 2\dot{r}\dot{\varphi} \sin \theta + 2r\dot{\theta}\dot{\varphi} \cos \theta \end{aligned}$$



Mechanics

$$\begin{aligned} \dot{\alpha} &= \frac{d\alpha}{dt}(\beta, t) = \frac{\partial \alpha}{\partial \beta} \dot{\beta} + \frac{\partial \alpha}{\partial t} & \vec{L} &= \vec{R} \times M \dot{\vec{R}} + (\vec{r}_i - \vec{R}) \times m_i (\dot{\vec{r}}_i - \dot{\vec{R}}) & \frac{\partial}{\partial \epsilon} S[q + \epsilon] \Big|_{\epsilon=0}^{\epsilon(t_1)=\epsilon(t_2)=0} &= 0 & \{u, v\} &= \frac{\partial u}{\partial q} \frac{\partial v}{\partial p} - \frac{\partial u}{\partial p} \frac{\partial v}{\partial q} \\ \vec{p} &:= m \dot{\vec{r}}; \vec{F} = \dot{\vec{p}}; \frac{d(mT)}{dt} = \vec{F} \cdot \vec{p} & \vec{\tau}_O &= \dot{\vec{L}}_O + \vec{v}_O \times \vec{p} & p &:= \frac{\partial \mathcal{L}}{\partial \dot{q}}; \dot{p} = \frac{\partial \mathcal{L}}{\partial q} & \frac{du}{dt} &= \{u, \mathcal{H}\} + \frac{\partial u}{\partial t} \\ M &:= \sum_i m_i; \vec{R} := \frac{m_i \vec{r}_i}{M} & \tau_1 &= I_1 \dot{\omega}_1 + (I_3 - I_2) \omega_3 \omega_2 & \mathcal{H}(q, p, t) &= \dot{q} p - \mathcal{L} & \eta &= (q, p); \Gamma = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \\ T &= \frac{1}{2} M \dot{\vec{R}}^2 + \frac{1}{2} m_i (\dot{\vec{r}}_i - \dot{\vec{R}})^2 & \mathcal{L}(q, \dot{q}, t) &= T - V + \frac{d}{dt} f(q, t) & \dot{q} &= \frac{\partial \mathcal{H}}{\partial p}; \dot{p} = -\frac{\partial \mathcal{H}}{\partial q} & \dot{\eta} &= \Gamma \frac{\partial \mathcal{H}}{\partial \eta}; \{u, v\} = \frac{\partial u}{\partial \eta} \Gamma \frac{\partial v}{\partial \eta} \\ & & S[q] &= \int_{t_1}^{t_2} \mathcal{L}(q, \dot{q}, t) dt & \frac{d\mathcal{H}}{dt} &= \frac{\partial \mathcal{H}}{\partial t} = -\frac{\partial \mathcal{L}}{\partial t} \end{aligned}$$

Inertia

$$\begin{aligned} \text{point: } &mr^2 & \text{rod: } &\frac{1}{12} mL^2 & \text{octahedron: } &\frac{1}{10} ms^2 & \text{cone: } &\frac{3}{10} mr^2 & \text{rectangulus: } &\frac{1}{12} m(a^2 + b^2) \\ \text{two points: } &\mu d^2 & \text{disk: } &\frac{1}{2} mr^2 & \text{sphere: } &\frac{2}{3} mr^2 & \text{torus: } &m(R^2 + \frac{3}{4} r^2) \\ \text{tetrahedron: } &\frac{1}{20} ms^2 & \text{ball: } &\frac{2}{5} mr^2 & \text{ellipsoid: } &I_a = \frac{1}{5} m(b^2 + c^2) \end{aligned}$$

Kepler

$$\begin{aligned} \langle U \rangle &= -2 \langle T \rangle & \vec{L} &= \vec{R} \times M \dot{\vec{R}} + \vec{r} \times \mu \dot{\vec{r}} & r &= \frac{k}{1 + \varepsilon \cos \theta} & \vec{A} &= \mu \dot{\vec{r}} \times \vec{L} - \mu \alpha \hat{r}, \dot{\vec{A}} = 0 \\ U_{\text{eff}} &= U + \frac{L^2}{2mr^2} & k &= \frac{L^2}{\mu \alpha}, \varepsilon = \sqrt{1 + \frac{2EL^2}{\mu \alpha^2}} & a &= \frac{k}{|1 - \varepsilon^2|} = \frac{\alpha}{2|E|} \\ T &= \frac{1}{2} M \dot{\vec{R}}^2 + \frac{1}{2} \mu \dot{\vec{r}}^2 & a^3 \omega^2 &= G(m_1 + m_2) = \frac{\alpha}{\mu} \end{aligned}$$

Relativity

$$\begin{aligned} \beta &= \frac{v}{c} = \tanh \chi & V'_\perp &= \frac{1}{\gamma} \frac{V_\perp}{1 - \frac{v V_\parallel}{c^2}} & x^\mu &= (ct, \vec{x}) & \partial_\mu &= \frac{\partial}{\partial x^\mu} = \left(\frac{1}{c} \frac{\partial}{\partial t}, \vec{\nabla} \right) \\ \gamma &= \frac{1}{\sqrt{1 - \beta^2}} = \cosh \chi & \frac{V'}{c} &= 1 - \frac{(1 - \frac{V^2}{c^2})(1 - \frac{v^2}{c^2})}{(1 - \frac{v V_\parallel}{c^2})^2} & v^\mu &= \frac{dx^\mu}{d\tau} = \gamma(c, \vec{v}) & g_{\mu\nu} &= \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & -1 & 0 & 0 \\ 0 & 0 & -1 & 0 \\ 0 & 0 & 0 & -1 \end{pmatrix} \\ \vec{p} &= \gamma m \vec{v}; \mathcal{E} = \gamma mc^2 & \chi'' &= \chi' + \chi & a^\mu &= \frac{dv^\mu}{d\tau} = \gamma \left(\frac{d\gamma}{dt} c, \frac{d(\gamma \vec{v})}{dt} \right) & x_\mu &= g_{\mu\nu} x^\nu \\ \text{free particle: } \mathcal{L} &= \frac{mc^2}{\gamma} & V'_\parallel &= \frac{V_\parallel - v}{1 - \frac{v V_\parallel}{c^2}} & p^\mu &= mv^\mu = \left(\frac{\mathcal{E}}{c}, \vec{p} \right) & \partial_\mu \partial^\mu &= \square \\ & & d\tau &= \frac{1}{\gamma} dt & \frac{dp^\mu}{d\tau} &= \gamma \left(\frac{v}{c} \frac{d\vec{p}}{dt}, \frac{d\vec{p}}{dt} \right) \end{aligned}$$

$$\begin{array}{llll} p^\mu p_\mu = (mc)^2 & \text{doppler: } \sqrt{\frac{1+\beta}{1-\beta}} \approx 1 + \beta & (\Lambda^0_0)^2 \geq 1 & M \rightarrow \sum_i m_i \quad E_A^{\min} = \frac{(\sum_i m_i)^2 - m_A^2 - m_B^2}{2m_B} \\ v^\mu a_\mu = 0 & \text{SO}_{1,3} = \left\{ \Lambda \left| \begin{array}{c} \Lambda^\top g \Lambda = g \\ \det \Lambda \geq 0 \end{array} \right. \right\} & \Lambda = \begin{pmatrix} \gamma & -\gamma \vec{\beta} \\ -\gamma \vec{\beta} & I + \frac{\gamma-1}{\beta^2} \vec{\beta} \otimes \vec{\beta} \end{pmatrix} & E_1^{\max} = \frac{M^2 + m_1^2 - \sum_{i \neq 1} m_i^2}{2M} \quad m, M_{\text{still}} \text{ 1D coll.} \\ & & & A + B_{\text{still}} \rightarrow \sum_i m_i \quad E'_m = \frac{(M+m)^2 E_m + 2Mm^2}{M^2 + m^2 + 2ME_m} \end{array}$$

Thermodynamics

$$\begin{array}{llll} \text{d}Q = T\text{d}S = \text{d}U + \text{d}L = \text{d}U + p\text{d}V - \mu\text{d}N & \kappa, C \geq 0 & \Omega(T,V,\mu) := U - TS - \mu N = -pV & \\ C_V := T \frac{\partial S}{\partial T} \Big|_V = \frac{\partial U}{\partial T} \Big|_V & \text{Euler: } U = ST - pV + \mu N & & \begin{array}{c} V \\ \swarrow \quad \searrow \\ U \quad \quad T \\ \swarrow \quad \searrow \\ S \quad \quad p \end{array} \quad \frac{\partial}{\partial T} \frac{G}{T} \Big|_p = -\frac{H}{T^2} \\ C_p := T \frac{\partial S}{\partial T} \Big|_p = \frac{\partial U}{\partial T} \Big|_p + p \frac{\partial V}{\partial T} \Big|_p = \frac{\partial H}{\partial T} \Big|_p & \text{Gibbs-Duhem: } S\text{d}T - V\text{d}p + N\text{d}\mu = 0 & & \frac{\partial}{\partial T} \frac{F}{T} \Big|_V = -\frac{U}{T^2} \\ \mu_J := \frac{\partial T}{\partial V} \Big|_{U,N}; \; \alpha := \frac{1}{V} \frac{\partial V}{\partial T} \Big|_p & \text{Fix } S, V, N : \min U \text{ at equilibrium} & & \\ \kappa_T := -\frac{1}{V} \frac{\partial V}{\partial p} \Big|_T; \; \kappa_S := -\frac{1}{V} \frac{\partial V}{\partial p} \Big|_S & \text{Fix } T, V, N : \min F := U - TS & & \delta E \leq T\delta S - p\delta V + H\delta \int M\text{d}V + \mu\delta N \\ \gamma := \frac{C_p}{C_V} = \frac{\kappa_T}{\kappa_S}; \; C_p - C_V = \frac{\alpha^2 VT}{\kappa_T} & \text{Fix } T, p, N : \min G := F + pV = \mu N & & \text{d}U = T\text{d}S + H\text{d}\langle M \rangle_V; \; \chi := \frac{\partial M}{\partial H} \\ & \text{Fix } S, p, N : \min H := U + pV & & \text{Clausius-Clapeyron: } \frac{\text{d}p}{\text{d}T} = \frac{s_2 - s_1}{v_2 - v_1} \end{array}$$

TM examples

$$\begin{array}{ll} \text{IDEAL GAS: } pV = Nk_{\text{B}}T & c_V, c_p = \frac{C_V, C_p}{n}, \; c_V = \frac{\text{dof}}{2}R, \; c_p = c_V + R \quad \text{d}Q = 0 : pV^\gamma, TV^{\gamma-1}, p^{\frac{1}{\gamma}-1}T \text{ const.} \\ & c_V = \frac{R}{\gamma-1}, \; c_p = \frac{\gamma}{\gamma-1}R, \; \gamma = 1 + \frac{2}{\text{dof}} \quad \text{V.D.WAALS: } \left(p - \frac{N^2a^2}{V^2}\right)(V - Nb) = Nk_{\text{B}}T \end{array}$$

Statistical mechanics

$$\begin{array}{llll} \text{MICROCANONICAL} & \left\langle x_i \frac{\partial \mathcal{H}}{\partial x_j} \right\rangle = \delta_{ij} kT & N = z \frac{\partial}{\partial z} \log \mathcal{Z} \Big|_{\beta, V} & B_1 = 1, \; B_2 = -b_2, \; B_3 = 4b_2^2 - 2b_3 \\ \rho := \frac{1}{N!} \begin{cases} \text{const.} & E < \mathcal{H} < E + \Delta \\ 0 & \text{otherwise} \end{cases} & \text{CANONICAL} & U = -\frac{\partial}{\partial \beta} \log \mathcal{Z} \Big|_{z, V}; \; \Omega = -\frac{\log \mathcal{Z}}{\beta} & f_{ij} := e^{-\beta v(|\vec{r}_i - \vec{r}_j|)} - 1 \\ \Gamma(E) := \frac{1}{N!} \int_{E < \mathcal{H} < E + \Delta} \text{d}p\text{d}q & \rho = \frac{1}{h^{3N} N!} e^{-\beta \mathcal{H}}; \; \beta := \frac{1}{k_{\text{B}}T} & \lambda_T^2 := \frac{h^2 \beta}{2\pi m}; \; z \ll 1 \rightarrow z \approx \frac{N \lambda_T^3}{V} & \mathcal{U}_l := \sum_{\substack{\text{connected} \\ l \text{ vertices}}} \prod_{\substack{\text{connections} \\ \langle i, j \rangle}} f_{ij} \\ S := k_{\text{B}} \log \Gamma(E) & Z := \frac{1}{h^{3N} N!} \int \text{d}p\text{d}q e^{-\beta \mathcal{H}} & \text{VIRIAL EXPANSION} & b_l = \frac{1}{l!} \frac{1}{V} \int \text{d}\vec{r}_1 \cdots \text{d}\vec{r}_l \mathcal{U}_l \\ T := \frac{\partial E}{\partial S} \Big|_V; \; p := -\frac{\partial E}{\partial V} \Big|_S & F = -\frac{\log Z}{\beta}; \; U = -\frac{\partial \log Z}{\partial \beta} & \mathcal{H} = \sum_i \frac{p_i^2}{2m} + \sum_{i < j} v(|\vec{r}_i - \vec{r}_j|) & \mathcal{W}_l := \sum_{\substack{\text{conn. } l \text{ vert.} \\ \text{that remain conn.} \\ \text{removing a point}}} \prod_{\langle i, j \rangle} f_{ij} \\ S = -k_{\text{B}} \langle \log \rho \rangle & \text{GRAND CANONICAL} & \frac{pV}{Nk_{\text{B}}T} = \sum_{l=1}^\infty B_l(T) \left(\frac{N}{V}\right)^{l-1} & B_l = -\frac{1}{l(l-2)!} \frac{1}{V} \int \text{d}\vec{r}_1 \cdots \text{d}\vec{r}_l \mathcal{W}_l \\ \mathcal{Z} := \sum_{N=0}^\infty z^N Z; \; z := e^{\beta \mu} & \rho = \frac{1}{h^{3N} N!} e^{-\beta(\mathcal{H} - \mu N + pV)} & \Omega = -Vk_{\text{B}}T \sum_{l=1}^\infty \left(\frac{z}{\lambda_T^3}\right)^l b_l(T) \end{array}$$

Statistical QM

$$\begin{array}{llll} \text{FERMIONS} & \int_{-\infty}^\infty \text{d}\epsilon f(\epsilon - \mu) \phi(\epsilon) = \frac{\pi D}{\sin(\pi D)} \Phi(\mu) & \Omega(T, V, \mu) = \frac{1}{\beta} \sum_\alpha \log(1 - e^{-\beta(\epsilon_\alpha - \mu)}) & \\ \Omega(T, V, \mu) = -\frac{1}{\beta} \sum_\alpha \log(1 + e^{-\beta(\epsilon_\alpha - \mu)}) & \Phi(\mu) := \int_{-\infty}^\mu \text{d}\epsilon \phi(\epsilon); \; D := \frac{1}{\beta} \frac{\text{d}}{\text{d}\mu} & \epsilon \propto p^2 : n(T, z) = n_0 + g\lambda_T^{-d} \text{Li}_{\frac{d}{2}}(z) & \\ f(\epsilon) := \frac{1}{1 + e^{\beta \epsilon}} & \frac{pV}{NkT} = 1 + 2^{-5/2} \frac{N\lambda_T^3}{V} + O((N\lambda_T^3/V)^2) & \epsilon \propto p^2 : p(T, z) = gk_{\text{B}}T \lambda_T^{-d} \text{Li}_{\frac{d}{2}+1}(z) & \\ n(\varepsilon, T, \mu) = f(\varepsilon - \mu) & \epsilon = \frac{p^2}{2m} \rightarrow \epsilon_F = \frac{h^2}{2m} \left(\frac{N}{V}\right)^{2/3} \left(\frac{3}{4\pi g}\right)^{2/3} & \frac{pV}{NkT} = 1 - 2^{-5/2} \frac{N\lambda_T^3}{V} + O((N\lambda_T^3/V)^2) & \end{array}$$

BOSONS

Electronics (MKS)

$$\begin{array}{llll} \begin{pmatrix} V \\ I \end{pmatrix} = \begin{pmatrix} V_0 \\ I_0 \end{pmatrix} e^{i\omega t}, \; Z = \frac{V}{I} & Z_{\text{series}} = \sum_k Z_k, \; \frac{1}{Z_{\text{parallel}}} = \sum_k \frac{1}{Z_k} & I_{A \rightarrow C} = I_0 (e^{\frac{V_{\text{AC}}}{V_T}} - 1), \; V_T = \eta \frac{k_{\text{B}}T}{q_e} & \\ Z_R = R, \; Z_C = -i\frac{1}{\omega C}, \; Z_L = i\omega L & \sum_{\text{loop}} V_k = 0, \; \sum_{\text{node}} I_k = 0 & I_{E, \text{out}} = I_0^E (e^{\frac{V_{\text{BE}}}{V_T}} - 1) - \alpha_R I_0^C (e^{\frac{V_{\text{BC}}}{V_T}} - 1) & \\ & \mathcal{E} = -L\dot{I}, \; L = \frac{\Phi_B}{I} & I_{C, \text{in}} = -I_0^C (e^{\frac{V_{\text{BC}}}{V_T}} - 1) + \alpha_F I_0^E (e^{\frac{V_{\text{BE}}}{V_T}} - 1) & \end{array}$$

Chemistry

$$\begin{array}{llll} H = U + pV & \exists k, (m_i) : v_{\text{r}} = k[\text{A}_i]^{m_i} & K_\chi = \frac{\prod x_{\text{B}_j}^{b_j}}{\prod x_{\text{A}_i}^{a_i}}, \; \chi = \frac{n}{n_{\text{tot}}} & \Delta G = RT \ln \frac{Q}{K} \\ \text{d}p = 0 \rightarrow \Delta H = \text{heat transfer} & k = Ae^{-\frac{E_a}{RT}} \text{ (Arrhenius)} & K_c = K_p(RT)^{\sum a_i - \sum b_j} & \ln \frac{K_2}{K_1} = -\frac{\Delta H^\circ}{R} \left(\frac{1}{T_2} - \frac{1}{T_1}\right) \\ G = H - TS & a_{(\ell)} = \gamma \frac{[\text{X}]}{[\text{X}]_0}, \; [\text{X}]_0 = 1 \frac{\text{mol}}{\text{L}} & K_c = K_n V^{\sum a_i - \sum b_j} & K_{\text{w}} = [\text{H}_3\text{O}^+][\text{OH}^-] = 10^{-14} \\ a_i \text{A}_i \rightarrow b_j \text{B}_j & a_{(g)} = \gamma \frac{p}{p_0}, \; p_0 = 1 \text{ atm} & K_\chi = K_n n_{\text{tot}}^{\sum a_i - \sum b_j} & \Delta E = \Delta E^\circ - \frac{RT}{n_e N_A q_e} \ln Q \text{ (Nerst)} \\ \Delta H_{\text{r}}^\circ = b_j \Delta H_{\text{f}}^\circ(\text{B}_j) - a_i \Delta H_{\text{f}}^\circ(\text{A}_i) & K = \frac{\prod a_{\text{B}_j}^{b_j}}{\prod a_{\text{A}_i}^{a_i}}, \; K_c = \frac{\prod [\text{B}_j]^{b_j}}{\prod [\text{A}_i]^{a_i}} & \Delta G_{\text{r}}^\circ = -RT \ln K & (\text{std}) \; \Delta E = \Delta E^\circ - \frac{0.059}{n_e} \log_{10} Q \\ \forall i, j : v_{\text{r}} = -\frac{1}{a_i} \frac{\Delta[\text{A}_i]}{\Delta t} = \frac{1}{b_j} \frac{\Delta[\text{B}_j]}{\Delta t} & K_p = \frac{\prod p_{\text{B}_j}^{b_j}}{\prod p_{\text{A}_i}^{a_i}}, \; K_n = \frac{\prod n_{\text{B}_j}^{b_j}}{\prod n_{\text{A}_i}^{a_i}} & Q = K(t) = \frac{\prod a_{\text{B}_j}^{b_j}(t)}{\prod a_{\text{A}_i}^{a_i}(t)} & \text{pH} = -\log_{10} [\text{H}_3\text{O}^+] \\ & & & K_a = \frac{[\text{A}^-][\text{H}_3\text{O}^+]}{[\text{AH}]}\end{array}$$

CGS→MKS

$$\begin{array}{llllllll} \text{Substitutions:} & \vec{E}, V \times \sqrt{4\pi\epsilon_0} & \rho, \vec{J}, I, \vec{P}/\sqrt{4\pi\epsilon_0} & \vec{H} \times \sqrt{4\pi\mu_0} & \sigma \text{ (cond.)}/4\pi\epsilon_0 & \mu/\mu_0 & L \times 4\pi\epsilon_0 & \\ c \mapsto \frac{1}{\sqrt{\epsilon_0\mu_0}} & \vec{D} \times \sqrt{\frac{4\pi}{\epsilon_0}} & \vec{B}, \vec{A} \times \sqrt{\frac{4\pi}{\mu_0}} & \vec{M} \times \sqrt{\frac{\mu_0}{4\pi}} & \varepsilon/\varepsilon_0 & R, Z \times 4\pi\epsilon_0 & C/4\pi\epsilon_0 & \end{array}$$

MKS→CGS

$$\begin{array}{llllllll} \mu_0 \mapsto \frac{4\pi}{c^2} & \vec{D}/4\pi & \vec{B}, \vec{A}/c & \vec{M} \times c & \varepsilon/4\pi & R, Z \times 1 & C \times 1 & \\ \varepsilon_0 \mapsto \frac{1}{4\pi} & \vec{E}, V \times 1 & \rho, \vec{J}, I, \vec{P} \times 1 & \vec{H} \times \frac{c}{4\pi} & \sigma \text{ (cond.)} \times 1 & \mu \times \frac{4\pi}{c^2} & L \times 1 & \end{array}$$

Electrostatics (CGS)

$$\begin{array}{ll} \vec{F}_{12} = q_1 q_2 \frac{\vec{r}_2 - \vec{r}_1}{|\vec{r}_1 - \vec{r}_2|^3}; \; \vec{E}_1 = \frac{\vec{F}_{12}}{q_2}; \; V(\vec{r}) = \int \text{d}^3r' \frac{\rho(\vec{r}')}{|\vec{r} - \vec{r}'|}; \; \rho_q = \delta(\vec{r} - \vec{r}_q) & U = \frac{1}{8\pi} \int E^2 \text{d}^3x; \; \tilde{U} = \frac{1}{2} \frac{q_i q_j}{|\vec{r}_i - \vec{r}_j|} = \frac{1}{8\pi} \sum_{i \neq j} \int \vec{E}_i \vec{E}_j \text{d}^3x \\ \oint \vec{E} \text{d}\vec{S} = 4\pi \int \rho \text{d}^3x; \; -\nabla^2 V = \vec{\nabla} \vec{E} = 4\pi \rho; \; \vec{\nabla} \times \vec{E} = 0 & V(\vec{r}) = \int \rho G_{\text{D}}(\vec{r}) \text{d}^3x - \frac{1}{4\pi} \oint_S V \frac{\partial G_{\text{D}}}{\partial n} \text{d}S \\ & V(\vec{r}) = \langle V \rangle_S + \int \rho G_{\text{N}}(\vec{r}) \text{d}^3x + \frac{1}{4\pi} \oint_S \frac{\partial V}{\partial n} G_{\text{N}}(\vec{r}) \text{d}S \end{array}$$

$$\nabla_y^2 G(\vec{x}, \vec{y}) = -4\pi\delta(\vec{x} - \vec{y}); G_D(\vec{x}, \vec{y})|_{\vec{y} \in S} = 0; \frac{\partial G_N}{\partial n}|_{\vec{y} \in S} = -\frac{4\pi}{S}$$

$$U_{\text{sphere}} = \frac{3}{5} \frac{Q^2}{R}; \vec{p} = \int d^3r \rho \vec{r}; \vec{E}_{\text{dip}} = \frac{3(\vec{p}\vec{r})\vec{r} - \vec{p}}{r^3}; V_{\text{dip}} = \frac{\vec{p}\vec{r}}{r^2}$$

$$\text{force on a dipole: } \vec{F}_{\text{dip}} = (\vec{p}\vec{\nabla})\vec{E}$$

$$Q_{ij} = \int d^3r \rho(\vec{r})(3r_i r_j - \delta_{ij} r^2); V_{\text{quad}} = \frac{1}{6r^5} Q_{ij} (3r_i r_j - \delta_{ij} r^2)$$

Magnetostatics (CGS)

$$\vec{\nabla} \vec{J} = -\frac{\partial \rho}{\partial t} = 0; I = \int \vec{J} d\vec{S}$$

$$\text{solenoid: } B = 4\pi \frac{I_s}{c}; \text{wires: } \frac{dF}{dl} = \frac{2}{c^2} \frac{I_1 I_2}{d}$$

$$d\vec{F} = \frac{I d\vec{l}}{c} \times \vec{B} = d^3x \frac{\vec{J}}{c} \times \vec{B}; \vec{F}_q = q \frac{\vec{r}}{c} \times \vec{B}$$

$$d\vec{B} = \frac{I d\vec{l}}{c} \times \frac{\vec{r}}{r^3}; \vec{B}_q = q \frac{\vec{r}}{c} \times \frac{\vec{r}}{r^3}$$

Electromagnetism (CGS)

$$\text{Faraday: } \mathcal{E} = -\frac{1}{c} \frac{d\Phi_B}{dt}; \int d^3x \vec{J} = \dot{\vec{p}}$$

$$\vec{\nabla} \times \vec{E} = -\frac{1}{c} \frac{\partial \vec{B}}{\partial t}; \vec{\nabla} \vec{E} = 4\pi \rho; \vec{\nabla} \vec{J} = -\frac{\partial \rho}{\partial t}$$

$$\vec{\nabla} \times \vec{B} = 4\pi \frac{\vec{J}}{c} + \frac{1}{c} \frac{\partial \vec{E}}{\partial t}; \vec{\nabla} \vec{B} = 0$$

$$d\vec{F} = d^3x (\rho \vec{E} + \frac{\vec{J}}{c} \times \vec{B}); \vec{F}_q = q(\vec{E} + \frac{\vec{r}}{c} \times \vec{B})$$

$$u = \frac{E^2 + B^2}{8\pi}; \vec{S} = \frac{c}{4\pi} \vec{E} \times \vec{B}; \vec{g} = \frac{\vec{S}}{c^2}$$

$$\mathbf{T}^E = \frac{1}{4\pi} (\vec{E} \otimes \vec{E} - \frac{1}{2} E^2); \mathbf{T} = \mathbf{T}^E + \mathbf{T}^B$$

$$-\frac{\partial u}{\partial t} = \vec{J} \vec{E} + \vec{\nabla} \vec{S}; -\frac{\partial \vec{g}}{\partial t} = \rho \vec{E} + \frac{\vec{J}}{c} \times \vec{B} - \vec{\nabla} \mathbf{T}$$

$$\vec{B} = \vec{\nabla} \times \vec{A}; \vec{E} = -\vec{\nabla} \phi - \frac{1}{c} \frac{\partial \vec{A}}{\partial t}$$

$$-\nabla^2 \phi - \frac{1}{c} \frac{\partial}{\partial t} \vec{\nabla} \vec{A} = 4\pi \rho$$

$$\vec{\nabla} (\vec{\nabla} \vec{A} + \frac{1}{c} \frac{\partial \phi}{\partial t}) - \nabla^2 \vec{A} + \frac{1}{c} \frac{\partial^2 \vec{A}}{\partial t^2} = 4\pi \frac{\vec{J}}{c}$$

$$(\phi, \vec{A}) \cong (\phi - \frac{1}{c} \frac{\partial \Lambda}{\partial t}, \vec{A} + \vec{\nabla} \Lambda)$$

$$(\phi, \vec{A}) = \int d^3r' \frac{(\rho, \frac{\vec{J}}{c})(\vec{r}', t - \frac{1}{c}|\vec{r} - \vec{r}'|)}{|\vec{r} - \vec{r}'|}$$

$$\vec{\nabla} \vec{A} = 0 \rightarrow \square \vec{A} = \frac{4\pi}{c} (\vec{J} - \vec{J}_L) =: \frac{4\pi}{c} \vec{J}_T$$

E.M. in matter (CGS)

$$\vec{\nabla} \vec{D} = 4\pi \rho_{\text{ext}}; \vec{\nabla} \times \vec{E} = -\frac{1}{c} \frac{\partial \vec{B}}{\partial t}$$

$$\vec{\nabla} \vec{B} = 0; \vec{\nabla} \times \vec{H} = 4\pi \frac{\vec{J}_{\text{ext}}}{c} + \frac{1}{c} \frac{\partial \vec{D}}{\partial t}$$

$$\vec{P} = \frac{d\langle \vec{p} \rangle}{dV}; \vec{M} = \frac{d\langle \vec{m} \rangle}{dV}$$

$$\rho_{\text{pol}} = -\vec{\nabla} \vec{P}; \sigma_{\text{pol}} = \hat{n} \vec{P}; \frac{\vec{J}_{\text{mag}}}{c} = \vec{\nabla} \times \vec{M}$$

$$\vec{D}_{\text{pol}} = \vec{E} + 4\pi \vec{P}; \vec{H}_{\text{mag}} = \vec{B} - 4\pi \vec{M}$$

$$\text{static linear isotropic: } \vec{P} = \chi \vec{E}$$

$$\text{static linear: } P_i = \chi_{ij} E_j$$

$$\text{static linear: } \varepsilon = 1 + 4\pi \chi$$

$$\text{static: } \Delta D_{\perp} = 4\pi \sigma_{\text{ext}}; \Delta E_{\parallel} = 0$$

$$\text{static linear: } u = \frac{1}{8\pi} \vec{E} \vec{D}$$

$$\Delta U_{\text{dielectric}} = -\frac{1}{2} \int d^3r \vec{P} \vec{E}_0$$

$$\text{plane capacitor: } C = \frac{\varepsilon}{4\pi} \frac{S}{d}$$

$$\text{cilindric capacitor: } C = \frac{L}{2 \log \frac{R}{r}}$$

$$\text{atomic polarizability: } \vec{p} = \alpha \vec{E}_{\text{loc}}$$

Quantum mechanics (CGS)

$$r_e = \frac{e^2}{m_e c^2}; \alpha = \frac{e^2}{\hbar c}; \mu_B = \frac{e \hbar}{2 m_e c}$$

$$\lambda_{\text{Broglie}} = \frac{h}{p}$$

$$\text{Planck: } \frac{8\pi \hbar}{c^3} \frac{\nu^3}{e^{kT} - 1} d\nu$$

$$i \hbar \frac{\partial \mathcal{U}}{\partial t} = \mathcal{H} \mathcal{U}; \frac{\partial \mathcal{H}}{\partial t} = 0 \Rightarrow \mathcal{U}(t) = e^{-\frac{i \mathcal{H} t}{\hbar}}$$

$$V(r, \theta) = \sum_{l=0}^{\infty} (A_l r^l + \frac{B_l}{r^{l+1}}) P_l(\cos \theta)$$

$$V(r, \theta, \varphi) = \sum_{l=0}^{\infty} \sum_{m=-l}^l (A_{lm} r^l + \frac{B_{lm}}{r^{l+1}}) Y_{lm}(\theta, \varphi)$$

$$V(r > \text{diam supp } \rho, \theta, \varphi) = \sum_{l=0}^{\infty} \frac{4\pi}{2l+1} \sum_{m=-l}^l q_{lm}[\rho] \frac{Y_{lm}(\theta, \varphi)}{r^{l+1}}$$

$$q_{lm}[\rho] = \int_0^{\infty} r^2 dr \int_0^{2\pi} d\varphi \int_0^{\pi} \sin \theta d\theta r^l \rho(r, \theta, \varphi) Y_{lm}^*(\theta, \varphi)$$

$$\vec{B} = \vec{\nabla} \times \vec{A}; \vec{A} = \int d^3r' \frac{\vec{J}}{c} \frac{1}{|\vec{r} - \vec{r}'|} + \vec{\nabla} \Lambda$$

$$\vec{B} = \int d^3r' \frac{\vec{J}}{c} \times \frac{\vec{r} - \vec{r}'}{|\vec{r} - \vec{r}'|^3}$$

$$\varphi = \frac{I}{c} \Omega, \vec{B} = -\vec{\nabla} \varphi$$

$$\vec{\nabla} \vec{A} = 0 \rightarrow \nabla^2 \vec{A} = -4\pi \frac{\vec{J}}{c}$$

$$\vec{J}_L = \frac{1}{4\pi} \vec{\nabla} \frac{\partial \phi}{\partial t} = -\frac{1}{4\pi} \vec{\nabla} \int \frac{\vec{\nabla}' \vec{J}'}{|\vec{r} - \vec{r}'|} d^3r'$$

$$\vec{E}'_{\parallel} = \vec{E}_{\parallel}; \vec{B}'_{\parallel} = \vec{B}_{\parallel}$$

$$\vec{E}'_{\perp} = \gamma (\vec{E}_{\perp} + \frac{\vec{v}}{c} \times \vec{B})$$

$$\vec{B}'_{\perp} = \gamma (\vec{B}_{\perp} - \frac{\vec{v}}{c} \times \vec{E})$$

$$\vec{\omega}_{\text{Larmor}} = -\frac{q}{2mc} \vec{B}$$

$$\text{plane wave: } \begin{cases} \vec{E} = \vec{E}_0 e^{i(\vec{k}\vec{r} - \omega t)} \\ \vec{B} = \hat{k} \times \vec{E} \\ \omega = ck \end{cases}$$

$$\vec{B}_{\text{diprad}} = \frac{1}{c^2} \frac{\ddot{\vec{p}} \times \hat{r}}{r} \Big|_{t_{\text{rit}}}; \vec{E}_{\text{diprad}} = \vec{B}_{\text{diprad}} \times \hat{r}$$

$$\text{Larmor: } P = \frac{2}{3c^3} |\ddot{\vec{p}}|^2$$

$$\text{Rel. Larmor: } P = \frac{2}{3c^3} q^2 \gamma^6 (a^2 - (\vec{a} \times \vec{\beta})^2)$$

$$\vec{A}_{\text{dm}} = \frac{1}{c} \frac{\dot{\vec{m}} \times \hat{r}}{r} \Big|_{t_{\text{rit}}}$$

$$\text{L.W.: } (\phi, \vec{A}) = \frac{q(1, \frac{\vec{v}}{c})}{[r - \frac{\vec{v} \cdot \vec{r}}{c}]_{t_{\text{rit}}}}; t_{\text{rit}} = t - \frac{r}{c} \Big|_{t_{\text{rit}}}$$

$$A^{\mu} = (\phi, \vec{A}); J^{\mu} = (c\rho, \vec{J})$$

$$\text{non-interacting gas: } \vec{p} = \alpha \vec{E}_0; \chi = n\alpha$$

$$\text{hom. cubic isotropic: } \chi = \frac{1}{\frac{1}{n\alpha} - \frac{4\pi}{3}}$$

$$\text{Clausius-Mossotti: } \frac{\varepsilon - 1}{\varepsilon + 2} = \frac{4\pi}{3} n\alpha$$

$$\text{perm. dipole: } \chi = \frac{1}{3} \frac{n p_0^2}{kT}$$

$$\text{local field: } \vec{E}_{\text{loc}} = \vec{E} + \frac{4\pi}{3} \vec{P}$$

$$\vec{J} \vec{E} = -\vec{\nabla} \left(\frac{c}{4\pi} \vec{E} \times \vec{H} \right) - \frac{1}{4\pi} \left(\vec{E} \frac{\partial \vec{D}}{\partial t} + \vec{H} \frac{\partial \vec{B}}{\partial t} \right)$$

$$n = \sqrt{\varepsilon \mu}; k = n \frac{\omega}{c}$$

$$\text{plane wave: } B = nE$$

$$\vec{J}_c = \sigma \vec{E}; \varepsilon_{\sigma} = 1 + i \frac{4\pi \sigma}{\omega}$$

$$\omega_p^2 = 4\pi \frac{n_{\text{vol}} q^2}{m}; \omega_{\text{cyclo}} = \frac{qB}{mc}$$

$$\text{I: } u = \frac{1}{8\pi} (\vec{E} \vec{D} + \vec{H} \vec{B})$$

$$\text{I: } \langle S_z \rangle = \frac{c}{n} \langle u \rangle$$

$$\text{II: } u = \frac{1}{8\pi} \left(\frac{\partial}{\partial \omega} (\varepsilon \omega) E^2 + \frac{\partial}{\partial \omega} (\mu \omega) H^2 \right)$$

$$\text{II: } \langle S_z \rangle = v_g \langle u \rangle; v_g = \frac{\partial \omega}{\partial k} = \frac{c}{n + \omega \frac{\partial n}{\partial \omega}}$$

$$\text{III: } \langle W \rangle = \frac{\omega}{4\pi} (\text{Im } \varepsilon \langle E^2 \rangle + \text{Im } \mu \langle H^2 \rangle)$$

$$[\mathcal{H}(t), \mathcal{H}(t')] = 0 \Rightarrow \mathcal{U}(t) = e^{-\frac{i \int_0^t dt' \mathcal{H}(t')}{\hbar}}$$

$$\mathcal{U} = \left(\frac{-i}{\hbar} \right)^k \int_0^t dt_1 \dots \int_0^{t_{k-1}} dt_k \mathcal{H}(t_1) \dots \mathcal{H}(t_k)$$

$$A_H(t) = \mathcal{U}(t)^{\dagger} A \mathcal{U}(t)$$

$$\frac{\partial \mathcal{H}}{\partial t} = 0 \Rightarrow \frac{dA_H}{dt} = \frac{[A_H, \mathcal{H}]}{i\hbar}$$

$$(A \otimes B)(|a\rangle \otimes |b\rangle) = A|a\rangle \otimes B|b\rangle$$

$$\text{Lorenz gauge: } \partial_{\alpha} A^{\alpha} = 0$$

$$\text{Temporal gauge: } \phi = 0$$

$$\text{Axial gauge: } A_3 = 0$$

$$\text{Coulomb gauge: } \vec{\nabla} \vec{A} = 0$$

$$F^{\mu\nu} = \partial^{\mu} A^{\nu} - \partial^{\nu} A^{\mu}; \mathcal{F}^{\mu\nu} = \frac{1}{2} \varepsilon^{\mu\nu\alpha\beta} F_{\alpha\beta}$$

$$F^{\mu\nu} = \begin{pmatrix} 0 & -E_x & -E_y & -E_z \\ E_x & 0 & -B_z & B_y \\ E_y & B_z & 0 & -B_x \\ E_z & -B_y & B_x & 0 \end{pmatrix}$$

$$\partial_{\alpha} F^{\alpha\nu} = 4\pi \frac{J^{\nu}}{c}; \partial_{\alpha} \mathcal{F}^{\alpha\nu} = 0; \frac{dp^{\mu}}{d\tau} = q F^{\mu\alpha} \frac{v_{\alpha}}{c}$$

$$\partial_{\mu} F_{\nu\sigma} + \partial_{\nu} F_{\sigma\mu} + \partial_{\sigma} F_{\mu\nu} = 0; \det F = (\vec{E} \vec{B})^2$$

$$F^{\alpha\beta} F_{\alpha\beta} = 2(B^2 - E^2); F^{\alpha\beta} \mathcal{F}_{\alpha\beta} = 4\vec{E} \vec{B}$$

$$\Theta^{\mu\nu} = \frac{1}{4\pi} (F^{\mu}_{\alpha} F^{\alpha\nu} + \frac{1}{4} g^{\mu\nu} F_{\alpha\beta} F^{\alpha\beta})$$

$$\Theta^{\mu\nu} = \left(\frac{u}{c\vec{g}} - \mathbf{T} \right); \partial_{\alpha} \Theta^{\alpha\nu} = \frac{J_{\alpha}}{c} F^{\alpha\nu} = -G^{\nu}$$

$$\mathcal{L} = \frac{mc^2}{\gamma} - q \vec{A} \frac{\vec{v}}{c} + q\phi; \mathcal{H} = \frac{1}{2m} \left(\vec{p} - \frac{q\vec{A}}{c} \right)^2 + q\phi$$

$$\text{plane wave: } \mathbf{T} = -u \hat{k} \otimes \hat{k}; \Theta^{\mu\nu} = u \hat{k}^{\mu} \hat{k}^{\nu}$$

$$\text{Fresnel TE (S): } \frac{E_t}{E_i} = \frac{2}{1 + \frac{k_{tz}}{k_{iz}}}; \frac{E_r}{E_i} = \frac{1 - \frac{k_{tz}}{k_{iz}}}{1 + \frac{k_{tz}}{k_{iz}}}$$

$$\text{TM (P): } \frac{E_t}{E_i} = \frac{2}{\frac{n_2}{n_1} + \frac{n_1}{n_2} \frac{k_{tz}}{k_{iz}}}; \frac{E_r}{E_i} = \frac{\frac{n_2}{n_1} - \frac{n_1}{n_2} \frac{k_{tz}}{k_{iz}}}{\frac{n_2}{n_1} + \frac{n_1}{n_2} \frac{k_{tz}}{k_{iz}}}$$

$$\text{Fresnel: } k_{tz} = \pm \sqrt{\varepsilon_2 \left(\frac{\omega}{c} \right)^2 - k_x^2}, \text{Im } k_{tz} > 0$$

$$\text{Drüde-Lorentz: } \varepsilon = 1 - \frac{\omega_p^2}{\omega^2 + i\gamma\omega - \omega_0^2}$$

$$P(t) = \int_{-\infty}^{\infty} g(t - t') E(t') dt'$$

$$P(\omega) = \chi(\omega) E(\omega)$$

$$\chi(\omega) = \int_{-\infty}^{\infty} e^{i\omega t} g(t) dt; \chi(-\omega) = \chi^*(\omega)$$

$$g(t < 0) = 0 \Rightarrow$$

$$\text{Re } \varepsilon(\omega) = 1 + \frac{2}{\pi} \int_0^{\infty} \frac{\omega' (\text{Im } \varepsilon(\omega') - \frac{4\pi\sigma_0}{\omega'})}{\omega'^2 - \omega^2} d\omega'$$

$$\text{Im } \varepsilon(\omega) = -\frac{2\omega}{\pi} \int_0^{\infty} \frac{\text{Re } \varepsilon(\omega') - 1}{\omega'^2 - \omega^2} d\omega' + \frac{4\pi\sigma_0}{\omega}$$

$$\text{sum rule: } \frac{\pi}{2} \omega_p^2 = \int_0^{\infty} \omega \text{Im } \varepsilon d\omega$$

$$\text{sum rule: } 2\pi^2 \sigma_0 = \int_0^{\infty} (1 - \text{Re } \varepsilon) d\omega$$

$$\text{sum rule: } \int_0^{\infty} (\text{Re } n - 1) d\omega = 0$$

$$\text{Miller rule: } \chi^{(2)}(\omega, \omega) \propto \chi^{(1)}(\omega)^2 \chi^{(1)}(2\omega)$$

$$(\langle a| \otimes \langle b|)(|c\rangle \otimes |d\rangle) = \langle a|c\rangle \langle b|d\rangle$$

$$A^{(1)} + B^{(2)} = A^{(1)} \otimes \mathbb{1}^{(2)} + \mathbb{1}^{(1)} \otimes B^{(2)}$$

$$[A, BC] = [A, B]C + B[A, C]$$

$$[A, [B, C]] + [B, [C, A]] + [C, [A, B]] = 0$$

$$[X, P] = i\hbar; \langle x|p\rangle = \frac{1}{\sqrt{2\pi\hbar}} e^{\frac{ipx}{\hbar}}$$

$$\begin{aligned}\langle x|X|\psi\rangle &= x\,\langle x|\psi\rangle\,; \, \langle x|P|\psi\rangle = \frac{\hbar}{i}\frac{\partial}{\partial x}\,\langle x|\psi\rangle \\ \langle (A-\langle A\rangle)^2\rangle \langle (B-\langle B\rangle)^2\rangle &\geq \frac{1}{4}|\langle [A,B]\rangle|^2 \\ e^BAe^{-B} &= A + [B,A] + \tfrac{1}{2!}[B,[B,A]] + \cdots \\ [A,B] \propto 1 &\Rightarrow [A,f(B)] = [A,B]f'(B) \\ [A,B] \propto 1 &\Rightarrow e^Ae^B = e^{A+B+\frac{1}{2}[A,B]} \\ e^{\frac{ip'x}{\hbar}}|p\rangle &= |p+p'\rangle\,; \, e^{-\frac{iPx'}{\hbar}}|x\rangle = |x+x'\rangle \\ \psi(x) &= \langle x|\psi\rangle\,; \, \rho = |\psi|^2\,; \, \psi = \sqrt{\rho}e^{\frac{iS}{\hbar}} \\ \mathcal{H} = \frac{\vec{P}^2}{2m} + V(\vec{X}) : \vec{j} &= \frac{\hbar}{m}\operatorname{Im}(\psi^*\vec{\nabla}\psi) = \frac{\rho\vec{\nabla}S}{m} \\ \frac{\partial\rho}{\partial t} &= -\vec{\nabla}\vec{j}\,; \, \int\mathrm{d}^3x\vec{j} = \frac{\langle\vec{P}\rangle}{m}\end{aligned}$$

QM perturbative

$$\begin{aligned}\mathcal{H} &= \mathcal{H}_0 + V \\ E &= E_0 + \epsilon_1 + \epsilon_2 + \cdots \\ V \mapsto kV &\Rightarrow \epsilon_n \mapsto k^n \epsilon_n \\ |\psi\rangle &= |\psi_0\rangle + |\psi_1\rangle + \cdots\,; \, \langle\psi_0|\psi\rangle = 1 \\ \epsilon_1 &= \langle\psi_0|V|\psi_0\rangle \\ |\psi_1\rangle &= \sum_{\alpha\neq\psi_0}|\alpha\rangle\,\frac{\langle\alpha|V|\psi_0\rangle}{E_0-E_\alpha} \\ \epsilon_2 &= \sum_{\alpha\neq\psi_0}\frac{|\langle\alpha|V|\psi_0\rangle|^2}{E_0-E_\alpha} = \langle\psi_0|V|\psi_1\rangle \\ \epsilon_n &= \langle\psi_0|V|\psi_{n-1}\rangle\end{aligned}$$

QM scattering

$$\begin{aligned}\frac{\mathrm{d}\sigma}{\mathrm{d}\Omega} &= |f|^2\,; \, \text{elastic: } f = \frac{k}{2\pi\hbar v}\left\langle\vec{k}_f\right|T\left|\vec{k}_i\right\rangle \\ \text{anelastic: } f &= \sqrt{\frac{k'^2}{4\pi^2\hbar^2vv'}}\left\langle\vec{k}',b\right|T\left|\vec{k},a\right\rangle\end{aligned}$$

QM rotations

$$\begin{aligned}\sigma_1 &= \left(\begin{smallmatrix} 0 & 1 \\ 1 & 0 \end{smallmatrix}\right); \, \sigma_2 = \left(\begin{smallmatrix} 0 & -i \\ i & 0 \end{smallmatrix}\right); \, \sigma_3 = \left(\begin{smallmatrix} 1 & 0 \\ 0 & -1 \end{smallmatrix}\right) \\ \sigma_i\sigma_j &= \delta_{ij} + i\varepsilon_{ijk}\sigma_k\,; \, \operatorname{tr}(\hat{n}\vec{\sigma}) = 0 \\ [\sigma_i,\sigma_j] &= 2i\varepsilon_{ijk}\sigma_k\,; \, \{\sigma_i,\sigma_j\} = 2\delta_{ij} \\ (\vec{\sigma}\vec{a})(\vec{\sigma}\vec{b}) &= \vec{a}\vec{b} + i\vec{\sigma}(\vec{a}\times\vec{b}) \\ e^{i\vec{\sigma}\hat{n}\alpha} &= \cos\alpha + i(\vec{\sigma}\hat{n})\sin\alpha \\ |\vec{\sigma}\hat{n},1\rangle &= \cos\frac{\theta}{2}|\sigma_3,1\rangle + e^{i\varphi}\sin\frac{\theta}{2}|\sigma_3,-1\rangle \\ M &= \vec{v}\vec{\sigma} \Rightarrow \lambda_M = \pm|\vec{v}| \\ U(R_{\hat{n},\phi}) &= \exp\Big(-i\vec{J}\hat{n}\phi\Big) \\ [J_i,J_j] &= i\varepsilon_{ijk}J_k\,; \, J_{\pm} := J_x \pm iJ_y \\ [J_+,J_-] &= 2J_z\,; \, [J_z,J_{\pm}] = \pm J_{\pm} \\ [J^2,J_{\pm}] &= [J^2,J_z] = 0 \\ J^2|j,m\rangle &= j(j+1)|j,m\rangle \\ J_z|j,m\rangle &= m|j,m\rangle \\ J_{\pm}|j,m\rangle &= \sqrt{j(j+1)-m(m\pm1)}|j,m\pm1\rangle \\ m &= -j,j-1,\ldots,j\,; \, 2j\in\mathbb{N}\end{aligned}$$

QM solutions

$$\begin{aligned}\text{BOX: } \mathcal{H} &= \frac{P^2}{2m} + \begin{cases} 0 & 0 < x < L \\ \infty & \text{otherwise} \end{cases} \\ E_n &= \frac{n^2\pi^2\hbar^2}{2mL^2}, \, n \geq 1 \\ \psi_n(x) &= \sqrt{\frac{2}{L}}\sin(n\pi\frac{x}{L}) = \sqrt{\frac{2}{L}}\sin\Big(\sqrt{\frac{2mE}{\hbar^2}}x\Big) \\ \Delta x^2 &= L^2\Big(\frac{1}{12} - \frac{1}{2n^2\pi^2}\Big)\,; \, \Delta p = \frac{\hbar n\pi}{L} \\ \text{HARM.: } \mathcal{H} &= \frac{P^2}{2m} + \frac{m\omega^2X^2}{2}\,; \, x_0 := \sqrt{\frac{\hbar}{m\omega}}\end{aligned}$$

$$\begin{aligned}\psi(x,t) &= \int\mathrm{d}x'K(x,t;x')\psi(x',t=0) \\ K(x,t;x') &= \sum_E\psi_E(x')^*\psi_E(x)e^{-\frac{iEt}{\hbar}} = \\ &= \langle x|e^{-\frac{i\mathcal{H}t}{\hbar}}|x'\rangle \\ (\mathcal{H}-i\hbar\frac{\partial}{\partial t})K(x,t;x') &= -i\hbar\delta(x-x')\delta(t) \\ \rho[|\alpha_i\rangle,w_i] &:= \sum_iw_i|\alpha_i\rangle\langle\alpha_i| \\ \operatorname{tr}\rho &= 1\,; \, [A] := \operatorname{tr}(\rho A) \\ \#\{w_i>0\} &= 1 \iff \operatorname{tr}(\rho^2) = 1 \\ \#\{w_i>0\}>1 &\iff 0<\operatorname{tr}(\rho^2)<1 \\ \frac{\partial\rho}{\partial t} &= -\frac{[\rho,\mathcal{H}]}{i\hbar} \\ W_\psi(x,p) &= \int\frac{\mathrm{d}y}{2\pi\hbar}\langle x+\frac{y}{2}|\psi\rangle\langle\psi|x-\frac{y}{2}\rangle e^{-\frac{ipy}{2}}\end{aligned}$$

$$\begin{aligned}Q &:= \sum_{\alpha\neq\psi_0}|\alpha\rangle\langle\alpha|\,; \, G_Q := Q\frac{1}{E_0-\mathcal{H}_0}Q \\ |\psi_n\rangle &= G_QV|\psi_{n-1}\rangle - \sum_{s=1}^{n-1}\epsilon_sG_Q|\psi_{n-s}\rangle \\ \mathcal{H} &= \mathcal{H}_0 + V(t) \\ |\psi(t)\rangle &= \sum_k a_k(t)e^{-iE_kt/\hbar}|k\rangle \\ a_k &= a_k^{(0)} + a_k^{(1)} + \cdots \\ \hbar\omega_{ab} &:= E_a - E_b\,; \, V_{ab} := \langle a|V|b\rangle \\ a_k^{(n+1)}(t) &= \frac{i}{\hbar}\sum_s\int_0^t\mathrm{d}\tau e^{i\omega_{ks}\tau}V_{ks}(\tau)a_s^{(n)}(\tau) \\ P_{i\rightarrow f}(t) &= \left|\frac{1}{\hbar}\int_0^t\mathrm{d}\tau e^{i\omega_{fi}\tau}V_{fi}(\tau)\right|^2\end{aligned}$$

$$\begin{aligned}\text{Born: } \langle|T|\rangle &\approx \langle|V|\rangle \\ \vec{q} &:= \vec{k}' - \vec{k}\,; \, \left\langle\vec{k}'\right|V\left|\vec{k}\right\rangle = \int\mathrm{d}\vec{x}V(\vec{x})e^{-i\vec{q}\vec{x}} \\ V = V(r) : \left\langle\vec{k}'\right|V\left|\vec{k}\right\rangle &= 4\pi\int_0^\infty\mathrm{d}r\,rV\frac{\sin(qr)}{q}\end{aligned}$$

$$\begin{aligned}U\in\mathrm{SU}_2: U &= \left(\begin{smallmatrix} a & b \\ -b^* & a^* \end{smallmatrix}\right), \, |a|^2 + |b|^2 = 1 \\ U &= e^{-\frac{i\sigma_z\alpha}{2}}e^{-\frac{i\sigma_y\beta}{2}}e^{-\frac{i\sigma_z\gamma}{2}} \\ a &= \cos\frac{\phi}{2} - in_z\sin\frac{\phi}{2} = e^{-i\frac{\alpha+\gamma}{2}}\cos\frac{\beta}{2} \\ b &= -\sin\frac{\phi}{2}(n_y+in_x) = -e^{-i\frac{\alpha-\gamma}{2}}\sin\frac{\beta}{2} \\ U(\vec{v}\vec{\sigma})U^\dagger &= (R_{\hat{n},\phi}\vec{v})\vec{\sigma} \\ \vec{L} = \vec{X}\times\vec{P}; \, L_z &= \frac{\hbar}{i}\frac{\partial}{\partial\varphi} \\ L_{\pm} &= \hbar e^{\pm i\varphi}\Big(\pm\frac{\partial}{\partial\theta} + i\cot\theta\frac{\partial}{\partial\varphi}\Big) \\ \vec{L}^2 &= -\hbar^2\Big(\frac{1}{\sin^2\theta}\frac{\partial^2}{\partial\varphi^2} + \frac{1}{\sin\theta}\frac{\partial}{\partial\theta}\Big(\frac{1}{\sin\theta}\frac{\partial}{\partial\theta}\Big)\Big) \\ \vec{L}^2Y_{lm} &= \hbar l(l+1)Y_{lm}\,; \, L_zY_{lm} = \hbar mY_{lm} \\ A = \vec{A} : \leftrightarrow [J_i,A_j] &= i\varepsilon_{ijk}A_k \\ T = \mathbf{T} : \leftrightarrow [J_z,T_q] &= qT_q, \\ \Big[J_{\pm},T_q^{(k)}\Big] &= \sqrt{k(k+1)-q(q\pm1)}T_{q\pm1}^{(k)} \\ U(R)^\dagger\vec{A}U(R) &= R\vec{A} \\ \Big[\vec{\theta}\vec{J},\vec{A}\Big] &= -i\vec{\theta}\times\vec{A}\end{aligned}$$

$$\begin{aligned}A &:= \frac{1}{\sqrt{2}x_0}\Big(X + \frac{iP}{m\omega}\Big)\,; \, A^\dagger = \frac{1}{\sqrt{2}x_0}\Big(X - \frac{iP}{m\omega}\Big) \\ N &:= A^\dagger A = \frac{\mathcal{H}}{\hbar\omega} - \frac{1}{2}\,; \, \mathcal{H} = \hbar\omega\Big(N + \frac{1}{2}\Big) \\ [A,A^\dagger] &= 1\,; \, [N,A] = -A\,; \, [N,A^\dagger] = A^\dagger \\ A^\dagger|n\rangle &= \sqrt{n+1}|n+1\rangle\,; \, A|n\rangle = \sqrt{n}|n-1\rangle \\ |n\rangle &= \frac{(A^\dagger)^n}{\sqrt{n!}}|0\rangle, \, n = 0,1,\ldots \\ \psi_n(x) &= \frac{1}{\pi^{\frac{1}{4}}\sqrt{2^n n!}x_0}H_n\Big(\frac{x}{x_0}\Big)e^{-\frac{1}{2}\Big(\frac{x}{x_0}\Big)^2}\end{aligned}$$

$$\begin{aligned}V(x+a) = V(x) \Rightarrow \psi_{nk}(x) &= e^{ikx}u_{nk}(x), \\ u_{nk}(x+a) &= u_{nk}(x) \\ \text{trasm.} &\approx e^{-\int_0^d\frac{\sqrt{2m(V(x)-E)}}{\hbar}\mathrm{d}x} \\ \mathcal{H}_{\text{em}} &= \frac{1}{2m}\Big(\vec{p}-\frac{e}{c}\vec{A}\Big)^2 + e\phi(\vec{x}) - \vec{\mu}\vec{B} \\ \vec{j}_{\text{em}} &= \frac{e\hbar}{m}\operatorname{Im}(\psi^*\vec{\nabla}\psi) - \frac{e^2}{mc}|\psi|^2\vec{A} + c\vec{\nabla}\times(\psi^*\vec{\mu}\psi) \\ &= \frac{e}{m}\operatorname{Im}\Big(\psi^*\Big(\vec{p}-\frac{e}{c}\vec{A}\Big)\psi\Big) + c\vec{\nabla}\times(\psi^*\vec{\mu}\psi) \\ D_\mu &= \partial_\mu + i\frac{q}{\hbar c}A_\mu \\ \vec{A} \mapsto \vec{A} + \vec{\nabla}\Lambda \Rightarrow |\psi\rangle &\mapsto e^{\frac{iq\Lambda}{\hbar c}}|\psi\rangle \\ \vec{\mu} &= g\frac{e\hbar}{2mc}\vec{L}\end{aligned}$$

$$\begin{aligned}P_{i\rightarrow f} &= \left|\frac{1}{\hbar\omega_{fi}}\int_{-\infty}^\infty\mathrm{d}t\,e^{i\omega_{fi}t}V'_{fi}(t) + O(V^2)\right|^2 \\ |\psi(t)\rangle &=: e^{-i\mathcal{H}_0t/\hbar}|\tilde{\psi}(t)\rangle \\ |\tilde{\psi}(t)\rangle &=: \tilde{U}(t)|\tilde{\psi}(0)\rangle \\ i\hbar\frac{\partial\tilde{U}}{\partial t} &=: \tilde{\mathcal{H}}\tilde{U}\,; \, \tilde{\mathcal{H}}(t) = e^{i\mathcal{H}_0t/\hbar}V(t)e^{-i\mathcal{H}_0t/\hbar} \\ R_{i\rightarrow f} &\approx \frac{2\pi}{\hbar}\Big|V_{fi} + \sum_{s\neq i}\frac{V_{fs}V_{si}}{E_i-E_s+i\varepsilon}\Big|^2\delta(E_f-E_i) \\ V(t) &= Fe^{-i\omega t} + F^\dagger e^{i\omega t} \rightarrow \\ \rightarrow R_{i\rightarrow f} &\approx \frac{2\pi}{\hbar}|F_{fi}|^2\delta(E_f-E_i\pm\hbar\omega)\end{aligned}$$

$$\begin{aligned}\psi(\vec{x}) &= e^{i\vec{k}\vec{x}} - \frac{m}{2\pi\hbar^2}\int\mathrm{d}\vec{y}\frac{e^{i\vec{k}|\vec{x}-\vec{y}|}}{|\vec{x}-\vec{y}|}V(\vec{y})\psi(\vec{y}) \\ f &= -\frac{m}{2\pi\hbar^2}\int\mathrm{d}\vec{y}\,e^{-i\vec{k}'\vec{y}}V(\vec{y})\psi(\vec{y}) \\ \mathrm{d}\Omega &= 2\pi\frac{q\mathrm{d}q}{k^2}\,; \, \sigma_{\text{tot}} = \frac{4\pi}{k}\operatorname{Im}(f(0)) \\ \text{charge: } f &= -\frac{2m}{\hbar^2}\frac{q_{\text{part}}}{q^2}\Big(Q_{\text{tot}} - \int\mathrm{d}\vec{r}\,\rho(\vec{r})e^{-i\vec{q}\vec{r}}\Big) \\ \Big[\hat{n}\vec{J},T_q^{(k)}\Big] &= \sum_{q'}\langle kq'|\hat{n}\vec{J}|kq\rangle T_{q'}^{(k)} \\ \langle jm'|\,e^{-i\alpha J_z}\,e^{-i\beta J_y}\,e^{-i\gamma J_z}\,|jm\rangle &= \\ = e^{-i\alpha m'}\langle jm'|\,e^{-i\beta J_y}\,|jm\rangle\,e^{-i\gamma m} & \\ \langle jm'|\,e^{-i\beta J_y}\,|jm\rangle = & \\ = (-1)^{m'-m}\sqrt{\frac{(j+m')!(j-m')!}{(j+m)!(j-m)!}}\Big(\cos\frac{\beta}{2}\Big)^{2j}. & \\ \cdot\sum_\mu(-1)^\mu\binom{j+m}{\mu}\binom{j-m}{j-m'-\mu}\Big(\tan\frac{\beta}{2}\Big)^{m'-m+2\mu} &\end{aligned}$$

$$\begin{aligned}Y_{lk}(\theta,\varphi) &= \sqrt{\frac{2l+1}{4\pi}}\langle lk|e^{-i\theta J_y}|l0\rangle e^{ik\varphi} \\ C_{M;m_1m_2}^{J;j_1j_2} &:= \langle j_1m_1;j_2m_2|j_1j_2JM\rangle \\ C\neq 0 \Rightarrow M &= m_1+m_2\,; \, C\in\mathbb{R} \\ C\neq 0 \Rightarrow |j_1-j_2| &\leq J\leq j_1+j_2 \\ C_{M;m_1m_2}^{J;j_1j_2} &= (-1)^{J-j_1-j_2}C_{M;m_2m_1}^{J;j_2j_1} \\ \langle\alpha jm|T_q^{(k)}|\beta j'm'\rangle &= \langle\alpha j\|T^{(k)}\|\beta j'\rangle C_{m;m'_q}^{j;j'k} \\ J_x^{(1)} &= \frac{1}{\sqrt{2}}\begin{pmatrix} 1 & 1 \\ 1 & 1 \end{pmatrix}\end{aligned}$$

$$\begin{aligned}\langle l+1,m|\cos\theta|lm\rangle &= \sqrt{\frac{(l+1)^2-m^2}{(2l+1)(2l+3)}} \\ \sum_{n=0}^\infty H_n(x)\frac{t^n}{n!} &= e^{-t^2+2tx} \\ H_n(x) &= e^{\frac{x^2}{2}}\Big(x-\frac{\mathrm{d}}{\mathrm{d}x}\Big)^ne^{-\frac{x^2}{2}} \\ H_n(x) &= (-1)^ne^{x^2}\frac{\mathrm{d}^n}{\mathrm{d}x^n}e^{-x^2} \\ H_n(-x) &= (-1)^nH_n(x) \\ n\text{ even: } H_n(0) &= (-1)^{\frac{n}{2}}\frac{n!}{(n/2)!} \\ H'_n(x) &= 2nH_{n-1}(x)\,; \, H_0 = 1 \\ H_1 &= 2x\,; \, H_2 = 4x^2-2\,; \, H_3 = 8x^3-12x\end{aligned}$$

$$\begin{aligned} H_{n+1}(x) &= 2xH_n(x) - 2nH_{n-1}(x) \\ H_n''(x) &= 2xH_n'(x) - 2nH_n(x) \\ \int_{-\infty}^{\infty} \mathrm{d} x H_n(x) H_m(x) e^{-x^2} &= \sqrt{\pi} 2^n n! \delta_{nm} \\ A_H(t) &= A e^{-i\omega t} \\ A\left|\alpha\right\rangle &= \alpha\left|\alpha\right\rangle, \left|\alpha\right\rangle = \sum_{n=0}^{\infty} \frac{\alpha^n}{\sqrt{n!}} e^{-\frac{1}{2}|\alpha|^2} \left|n\right\rangle \\ D(\alpha) &:= e^{\alpha A^{\dagger} - \alpha^* A}; \; D(\alpha)\left|0\right\rangle = \left|\alpha\right\rangle \\ \text{DELTA: } \mathcal{H} &= \frac{P^2}{2m} - \lambda \delta(x), \; \lambda > 0 \\ x_0 &:= \frac{\hbar^2}{\lambda m}; \; \beta := \frac{m\lambda}{\hbar^2}; \; k^2 := \frac{2mE}{\hbar^2} \\ \psi_{\text{bounded}}(x) &= \frac{1}{\sqrt{x_0}} e^{-\frac{|x|}{x_0}}, \; E_{\text{bounded}} = -\frac{\lambda}{2x_0} \\ R &= \frac{1}{-1+i\frac{k}{\beta}}; \; T = \frac{1}{1+i\frac{\beta}{k}} = R+1 \end{aligned}$$

$$\begin{aligned} \text{STEP: } \mathcal{H} &= \frac{P^2}{2m} + \begin{cases} 0 & x < 0 \\ V_0 > 0 & x > 0 \end{cases} \\ k^2 &:= \frac{2mE}{\hbar^2}, \; q^2 := \frac{2m(E-V_0)}{\hbar^2} \\ \psi_{\text{right}}(x) &\propto \begin{cases} e^{ikx} + \frac{k-q}{k+q} e^{-ikx} & x < 0 \\ \frac{2k}{k+q} e^{iqx} & x > 0 \end{cases} \end{aligned}$$

Particle physics

$$\begin{aligned} M(A,Z) &= Zm_{\rm p} + (A-Z)m_{\rm n} - B(A,Z) \\ B(A,Z) &= a_v A - a_s A^{2/3} - a_c \frac{Z(Z-1)}{A^{1/3}} - a_{\rm sym} \frac{(A-2Z)^2}{A} + a_p A^{-3/4} \Delta \\ \Delta &= \begin{cases} 0 & A \text{ odd} \\ \begin{matrix} 1 & Z \text{ even} \\ -1 & Z \text{ odd} \end{matrix} & A \text{ even} \end{cases} \\ a_v &= 15.5; \; a_s = 16.8; \; a_c = 0.72; \; a_{\rm sym} = 23; \; a_p = 34 \text{ [MeV]} \\ \frac{\partial M}{\partial Z} &= 0 : Z = \frac{m_{\rm n} - m_{\rm p} + 4a_{\rm sym}}{\frac{2a_c}{A^{1/3}} + \frac{8a_{\rm sym}}{A}}; \; r_{\rm nuc} \approx 1.5 \sqrt[3]{A} \text{ fm} \\ s_{ab} &:= (p_a + p_b)^2; \; M \rightarrow abc : (m_a + m_b)^2 \leq s_{ab} \leq (M - m_c)^2 \end{aligned}$$

QFT fields

$$\begin{aligned} \varphi(x) &= \int \frac{\mathrm{d}^3p}{(2\pi)^3} \frac{1}{\sqrt{2p^0}} \left(a(\vec{p}) e^{-ip\cdot x} + a^\dagger(\vec{p}) e^{ip\cdot x} \right) \\ \phi(x) &= \int \frac{\mathrm{d}^3p}{(2\pi)^3} \frac{1}{\sqrt{2p^0}} \left(a_+(\vec{p}) e^{-ip\cdot x} + a_-^\dagger(\vec{p}) e^{ip\cdot x} \right) \\ \phi^*(x) &= \int \frac{\mathrm{d}^3p}{(2\pi)^3} \frac{1}{\sqrt{2p^0}} \left(a_+^\dagger(\vec{p}) e^{ip\cdot x} + a_-(\vec{p}) e^{-ip\cdot x} \right) \end{aligned}$$

QFT ($\hbar=c=1$)

$$\begin{aligned} -\sigma^2 \vec{\sigma} \sigma^2 &= \vec{\sigma}^\top \\ -\sigma^2 \vec{\sigma}^* \sigma^2 &= \vec{\sigma} \\ [J^{\mu\nu}, J^{\rho\sigma}] &= i(\eta^{\mu\sigma} J^{\nu\rho} + \eta^{\nu\rho} J^{\mu\sigma} - \eta^{\mu\rho} J^{\nu\sigma} - \eta^{\nu\sigma} J^{\mu\rho}) \\ J^i &:= \tfrac{1}{2} \varepsilon^{ijk} J^{jk} \\ K^i &:= \tfrac{1}{2} J^{0i} \\ \vec{J}_\pm &:= \tfrac{1}{2} (\vec{J} \pm i \vec{K}) \\ [J^i, J^j] &= i \varepsilon^{ijk} J^k \\ J^{ij} &= \varepsilon^{ijk} J^k \\ [K^i, K^j] &= -i J^{ij} \\ [J^i, K^j] &= i \varepsilon^{ijk} K^k \\ [J_+, J_-] &= 0 \\ [J_\pm^i, J_\pm^j] &= i \varepsilon^{ijk} J_\pm^k \\ \text{fermions: } \{a, a^\dagger\} &= 1; \{a, a\} = \{a^\dagger, a^\dagger\} = 0 \\ b_\alpha &:= \eta_\alpha a_\alpha \\ \eta_\alpha &:= \prod_{\beta=1}^{\alpha-1} \begin{pmatrix} -1 & \\ & 1 \end{pmatrix}_\beta \end{aligned}$$

$$\begin{aligned} \text{HYDROGEN: } \mathcal{H} &= \frac{\vec{p}^2}{2M} - \frac{e^2}{X} \\ a := r_B &:= \frac{\hbar^2}{Me^2}; \; \text{Rydberg} := \frac{e^2}{2a} \\ E_n &= -\frac{1}{n^2} \frac{e^2}{2a}; \; \text{degen.} = n^2 \\ \psi_{nlm} &= R_{nl} Y_{lm}; \; \vec{j} = \frac{\hbar}{M} \hat{\varphi} \frac{m}{r \sin \theta} |\psi|^2 \\ R_{nl} &= 2 \sqrt{\frac{(n-l-1)!}{a^3 n^4 (n+l)!}} e^{-\frac{r}{na}} \left(\frac{2r}{na}\right)^l L_{n+l}^{2l+1} \left(\frac{2r}{na}\right) \\ L_n^{(j)}(x) &= \sum_{m=0}^{n-j} (-1)^m \binom{n}{n-j-m} \frac{x^m}{m!} \\ L_k(x) &= e^x \frac{\mathrm{d}^k}{\mathrm{d} x^k} (x^k e^{-x}) \\ L_k^{(j)} &= (-1)^j \frac{\mathrm{d}^j}{\mathrm{d} x^j} L_k(x) \\ R_{1s} &= 2 a^{-\frac{3}{2}} e^{-\frac{r}{a}} \\ R_{2s} &= \frac{1}{\sqrt{2}} a^{-\frac{3}{2}} \left(1 - \frac{r}{2a}\right) e^{-\frac{r}{2a}} \\ R_{2p} &= \frac{1}{2\sqrt{6}} a^{-\frac{3}{2}} \frac{r}{a} e^{-\frac{r}{2a}} \\ \Delta \mathcal{H}_{\text{f.s.}} &= -\frac{1}{8} \frac{P^4}{m^3 c^2} + \pi K \delta(\vec{X}) + (g-1) K \frac{\vec{L} \vec{S}}{X^3} \\ K &:= \frac{Ze^2 \hbar^2}{2m^2 c^2}; \; \vec{J} = \vec{L} + \vec{S} \\ \Delta E_{\text{f.s.}}(n,l,j) &= -Z^4 \frac{\alpha^2}{2n^3} \left(\frac{1}{j+\frac{1}{2}} - \frac{3}{4n}\right) \end{aligned}$$

$$\begin{aligned} \Delta \mathcal{H}_{\text{h.s.}} &= C \vec{I} \vec{J} \\ C_{nlj} &= 2 g_N \mu_N \mu_B \begin{cases} \frac{8\pi}{3} |\psi(0)|^2 & l=0 \\ \langle \frac{1}{R^3} \rangle \frac{l(l+1)}{j(j+1)} & l>0 \end{cases} \\ \Delta E_{jm}^{\text{Zeeman}} &= \mu_B B m \Big(1 + (g-1) \cdot \frac{j(j+1) + s(s+1) - l(l+1)}{2j(j+1)}\Big) \\ \vec{\omega}_{\text{Thomas}} &= \frac{\vec{a} \times \vec{v}}{2c^2} \\ \text{RUTHERFORD: } V(r) &= \frac{zZe^2}{r} \\ \frac{\mathrm{d}\sigma}{\mathrm{d}\Omega} &= \left(\frac{zZe^2}{2mv^2}\right)^2 \left(\sin^4(\theta/2)\right)^{-1} \\ \text{YUKAWA: } V(r) &= \frac{\alpha}{r} e^{-\mu r} \\ \frac{\mathrm{d}\sigma}{\mathrm{d}\Omega} &= \left(\frac{2m\alpha}{\hbar^2}\right)^2 (|\vec{k}' - \vec{k}|^2 + \mu^2)^{-2} \\ q &:= \vec{k}' - \vec{k}; \; \frac{\mathrm{d}\sigma}{\mathrm{d}q} = \frac{2\pi q}{k^2} \frac{\mathrm{d}\sigma}{\mathrm{d}\Omega} \\ \sigma &= 16\pi \left(\frac{m\alpha}{\hbar^2}\right)^2 \frac{1}{\mu^2} \frac{1}{4k^2 + \mu^2} \\ \text{CHARGED SHERE } e\rho(r) & \\ \left\langle \vec{k}' \left| V \right| \vec{k} \right\rangle &= \frac{4\pi}{q^2} z e^2 \int \mathrm{d}\vec{r} \, \rho(r) e^{-i\vec{q}\vec{r}} \end{aligned}$$

$$M \rightarrow abc : s_{ab} + s_{bc} + s_{ac} = M^2 + m_a^2 + m_b^2 + m_c^2$$

$$a_i A_i \rightarrow b_j B_j : Q := a_i m_{A_i} - b_j m_{B_j}$$

$$p=qBR$$

$$\frac{\mathrm{d}^3\vec{p}}{2E} = \mathrm{d}^4p \delta(p^2 - m^2) \theta(p_0)$$

$$\mathrm{d}L_p = \Big(\prod_n \frac{\mathrm{d}^3\vec{p}_n}{2E_n}\Big) \delta^4(p_{\text{in}} - \sum_n p_n); \; \mathrm{d}\sigma = f_{\text{coll}}(p_1,\dots,p_n) \mathrm{d}L_p$$

$$\text{two body: } \frac{\mathrm{d}\sigma}{\mathrm{d}\Omega_1} = f(\Omega_1) \frac{p_1}{4\sqrt{s}}; \; \sqrt{s} = \text{c.m. energy}$$

$$\text{Rutherford: } \tan\frac{\theta}{2} = \frac{1}{4\pi\epsilon_0} \frac{Qqm}{p^2b}; \; \frac{\mathrm{d}\sigma}{\mathrm{d}\Omega} = \left|\frac{b}{\sin\theta} \frac{\mathrm{d}b}{\mathrm{d}\theta}\right|; \; \frac{\mathrm{d}\sigma}{\mathrm{d}\Omega} = \frac{d_{\text{min}}^2}{16} \frac{1}{\sin^4\frac{\theta}{2}}$$

$$\begin{aligned} \frac{\mathrm{d}\sigma}{\mathrm{d}\Omega}\big|_{\text{Mott}} &= \frac{\mathrm{d}\sigma}{\mathrm{d}\Omega}\big|_{\text{Rutherford}} \cdot \cos^2\frac{\theta}{2} \\ \text{mass defect} &:= M - A \cdot \text{amu} \end{aligned}$$

$$B_\mu(x) = \sum_b \int \frac{\mathrm{d}^3p}{(2\pi)^3} \frac{1}{\sqrt{2p^0}} \left(a_b(\vec{p}) \varepsilon_\mu(\vec{p},b) e^{-ip\cdot x} + a_b^\dagger(\vec{p}) \varepsilon_\mu^*(\vec{p},b) e^{ip\cdot x} \right)$$

$$A_\mu(x) = \sum_\lambda \int \frac{\mathrm{d}^3p}{(2\pi)^3} \frac{1}{\sqrt{2|\vec{p}|}} \left(a_\lambda(\vec{p}) \varepsilon_\mu(\vec{p},\lambda) e^{-ip\cdot x} + a_\lambda^\dagger(\vec{p}) \varepsilon_\mu^*(\vec{p},\lambda) e^{ip\cdot x} \right)$$

$$\psi(x) = \sum_r \int \frac{\mathrm{d}^3p}{(2\pi)^3} \frac{1}{\sqrt{2p^0}} \left(b_r(\vec{p}) u_r(\vec{p}) e^{-ip\cdot x} + d_r^\dagger(\vec{p}) v_r(\vec{p}) e^{ip\cdot x} \right)$$

$$\bar{\psi}(x) = \sum_r \int \frac{\mathrm{d}^3p}{(2\pi)^3} \frac{1}{\sqrt{2p^0}} \left(b_r^\dagger(\vec{p}) \bar{u}_r(\vec{p}) e^{ip\cdot x} + d_r(\vec{p}) \bar{v}_r(\vec{p}) e^{-ip\cdot x} \right)$$

$$\langle \vec{x} | \vec{p} \rangle := \frac{1}{\sqrt{V}} e^{i\vec{p}\vec{x}} \; (\text{then } V \rightarrow \infty)$$

$$\langle \vec{p}_1 | \vec{p}_2 \rangle = \frac{(2\pi)^3}{V} \delta(\vec{p}_1 - \vec{p}_2)$$

$$\tilde{U}(\infty,-\infty)=T_W\exp(i\int\mathrm{d}^4x\mathcal{L}_I[\Phi](x))$$

$$S[\Phi]=:\int\mathrm{d}^4x\mathcal{L}[\Phi](x)=:\int\mathrm{d}^4x(\mathcal{L}_0[\Phi](x)+\mathcal{L}_I[\Phi](x))$$

$$\begin{aligned} T_W\Phi_1(x)\Phi_2(y) &:= \theta(x^0-y^0)\Phi_1(x)\Phi_2(y) \pm \theta(y^0-x^0)\Phi_2(y)\Phi_1(x), \\ &\quad - \text{ iff } \Phi_1 \text{ and } \Phi_2 \text{ fermions} \end{aligned}$$

$$T_W\frac{\partial\Phi_1}{\partial x}\frac{\partial\Phi_2}{\partial y}:=\frac{\partial}{\partial x}\frac{\partial}{\partial y}T_W\Phi_1(x)\Phi_2(y)$$

$$S[\varphi]=S[x\mapsto\varphi(x+\lambda)]+\mathrm{O}(x^{(\cdot)}\lambda^2)\rightarrow$$

$$\rightarrow T^{\mu\nu} = \partial^\mu \varphi \partial^\nu \varphi - \tfrac{1}{2} \eta^{\mu\nu} \mathcal{L}$$

$$\text{standard: } \gamma^0_S = \left(\begin{smallmatrix} 1 & \\ & -1 \end{smallmatrix} \right)$$

$$\text{standard: } \gamma^i_S = \left(\begin{smallmatrix} & \sigma^i \\ -\sigma^i & \end{smallmatrix} \right)$$

$$\bar{u} := u^\dagger \gamma^0 \underset{(\text{standard})}{=} \left(\begin{smallmatrix} u_1^* & u_2^* & -u_3^* & -u_4^* \end{smallmatrix} \right)$$

$$\gamma^5 = i\gamma^0\gamma^1\gamma^2\gamma^3 \underset{(\text{standard})}{=} \left(\begin{smallmatrix} 1 & \\ & -1 \end{smallmatrix} \right)$$

$$\bar{a}b \text{ scalar}$$

$$\bar{a}\gamma^5b \text{ pseudoscalar}$$

$$\bar{a}\gamma^\mu b \text{ vector}$$

$$\bar{a}\gamma^\mu\gamma^5b \text{ pseudovector}$$

$$\bar{a}\sigma^{\mu\nu}b \text{ tensor}$$

$$:a^\dagger\cdots a\cdots a^\dagger::=a^\dagger a^\dagger\cdots aa+\ldots \text{ (creation to the left)}$$

$$\text{real scalar } \mathcal{L}_0 = \frac{1}{2}(\partial_\mu\varphi\partial^\mu\varphi - m^2\varphi^2), \ (\partial^\mu\partial_\mu + m^2)\varphi = 0$$

$$\text{complex scalar } \mathcal{L}_0 = \partial_\mu\phi^*\partial^\mu\phi - m^2\phi^*\phi, \ (\partial_\mu\partial^\mu + m^2)\left(\begin{smallmatrix}\phi\\ \phi^*\end{smallmatrix}\right) = 0$$

$$\text{real vector } \mathcal{L}_0 = -\frac{1}{4}F^{\mu\nu}F_{\mu\nu} + \frac{m^2}{2}B^\mu B_\mu, \ F^{\mu\nu} := \partial^\mu B^\nu - \partial^\nu B^\mu, \\ (\partial_\mu\partial^\mu + m^2)B^\nu = 0, \ \partial_\mu B^\mu = 0$$

$$\text{Dirac spinor } \mathcal{L}_0 = \bar{\psi}(i\gamma^\mu\partial_\mu - m)\psi, \ (i\gamma^\mu\partial_\mu - m)\psi = 0$$

$$\{\gamma^\mu,\gamma^\nu\}=2\eta^{\mu\nu}$$

$$\Lambda(\omega)=e^{\frac{i}{2}\omega_{\mu\nu}J^{\mu\nu}}$$

$$R(\omega)=e^{\frac{i}{4}\omega_{\mu\nu}\Sigma^{\mu\nu}}$$

$$R^{-1}(\omega)\gamma^\mu R(\omega)=\lambda^\mu_\nu(\omega)\gamma^\nu$$

$$R^{-1}(\omega)=\gamma^0R^\dagger(\omega)\gamma^0$$

$$\bar{u}_r(\vec{p})u_s(\vec{p})=-\bar{v}_r(\vec{p})v_s(\vec{p})=2m\delta_{rs}$$

$$u_r^\dagger(\vec{p})u_s(\vec{p})=v_r^\dagger(\vec{p})v_s(\vec{p})=2E(\vec{p})\delta_{rs}$$

$$\bar{v}_r(\vec{p})u_s(\vec{p})=\bar{u}_r(\vec{p})v_s(\vec{p})=0$$

$$v_r^\dagger(\vec{p})u_s(-\vec{p})=u_r^\dagger(\vec{p})v_s(-\vec{p})=0$$

$$u_r^\alpha(\vec{p})\bar{u}_r^\beta(\vec{p})-v_r^\alpha(\vec{p})\bar{v}_r^\beta(\vec{p})=2m\delta^{\alpha\beta}$$

$$u_r^\alpha(\vec{p})\bar{u}_r^\beta(\vec{p})=(\not{p}+m)^{\alpha\beta}$$

$$v_r^\alpha(\vec{p})\bar{v}_r^\beta(\vec{p})=(\not{p}-m)^{\alpha\beta}$$

$$(\gamma^0)^\dagger=\gamma^0$$

$$(\gamma^i)^\dagger=-\gamma^i$$

$$\gamma^0(\gamma^i)^\dagger\gamma^0=\gamma^i$$

$$(\gamma^5)^\dagger=\gamma^5$$

$$\{\gamma^5,\gamma^\mu\}=0$$

$$[\gamma^5,\Sigma^{\mu\nu}]=0$$

$$\mathrm{tr}\,\gamma^5=0$$

$$\not{p}:=\gamma^\mu p_\mu$$

$$\{\gamma^5,\gamma^\mu\}=0$$

$$\not{p}\not{q}=pq-i\Sigma_{\mu\nu}p^\mu q^\nu$$

$$\gamma_\mu\not{p}\gamma^\mu=-2\not{p}$$

$$\gamma_\mu\not{p}\not{q}\not{k}\gamma^\mu=-2\not{k}\not{q}\not{p}$$

$$\gamma_\mu\not{p}\not{q}\gamma^\mu=4pq$$

$$\mathrm{tr}\underbrace{(\gamma^\mu\gamma^\nu\cdots\gamma^\sigma)}_{\mathrm{dispari}}=0$$

$$\mathrm{tr}(\gamma^5\not{p}\not{q})=0$$

$$\mathrm{tr}(\not{p}\not{q})=4pq$$

$$\mathrm{tr}(\gamma^5\not{p}\not{q}\not{k}\not{l})=4i\varepsilon_{\mu\nu\rho\sigma}p^\mu q^\nu k^\rho l^\sigma$$

$$\mathrm{tr}(\not{p}\not{q}\not{k}\not{l})=4((pq)(kl)-(pk)(ql)+(pl)(qk))$$

$$\text{base } \{1,\gamma^5,\gamma^\mu,\gamma^\mu\gamma^5,\Sigma^{\mu\nu}\}$$

$$E(\vec{p}):=\sqrt{m^2+\vec{p}^2}$$

$$\{b_r(\vec{p}),b_s^\dagger(\vec{k})\}=\{d_r(\vec{p}),d_s^\dagger(\vec{k})\}=\frac{(2\pi)^3}{V}\delta_{rs}\delta^3(\vec{p}-\vec{k})$$

$$\{a(\vec{p}),a^\dagger(\vec{k})\}=\frac{(2\pi)^3}{V}\delta^3(\vec{p}-\vec{k})$$

$$\theta^{\mu\nu}:=i\bar{\psi}\gamma^\mu\partial^\nu\psi$$

$$P^\mu:=\int\mathrm{d}^3x\theta^{0\nu}=\int\mathrm{d}^3x\psi^\dagger i\partial^\nu\psi$$

$$\theta_a:=\frac{1}{2}\varepsilon_{abc}\omega^{bc}$$

$$\eta_a:=\omega^{0a}$$

$$\Lambda(\omega)=e^{i(\theta_aJ^a+\eta_aK^a)}$$

$$J^aK^a=0$$

$$J^aJ^a-K^aK^a=3$$

$$J_R=J_-$$

$$J_L=J_+$$

$$\psi_R=\frac{1+\gamma^5}{2}\psi\underset{\text{(standard)}}{=}\begin{pmatrix}\xi_R\\ \xi_R\end{pmatrix},\ \xi_R=\frac{1}{2}\begin{pmatrix}\psi_1+\psi_3\\ \psi_2+\psi_4\end{pmatrix}$$

$$\psi_L=\frac{1-\gamma^5}{2}\psi\underset{\text{(standard)}}{=}\begin{pmatrix}\xi_L\\ -\xi_L\end{pmatrix},\ \xi_L=\frac{1}{2}\begin{pmatrix}\psi_1-\psi_3\\ \psi_2-\psi_4\end{pmatrix}$$

$$\Sigma_S^{0a}=i\left(\begin{smallmatrix} & \sigma^a \\ \sigma^a & \end{smallmatrix}\right)$$

$$\Sigma_S^{ab}=\varepsilon^{abc}\left(\begin{smallmatrix} \sigma^c & \\ & \sigma^c \end{smallmatrix}\right)$$

$$\frac{i}{4}\omega_{\mu\nu}\Sigma^{\mu\nu}=\frac{1}{2}\left(\begin{smallmatrix} i\theta_a\sigma^a & -\eta_a\sigma^a \\ -\eta_a\sigma^a & i\theta_a\sigma^a \end{smallmatrix}\right)$$

$$\text{chiral: } \gamma_C^0 = \left(\begin{smallmatrix} & 1 \\ 1 & \end{smallmatrix}\right)$$

$$\text{chiral: } \gamma_C^i = \left(\begin{smallmatrix} & -\sigma^i \\ \sigma^i & \end{smallmatrix}\right)$$

$$\gamma_C^5 = \left(\begin{smallmatrix} 1 & \\ & -1 \end{smallmatrix}\right)$$

$$\gamma_C^0\gamma_C^\mu=\left(\begin{smallmatrix} \sigma^\mu & \\ & \sigma_\mu \end{smallmatrix}\right)$$

$$\frac{1}{2}(\sigma_\mu\sigma^\nu+\sigma_\nu\sigma^\mu)=\frac{1}{2}(\sigma^\mu\sigma_\nu+\sigma^\nu\sigma_\mu)=\eta^{\mu\nu}$$

$$\sigma_\mu\sigma^\nu\sigma^\mu=-2\sigma_\nu$$

$$\sigma^\mu\sigma_\nu\sigma^\mu=-2\sigma^\nu$$

$$\sigma^{\mu\nu}:=\frac{i}{2}(\sigma_\mu\sigma^\nu-\sigma_\nu\sigma^\mu)$$

$$\bar{\sigma}^{\mu\nu}:=\frac{i}{2}(\sigma^\mu\sigma_\nu-\sigma^\nu\sigma_\mu)$$

$$\sigma^{0j}=i\sigma^j$$

$$\bar{\sigma}^{0j}=-i\sigma^j$$

$$\sigma^{ij}=\bar{\sigma}^{ij}=\varepsilon^{ijk}\sigma^k$$

$$\bar{\sigma}^{\mu\nu}=-\frac{1}{2}\varepsilon^{\mu\nu\rho\sigma}\sigma_{\rho\sigma}$$

$$\mathrm{tr}(\sigma_\mu\sigma^\nu)=2\eta^{\mu\nu}$$

$$\mathrm{tr}(\sigma_\mu\sigma^\nu\sigma_\rho\sigma^\sigma)=2(\eta^{\mu\nu}\eta^{\rho\sigma}-\eta^{\mu\rho}\eta^{\nu\sigma}+\eta^{\mu\sigma}\eta^{\nu\rho}-i\varepsilon^{\mu\nu\rho\sigma})$$

$$\sigma_C^{\mu\nu}=\left(\begin{smallmatrix} \sigma^{\mu\nu} & \\ & \bar{\sigma}^{\mu\nu} \end{smallmatrix}\right)$$

$$\gamma_C^\mu=W^{-1}\gamma_S^\mu W$$

$$W=\frac{1}{\sqrt{2}}\left(\begin{smallmatrix} 1 & 1 \\ 1 & -1 \end{smallmatrix}\right)$$

$$\text{parity: } \psi(x) \stackrel{P}{\mapsto} \gamma^0\psi(Px)$$

$$\text{charge: } \psi \stackrel{C}{\mapsto} i\gamma^2\psi^\dagger$$

$$\bar{\psi} \stackrel{C}{\mapsto} -\psi(i\gamma^2)^{-1}$$

$$\text{time: } \psi(x) \stackrel{T}{\mapsto} i\gamma^1\gamma^3\psi(Tx)$$

$$\tilde{U}(\infty,-\infty)=T_W e^{iS_I}=:e^{i\theta_0}\mathbb{S}=:e^{i\theta_0}(1+i\mathbb{T})$$

$$A(\{\vec{k}_{\mathrm{out}}\},\{\vec{k}_{\mathrm{in}}\})=(2\pi)^4\delta^4(\sum k_{\mathrm{out}}-\sum k_{\mathrm{in}})\cdot$$

$$\cdot \prod \frac{1}{\sqrt{2E(\vec{k}_{\mathrm{out}})}} \prod \frac{1}{\sqrt{2E(\vec{k}_{\mathrm{in}})}} V^{-(\#\mathrm{in}+\#\mathrm{out})/2} \mathcal{M}(\{\vec{k}_{\mathrm{out}}\},\{\vec{k}_{\mathrm{in}}\})$$

$$\frac{\mathrm{d}P}{\mathrm{d}t}=\frac{(2\pi)^4\delta^4(\sum k_{\mathrm{out}}-\sum k_{\mathrm{in}})}{V^{\#\mathrm{in}-1}\prod 2E(k_{\mathrm{in}})}|\mathcal{M}|^2\prod\frac{\mathrm{d}^3k_{\mathrm{out}}}{(2\pi)^32E(k_{\mathrm{out}})}$$

$$\mathrm{d}\sigma_{12\rightarrow\mathrm{out}}=\frac{1}{4\sqrt{(k_1k_2)^2-m_1^2m_2^2}}\prod\frac{\mathrm{d}^3k_{\mathrm{out}}}{(2\pi)^32E(k_{\mathrm{out}})}\cdot$$

$$\cdot |\mathcal{M}|^2 (2\pi)^4 \delta^4(k_1+k_2-\sum k_{\mathrm{out}})$$

$$\overline{\prod_\mu \gamma^\mu}:=\gamma^0(\prod_\mu \gamma^\mu)^\dagger\gamma^0=\prod_{\text{rev. }\mu} \bar{\gamma}^\mu$$

$$\bar{\gamma}^\mu=\gamma^\mu$$

$$\bar{\gamma}^5=-\gamma^5$$

$$\sum_{rs} \left| \bar{u}_s(\vec{k}) \prod_\mu \gamma^\mu v_s(\vec{p}) \right|^2 = \mathrm{tr} \left((\not{k} + m) \prod_\mu \gamma^\mu (\not{p} - m) \overline{\prod_\mu \gamma^\mu} \right)$$