

Disputation i Teknisk Mekanik

Fredag 2020-03-27, kl 10.00

Respondent: Zhouyang Ge

Titel: Droppinteraktioner och suspensionsflöden

Handledare: Prof. Luca Brandt

Asso. Prof. Outi Tammisola

Fakultetsopponent: Asso. Prof. François Gallaire, École polytechnique fédérale de Lausanne, Schweiz

Betygsnämnd: Dr. Elisabeth Lemaire, Institut de Physique de Nice, Frankrike

Dr. Martin Trulsson, Lund Universitet, Sverige

Dr. Gustaf Mårtensson, Mycronic AB, Sverige

Ordförande: Prof. Fredrik Lundell

Sponsorer: EU Horizon 2020 (MICROFLUSA), VR

Procedure

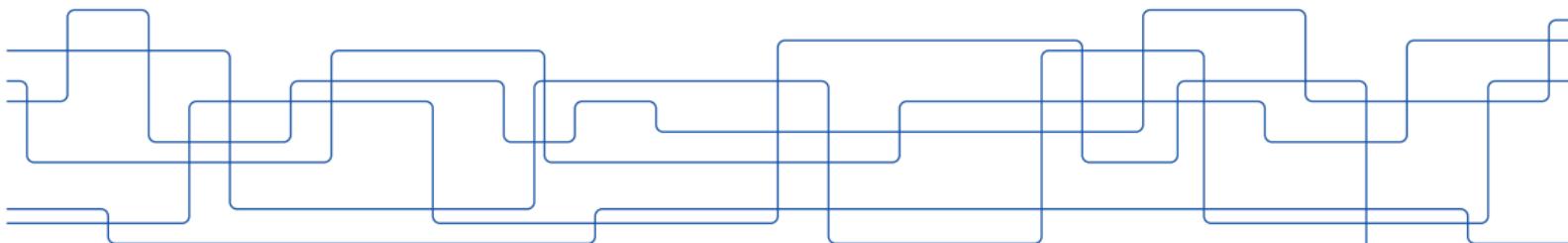
- ▶ The respondent will present the thesis
- ▶ The opponent will discuss the thesis
- ▶ The grading committee will ask questions
- ▶ The audience may ask questions
- ▶ The public part of the defence will be closed
- ▶ The result will be announced at Osquars Backe 18, floor 6



On Droplet Interactions and Suspension Flow

Zhouyang Ge

*Department of Engineering Mechanics,
KTH Royal Institute of Technology, Stockholm, Sweden*



But what are *droplets*?



Droplets are micron to millemetre sized liquid balls formed under surface tension.*

*Photo by Martin Brechtl on Unsplash.

†Source: <https://youtu.be/P1ww1IXRfTA>.

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BBC interview of Richard Feynman (1983).†

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Let's have some fun!

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Experiments, strategy and questions

A simple q2D dipolar model

3D numerical simulations

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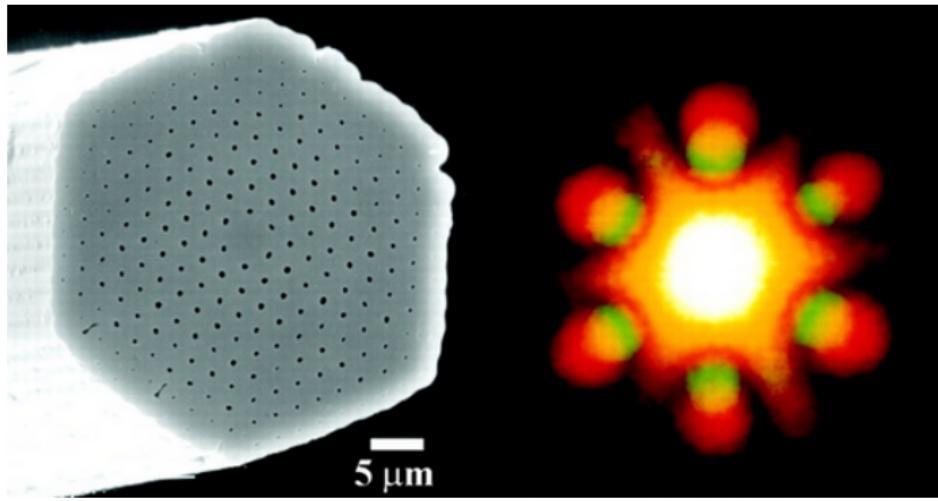
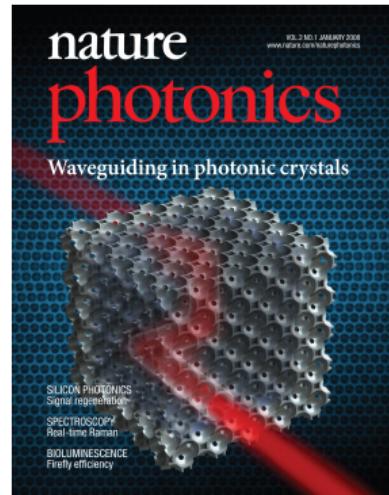


Figure 1: (a) A scanning electron micrograph of a solid-core photonic crystal fibre and its far-field optical pattern.*
© AAAS Science. (b) Cover photo of the January 2008 issue of Nature Photonics, showing an artistic rendering of a light beam passing through a 3D photonic crystal. © Springer Nature.



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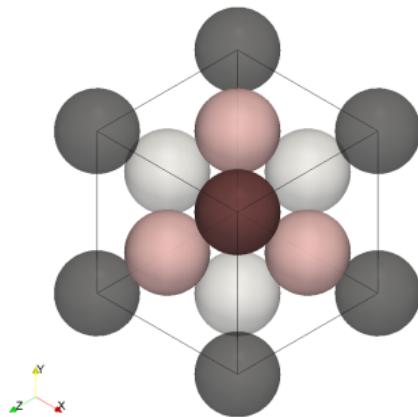
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Despite the theoretical promise, fabricating photonic crystals with 3D bandgaps is challenging in practice.

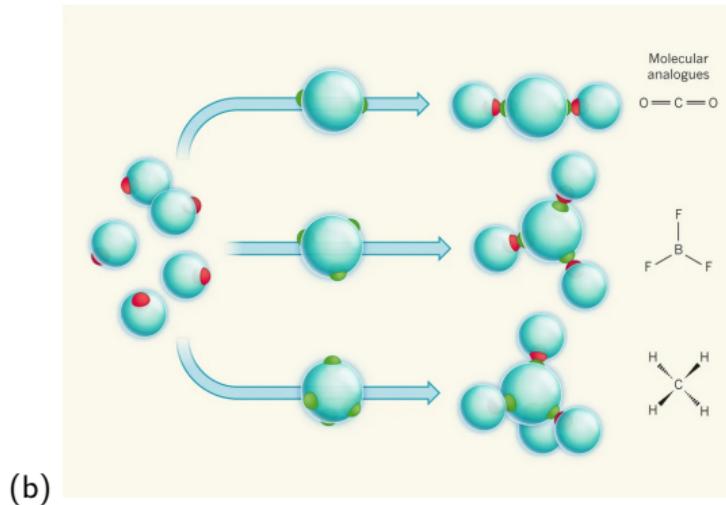
*Y. Wang *et al.* Nature 491, 51–55 (2012).

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(a)



(b)

Figure 2: (a) A diamond cubic (view angle 111). (b) Self-assembly of micron-sized patchy particles, using single-stranded DNAs as selective “glues” to bind to nearby colloids.* The directional bindings mimic the arrangements of bonds around atoms, and the obtained cluster may serve as basic **building blocks** for more complex structures. However, the process relies on **Brownian motion** to bring complementary particles into their range of attraction. Therefore, even if they can self-assemble into a diamond lattice, the rate of production will be **too slow** to meet practical needs. © Springer Nature.

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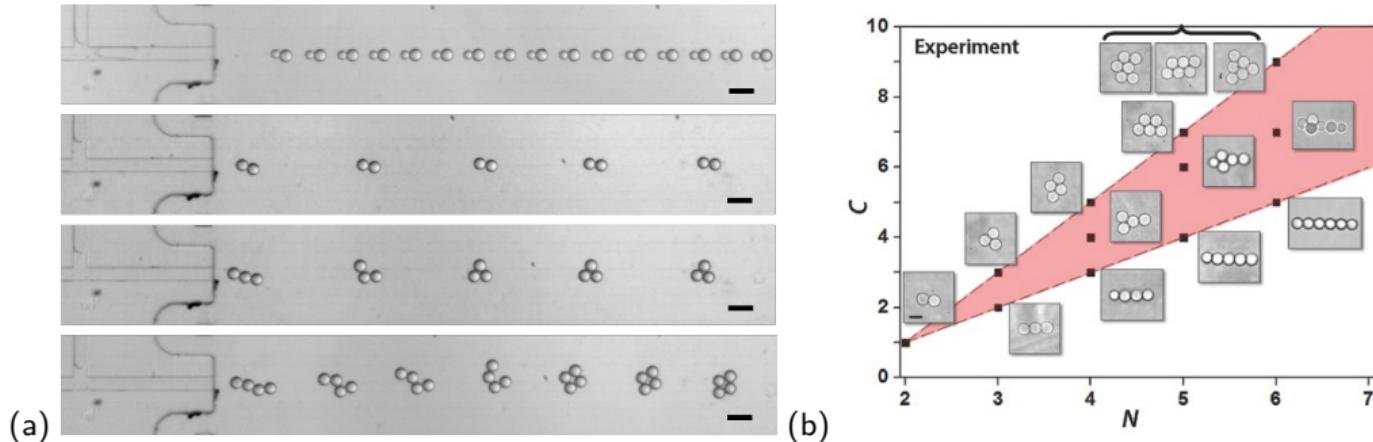


Figure 3: (a) Self-assembly of two to four droplets in a **microfluidic** chip; scale bars, $100 \mu\text{m}$. (b) Observed cluster morphologies mapped on a $C - N$ diagram, where C represents the number of droplet-droplet contacts and N is the number of droplets per cluster.* Scale bar is $5 \mu\text{m}$. Images courtesy of Dr. Bingqing Shen.

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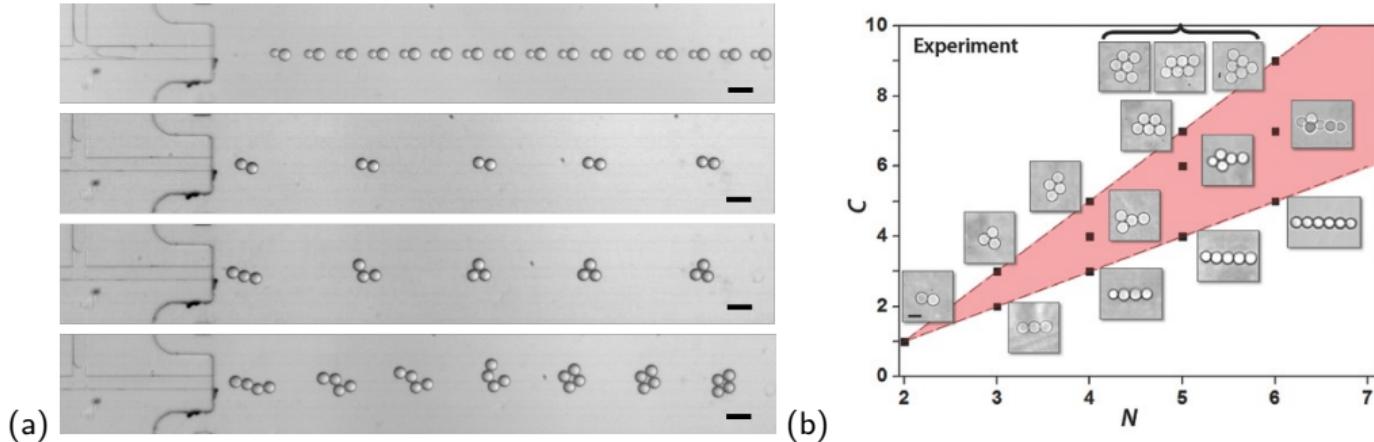


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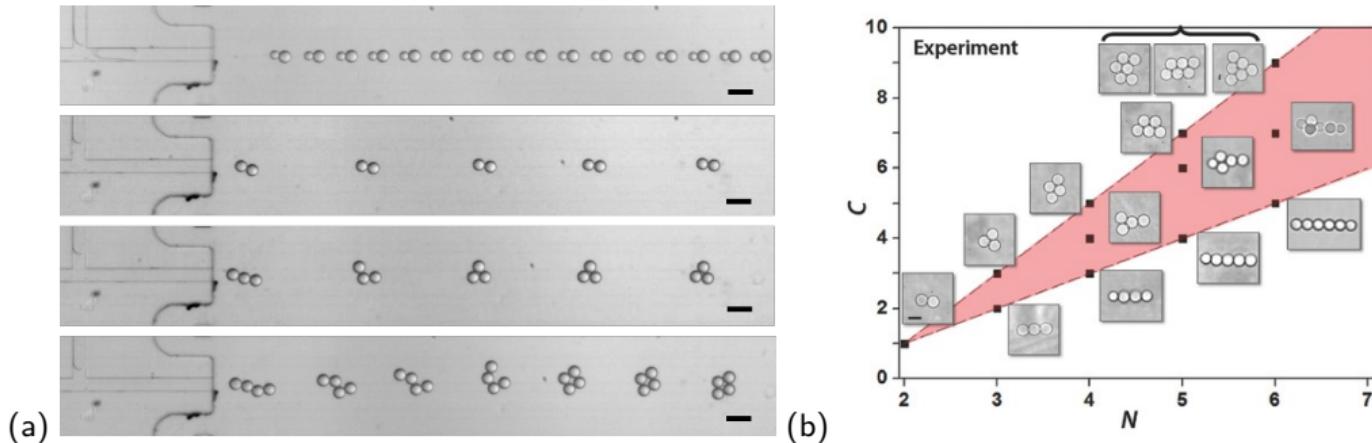


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The questions are, why do they organize into such patterns? What is the hydrodynamic mechanism?
Can we further optimize it?

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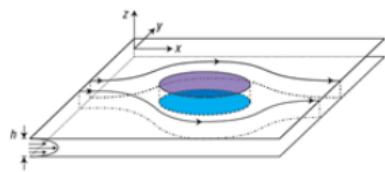
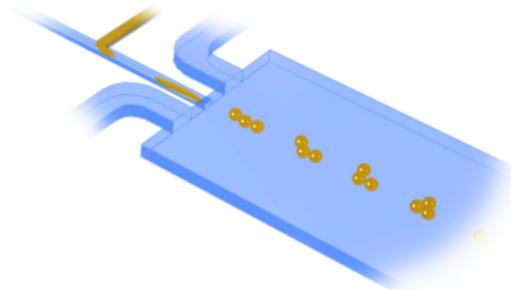
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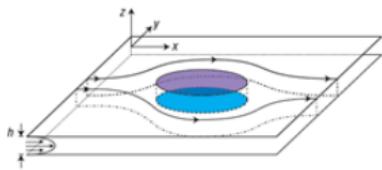
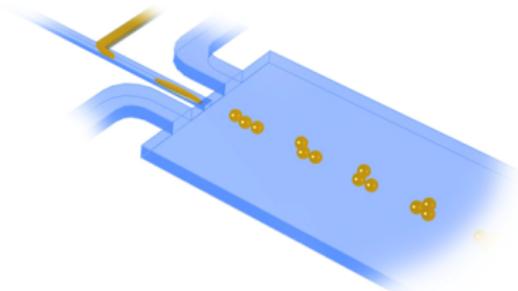
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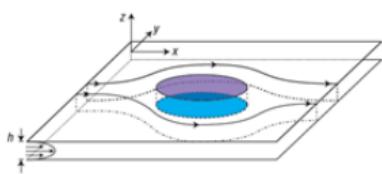
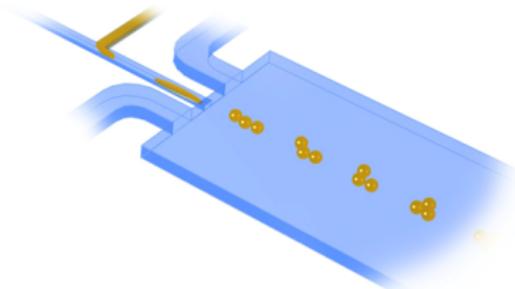


Obersvations & hypotheses

- ▶ Droplets are in quasi-two-dimensional (q2D) space.
- ▶ Depletion force attracts nearby drops.
- ▶ Hydrodynamic interactions (HI) lead to the rearrangement dynamics.

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Equations of motion (due to HI)

$$\mathbf{u}^\infty(x, y, z) \approx -\frac{z(H-z)}{2\mu} \nabla p(x, y),$$

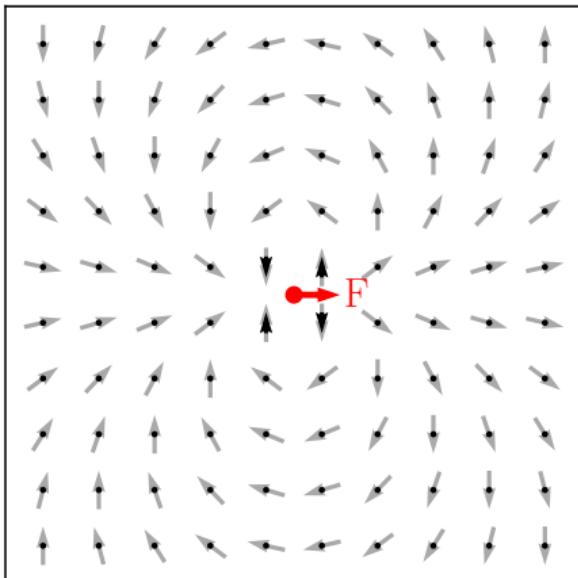
$$\delta\mathbf{u}_{ij} = \mathbf{B}_{ij}\mathbf{F}_j, \quad \mathbf{U}_i = \mathbf{u}_i^\infty + \sum_{j \neq i} \delta\mathbf{u}_{ij},$$

where

$$\mathbf{B}(\mathbf{x}) \approx -\frac{\alpha H}{\mu} \left(\frac{\mathbf{I}}{|\mathbf{x}|^2} - \frac{2\mathbf{x}\mathbf{x}}{|\mathbf{x}|^4} \right).$$

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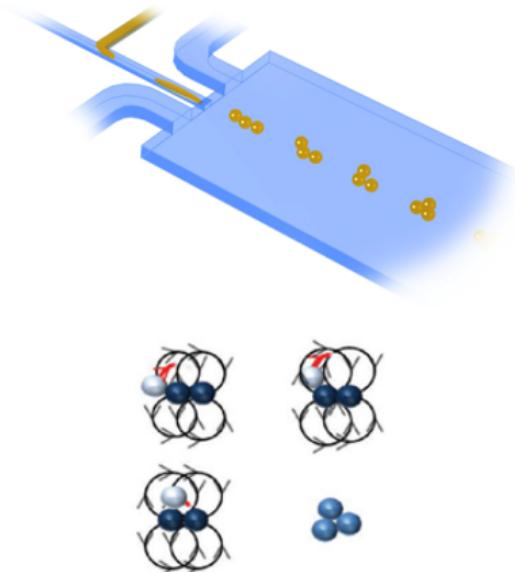
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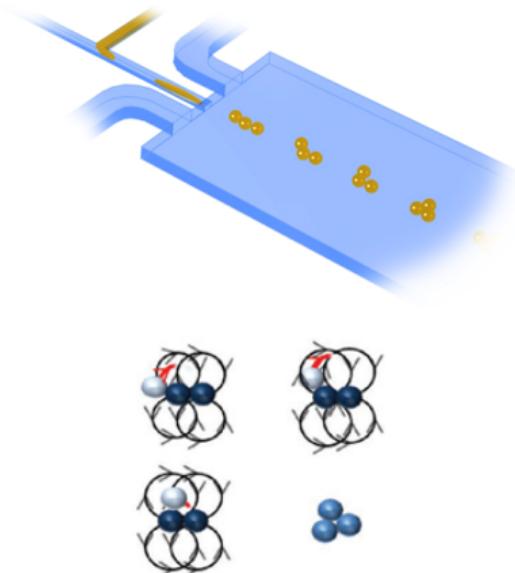
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Is this picture really correct?

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Mathematical formulations

The incompressible Navier-Stokes (NS) equations read,

$$\nabla \cdot \mathbf{u} = 0, \quad (1a)$$

$$\rho \left(\frac{\partial \mathbf{u}}{\partial t} + \mathbf{u} \cdot \nabla \mathbf{u} \right) = \nabla \cdot \boldsymbol{\sigma} + \mathbf{f}, \quad (1b)$$

where $\boldsymbol{\sigma}$ is the viscous stress tensor, given as

$$\boldsymbol{\sigma} = -p\mathbf{I} + \mu \left(\nabla \mathbf{u} + (\nabla \mathbf{u})^T \right), \quad (2)$$

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In developing the numerical solver, no such assumptions are made.

Hypotheses for cascade directions

- ▶ In Batchelor we believe.
- ▶ Yes we still do.

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