

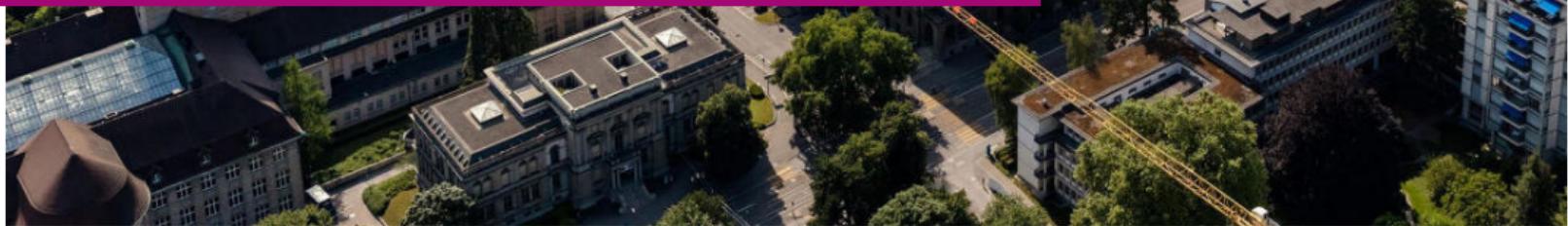


# Introducing PlesioGestroPy

a Python realization of the PG model

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Oct. 2023



# Outline

1. The PG model
2. The old implementation
3. Introducing PlesioGeostroPy
4. Future works

# PG, QG, and all the geostrophic models

Premise: the planet is rotating so fast that the Coriolis effect is dominant in the rotating frame  
( $=2\Omega \times \mathbf{u}$  is the dominant term in Navier-Stokes)

Geostrophy: pressure gradient balanced by Coriolis force  $2\Omega \times \mathbf{u} + \nabla p = \mathbf{0}$   
(this balance produces Taylor columns, i.e. invariance of velocity along the rotation axis)

Conclusion: the planet is "nearly" geostrophic  
(=the flow organized into columns)  
(=Variation  $|\frac{\partial \mathbf{u}}{\partial z}| \ll |\nabla_e \mathbf{u}|$ )

Premise somewhat confirmed by numerical computations (see e.g. right picture, from Jault and Finlay (2015)), at least outside the tangent cylinder.



# PG, QG, and all the geostrophic models

Strictly geostrophic

$$\mathbf{u} = \nabla \times (\Psi \hat{\mathbf{z}})$$

does not satisfy non-penetrating BC.

How to exploit the "near"-geostrophy?

We can choose an ansatz for the flow that is close to geostrophic flow.

Also other options

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does not satisfy divergence-free.

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Schaeffer and Cardin QG ansatz

$$\mathbf{u} = \frac{1}{H} \nabla \times (\Psi \hat{\mathbf{z}}) + \frac{z}{sH^2} \frac{dH}{ds} \frac{\partial \Psi}{\partial \phi} \hat{\mathbf{z}}$$

non-penetration + divergence-free.

# PG, QG, and all the geostrophic models

Plesio-Geostrophy (PG) (Jackson and Maffei 2020) is built on the QG ansatz by Schaeffer and Cardin. In the absence of magnetic diffusion, PG converts

Vector fields  $\mathbf{u}(\mathbf{r})$  and  $\mathbf{B}(\mathbf{r})$  living in  
3D space,  
along with their evolution equations  
(Navier-Stokes + magnetic  
induction eqs)

→ exactly to →

15 PG quantities:  $\Psi, \widetilde{B_s^2}, \widetilde{B_\phi^2}, \widetilde{B_s B_\phi},$   
 $\widetilde{B_s B_z}, \widetilde{B_\phi B_z}, \widetilde{z B_s^2}, \widetilde{z B_\phi^2}, \widetilde{z B_s B_\phi},$   
 $B_{es}, B_{e\phi}, B_{ez}, B_{es,z}, B_{e\phi,z}, B_r$   
all of which live in 2D space, along  
with their evolution equations.

# How did I end up here?

Since the start of my doctoral studies, I

- studied surface operators and tried to derive the boundary terms in the diffusive torsional oscillation (TO) equation for **2 months**  $\Rightarrow$  didn't lead anywhere;
- studied torsional Alfvén waves and calculated the 1-D eigenmodes of the torsional oscillation for **2 month**  $\Rightarrow$  didn't lead anywhere;
- studied anelastic approximation for **1 month** in order to work on the anelastic version of QuICC, which may help the Jupiter simulation  $\Rightarrow$  project was called off and deemed unpromising before I could start to think about the numerics;
- picked up the thread and studied the reflection of Alfvén waves at the fluid-solid interface for **2 months**  $\Rightarrow$  wrote a sixty-page document and yet had more questions than before;
- have been mainly working on the PG model since Aug 2023.

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# The old implementation - code *Daria*

Mathematica implementation (previous member of the group, Dr. Daria Holdenried-Chernoff)

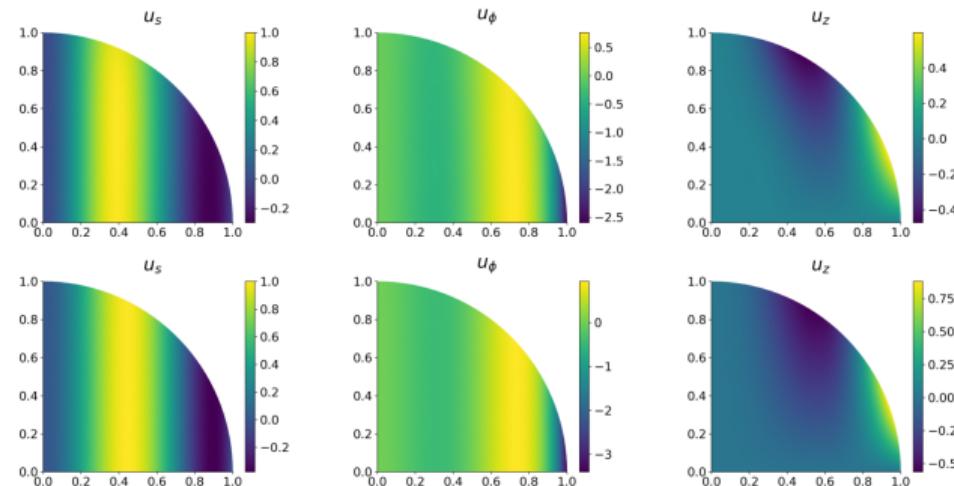


Figure: Eigenmodes calculated using 3-D code *Jiawen* (top) and PG code *Daria*.

Why reinvent the wheels when we already have a functioning implementation?

# The old implementation - code *Daria*

Citizens of Gotham, riddle me this... what is this code snippet doing?

```
In[=]= Les =  $\frac{1}{s} D[s \cdot Bs2[s, p, t], s] + \frac{s^2}{H[s]} (BsH[s, p, z, t] * BsH[s, p, z, t] + BsmH[s, p, z, t] * BsmH[s, p, z, t]) + \frac{1}{s} D[BsBp[s, p, t], p] - \frac{1}{s} Bp2[s, p, t] + BsH[s, p, z, t] * BzH[s, p, z, t] - BsmH[s, p, z, t] * BzmH[s, p, z, t];$   
  
In[=]= Lestmp =  
   $Le[s] / . Bs2[s, p, t] \rightarrow A0[s, p, t] + x * A[s, p, t] /. D[Bs2[s, p, t], s] \rightarrow D[A0[s, p, t] + x * A[s, p, t], s] /. D[Bs2[s, p, t], p] \rightarrow D[A0[s, p, t] + x * A[s, p, t], p] /.$   
   $Bp2[s, p, t] \rightarrow B0[s, p, t] + x * B[s, p, t] /. D[Bp2[s, p, t], s] \rightarrow D[B0[s, p, t] + x * B[s, p, t], s] /. D[Bp2[s, p, t], p] \rightarrow D[B0[s, p, t] + x * B[s, p, t], p] /.$   
   $BsBp[s, p, t] \rightarrow C0[s, p, t] + x * Csp[s, p, t] /. D[BsBp[s, p, t], s] \rightarrow D[C0[s, p, t] + x * Csp[s, p, t], s] /. D[BsBp[s, p, t], p] \rightarrow D[C0[s, p, t] + x * Csp[s, p, t], p] /.$   
   $BsmH[s, p, z, t] \rightarrow amh0[s, p, z, t] + x * amh[s, p, z, t] /. D[BsmH[s, p, z, t], s] \rightarrow D[amh0[s, p, z, t], s] + x * D[amh[s, p, z, t], s] /.$   
   $D[BsmH[s, p, z, t], p] \rightarrow D[amh0[s, p, z, t], p] + x * D[amh[s, p, z, t], p] /. Bsh[s, p, z, t] \rightarrow ah0[s, p, z, t] + x * ah[s, p, z, t] /.$   
   $D[Bsh[s, p, z, t], s] \rightarrow D[ah0[s, p, z, t], s] + x * D[ah[s, p, z, t], s] /. D[Bsh[s, p, z, t], p] \rightarrow D[ah0[s, p, z, t], p] + x * D[ah[s, p, z, t], p] /.$   
   $BpmH[s, p, z, t] \rightarrow bmh0[s, p, z, t] + x * bmh[s, p, z, t] /. D[BpmH[s, p, z, t], s] \rightarrow D[bmh0[s, p, z, t], s] + x * D[bmh[s, p, z, t], s] /.$   
   $D[BpmH[s, p, z, t], p] \rightarrow D[bmh0[s, p, z, t], p] + x * D[bmh[s, p, z, t], p] /. Bph[s, p, z, t] \rightarrow bh0[s, p, z, t] + x * bh[s, p, z, t] /.$   
   $D[Bph[s, p, z, t], s] \rightarrow D[bh0[s, p, z, t], s] + x * D[bh[s, p, z, t], s] /. D[Bph[s, p, z, t], p] \rightarrow D[bh0[s, p, z, t], p] + x * D[bh[s, p, z, t], p] /.$   
   $D[Bph[s, p, z, t], z] \rightarrow D[bh0[s, p, z, t], z] + x * D[bh[s, p, z, t], z] /. BzmH[s, p, z, t] \rightarrow cmh0[s, p, z, t] + x * cmh[s, p, z, t] /. BzH[s, p, z, t] \rightarrow ch0[s, p, z, t] + x * ch[s, p, z, t] /.$   
   $D[BzH[s, p, z, t], z] \rightarrow D[ch0[s, p, z, t], z] + x * D[ch[s, p, z, t], z] /. D[BzH[s, p, z, t], s] \rightarrow D[ch0[s, p, z, t], s] + x * D[ch[s, p, z, t], s] /.$   
   $D[BzmH[s, p, z, t], s] \rightarrow D[cmh0[s, p, z, t], s] + x * D[cmh[s, p, z, t], s];$ 
```

# The old implementation - code *Daria*

Citizens of Gotham, riddle me this... what is this code snippet doing?

```
In[=]= Les =  $\frac{1}{s} D[s \cdot Bs2[s, p, t], s] + \frac{s^2}{H[s]} (BsH[s, p, z, t] * BsH[s, p, z, t] + BsmH[s, p, z, t] * BsmH[s, p, z, t]) + \frac{1}{s} D[BsBp[s, p, t], p] - \frac{1}{s} Bp2[s, p, t] + BsH[s, p, z, t] * BzH[s, p, z, t] - BsmH[s, p, z, t] * BzmH[s, p, z, t];$   
  
In[=]= Lestmp =  
  Les /. Bs2[s, p, t] → A0[s, p, t] + x*A[s, p, t] /. D[Bs2[s, p, t], s] → D[A0[s, p, t] + x*A[s, p, t], s] /. D[Bs2[s, p, t], p] → D[A0[s, p, t] + x*A[s, p, t], p] /.  
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  BsBp[s, p, t] → C0[s, p, t] + x*Csp[s, p, t] /. D[BsBp[s, p, t], s] → D[C0[s, p, t] + x*Csp[s, p, t], s] /. D[BsBp[s, p, t], p] → D[C0[s, p, t] + x*Csp[s, p, t], p] /.  
  BsmH[s, p, z, t] → amh0[s, p, z, t] + x*amh[s, p, z, t] /. D[BsmH[s, p, z, t], s] → D[amh0[s, p, z, t], s] + x*D[amh[s, p, z, t], s] /.  
  D[BsmH[s, p, z, t], p] → D[amh0[s, p, z, t], p] + x*D[amh[s, p, z, t], p] /. Bsh[s, p, z, t] → ah0[s, p, z, t] + x*ah[s, p, z, t] /.  
  D[Bsh[s, p, z, t], s] → D[ah0[s, p, z, t], s] + x*D[ah[s, p, z, t], s] /. D[Bsh[s, p, z, t], p] → D[ah0[s, p, z, t], p] + x*D[ah[s, p, z, t], p] /.  
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  D[Bph[s, p, z, t], s] → D[bh0[s, p, z, t], s] + x*D[bh[s, p, z, t], s] /. D[Bph[s, p, z, t], p] → D[bh0[s, p, z, t], p] + x*D[bh[s, p, z, t], p] /.  
  D[Bph[s, p, z, t], z] → D[bh0[s, p, z, t], z] + x*D[bh[s, p, z, t], z] /. BzmH[s, p, z, t] → cmh0[s, p, z, t] + x*cmh[s, p, z, t] /. BzH[s, p, z, t] → ch0[s, p, z, t] + x*ch[s, p, z, t] /.  
  D[BzH[s, p, z, t], z] → D[ch0[s, p, z, t], z] + x*D[ch[s, p, z, t], z] /. D[BzH[s, p, z, t], s] → D[ch0[s, p, z, t], s] + x*D[ch[s, p, z, t], s] /.  
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```

This rewrites the magnetic quantities in Lorentz force  $\overline{L_\phi}$  as background + perturbation for linearization.  
e.g.  $Bp2[s, p, t] \rightarrow B0[s, p, t] + x*B[s, p, t]$  equals to  $\overline{B_\phi^2} = \overline{B_{\phi 0}^2} + \epsilon \overline{B_\phi^2}'$

## The old implementation - code *Daria*

Why is the code undesirable for **developers** and users alike?

- long, repetitive operations  $\Rightarrow$  hard to debug
  - namings: the background and perturbation fields are named from A to H  $\Rightarrow$  hard to debug or invoke.
  - At some point I even need a correspondence table to decipher the code.



# The old implementation - code *Daria*

```
In[1]:= psl[m_, n_] := (1 - s^2)^3/2 s^n JacobiP[n - 1, 3/2, m, 2 s^2 - 1]

In[2]:= Fm[m_] := Piecewise[{(1, m == 0), (s, m == 1), {s^Abs[n]-2, m > 1}}]

Beven[m_] := Piecewise[{(-1/2, m == 0), {1 - 1/2, m == 1}, {m - 5/2, m > 1}}]

In[3]:= BackgRegFactor = 1; (* +s (-1+s^2) ; *)

In[4]:= BsExp[m_, n_] := BackgRegFactor * (1 - s^2)^1/2 s^Abs[n]-1 JacobiP[n - 1, 1, Abs[m] - 3/2, 2 s^2 - 1]

BpExp[m_, n_] := BackgRegFactor * (1 - s^2)^1/2 s^Abs[n]+1 JacobiP[n - 1, 1, Abs[m] + 1/2, 2 s^2 - 1]

BzExp[m_, n_] := BackgRegFactor * (1 - s^2) s^Abs[n] JacobiP[n - 1, 2, Abs[m] - 1/2, 2 s^2 - 1]

Magnetic moments expansions for toroidal background field

In[5]:= Bsp[m_, n_] := (1 - s^2)^1/2 s^Abs[n-2] * JacobiP[n - 1, 1, Abs[m - 2] - 1/2, 2 s^2 - 1]

Bpp[m_, n_] := (1 - s^2)^1/2 * Fm[m] * s^2 * JacobiP[n - 1, 1, Beven[m] + 2, 2 s^2 - 1] (*Wherever Bpp

Bpz[m_, n_] := (1 - s^2) s^{n-1} JacobiP[n - 1, 2, m - 3/2, 2 s^2 - 1]

Bzpp[m_, n_] := (1 - s^2) Fm[m] * s^2 * JacobiP[n - 1, 2, Beven[m] + 2, 2 s^2 - 1] (*Wherever Bzpp
there needs to be a Bzsp contribution to ensure correct regularity. Note extra H from

Bzsp[m_, n_] := (1 - s^2) s^Abs[n-2] JacobiP[n - 1, 2, Abs[m - 2] - 1/2, 2 s^2 - 1]

Bs0[m_, n_] := (1 - s^2) s^Abs[n]-1 JacobiP[n - 1, 2, Abs[m] - 1/2, 2 s^2 - 1]

Bp0[m_, n_] := (1 - s^2) s^Abs[n]+1 JacobiP[n - 1, 2, Abs[m] + 1/2, 2 s^2 - 1] (*Factor of H ensu
```

Inflexible expansions: the expansion and the construction of matrices are hard coded in each notebook.

To change the expansion, one needs to

- produce a complete new Mathematica notebook;
- manually change the way the code collects the matrix elements.

Especially cumbersome to implement coupling, or if the bases do not coincide with the field quantities (more detail later).

# The old implementation

## Other problems

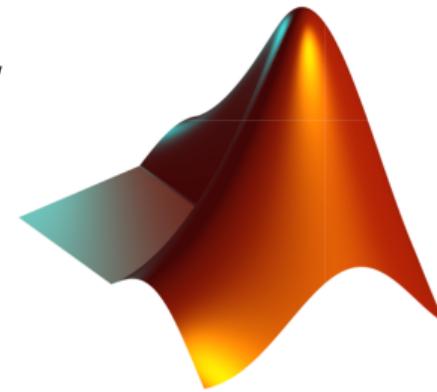
- The code is not efficient numerically, nor easily linked to efficient numerical libraries.
- Cryptic Mathematica syntax sugars

```
In[✓]:= sol = Solve[JacobiP[colN, α, β, x] == 0, x];  
rlocation = N[x /. # & /@ sol, prec];
```

- When something unexpected happens, difficult to understand the cause
- Or perhaps I just don't like Mathematica

## IMAGINE BEING A MATH PROGRAMMING ENVIRONMENT

THIS MEME MADE BY



THAT'S NOT SYMBOLIC BY DEFAULT

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# Introducing PlesioGestroPy

Make Plesio-Geostrophy model easy to use!

Symbolic manipulation  
[`sympy`]

- all PG equations stored, easily accessible
- linearization
- configure expansion
- collect matrix elements

Demo notebook available.

Numerical computation  
[`mpmath`,`numpy`/`scipy`]

- compute matrix elements
- quadrature
- eigenvalue solver
- time stepper
- comes in both std (`numpy+scipy` backend) and multi-precision (`mpmath`)

Post-processing  
[`matplotlib`]

# First batch of result: a glimpse of efficiency

Test problem: eigenvalue problem with  
Malkus background field, diffusionless.

A not-so-well-controlled test:

Assemble mass and stiffness matrices at  
the size of  $124 \times 124$  (PlesioGeostroPy)  
and  $73 \times 73$  (Mathematica)

Three versions:

- Code Daria, eval to 32nd digit
- PlesioGeostroPy, multiprecision eval (same precision)
- PlesioGeostroPy, scipy eval (double precision)

Table: Time lapse for computing mass and stiffness matrices

Code (version)	Time M	Time K
Code Daria Mathematica	5s	50s*
PlesioGeostroPy-mp32	16s	35s
PlesioGeostroPy-scipy	400ms	750ms

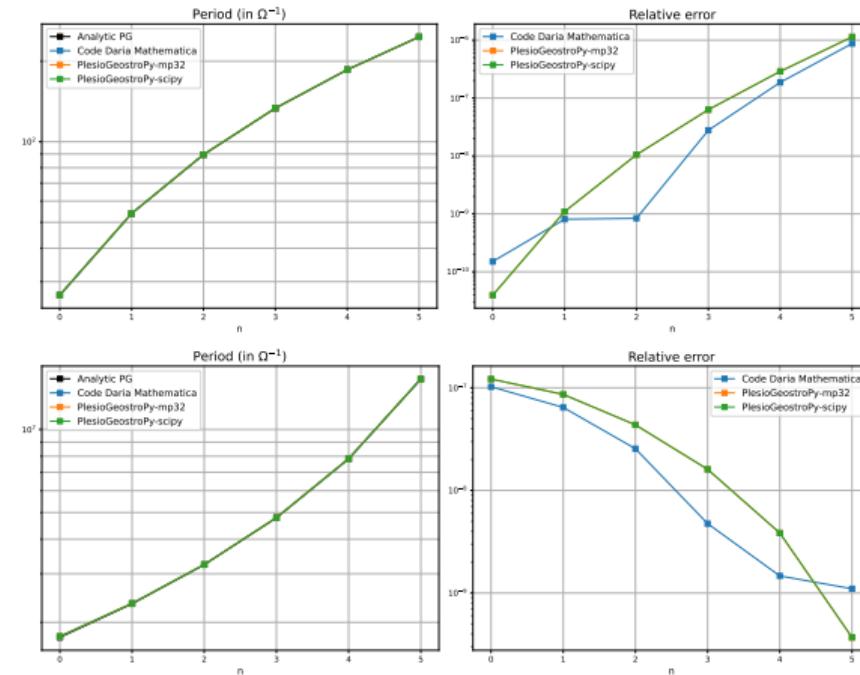
\*may not be representative as this involves other operations

# First batch of result: eigenvalues with Malkus background field

Test: eigenvalue problem with  
Malkus background field,  
diffusionless.

Eigenvalues (right) are  
computed for  $\text{Le} = 10^{-4}$ , the  
errors are measured relative to  
the analytical eigenvalue for the  
PG model.

*PlesioGeostroPy*  $\sim$  *Code Daria*  
 $\text{mp32} \approx \text{scipy}$ .



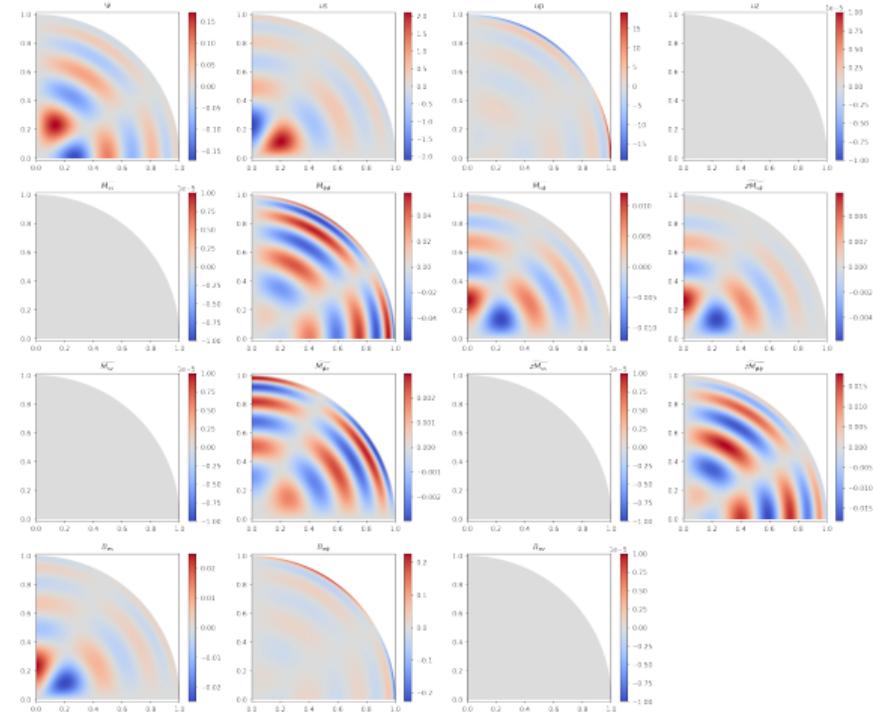
**Figure:** Periods and errors of the lowest 6 eigenmodes for the fast(top) and slow(bottom) branch.

# First batch of result: eigenmodes with Malkus background field

Test: eigenvalue problem with  
Malkus background field,  
diffusionless.

Eigenmodes are computed for  
 $\text{Le} = 10^{-4}$ .

Right fig. shows the equatorial  
sections of the 5-th eigenmode  
for  $m = 3$ .



# First batch of result: eigenmodes with Malkus background field

Test: eigenvalue problem with  
Malkus background field,  
diffusionless.

Eigenmodes are computed for  
 $\text{Le} = 10^{-4}$ .

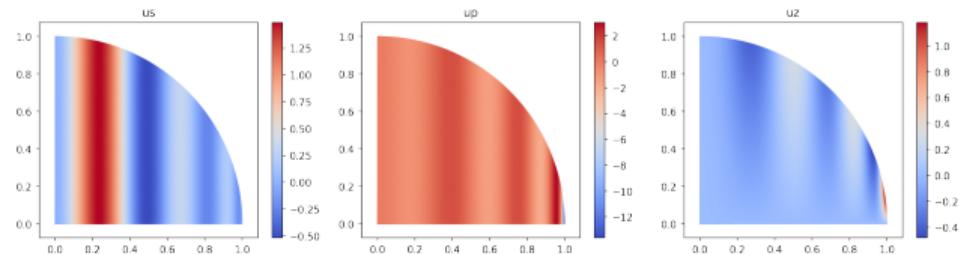


Figure: Meridional sections of the velocity field components for the 5-th eigenmode for  $m = 3$ .

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# To-do list

- Alternative expansions (even alternative evolution equations)
  - Additional coupling revealed (before group retreat), unlikely to implement all low-order coupling in the current way.
  - Need some thoughts + pen-and-paper derivation.
- Reduced systems of the eigenvalue problems
  - The eigenvalue problem with  $\mathbf{u}^0 = 0$  can always be reduced to a problem with two variables, instead of 15.
- Treatment of the  $B_r$ 
  - Circumvented in the eigenvalue problem.
  - Required for time stepping.
  - Possible to incorporate other than insulating BC?
- Time step the system