

VPI University Program

Photonics Curriculum Version 7.0

Lecture Series



Raman, Brillouin and
Rayleigh scattering

Fiber 4

Developed in cooperation with
Technische Universität Dresden,
Communications lab - RF engineering



Authors: M. Haas, N. Neumann

- **Fiber 1: Basics of Fiber Propagation**

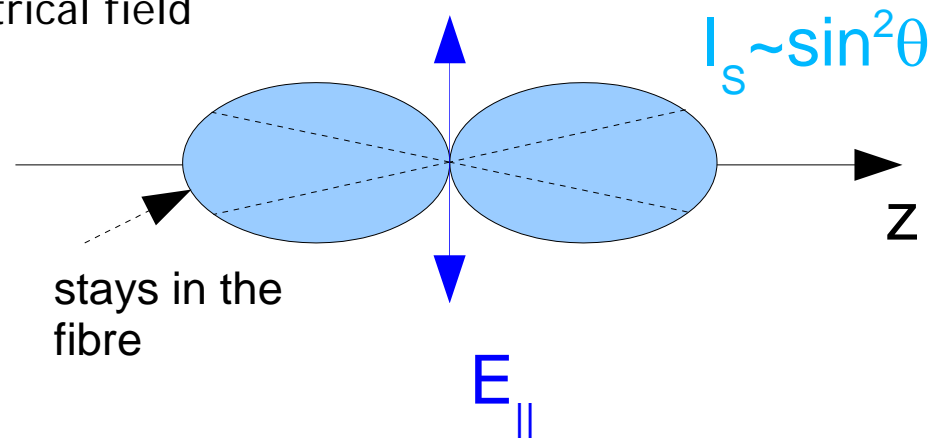
Module Objectives

- Rayleigh scattering
- Nonlinear effects in optical fibers
- Raman scattering
- Brillouin scattering
- Conclusion

Rayleigh scattering

Rayleigh scattering

- Optical field interacts with non-propagating density fluctuations in the fiber:
 - Microscopic defects in the amorphous structure of glass ($d \ll \lambda$)
 - Scattering of light ($P_r \sim \lambda^{-4}$)
 - Fiber attenuation (approx. 0.12 dB/km for SSMF @ 1550 nm)
 - Characterized by the fiber parameter: Rayleigh backscattering coefficient (-82 dB for SSMF @ 1550 nm and 1 ns pulse width)
 - Rayleigh scattering can be modeled as a radiating dipole
 - Radiation is rotationally symmetric to the polarization of the electrical field



Applications and System considerations

Applications:

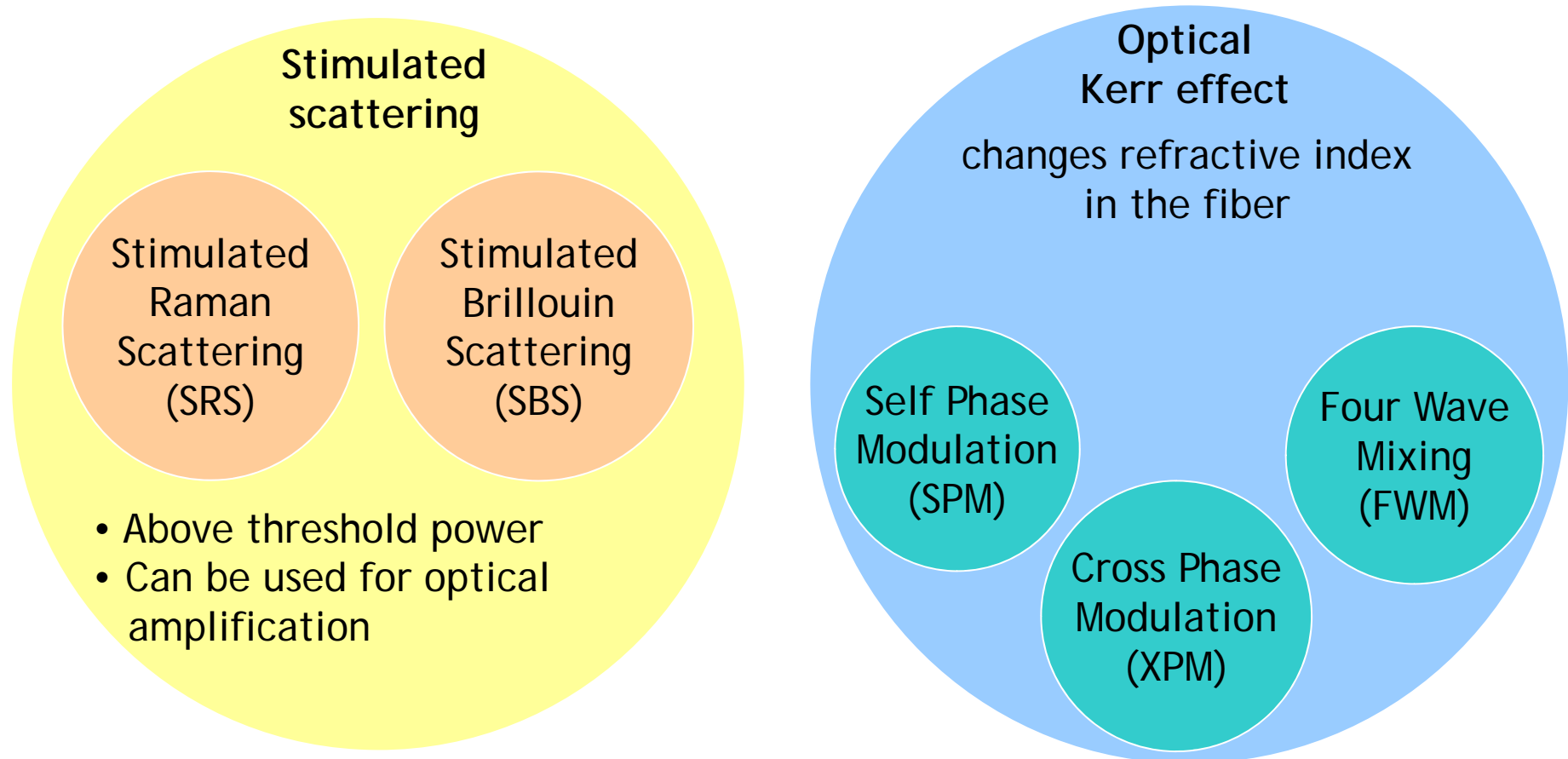
- Optical Time Domain Reflectometry (OTDR)
- Sensing

System considerations:

- Rayleigh scattering can degrade system performance of a fiber optical transmission system
 - Converts the phase noise of the source into amplitude noise at the receiver
 - Two possible system constellations:
 - Back scattered light in bidirectional optical links
 - Double back scattered light in unidirectional link
 - Ghost pulses in loop configurations

Nonlinear effects in optical fibers

Nonlinear effects



Covered in more detail in lecture “Fiber 2”

Nonlinear effects

- dependencies
 - Material
 - Field distribution inside waveguide → power density

- nonlinear coefficient γ
 n_2 : nonlinear-index coefficient
 A_{eff} : effective core Area

$$\gamma = \frac{k_0 n_2}{A_{eff}}$$

Common values for different fiber types:

Parameter/Fiber Type	SSMF	DSF	DCF
$\gamma / \text{W}^{-1} \text{ km}^{-1}$	1,2	1,76	8,3

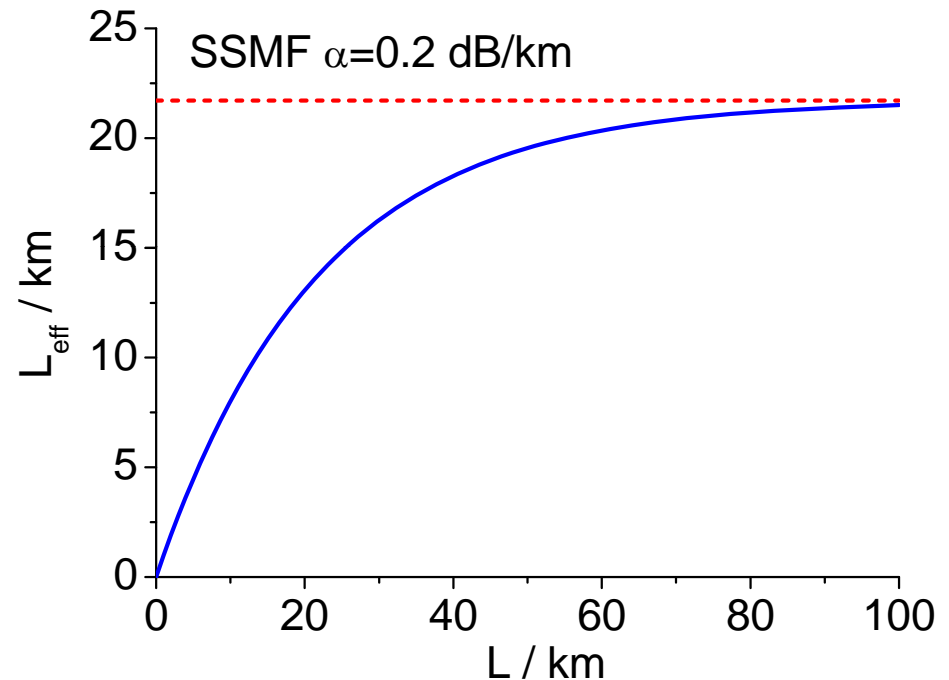
Nonlinear effects

- Effective length L_{eff}

$$L_{\text{eff}} = \frac{1 - e^{-\alpha' L}}{\alpha'},$$

$$\alpha' [km^{-1}] = \ln \left(10^{\frac{\alpha [dB/km]}{10}} \right),$$

$$\lim_{L \rightarrow \infty} L_{\text{eff}} = \frac{1}{\alpha'}$$



- A fiber without attenuation and an effective length of L_{eff} has the same nonlinear behavior as a fiber with attenuation α and length of L

Nonlinear effects

- Effective core Area:
 - “Reciprocal” overlap integral for the considered mode of the fiber

$$A_{eff} = \frac{\int_{-\infty}^{\infty} \int_{-\infty}^{\infty} |F(x, y)|^2 dx dy \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} |F(x, y)|^2 dx dy}{\int_{-\infty}^{\infty} \int_{-\infty}^{\infty} |F(x, y)|^2 |F(x, y)|^2 dx dy}$$

- $F(x, y)$ fiber mode field distribution
- Common values for different fiber types:

Parameter/Fiber Type	SSMF	DSF	DCF
$A_{eff} / \mu\text{m}^2$	80	54	21

Raman scattering

Raman scattering

- Spontaneous Raman scattering discovered by Raman
→ Nobel prize 1930
 - Optical field interacts with molecules:
 - Bound electrons oscillate at the frequency of the incident optical light
 - Induced oscillating dipole moment produces optical radiation with same frequency + phase shift
- phase causes refractive index of the medium

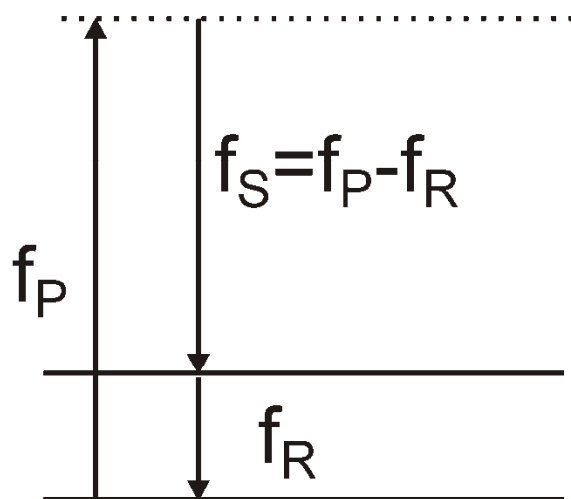
Raman scattering

- Molecular structure oscillates at various frequencies of molecular vibrations
 - Oscillation due to optical field contains sum and difference frequency terms between optical frequencies and vibrational frequencies
 - Sum and difference terms → Raman scattered light
- Inelastic scattering (energy transfer from/to lattice)
 - Optical Phonons are used to quantify loss/gain
 - Photon energy is lost → lattice is heated
 - Photon energy is gained → lattice is cooled

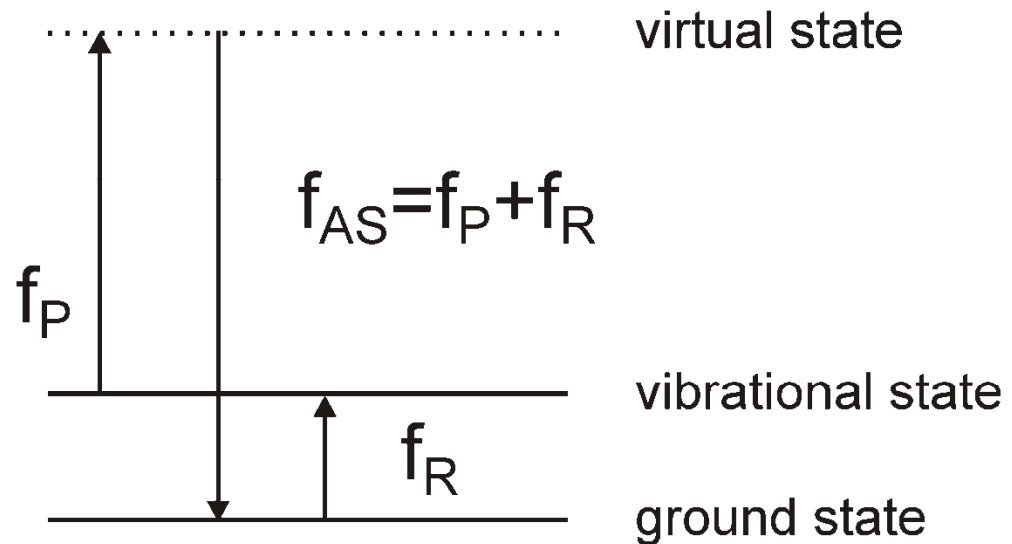
Stimulated Raman scattering Stokes, Anti-Stokes

- Scattered light shifted to lower frequencies
→ Stokes lines
- Scattered light shifted to higher frequencies
→ Anti-Stokes lines
- Anti-Stokes: orders of magnitude weaker than Stokes
→ not relevant for communications applications

Stimulated Raman scattering Stokes, Anti-Stokes



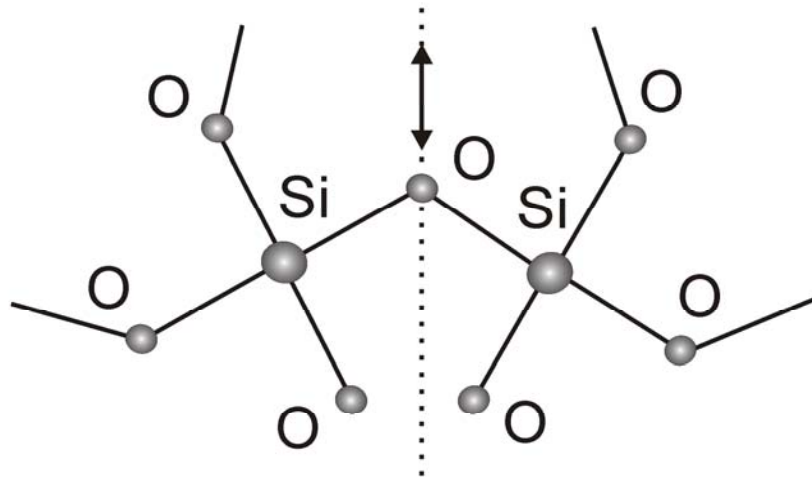
Stokes-wave
generation



Anti-Stokes-wave
generation

Stimulated Raman scattering (Anti-)Stokes frequency shift

- Frequency shift is determined by the frequency of the lattice phonon
- Silica glass:
dominant Raman line due to bending motion of Si-O-Si bond

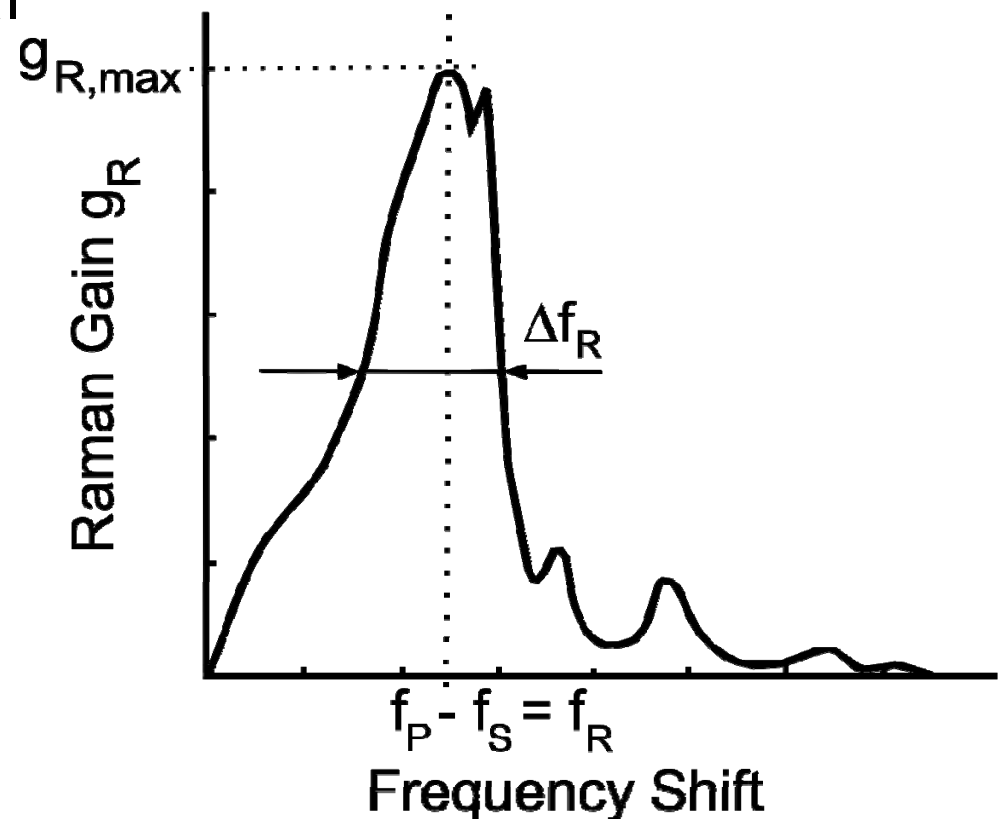


Stimulated Raman scattering

- Stimulation by pump with appropriate frequency shift with respect to signal light (determined by material)
→ amplification
- Pump and signal are coherently coupled by Raman process

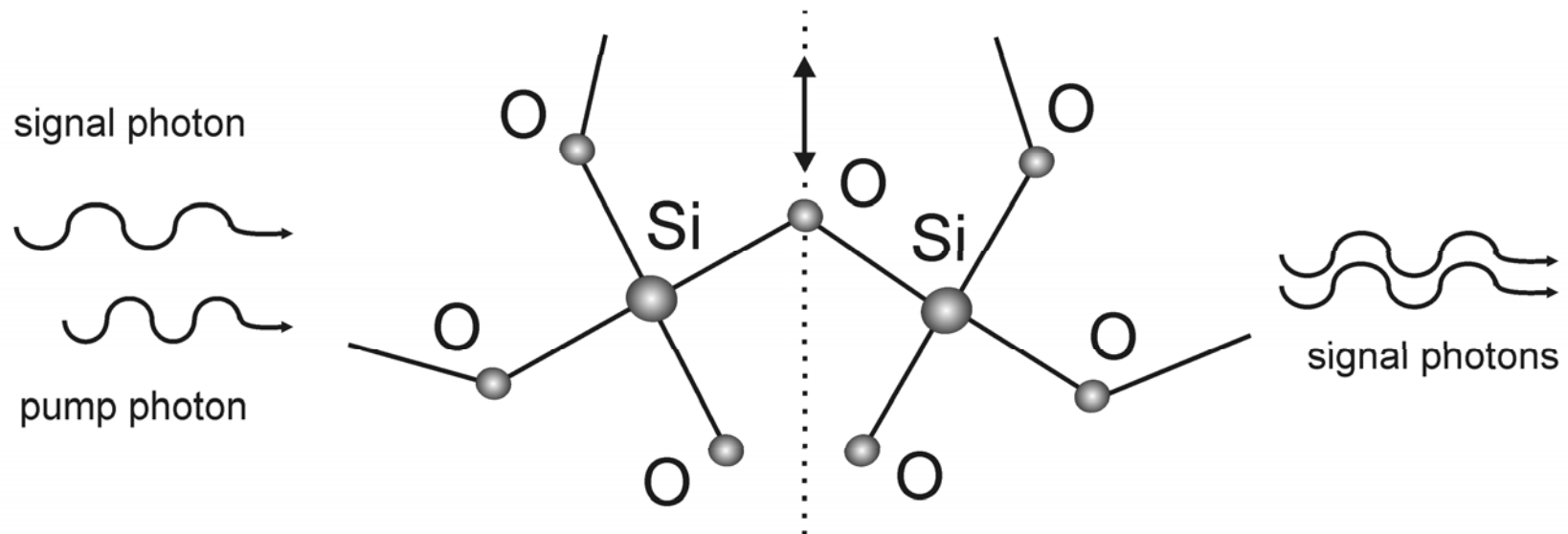
Stimulated Raman scattering

- Raman gain g_R is determined by frequency separation between pump and signal
 - Silica fibers:
 - $g_R = 1 \times 10^{-13} \text{ m/W}$
 - $f_R = 13 \text{ THz}$, $\Delta f_R = 6 \text{ THz}$
 - @ $\lambda = 1 \text{ }\mu\text{m}$
 - SSMF:
 - $g_R = 3.1 \times 10^{-14} \text{ m/W}$
 - @ $\lambda = 1.55 \text{ }\mu\text{m}$



Values from: T. Schneider, **Nonlinear Optics in Telecommunications**, Springer-Verlag Berlin, 2004
 W. Glaser, **Photonik für Ingenieure**, Verlag Technik Berlin, 1997

Stimulated Raman scattering



- Inelastic scattering (energy transfer from / to lattice)
 - Photon energy is lost (pump photon) → lattice is heated, bending motion of Si-O-Si bonds
 - Photon energy is gained (signal photon) → lattice is cooled

Stimulated Raman scattering

- SRS is non-resonant
 - Upper state is virtual
 - No long upper-state lifetimes
 - No buffering effect for pump fluctuations but averaging over fiber length (several kms for efficiency reasons) possible
 - Fast process: sub-picosecond
- Polarization dependent
 - Peak coupling strength is an order of magnitude stronger when co-polarized than if orthogonally polarized
- Application → Raman amplifiers

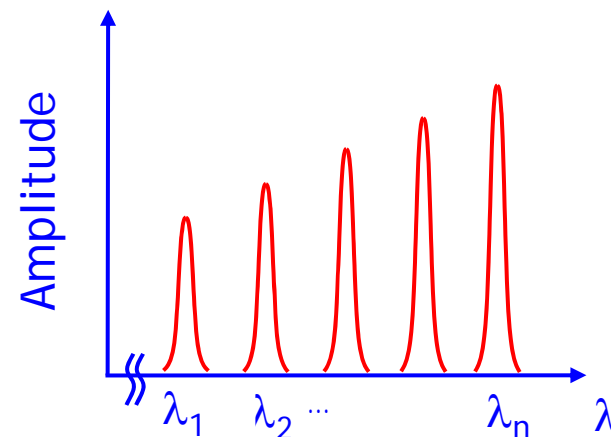
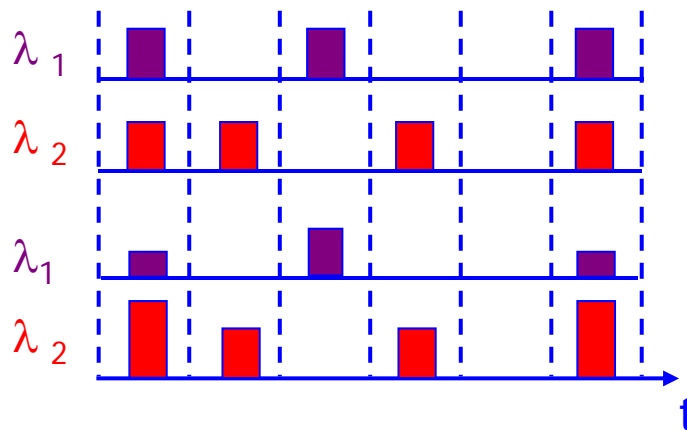
Applications

- Stimulated Raman Scattering
 - Raman lasers
 - Raman amplifiers
- Spontaneous Raman Scattering
 - Raman spectroscopy
(sometimes also using Stimulated Raman Scattering)

System considerations

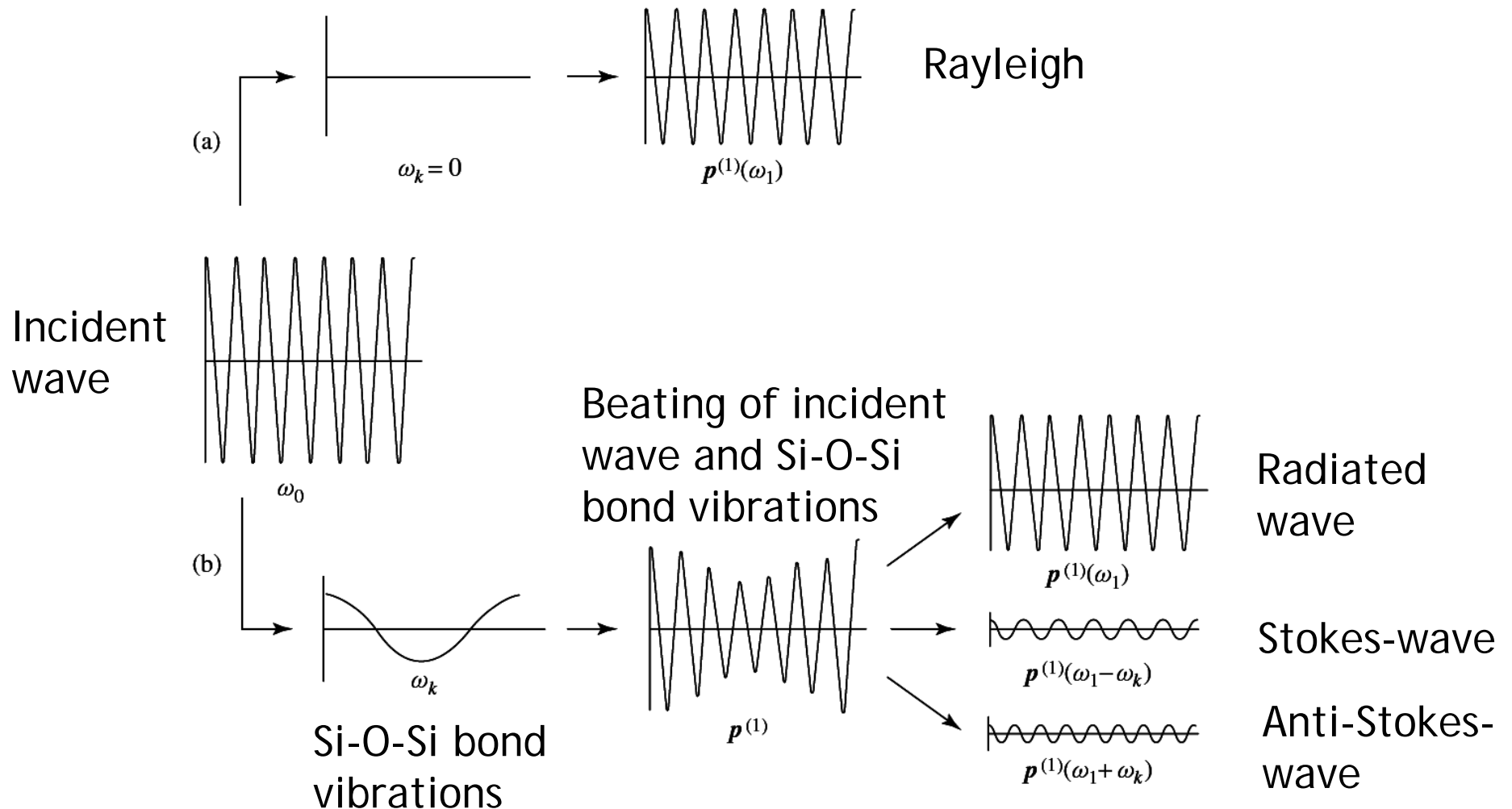
Raman tilt

- Amplification of higher wavelength channels at the expense of lower wavelength channels due to Raman scattering
- Time domain:
amplification or attenuation of single bits
- Frequency domain:
tilted WDM spectrum



References: T. Schneider, **Nonlinear Optics in Telecommunications**, Springer-Verlag Berlin, 2004
 G.P. Agrawal, **Nonlinear Fiber Optics**, Academic Press, 2007
 J. Bromage, **Raman Amplification for Fiber Communications Systems**, JLT, 2004

Raman and Rayleigh – a wave approach



Reference: Giordano Bruno, **Classical Theory of Rayleigh and Raman Scattering**

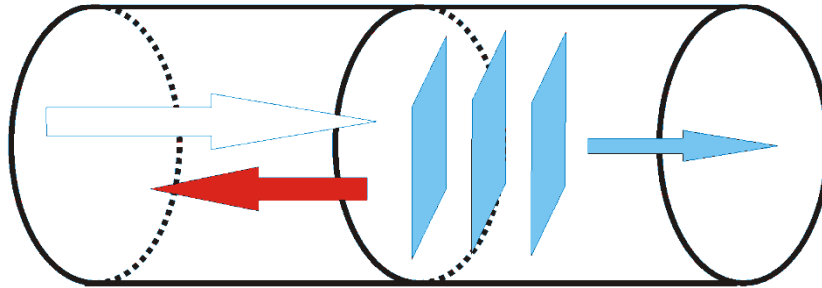
Brillouin scattering

Brillouin scattering

- Brillouin scattering discovered by Léon Brillouin
 - Optical field interacts with propagating density fluctuations in the fiber:
 - Thermal fluctuations of fiber density due to thermo-elastic motions of the molecules
 - Density fluctuations can be regarded as acoustic waves
 - Travel through the fiber at the speed of sound
 - Results in modulation of the refractive index of the fiber
- ➔ Scattering of light on the resulting refractive index grating

Brillouin scattering

pump wave
 E_P, ω_P, k_P



doppler shifted
 scattered wave
 E_S, ω_S, k_S

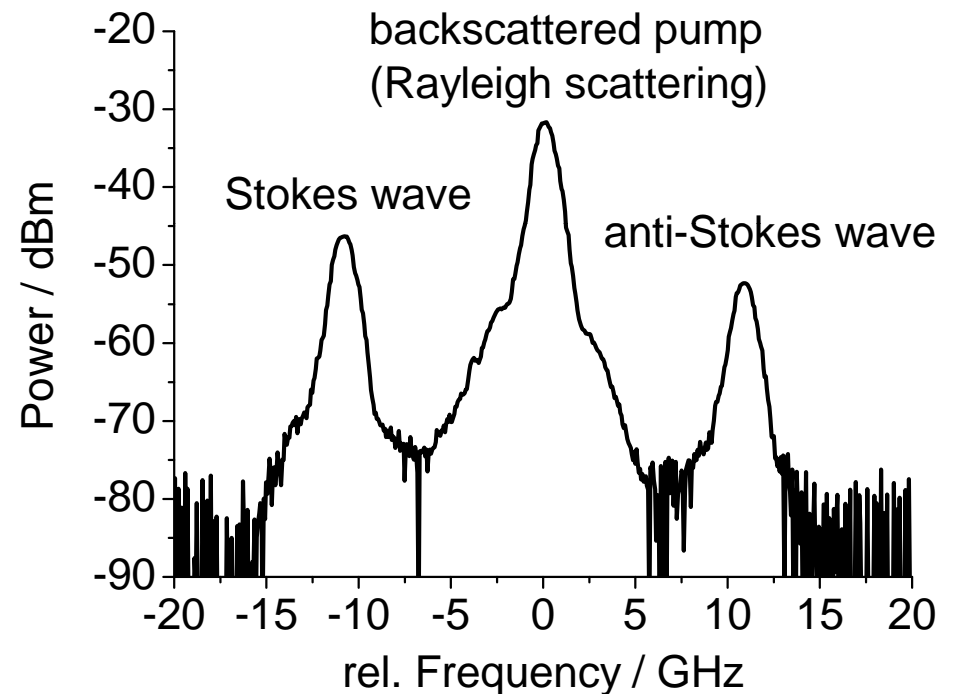
density (acoustic) wave
 v_a, ω_B, k_B

- Frequency shift of the back scattered wave for
 - Pump wavelength $\lambda_P = 1550$ nm
 - Refractive index $n = 1.45$
 - Sound velocity $v_a = 5.96$ km/s

$$f_B = \frac{2v_a n}{\lambda_P} = 11.15 \text{ GHz}$$

Brillouin scattering Stokes, Anti-Stokes

- Down shifted → Stokes wave
 - $f_S = f_P - f_B$
 - Acoustic wave propagates in the direction of the incident wave
- Up shifted → anti-Stokes wave
 - $f_S = f_P + f_B$
 - Acoustic wave propagates in the opposite direction of the incident wave



Brillouin scattering

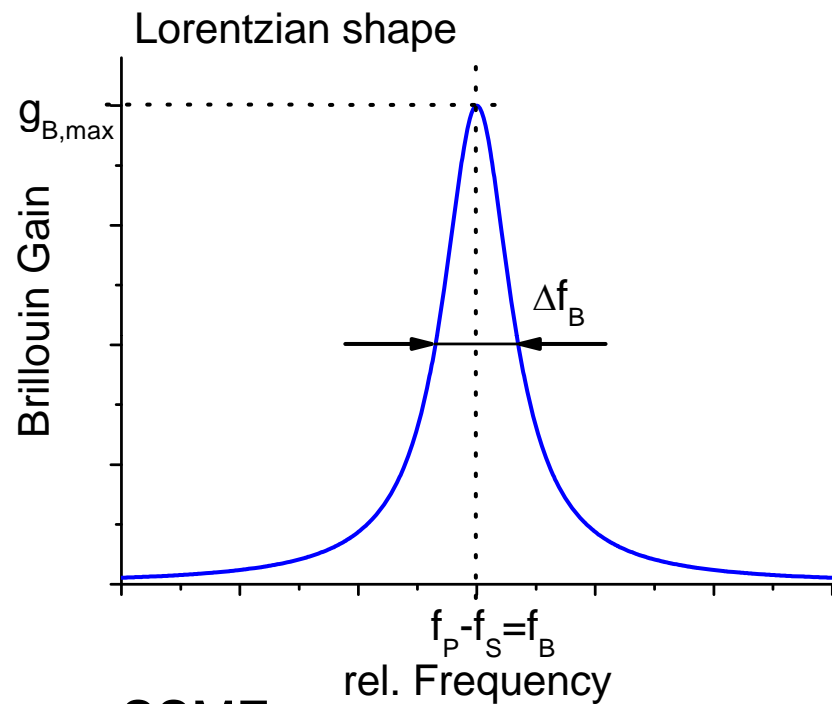
Brillouin gain:

- Lorentzian shape
(low gain region)

$$g_B(f) = \frac{1}{1 + \left[(f - f_B) / (\Delta f_B / 2) \right]^2} g_{B,\max}$$

- Gaussian shape
(high gain region)

$$g_B(f) = g_{B,\max} e^{-\ln(2) \frac{(f - f_B)^2}{(\Delta f_B / 2)^2}}$$

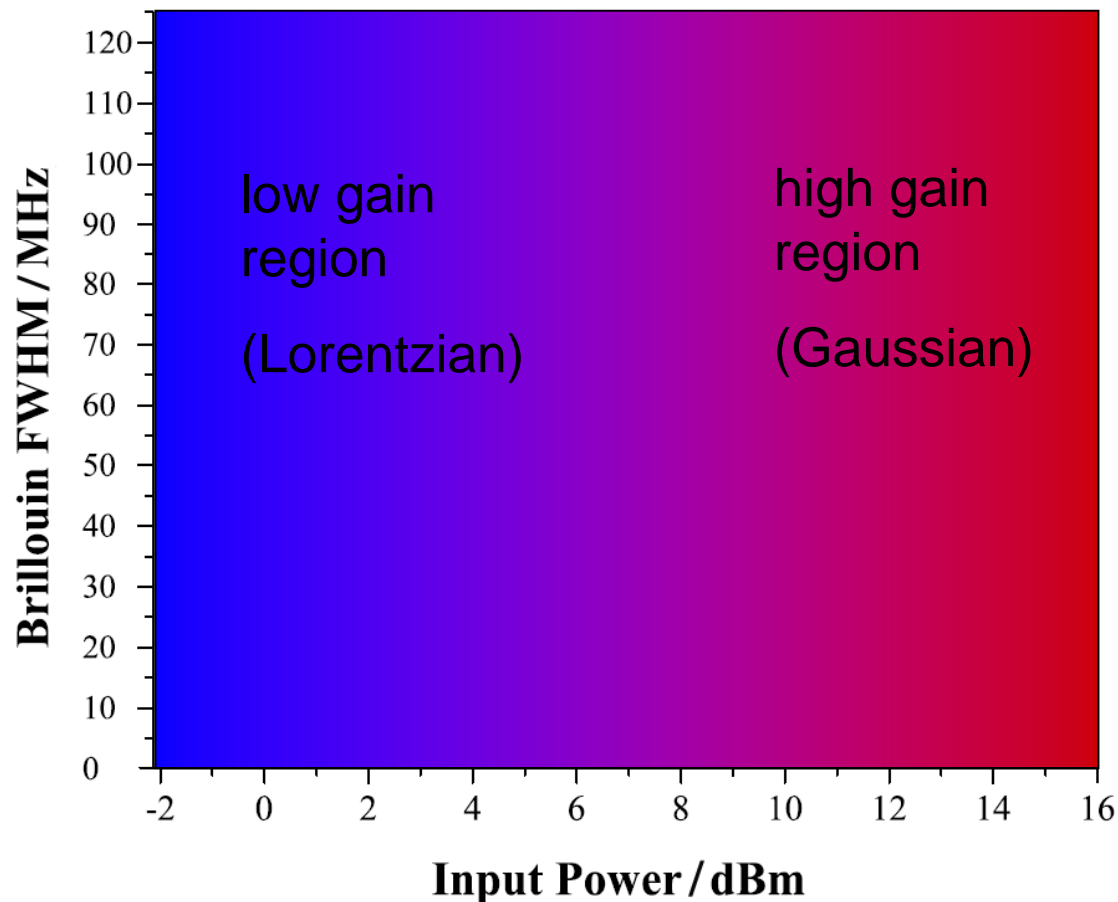


SSMF:

- $\lambda = 1.55 \mu\text{m}$
- $\Delta f_B = 35 \text{ MHz}$
- $g_B = 2 \times 10^{-11} \text{ m/W}$
- $f_B = f_P - f_S = 11.15 \text{ GHz}$

Values from: T. Schneider, **Nonlinear Optics in Telecommunications**, Springer-Verlag Berlin, 2004

- Full width half maximum (FWHM) depends on the pump power
- NZ-DSF fiber



Reference: A. Boh Ruffin, Stimulated Brillouin Scattering: An Overview of Measurements, System Impairments, and Applications

Stimulated Brillouin scattering

- Superposition between Stokes wave and incident wave (pump wave) leads to a fading with the beating frequency $\Delta f = f_P - f_S = f_B$
- Amplification of the propagating refractive index grating due to electrostriction
- Increase of the scattered Stokes wave power
- Exponential growth of the scattered power → stimulated scattering

Stimulated Brillouin scattering

Threshold power:

- $L \gg L_{eff}$
- Neglected pump depletion

$$P_{th} \approx 21 \frac{K \cdot A_{eff}}{g_B \cdot L_{eff}} \frac{\Delta f_B \otimes \Delta f_P}{\Delta f_B}$$

- K - polarization adjustment factor
- Convolution \otimes between the pulses

$$\Delta f_B \otimes \Delta f_P = \begin{cases} \sqrt{\Delta f_B^2 + \Delta f_P^2} & \text{Gaussian shape} \\ \Delta f_B + \Delta f_P & \text{Lorentzian shape} \end{cases}$$



Applications

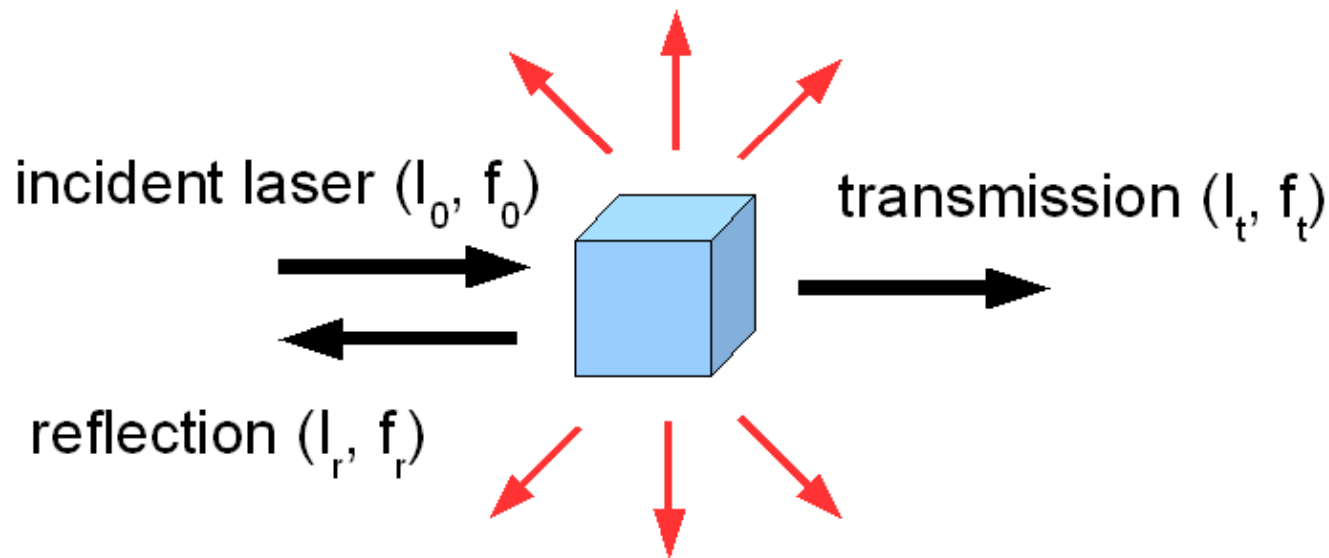
- Brillouin Amplifier
- Brillouin Laser
- Optical filter
- Brillouin optical spectrum analyzer (BOSA)
- Sensors (temperature, strain, ...)
- Slow and fast light

System considerations

- The signal transfers optical power to the backscattered wave → degradation of the SNR at the receiver
- Maximal optical power which can be transmitted through one channel in a fiber is limited (... but very high)
- Performance degradations in bidirectional optical links due to the backscattered wave

References: T. Schneider, **Nonlinear Optics in Telecommunications**, Springer-Verlag Berlin, 2004
G.P. Agrawal, **Nonlinear Fiber Optics**, Academic Press, 2007

Conclusion



scattered radiation

Rayleigh
(f_0)

- Particles
- Local refractive index variations

Brillouin
($f_0 \pm f_B$)

- Density fluctuations
- Acoustic waves
- $f_B \sim \text{GHz}$

Raman
($f_0 \pm f_R$)

- Molecules
- Electrons
- $f_R \sim \text{THz}$