

6-site Hubbard model

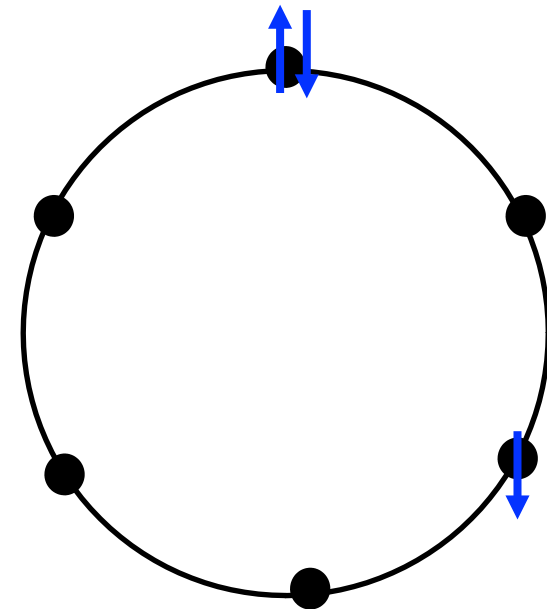
$$H = t \sum_{\langle ij \rangle, \sigma} c_{i\sigma}^\dagger c_{j\sigma} + U \sum_i n_{i\uparrow} n_{i\downarrow}$$

Large Fock space: $\dim 2^{12}$

Use conservation of S_z : (s_1, s_2) sectors of dim $\binom{6}{s_1} \binom{6}{s_2}$

For example a **basis** function from (1,2) sector:

in binary code (10000|101000)

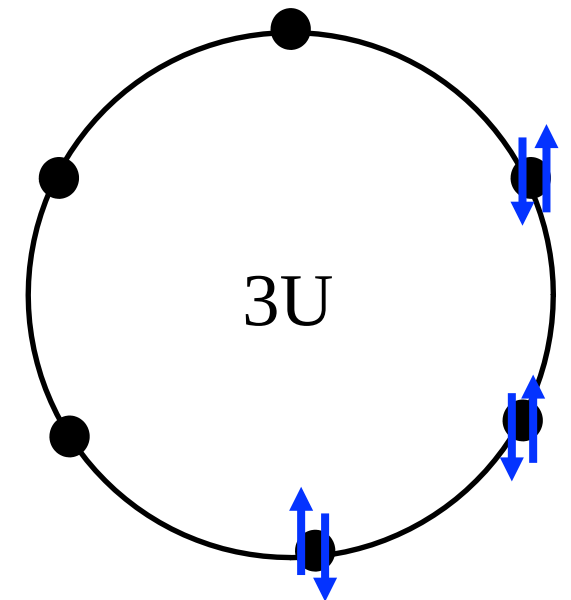
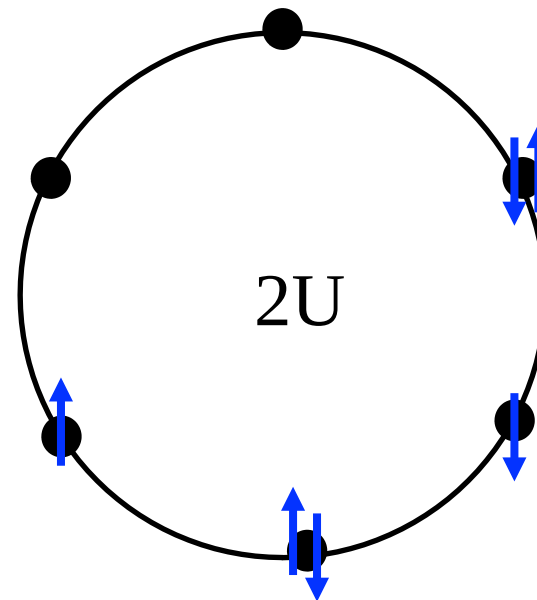
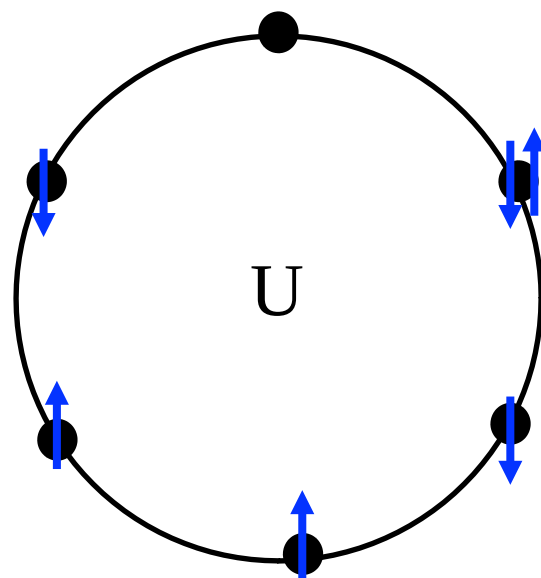
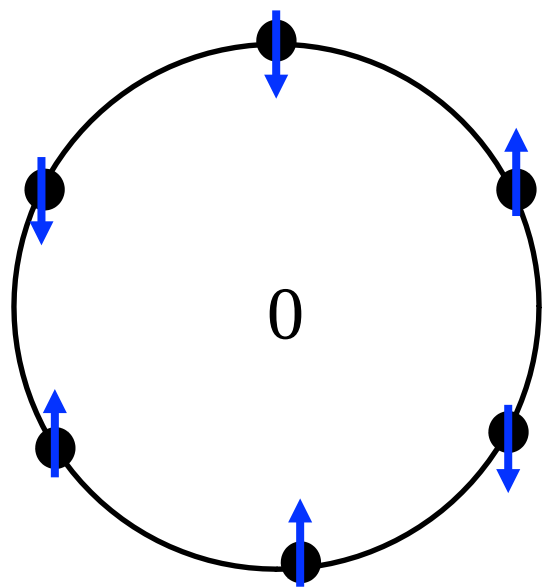


6-site Hubbard model

$$H = t \sum_{\langle ij \rangle, \sigma} c_{i\sigma}^\dagger c_{j\sigma} + U \sum_i n_{i\uparrow} n_{i\downarrow}$$

Matrix elements of the interaction part (diagonal in present basis):

Sector (3,3)



implementation of $\sum_i n_{i\uparrow} n_{i\downarrow}$

$(101000|100000) \rightarrow (101000) \cdot (100000) = 1$

implementation of S_{iz}

$(101000|100000) \rightarrow (101000) \cdot (100000) = (001000)$
i=2

6-site Hubbard model

$$H = t \sum_{\langle ij \rangle, \sigma} c_{i\sigma}^\dagger c_{j\sigma} + U \sum_i n_{i\uparrow} n_{i\downarrow}$$

Matrix elements of the hopping part:

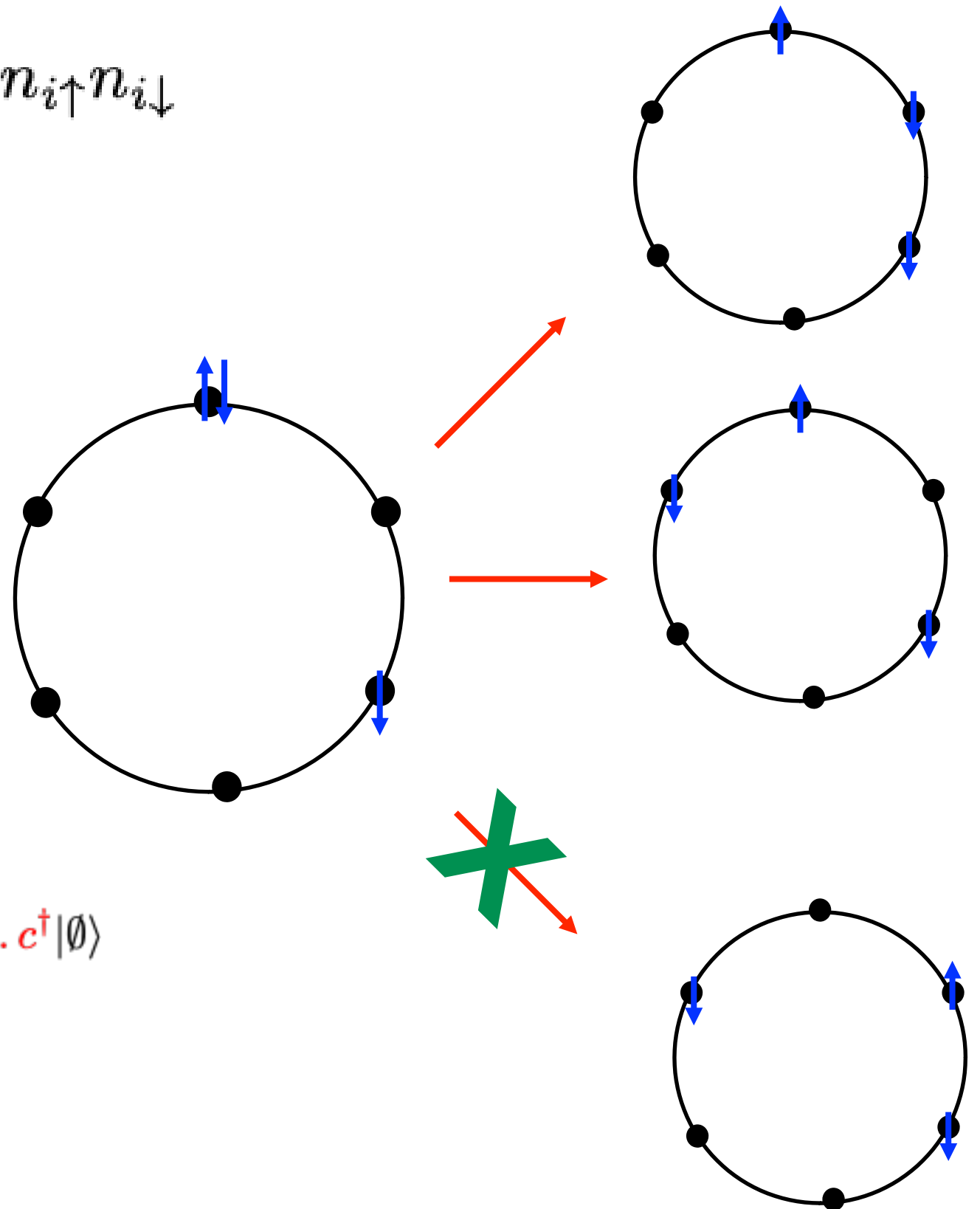
in binary code (10100|100000)

-> ~~(01100|100000)~~

✓ (11000|000000)

Signs:

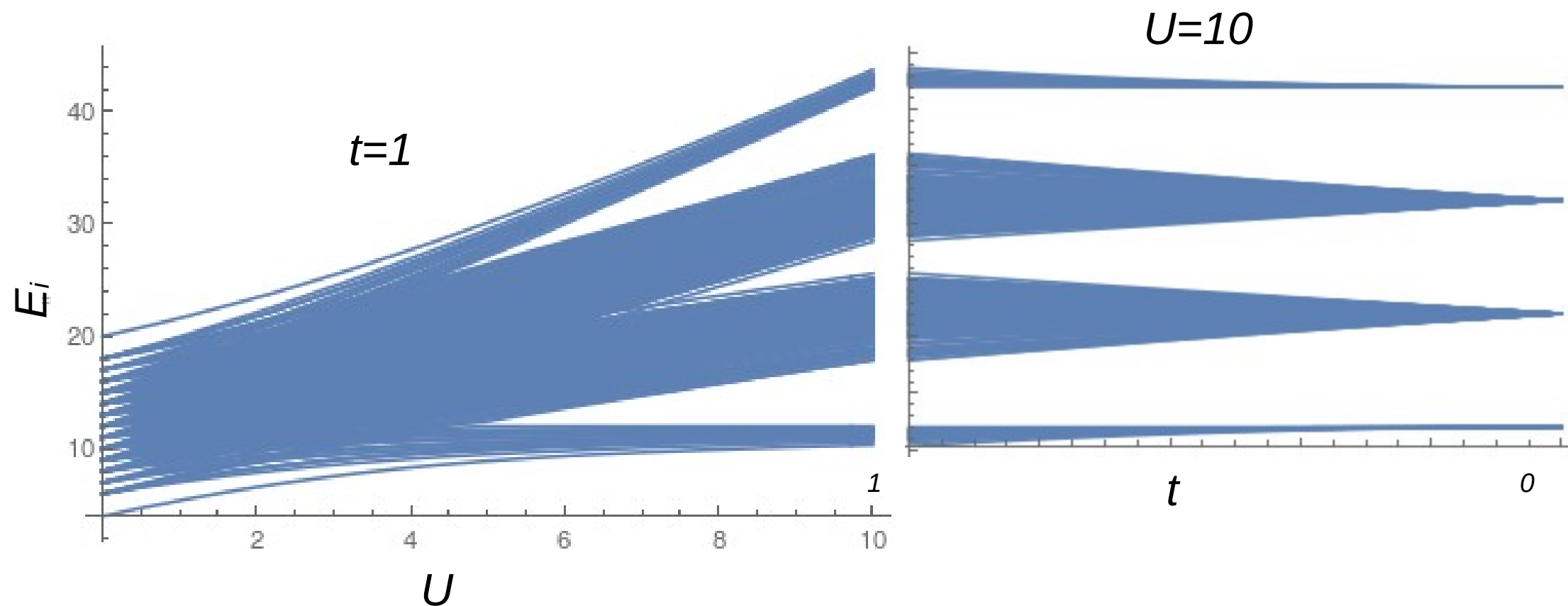
$$c_j | \textcolor{green}{c}^\dagger \dots \textcolor{green}{c}^\dagger c_j^\dagger \textcolor{red}{c}^\dagger \dots \textcolor{red}{c}^\dagger | \emptyset \rangle = (-1)^n \textcolor{green}{c}^\dagger \dots \textcolor{green}{c}^\dagger \textcolor{red}{c}^\dagger \dots \textcolor{red}{c}^\dagger | \emptyset \rangle$$



6-site Hubbard model

$$H = t \sum_{\langle ij \rangle, \sigma} c_{i\sigma}^\dagger c_{j\sigma} + U \sum_i n_{i\uparrow} n_{i\downarrow}$$

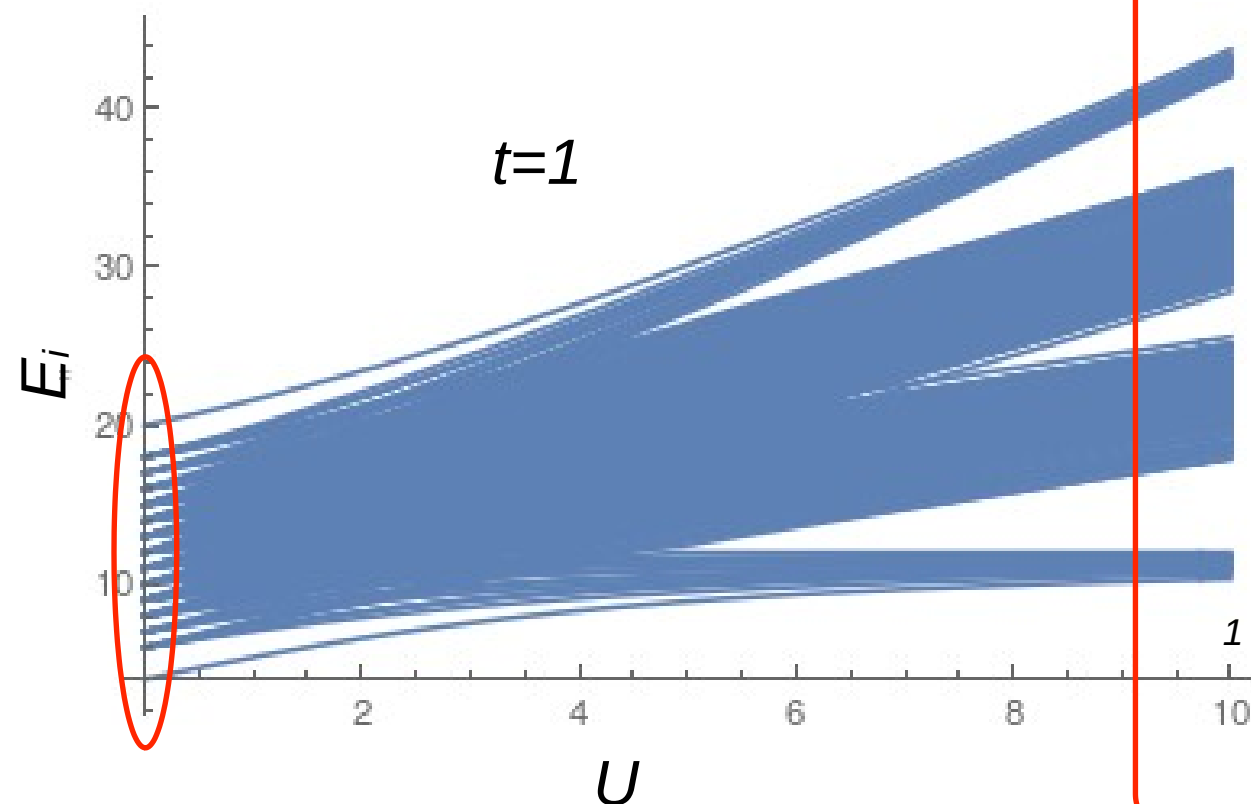
Spectrum of eigenenergies:
Sector (3,3)



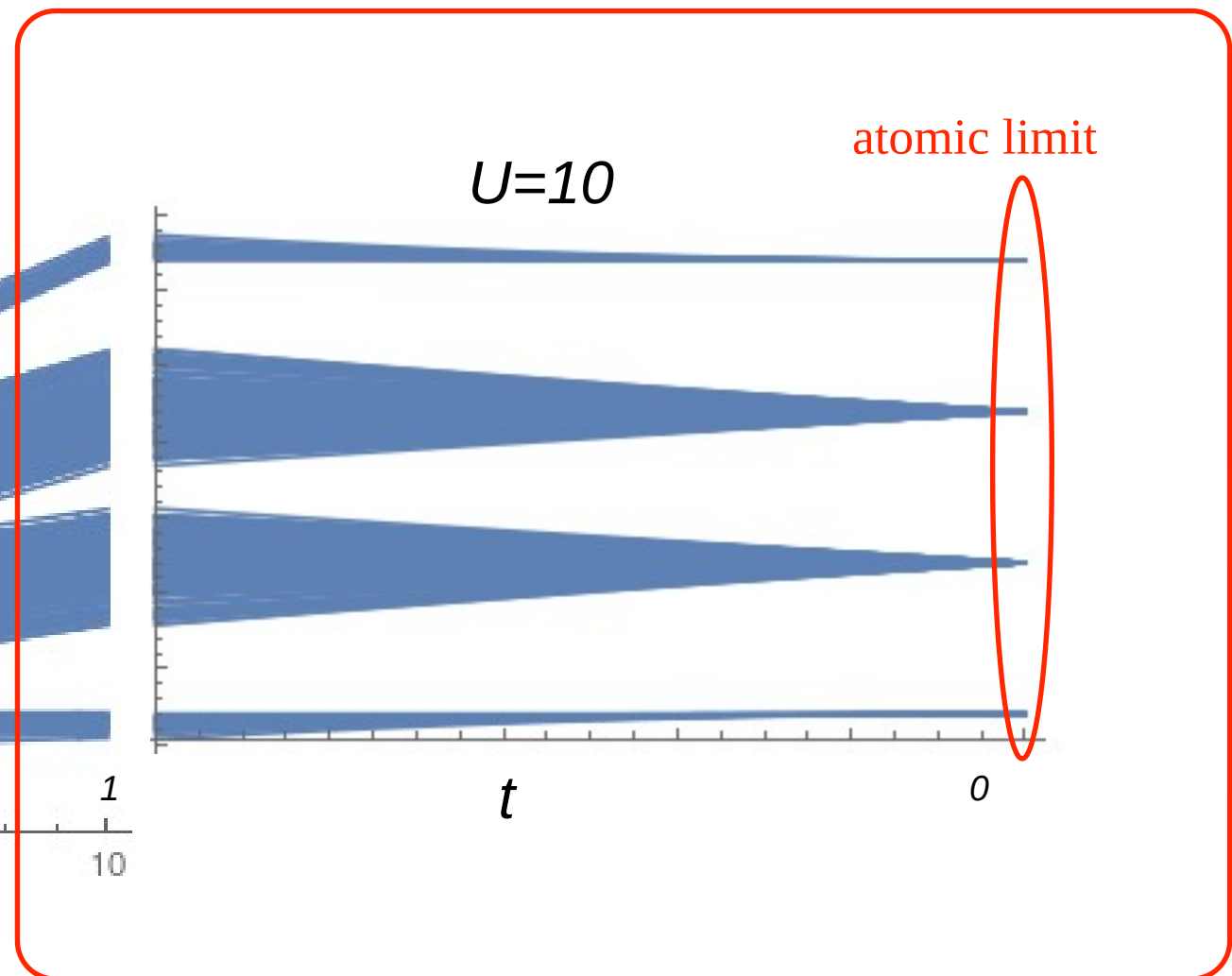
6-site Hubbard model

$$H = t \sum_{\langle ij \rangle, \sigma} c_{i\sigma}^\dagger c_{j\sigma} + U \sum_i n_{i\uparrow} n_{i\downarrow}$$

Spectrum of eigenenergies:
Sector (3,3)



Non-interacting limit
(weak-coupling expansion)



(strong-coupling expansion)

6-site Hubbard model

Non-interacting (canonical) bosons or fermions

=> We can find all eigenstates by diagonalizing the 1-p Hamiltonian (= hopping matrix)

$$H = \sum_{a,b} h_{ab} c_a^\dagger c_b$$

$$c_b = U_{bi} c_i, \quad (c_b^\dagger = U_{bi}^* c_i^\dagger = U_{ib}^\dagger c_i^\dagger)$$

$$\{c_i, c_j^\dagger\} = U_{ia}^\dagger \{c_a, c_b^\dagger\} U_{bj} = U_{ia}^\dagger \delta_{ab} U_{bj} = \delta_{ij}$$

$$H = \sum_i \epsilon_i c_i^\dagger c_i$$

$$|\phi\rangle = c_{i_1}^\dagger \dots c_{i_N}^\dagger |\text{vac}\rangle$$

$$H|\phi\rangle = \left(\sum_{k=1}^N \epsilon_{i_k} \right) |\phi\rangle$$

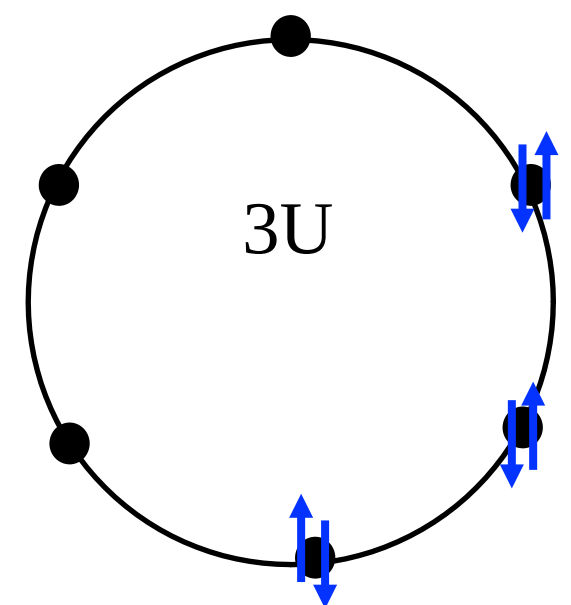
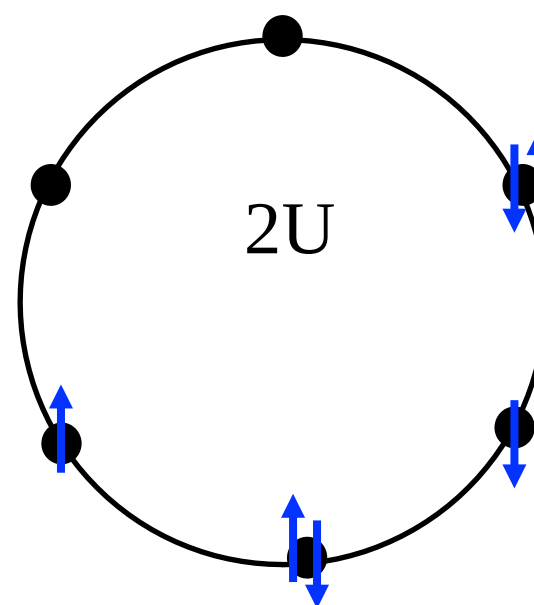
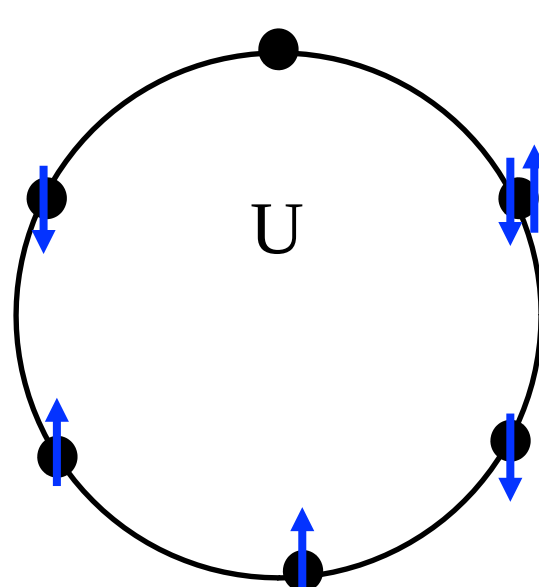
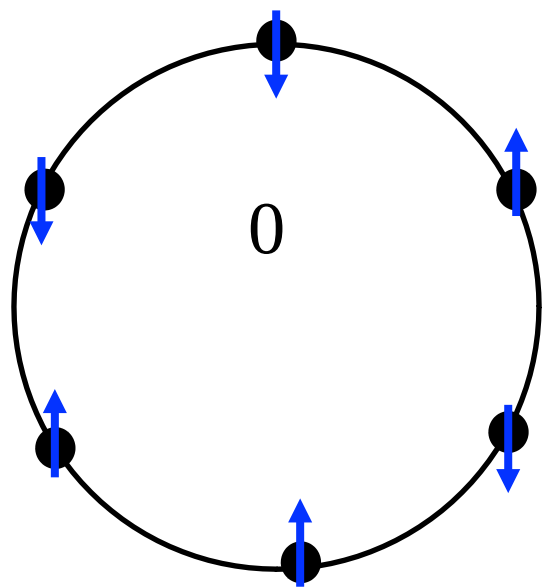
Canonical commutation relations!



6-site Hubbard model

$$H = t \sum_{\langle ij \rangle, \sigma} c_{i\sigma}^\dagger c_{j\sigma} + U \sum_i n_{i\uparrow} n_{i\downarrow}$$

Large $U \gg t$ limit:
Sector (3,3)



Degeneracy:

$$\binom{6}{3} = 20$$

$$6 \cdot 5 \cdot \binom{4}{2} = \cancel{120} \\ 180$$

$$\binom{6}{2} \cdot \binom{4}{2} \cdot 2 = \cancel{120} \\ 180$$

$$\binom{6}{3} = 20$$

6-site Hubbard model

$$H = t \sum_{\langle ij \rangle, \sigma} c_{i\sigma}^\dagger c_{j\sigma} + U \sum_i n_{i\uparrow} n_{i\downarrow}$$

Expectation values/correlation functions:

$$|\psi_g\rangle = \sum_l a_l |l\rangle$$

Simple form for operator diagonal in a given basis

$$\langle S_{iz} \rangle = \langle \psi_g | S_{iz} | \psi_g \rangle = \sum_l \langle l | S_{iz} | l \rangle |a_l|^2$$

The average value one gets when many measurement on site i are performed. Possible result of each individual measurement is 0, 1 and -1.

At half filling ($N=6$)

$$\langle S_{iz} \rangle = 0 \quad \langle n_{i\uparrow} \rangle = \frac{1}{2}$$

Fluctuations of S_z

$$\langle S_{iz}^2 \rangle - \langle S_{iz} \rangle^2 \neq 0$$

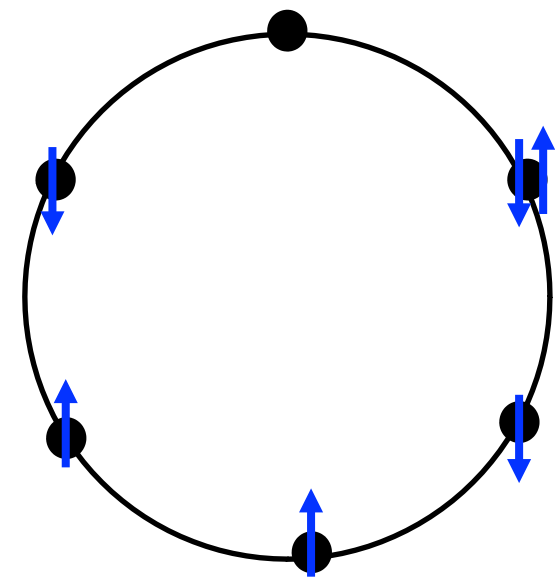
Total moment (occupation number):

$$S_z \equiv \sum_i S_{iz}$$

$$\langle S_z \rangle = 0 \quad \langle (\delta S_z)^2 \rangle \equiv \langle (S_z - \langle S_z \rangle)^2 \rangle = 0$$

$$\langle N \rangle = 6 \quad \langle (\delta N)^2 \rangle \equiv \langle (N - \langle N \rangle)^2 \rangle = 0$$

Conserved quantities (corresponding operators commute with Hamiltonian)



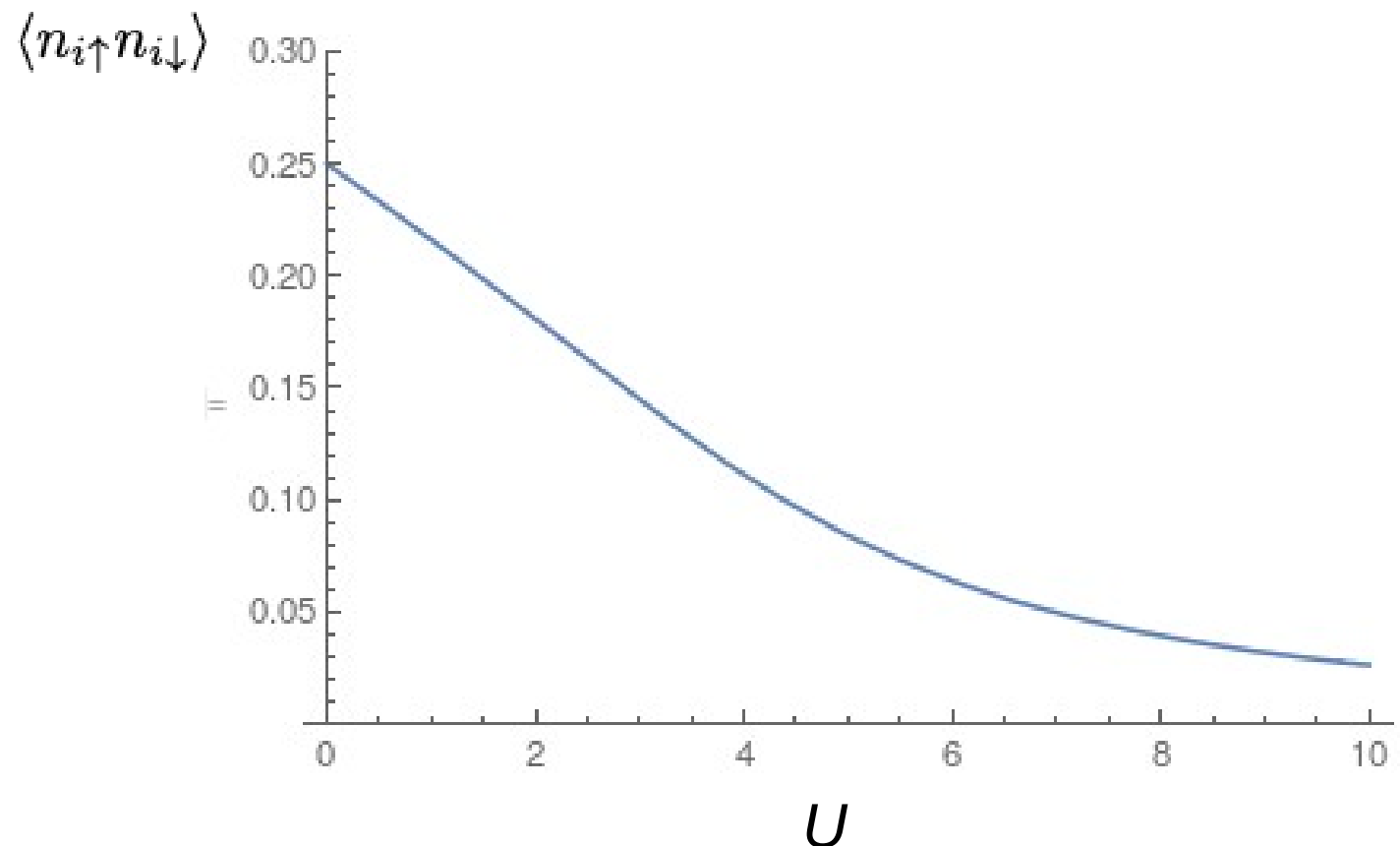
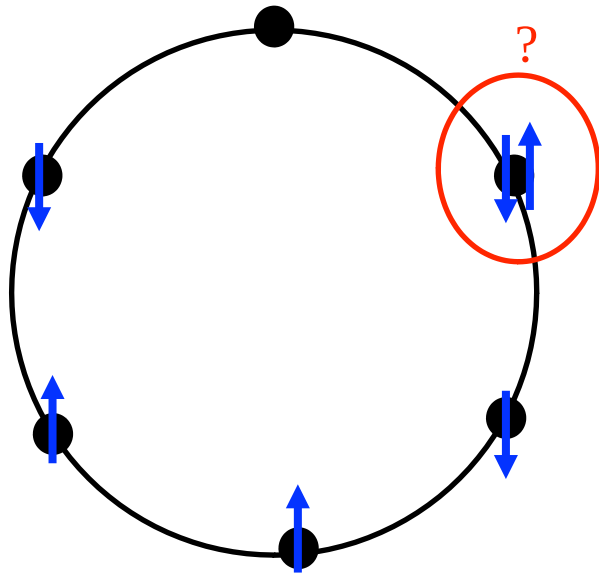
6-site Hubbard model

$$H = t \sum_{\langle ij \rangle, \sigma} c_{i\sigma}^\dagger c_{j\sigma} + U \sum_i n_{i\uparrow} n_{i\downarrow}$$

Expectation values/correlation functions:

Double occupancy: (probability to find two electrons in a given site)

$$\langle n_{i\uparrow} n_{i\downarrow} \rangle$$

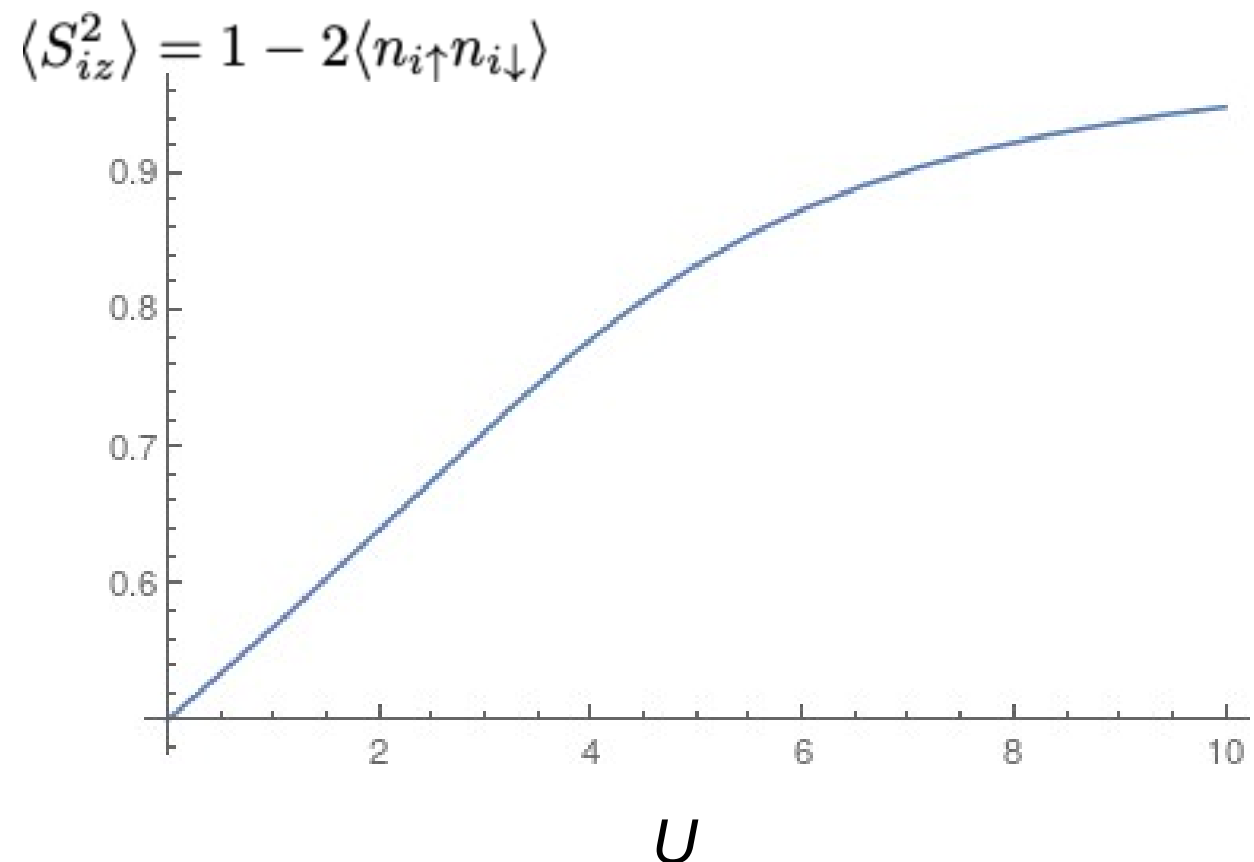
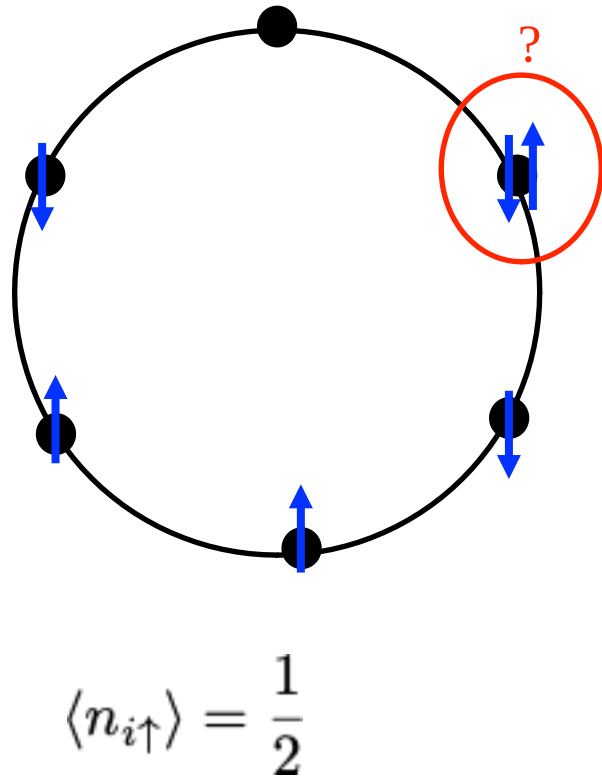


6-site Hubbard model

$$H = t \sum_{\langle ij \rangle, \sigma} c_{i\sigma}^\dagger c_{j\sigma} + U \sum_i n_{i\uparrow} n_{i\downarrow}$$

Expectation values/correlation functions:

(Fluctuating) local moment



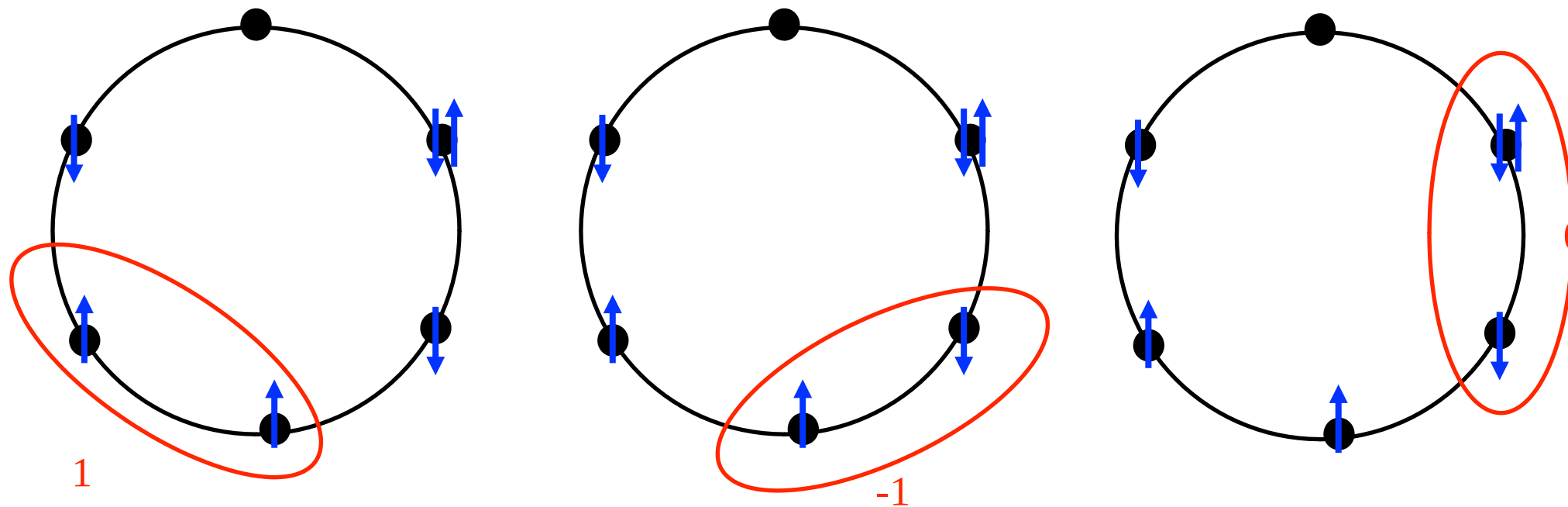
6-site Hubbard model

$$H = t \sum_{\langle ij \rangle, \sigma} c_{i\sigma}^\dagger c_{j\sigma} + U \sum_i n_{i\uparrow} n_{i\downarrow}$$

Expectation values/correlation functions:

Non-local spin-spin correlation function $\langle S_{iz} S_{jz} \rangle$

Weighted sum over configurations like

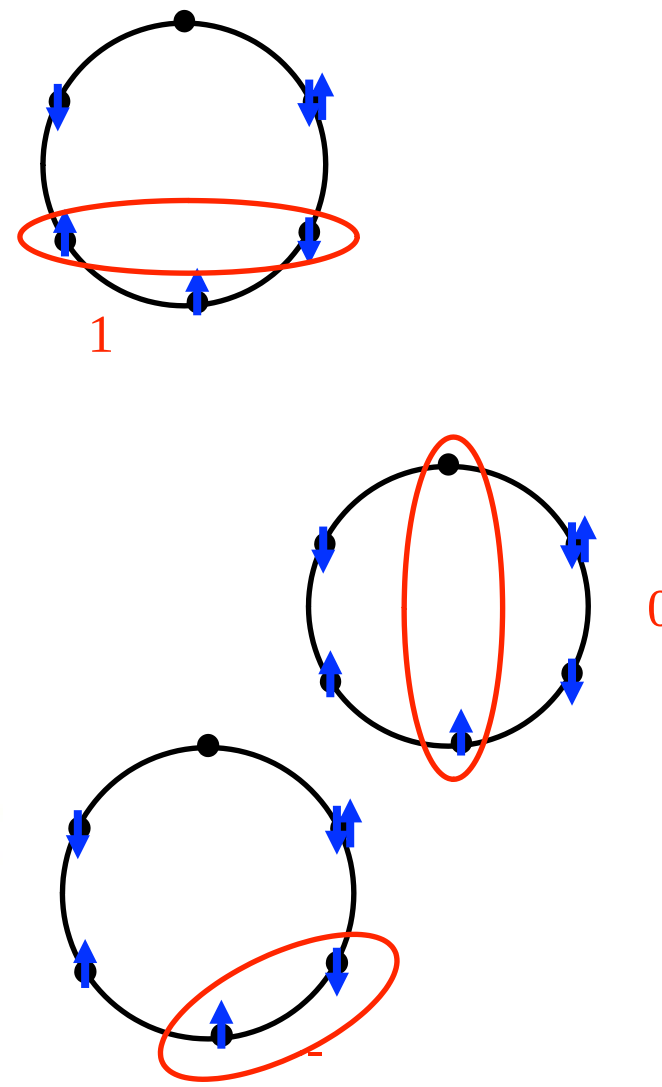
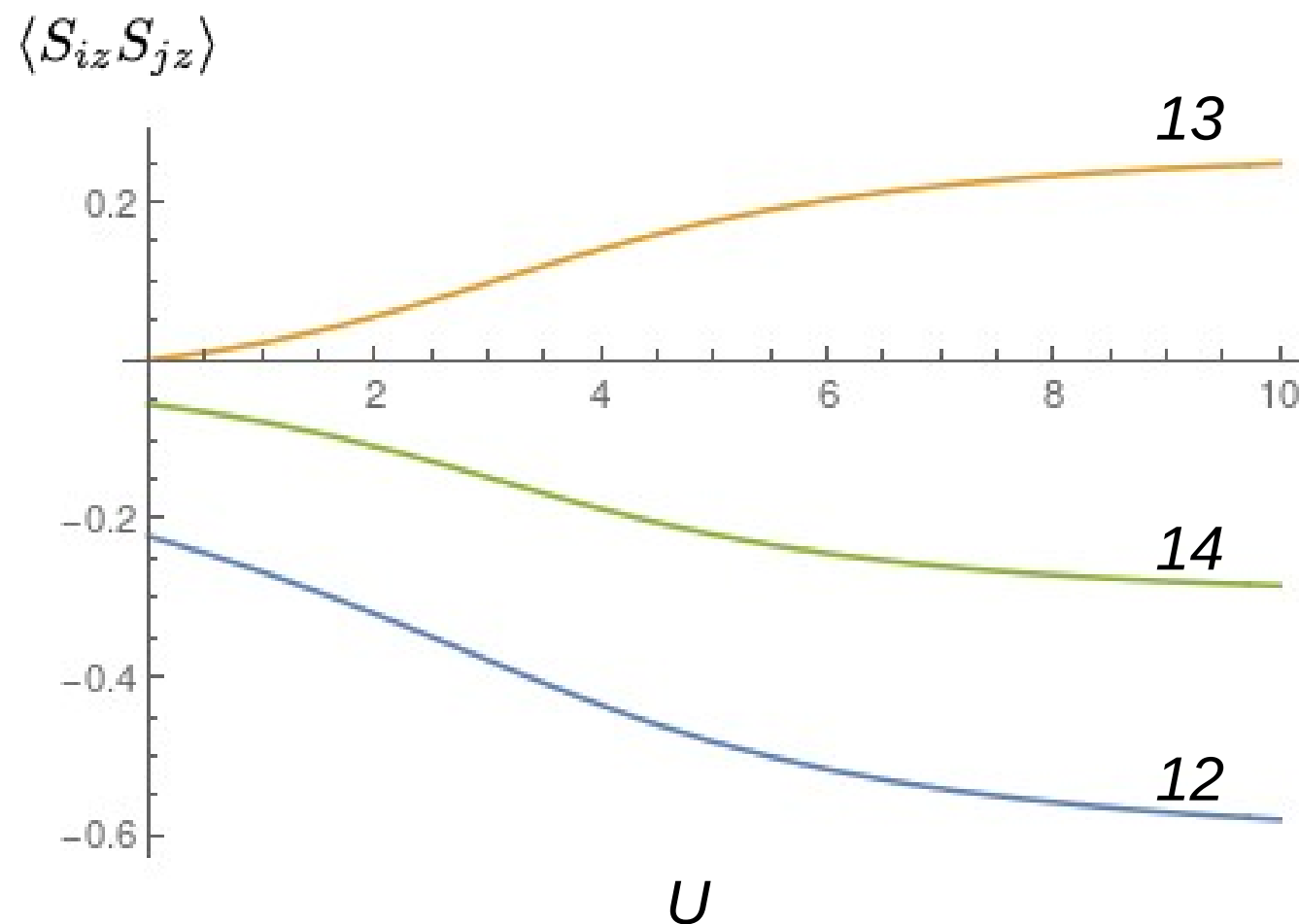


Which one has the largest weight?

6-site Hubbard model

$$H = t \sum_{\langle ij \rangle, \sigma} c_{i\sigma}^\dagger c_{j\sigma} + U \sum_i n_{i\uparrow} n_{i\downarrow}$$

Expectation values/correlation functions:

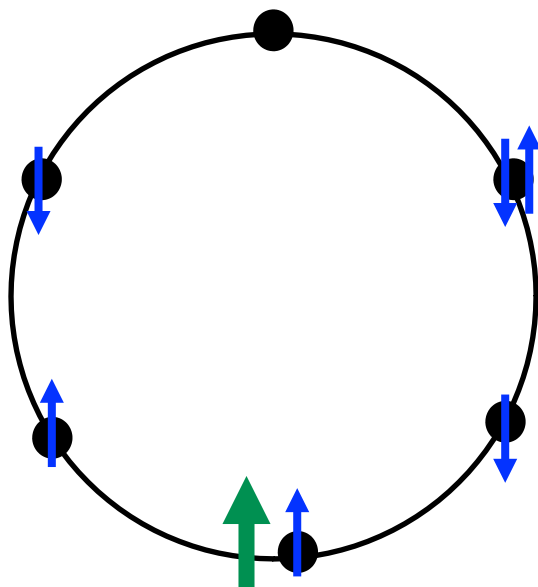


6-site Hubbard model

$$H = t \sum_{\langle ij \rangle, \sigma} c_{i\sigma}^\dagger c_{j\sigma} + U \sum_i n_{i\uparrow} n_{i\downarrow}$$

Why correlation functions?

- Contributions to interaction energy of the system $\langle n_{i\uparrow} n_{i\downarrow} \rangle$
- Response to small perturbations

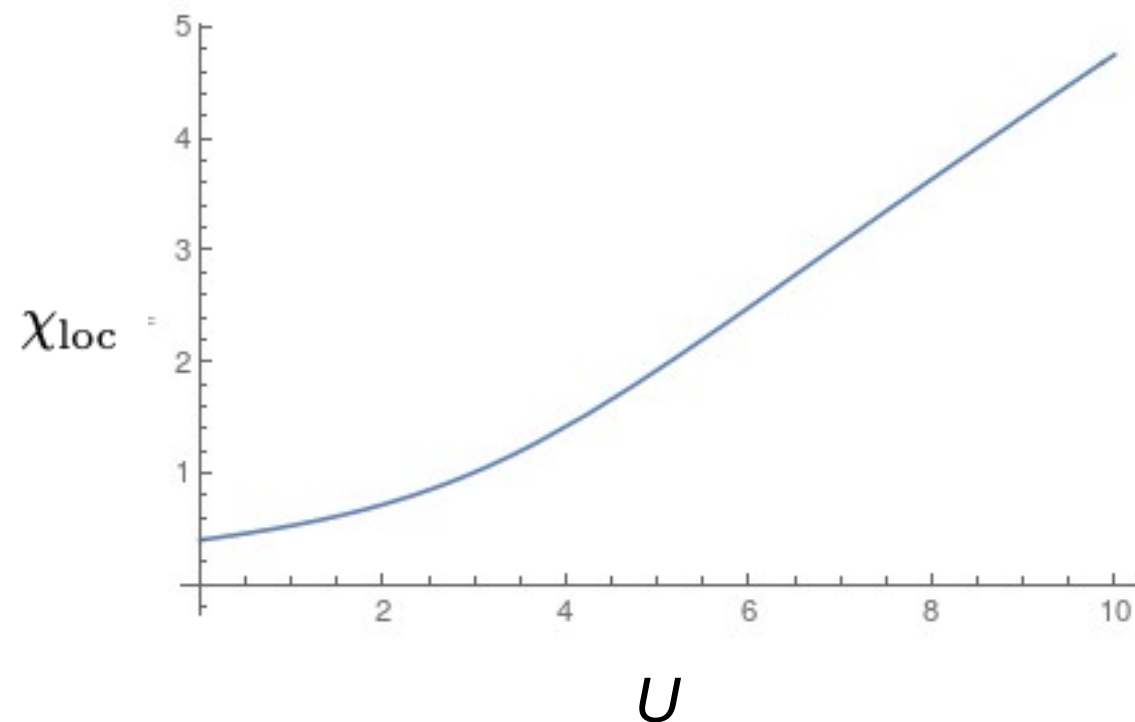


external local field

$$\delta \langle S_{iz} \rangle = \chi_{\text{loc}} \cdot \delta h$$

$$\chi_{\text{loc}} = 2 \sum_{n>g} \frac{|\langle \psi_n | S_{iz} | \psi_g \rangle|^2}{E_n - E_g}$$

Correction: Factor 2 missing in the recorded presentation

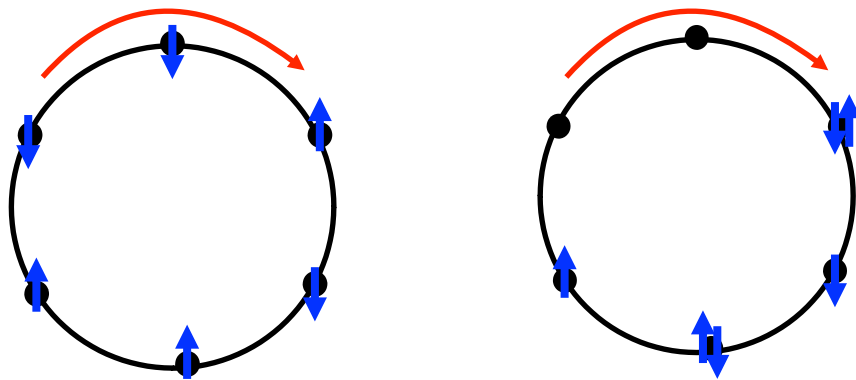


6-site Hubbard model

$$H = t \sum_{\langle ij \rangle, \sigma} c_{i\sigma}^\dagger c_{j\sigma} + U \sum_i n_{i\uparrow} n_{i\downarrow}$$

What about symmetry?

- We have used conservation of N and S_z when constructing the basis
- We did not use translation symmetry



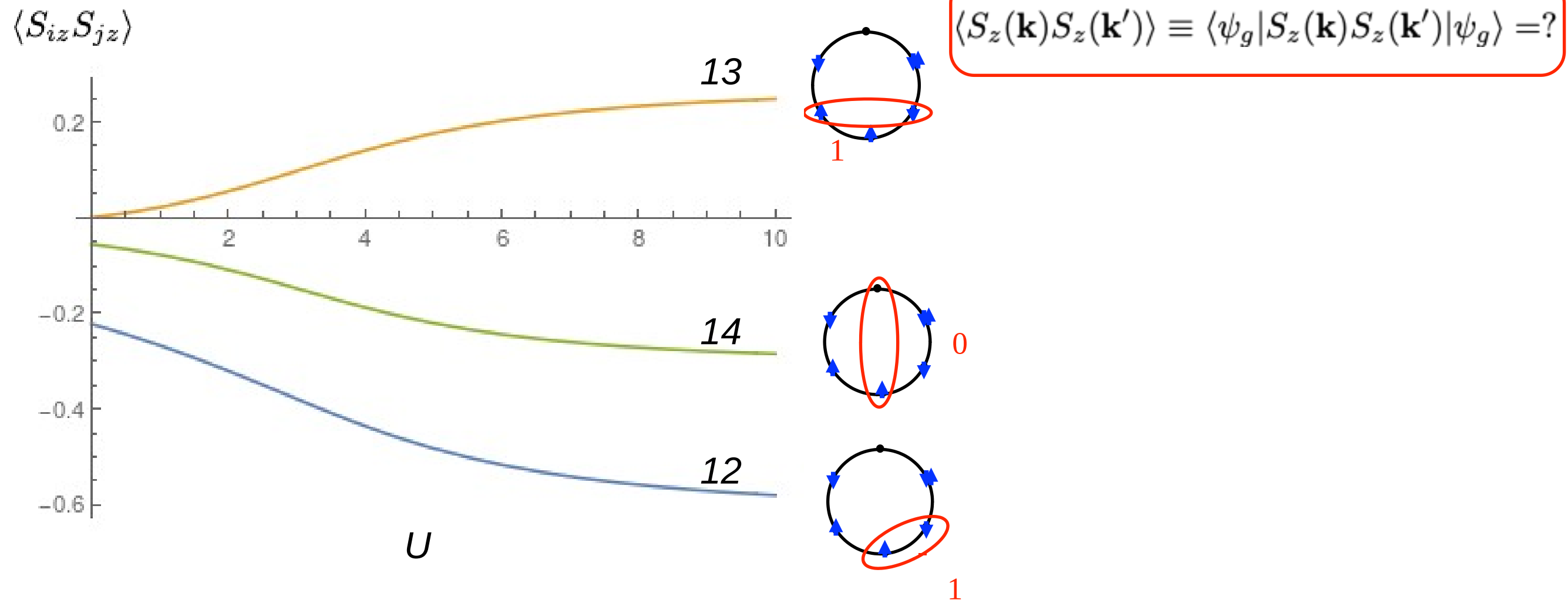
This would require a bit more 'brain' input

6-site Hubbard model

$$H = t \sum_{\langle ij \rangle, \sigma} c_{i\sigma}^\dagger c_{j\sigma} + U \sum_i n_{i\uparrow} n_{i\downarrow}$$

Translation symmetry is reflected in the correlation functions:

$$S_z(\mathbf{k}) = \sum_{\mathbf{R}} e^{i\mathbf{k} \cdot \mathbf{R}} S_{\mathbf{R}z}$$



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$$S_z(\mathbf{k}) = \sum_{\mathbf{R}} e^{i\mathbf{k} \cdot \mathbf{R}} S_{\mathbf{R}z}$$

$$\langle S_z(\mathbf{k}) S_z(\mathbf{k}') \rangle \equiv \langle \psi_g | S_z(\mathbf{k}) S_z(\mathbf{k}') | \psi_g \rangle = ?$$

invariant

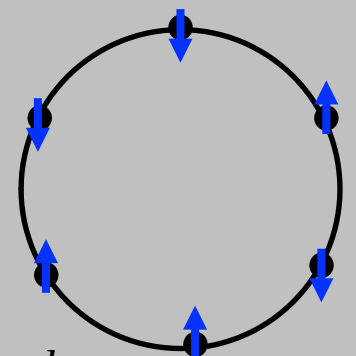
$$k + k' + (k'' - k'') = 0 \pmod{2\pi}$$

~~invariant. ($k''=0$)~~

Correction:

$$k'' = \pi$$

The ground state is dominated by the states of the type



The ground state phase does not matter as pointed out correctly.

$$\Rightarrow \langle S_z(\mathbf{k}) S_z(-\mathbf{k}') \rangle = \delta_{\mathbf{k}\mathbf{k}'}$$

6-site Hubbard model

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