

LECTURE 6: Information Theory, Entropy, Experiment Design

- The concept of **information theory** and **entropy** appears in many statistical problems.
- Here we will develop some basic theory and show how it can be applied to questions such as **how to compute statistical (in)dependence** and **how to design experiments** such that we achieve the goals
- In principle the development mirrors classical statistical mechanics, but there is no corresponding concept of quantum statistics

Information Theory

- Suppose we have a random variable with a discrete set of outcomes p_i , for $i=1, \dots, M$
- We construct a message from N independent outcomes of this variable
- We need $M \log_2 M$ bits to transmit this information
- But what if some are more likely than others: for large N we expect $N_i = N p_i$ events for each i

Information Theory

- Number of typical events is given by multinomial coefficient
 $g = N! / (\prod_{i=1}^M N_i!) \ll M^N$
- Remember Stirling formula $x! = x^x e^{-x} (2\pi x)^{1/2} + \dots$
- The number of bits needed to specify one of g events in large N limit is $\log_2 g = -N \sum_i p_i \log_2 p_i \ll N \log_2 M$: Shannon's theorem proves that in large N limit error with this number of bits vanishes
- Information content of p is $I(p_i) = \log_2 M + \sum_i p_i \log_2 p_i$

Entropy and information: discrete case

- Shannon information (Shannon 1948):

$$h(x) = -\log_2 p(x)$$
- Its average is called Shannon entropy:

$$H(X) = -\sum_i p_i \log_2 p_i$$
- $\log_2 g = \log_2 N! / (\prod_{i=1}^M N_i!)$ is known as entropy of mixing in the context of mixing of M components
- Example: English alphabet has information content of 4.7+4.1 bit
 $(p(x)=0: 0 \cdot \log_2 0 = 0) (\log_2 27 = 4.7)$
- Entropy is minimized at 0 for $p_i = \delta_{i,j}$ and maximized for $p_i = 1/M$: it is a measure of disorder

i	a_i	p_i	$h(p_i)$
1	a	.0575	4.1
2	b	.0128	6.3
3	c	.0263	5.2
4	d	.0285	5.1
5	e	.0913	3.5
6	f	.0173	5.9
7	g	.0133	6.2
8	h	.0313	5.0
9	i	.0599	4.1
10	j	.0006	10.7
11	k	.0084	6.9
12	l	.0335	4.9
13	m	.0235	5.4
14	n	.0596	4.1
15	o	.0689	3.9
16	p	.0192	5.7
17	q	.0008	10.3
18	r	.0508	4.3
19	s	.0567	4.1
20	t	.0706	3.8
21	u	.0334	4.9
22	v	.0069	7.2
23	w	.0119	6.4
24	x	.0073	7.1
25	y	.0164	5.9
26	z	.0007	10.4
27	-	.1928	2.4
$\sum_i p_i \log_2 \frac{1}{p_i}$			4.1

Relation between entropy and likelihood

- Instead of actual data likelihood we can replace it with its ensemble average
- Suppose we have N measurements x_i

$$L = \prod_i p(x_i)$$

$$\ln L = \sum_i \ln p(x_i)$$

$$\langle \ln L \rangle = \left\langle \sum_i \ln p(x_i) \right\rangle = N \langle \ln p(X) \rangle$$

$$\langle \ln p(X) \rangle \text{ (or } E\{\ln p(X)\}) = \int dx \cdot p(X) \ln p(X) = -H(X)$$

$$\langle \ln L \rangle = -NH(X)$$

Entropy for Continuous Distribution

$$H(X) = - \int p(x) \log p(x) dx = E\{-\log p(X)\}$$

- Not invariant under reparametrization: if we change x to $F(x)$ entropy changes by $\langle |F'(x)| \rangle$, so absolute value is meaningless. Not always positive definite. We will not distinguish \log_2 vs \log/\ln .
- In statistical mechanics this is solved by canonical conjugate pairs whose Jacobian is unity or if states are discretized (quantum statistics): no such concept in statistics

Entropy for Continuous Distribution

- Joint entropy of X and Y

$$H(X,Y) = - \int p(x,y) \log p(x,y) dx dy = E\{-\log p(X,Y)\}$$

- Conditional entropy of X given y

$$H(X|y) = - \int p(x|y) \log p(x|y) dx = E\{-\log p(X|Y) \mid Y = y\}$$

- Conditional of X given Y

$$\begin{aligned} H(X|Y) &: \int p(y) H(X|y) dy = - \int p(y) \int p(x|y) \log p(x|y) dx dy \\ &: - \int \int p(x,y) \log p(x|y) dx dy = E\{E\{-\log p(X|Y) \mid Y\}\} \end{aligned}$$

Maximum Entropy

- For a bounded interval $a < x < b$ find p with maximum entropy given the normalization constraint: use Lagrange multiplier method

$$H(p) \triangleq - \int_a^b p(x) \lg p(x) dx \quad \begin{array}{l} p(x) \geq 0 \\ \int_a^b p(x) dx = 1. \end{array}$$

$$J(p) \triangleq - \int_a^b p(x) \ln p(x) dx + \lambda_0 \left(\int_a^b p(x) dx - 1 \right)$$

$$\frac{\partial}{\partial p(x)} J(p) = -\ln p(x) - 1 + \lambda_0. \quad p(x) = e^{\lambda_0 - 1}. \quad \lambda_0 = 1 - \ln(b - a)$$

$$p(x) = \begin{cases} \frac{1}{b-a}, & a \leq x \leq b \\ 0, & \text{otherwise.} \end{cases}$$

Uniform distribution maximizes entropy

Maximum Entropy for Semi-unbounded Distributions

- If we are given mean on a semi-unbounded range from 0 to infinity, $p(x) = 0$ for $x < 0$

$$\int_{-\infty}^{\infty} x p(x) dx = \mu < \infty$$

$$J(p) \triangleq - \int_0^{\infty} p(x) \ln p(x) dx + \lambda_0 \left(\int_0^{\infty} p(x) dx - 1 \right) + \lambda_1 \left(\int_0^{\infty} x p(x) dx - \mu \right)$$

$$p(x) = \begin{cases} \frac{1}{\mu} e^{-x/\mu}, & x \geq 0 \\ 0, & \text{otherwise.} \end{cases}$$

Note the “Boltzmann” factor $e^{-\beta x}$

Maximum Entropy for Unbounded Distributions

- If we are given mean μ and variance s on an unbounded range from $-\infty$ to $+\infty$

$$J(p) \triangleq - \int_{-\infty}^{\infty} p(x) \ln p(x) dx + \lambda_0 \left(\int_{-\infty}^{\infty} p(x) dx - 1 \right) \\ + \lambda_1 \left(\int_{-\infty}^{\infty} x p(x) dx - \mu \right) + \lambda_2 \left(\int_{-\infty}^{\infty} x^2 p(x) dx - \sigma^2 \right)$$

- $p(x) = e^{(\lambda_0 - 1) + \lambda_1 x + \lambda_2 x^2} = \frac{1}{\sqrt{2\pi\sigma^2}} e^{-\frac{x^2}{2\sigma^2}} \quad \text{if } \mu=0$
- Can be further generalized: if we have constraints on first $n = 2k$ cumulants we obtain exponential of n -th order polynomial
- These “Boltzmann factors” have direct analogy with statistical mechanics

Kullback-Leibler (KL) divergence

- KL divergence is a relative entropy between two distributions (discrete or continuous)

$$D_{\text{KL}}(P\|Q) = \sum_i P(i) \log \frac{P(i)}{Q(i)}. \quad D_{\text{KL}}(P\|Q) = \int_{-\infty}^{\infty} p(x) \log \frac{p(x)}{q(x)} dx.$$

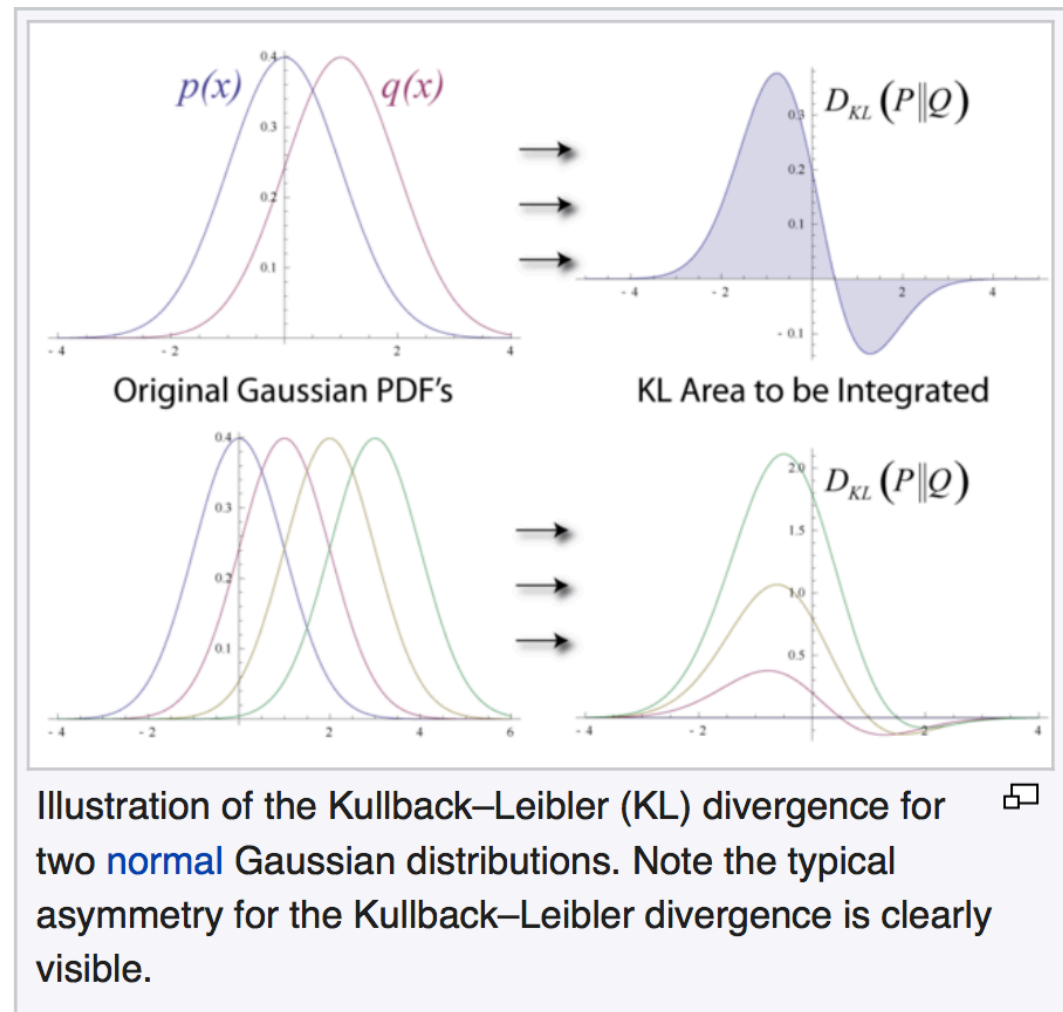
- Satisfies Gibbs inequality $\text{KL} \geq 0$: proof using Jensen inequality for convex functions (see e.g. MacKay 2.7) or:

$$\begin{aligned} \ln x \leq x - 1 \quad & - \sum_{i \in I} p_i \ln \frac{q_i}{p_i} \geq - \sum_{i \in I} p_i \left(\frac{q_i}{p_i} - 1 \right) \\ & = - \sum_{i \in I} q_i + \sum_{i \in I} p_i \end{aligned}$$

- This is 0 since probabilities are normalized
- It is not a distance: $\text{KL}(p, q)$ is not $\text{KL}(q, p)$

KL Divergence for Gaussians

- Always positive
- Increases as the two distributions differ from each other
- Only zero when the two distributions are equal
- Good way to probe how similar are two distributions: starting point for Variational Inference/Variational Bayes methods



Exercise: KL Divergence for Gaussians

- Assume $p = \text{gauss}(\mu_1, \sigma_1)$ and $q = \text{gauss}(\mu_2, \sigma_2)$
- Evaluate $\text{KL}(p||q)$ and show $\text{KL} > 0$
- Evaluate $\text{KL}(q||p)$ and show it differs from $\text{KL}(p||q)$

Solution: KL Divergence for Gaussians

$$KL(p||q) = \frac{1}{\sqrt{2\pi\sigma_1^2}} \int e^{\frac{-(x-\mu_1)^2}{2\sigma_1^2}} \left[\frac{(x-\mu_2)^2}{2\sigma_2^2} - \frac{(x-\mu_1)^2}{2\sigma_1^2} + \frac{1}{2} \ln 2\pi(\sigma_2^2 - \sigma_1^2) \right] dx$$

Let $u \equiv x - \mu_1$, $\Delta\mu = \mu_1 - \mu_2$

$$\begin{aligned} KL(p||q) &= \frac{1}{\sqrt{2\pi\sigma_1^2}} \int e^{\frac{-u^2}{2\sigma_1^2}} \left[\frac{(u - \Delta\mu)^2}{2\sigma_2^2} - \frac{u^2}{2\sigma_1^2} + \frac{1}{2} \ln 2\pi(\sigma_2^2 - \sigma_1^2) \right] dx \\ &= \frac{\sigma_1^2}{2\sigma_2^2} + \frac{\Delta\mu^2}{2\sigma_2^2} - \frac{1}{2} + \ln\sigma_2 - \ln\sigma_1 \end{aligned}$$

Minimized for $\Delta\mu = 0$

$$\frac{\sigma_1}{\sigma_2} = \alpha, \quad \frac{\partial}{\partial\alpha} \left(\frac{\alpha^2}{2} - \ln\alpha - \frac{1}{2} \right) = 0, \quad \alpha - \frac{1}{\alpha} = 0 \rightarrow \alpha = 1$$

Minimized for $\sigma_1 = \sigma_2$, $KL = 0$

$$KL(q||p) = \frac{\sigma_2^2}{2\sigma_1^2} + \frac{\Delta\mu^2}{2\sigma_1^2} - \frac{1}{2} + \ln\sigma_1 - \ln\sigma_2$$

KL Divergence and Negentropy

- **Negentropy**: KL divergence, i.e. relative entropy, against a Gaussian with equal variance

$$J(y) = H(y_G) - H(y) \geq 0$$

- Measures deviation of a distribution from gaussian. Can be approximated as

$$J(y) \approx \frac{1}{12}E\{y^3\}^2 + \frac{1}{48}kurt(y)^2 \quad kurt(y)=E(y^4)-3E(y^2)^2$$

- But other approximations may work better:

$$J(y) = [E\{G(y)\} - E\{G(g)\}]^2$$

$$G_1(y) = \frac{1}{a} \log \cosh (a y), \quad G_2(y) = -\exp(-y^2/2)$$

Mutual Information

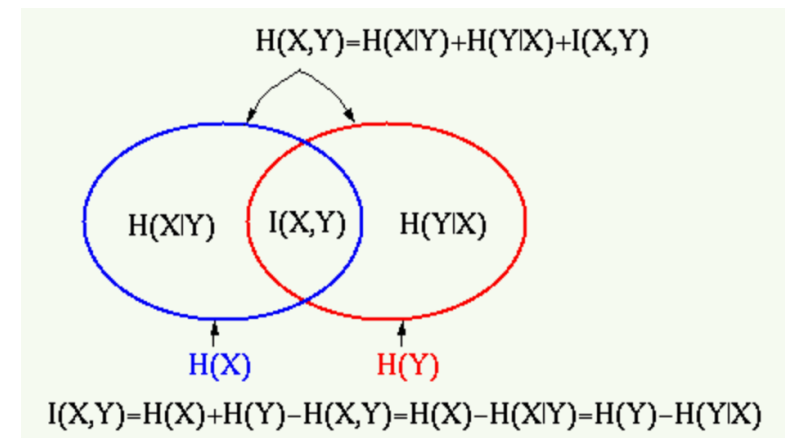
- Defined as amount of information shared between X and Y

$$\begin{aligned}
 I(X,Y) &: H(X) + H(Y) - H(X,Y) \\
 &: E\{-\log p(X)\} + E\{-\log p(Y)\} + E\{-\log p(X,Y)\} \\
 &: E\left\{\log \frac{p(X,Y)}{p(X)p(Y)}\right\}
 \end{aligned}$$

$$\begin{aligned}
 I(X,Y) &: E\left\{\log \frac{p(X,Y)}{p(X)p(Y)}\right\} \\
 &: E\left\{\log \frac{p(X|Y)}{p(X)}\right\} = H(X) - H(X|Y) \\
 &: E\left\{\log \frac{p(Y|X)}{p(Y)}\right\} = H(Y) - H(Y|X)
 \end{aligned}$$

$$I(X;Y) = D_{\text{KL}}(P(X,Y) \| P(X)P(Y))$$

- Minimizing $I(X,Y)$ is a good way to define independence: $I(X,Y)=0$ if $H(X|Y)=H(X)$ or $H(Y|X)=H(Y)$ and is positive (KL divergence)



Multi-Information

- Generalization of mutual $I(\mathbf{y}) = \int P(\mathbf{y}) \log_2 \frac{P(\mathbf{y})}{\prod_i P(y_i)} d\mathbf{y}$
- Information $I(X, Y)$ to several variables y (or s)
- Multi-information is 0 if y statistically independent

$$\begin{aligned} I(\hat{\mathbf{s}}) &= \sum_i H[(\mathbf{V}\mathbf{x}_w)_i] - H[\mathbf{V}\mathbf{x}_w] \\ &= \sum_i H[(\mathbf{V}\mathbf{x}_w)_i] - (H[\mathbf{x}_w] + \log_2 |\mathbf{V}|) \end{aligned}$$

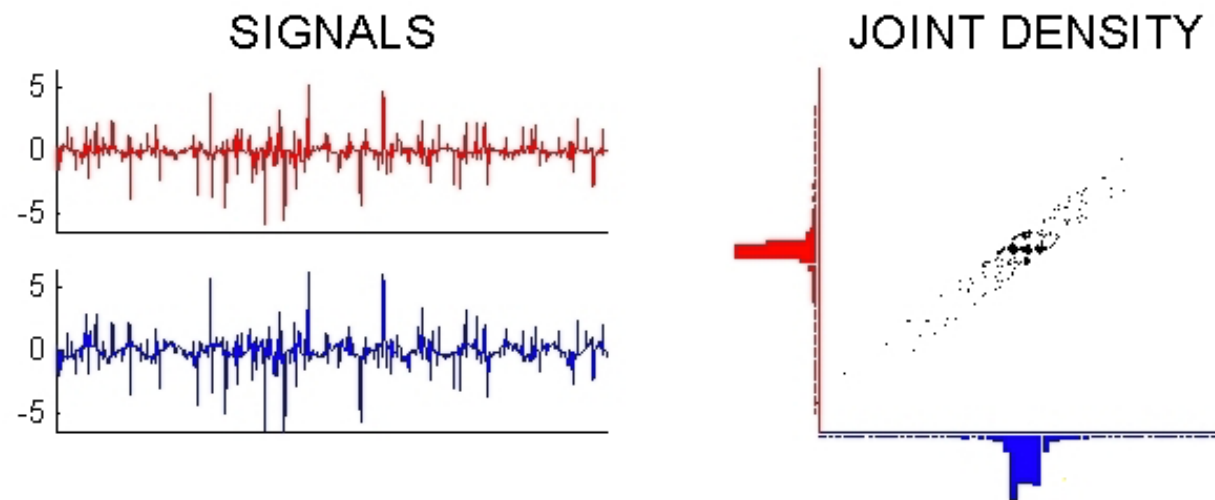
- ICA: we want to know \mathbf{V} that minimizes $I(s)$, $s = \mathbf{V}\mathbf{x}_w$

$$\mathbf{V} = \arg \min_{\mathbf{V}} \sum_i H[(\mathbf{V}\mathbf{x}_w)_i]$$

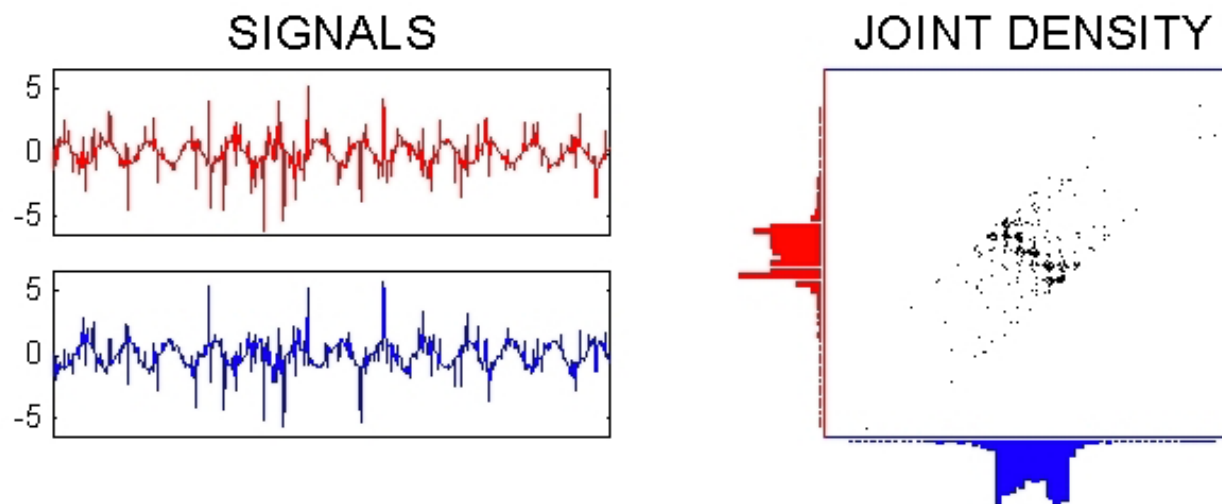
- This is equivalent to maximizing negentropy
 $J = \sum_i [H(s_{gi}) - H(s_i)]$ where $s = \mathbf{V}\mathbf{x}_w$
- We do not know $P(s)$ so we need to use some approximation to evaluate relative entropy

Fast ICA

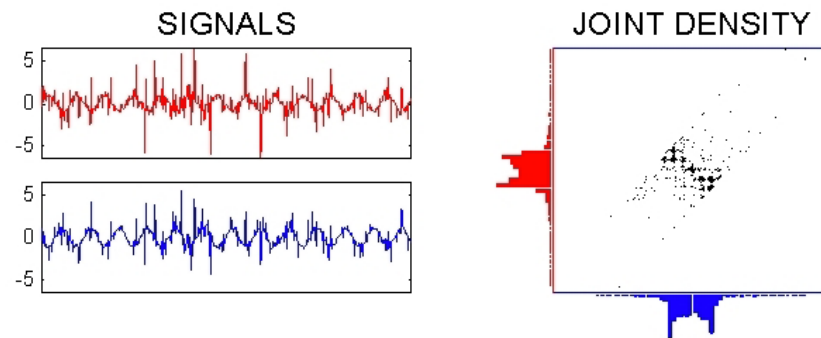
- Approximation for negentropy $J = [E(G(s)) - E(G(g))]^2$
- Maximize $\sum_i E(G(s_i)) = \sum_i E(G(V_i^t x_w)) = E(G(V^t x_w))$
- Subject to normalization for V : Lagrange multiplier β
 $O(V) = E(G(V^t x_w)) - \beta(V^T V - I)$
- This is optimization problem
- We will discuss how to solve optimization problems next, but typically this requires iterations, hence more complicated than linear algebra
- For large dimensions iterative methods are faster than linear algebra and even linear algebra problems are solved iteratively



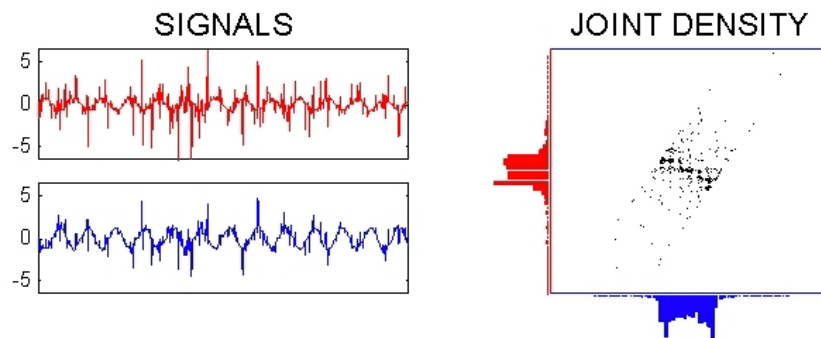
Input signals and density



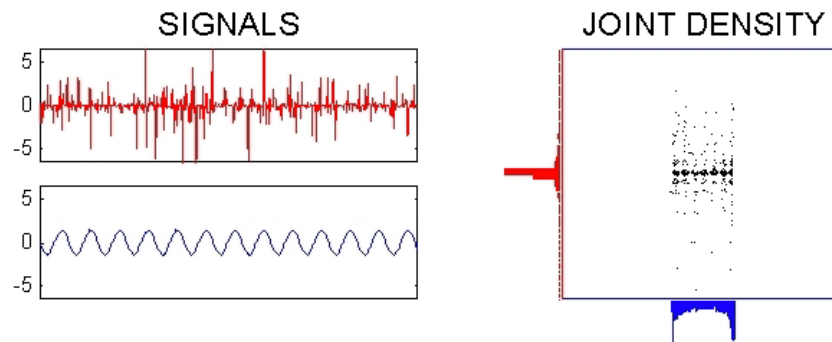
Whitened signals and density



Separated signals after 1 step of FastICA



Separated signals after 2 steps of FastICA



Separated signals after 4 steps of FastICA

Subsequent iterations rotate V until it decorrelates the two signals

Fisher Information Matrix (Metric)

- Quantify the power of future experiments
- Instead of actual data likelihood we can replace it with its ensemble average
- Suppose we have N measurements x_i

$$L = \prod_i p(x_i)$$

$$\ln L = \sum_i \ln p(x_i)$$

$$\langle \ln L \rangle = \left\langle \sum_i \ln p(x_i) \right\rangle = N \langle \ln p(X) \rangle$$

$$\langle \ln p(X) \rangle \text{ (or } E\{\ln p(X)\}) = \int dx \cdot p(X) \ln p(X) = -H(X)$$

Ensemble Averaging:

Precision matrix becomes Fisher matrix.

- we can Taylor expand around a fiducial model in terms of parameters Θ we wish to measure

$$\ln L(\vec{\theta}_{\text{fid}} + \Delta\vec{\theta}) = \ln L(\vec{\theta}_{\text{fid}}) + \sum_i \left. \frac{\partial \ln L}{\partial \theta_i} \right|_{\vec{\theta}_{\text{fid}}} \delta\theta_i + \frac{1}{2} \sum_{ij} \left. \frac{\partial^2 \ln L}{\partial \theta_i \partial \theta_j} \right|_{\vec{\theta}_{\text{fid}}} \delta\theta_i \delta\theta_j$$

MLE: $\left\langle \frac{\partial \ln L}{\partial \theta_i} \right\rangle = E\left(\frac{\partial \ln p}{\partial \theta_i}\right) = 0$ Maximized at fiducial model:
 $\theta_i = \theta_{i,\text{fid}}$

Fisher Matrix:

$$F_{ij} = -\left\langle \frac{\partial^2 \ln L}{\partial \theta_i \partial \theta_j} \right\rangle_{\theta_{\text{fid}}} = \frac{\partial^2 H}{\partial \theta_i \partial \theta_j}$$

$$F_{ij} = \int dx \frac{\partial \ln p(x, \vec{\theta})}{\partial \theta_i} \frac{\partial \ln p(x, \vec{\theta})}{\partial \theta_j} p(x, \vec{\theta})$$

$$- \int p(x, \vec{\theta}) dx = \int e^{-\ln p} dx = 1$$

$$\int e^{-\ln p} \frac{\partial \ln p}{\partial \theta_i} dx = 0$$

$$\int \left[e^{-\ln p} \frac{\partial^2 \ln p}{\partial \theta_i \partial \theta_j} - e^{-\ln p} \frac{\partial \ln p}{\partial \theta_i} \frac{\partial \ln p}{\partial \theta_j} \right] dx = 0$$

$$E\left(\frac{\partial^2 \ln p}{\partial \theta_i \partial \theta_j}\right) = \int p \frac{\partial^2 \ln p}{\partial \theta_i \partial \theta_j} dx = E\left(\frac{\partial \ln p}{\partial \theta_i} \frac{\partial \ln p}{\partial \theta_j}\right) = \int p \frac{\partial \ln p}{\partial \theta_i} \frac{\partial \ln p}{\partial \theta_j} dx$$

Back to Linear Least Squares

$$-\ln L = \sum_i \frac{(y_i - y(x_i|\vec{\theta}))^2}{2\sigma_i^2}$$

$$\left\langle \frac{\partial^2 \ln L}{\partial \theta_i \partial \theta_j} \right\rangle = \left\langle \sum_k \frac{\frac{\partial y(x_k)}{\partial \theta_i} \frac{\partial y(x_k)}{\partial \theta_j}}{2\sigma_i^2} \right\rangle = F_{ij}$$

Posterior: $p(\vec{\theta}) \propto e^{-\delta\theta_i F_{ij} \delta\theta_j / 2}$

Covariance Matrix: $\langle \theta_i \theta_j \rangle = \langle \theta_i \rangle \langle \theta_j \rangle = F_{ij}^{-1}$

Experiment Design

- When we design an experiment we may be able to choose several parameters: sampling of points x_i where we measure data y_i , noise level σ_i , number of data points x_i etc.
- At a given x_i information on parameter Θ_j is given by $(dy_i/d\Theta_j)^2/\sigma_i^2$: this suggests choosing x_i where this is maximized. Note that this can be computed at the fiducial model without actually taking any data
- If we have several parameters we need to break their degeneracies: this is not possible if we only observe at a single x_i : we need to compute full Fisher matrix and invert it to obtain the final error estimate
- By varying the design of the experiment we can predict what the expected error on any given parameter will be: this enables us to design experiment to reach the goals we wish to achieve

Literature

- D. Mackay, *Information Theory, Inference, and Learning Algorithms* (See course website), Chapter 2
- M. Kardar, *Statistical Physics of Particles*, Chapter 2
- ICA: J. Shlens, *A Tutorial on Independent Component Analysis*, <https://arxiv.org/pdf/1404.2986.pdf>