

13 Lecture 13 (23/03/2020)

13.1 Vibrations and oscillations part 2

Lecture 13 doesn't cover a whole lot of new material since the lectures covered chapter 22 earlier on (Lecture 05). It does have a little new material in the form of a problem involving rotations but these lecture notes will mainly just be example problems.

13.2 Ideal springs vs non-ideal springs

Before the example there will be a short subsection clearing up confusion between springs with a zero-length and ideal springs.

An ideal spring is a spring where for any elongation it will exert a spring force $F_v = ks$. This is illustrated in the figures below. Ideal springs are good for math since they're easy to work with but

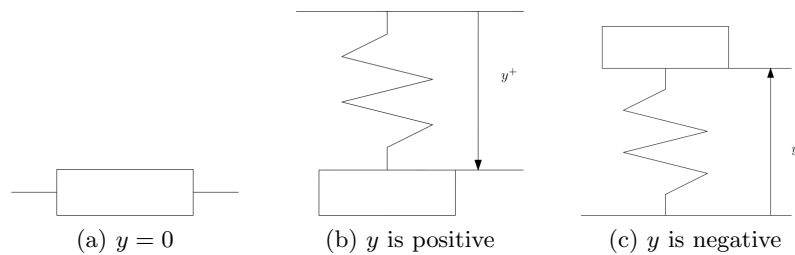


Figure 1: Force from an ideal spring

not so good for realistic results, since ideal springs don't actually exist. In general a real spring will have a zero-length, which is the length in which the spring exerts exactly no force. The principal of a spring with a zero-length is shown in the figure below.

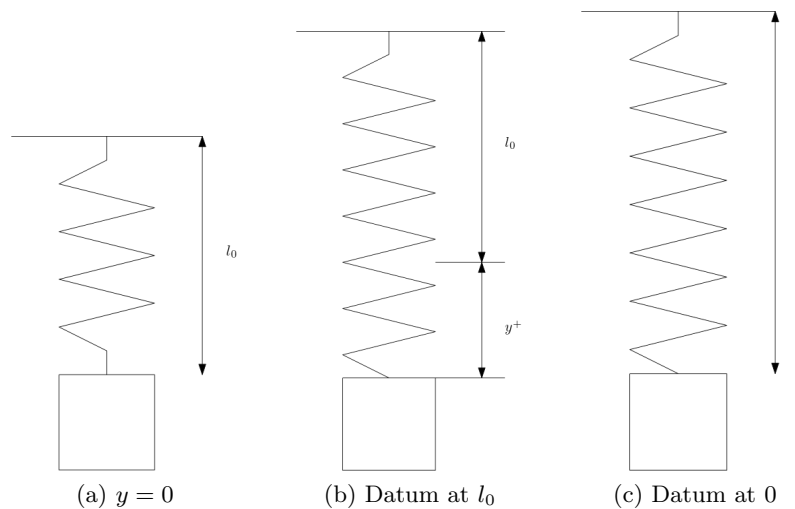


Figure 2: Force from a non-ideal spring

13.3 Spring with damper

Templates for images and figures

$$\Sigma F_x = F_{ext} - F_v - F_d = m\ddot{x} \quad (1)$$

$$F_v = kx \quad (2)$$

$$F_d = c\dot{x} \quad (3)$$

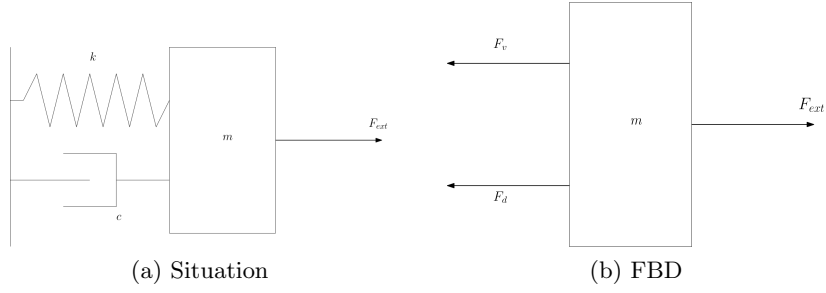


Figure 3: Oscillating spring with damper term

Substitute equations (2) and (3) into (1):

$$\begin{aligned} F_{ext} - kx - c\dot{x} &= m\ddot{x} \\ m\ddot{x} + c\dot{x} + kx &= F_{ext} \end{aligned} \quad (4)$$

Where equation (4) is a non-homogeneous 2nd order linear ODE¹.

13.4 Rotating oscillation problem with Torsion rod

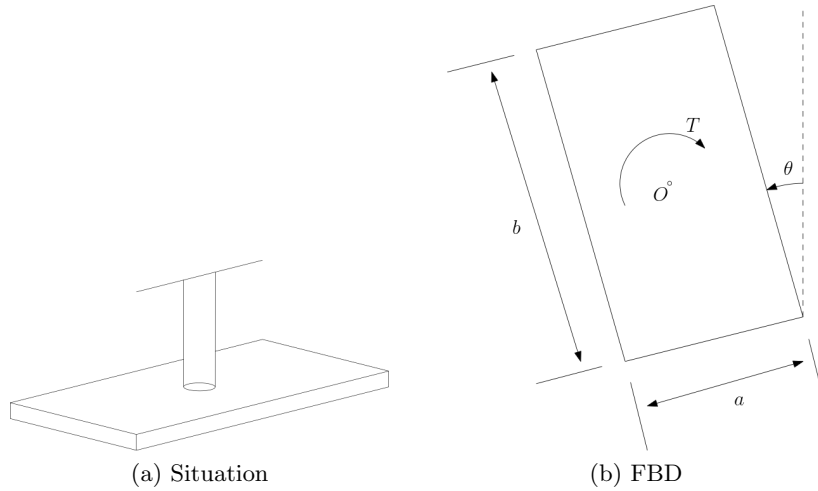


Figure 4: Oscillating plate on a torsion rod

$$\Sigma M_O = -T = I_O \ddot{\theta} \quad (1)$$

$$T = k\theta \quad (2)$$

$$I_O = \frac{1}{12}m(a^2 + b^2) \quad (3)$$

Substitute equation (2) and (3) into (1):

$$\begin{aligned} -k\theta &= \frac{1}{12}m(a^2 + b^2)\ddot{\theta} \\ \frac{1}{12}m(a^2 + b^2)\ddot{\theta} + k\theta &= 0 \end{aligned} \quad (4)$$

Where (4) is the homogeneous 2nd order linear ODE.

¹Ordinary differential equation

13.5 Complex spring systems example

There will be situation where computations will be done for spring systems where the oscillation will not be one single spring vibrating, but any random combination of springs in a series or parallel. The following example will cover such a system with 2 springs placed in series with each other. Since

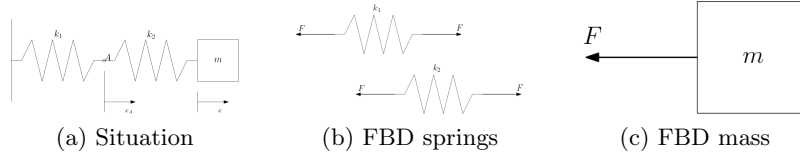


Figure 5: Spring system with springs in series rather than a single oscillating spring

Newton's 3rd law states Action = -Reaction:

$$\begin{aligned}
 F &= k_1 x_A \Rightarrow x_A = \frac{F}{k_1} \\
 F &= k_2(x - x_A) \Rightarrow F = k_2 x - \frac{k_2}{k_1} F \\
 F &= \frac{k_1 k_2}{k_1 + k_2} x
 \end{aligned} \tag{1}$$

Then when considering the FBD of the vibrating mass:

$$\begin{aligned}
 \Sigma F_x &= -F = m\ddot{x} \\
 -\frac{k_1 k_2}{k_1 + k_2} x &= m\ddot{x} \\
 m\ddot{x} + \frac{k_1 k_2}{k_1 + k_2} x &= 0
 \end{aligned} \tag{2}$$

Where $m\ddot{x} + \frac{k_1 k_2}{k_1 + k_2} x = 0$ would be the homogeneous Linear 2nd order ODE that needs to be solved for the answer.

13.6 Example with non-ideal spring and gravity

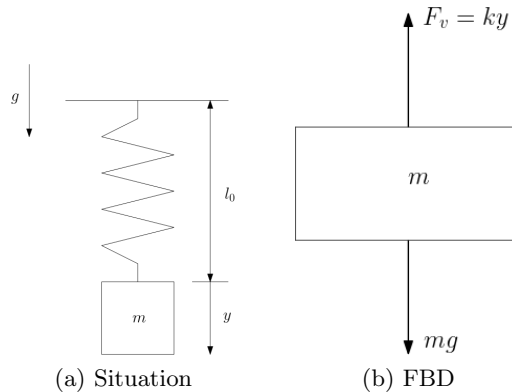


Figure 6: Oscillating non-ideal spring with gravity term

$$\Sigma F_y = mg - ky = m\ddot{y} \quad (1)$$

Equilibrium when $\ddot{y} = 0$

$$mg - ky = 0$$

$$y = y_e = \frac{mg}{k} \quad (2)$$

where y_e is the equilibrium position

$$m\ddot{y} + ky = mg \quad (3)$$

Equation (3) is a non-homogenous 2nd order linear ODE. The solution to this type of differential equation will take the form of $y = y_p + y_c$ where y_p is the particular solution and y_c is the characteristic equation. The particular solution is found by "geussing" the solution takes the same for as the right hand side of the equation. In this case this means that the solution is of the form $y_p = C$, where C is some constant. This gives the following:

$$\begin{cases} y_p = C \\ y'_p = 0 \\ y''_p = 0 \end{cases} \Rightarrow m \cdot 0 + kC = mg \quad (4)$$

$$C = \frac{mg}{k} = y_p \quad (5)$$

For the characteristic equation we guess the solution is of the form $y_c = e^{\lambda y}$. This gives the following:

$$\begin{cases} y = e^{\lambda y} \\ y' = \lambda e^{\lambda y} \\ y'' = \lambda^2 e^{\lambda y} \end{cases} \Rightarrow m\lambda^2 e^{\lambda y} + k\lambda e^{\lambda y} = 0 \quad (6)$$

$e^{\lambda y} \neq 0$, Thus:

$$m\lambda^2 + k\lambda = 0$$

$$\lambda(m\lambda + k) = 0$$

$$\lambda_1 = 0 \wedge \lambda_2 = -\frac{k}{m} \quad (7)$$

$$(8)$$

The total characteristic solution then becomes:

$$y_c = C_1 e^{\lambda_1 y} + C_2 e^{\lambda_2 y} \quad (9)$$

$$= C_1 + C_2 e^{-\frac{k}{m} y} \quad (10)$$

Where C_1 and C_2 are random constants which can be determined by solving for the initial conditions of the problem. The total solution for ht edifferential equation then becomes:

$$y = y_p + y_c = \frac{mg}{k} + C_1 + C_2 e^{-\frac{k}{m} y} \quad (11)$$